

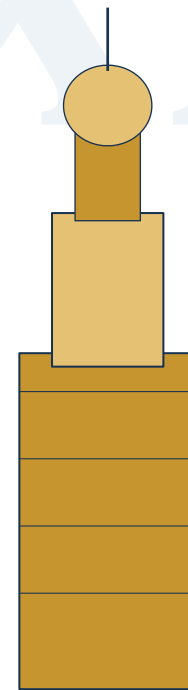
RSEF BIENNIAL

XL

ATLAS · EFT in top-quark production & decay

A Finite-Basis Framework for Fully-Differential EFT Measurements

Orthogonal basis expansions of the decay rate, and a simultaneous fit of EFT Wilson coefficients with the ATLAS detector.



Torre del Oro · Sevilla

Research project: CNS2022-135718

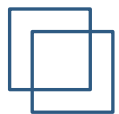


Mariam Chitishvili

ATLAS Collaboration

The XL RSEF Physics Biennial
Sevilla · 20–24 July 2026

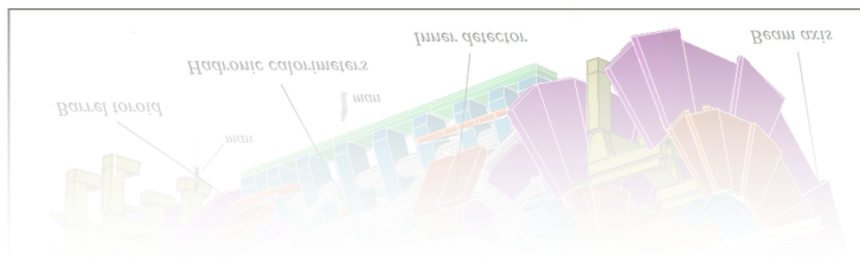
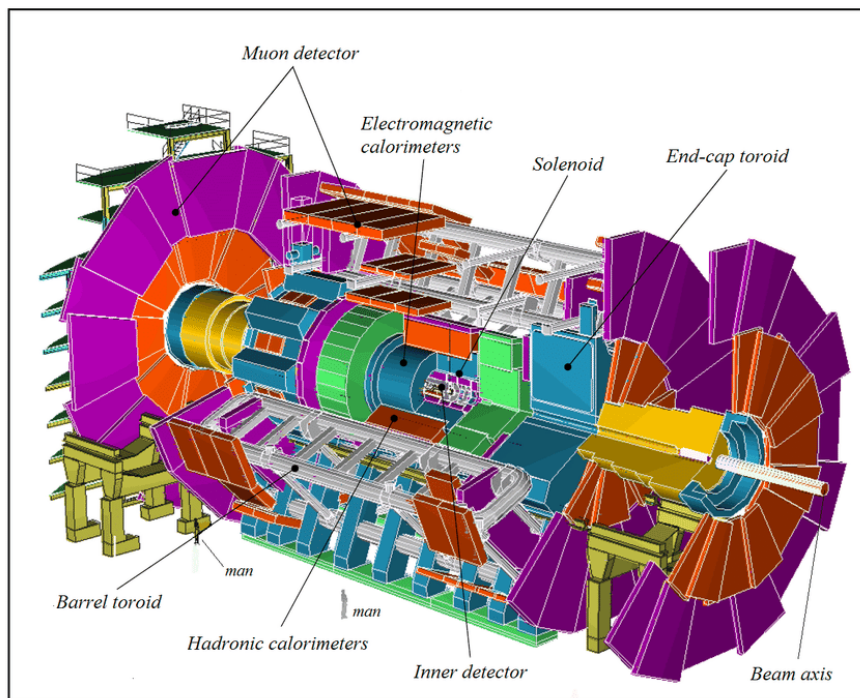




The ATLAS Experiment at the LHC



The largest general-purpose LHC detector — and a top-quark factory.



46 m

long · 25 m tall · ~7000 tonnes of instrumentation

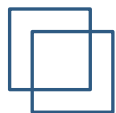
140 fb^{-1}

pp collisions at $\sqrt{s} = 13 \text{ TeV}$ (Run 2, 2015–2018)

$\sim 10^8$

top-quark pairs produced — precision statistics

Millions of read-out channels reconstruct every charged lepton, jet and missing-energy vector — the raw material for precision top-quark physics.



New Physics with the SMEFT



If heavy new physics lives at a scale $\Lambda \gg v_{EW}$, it imprints on the SM as higher-dimensional operators.

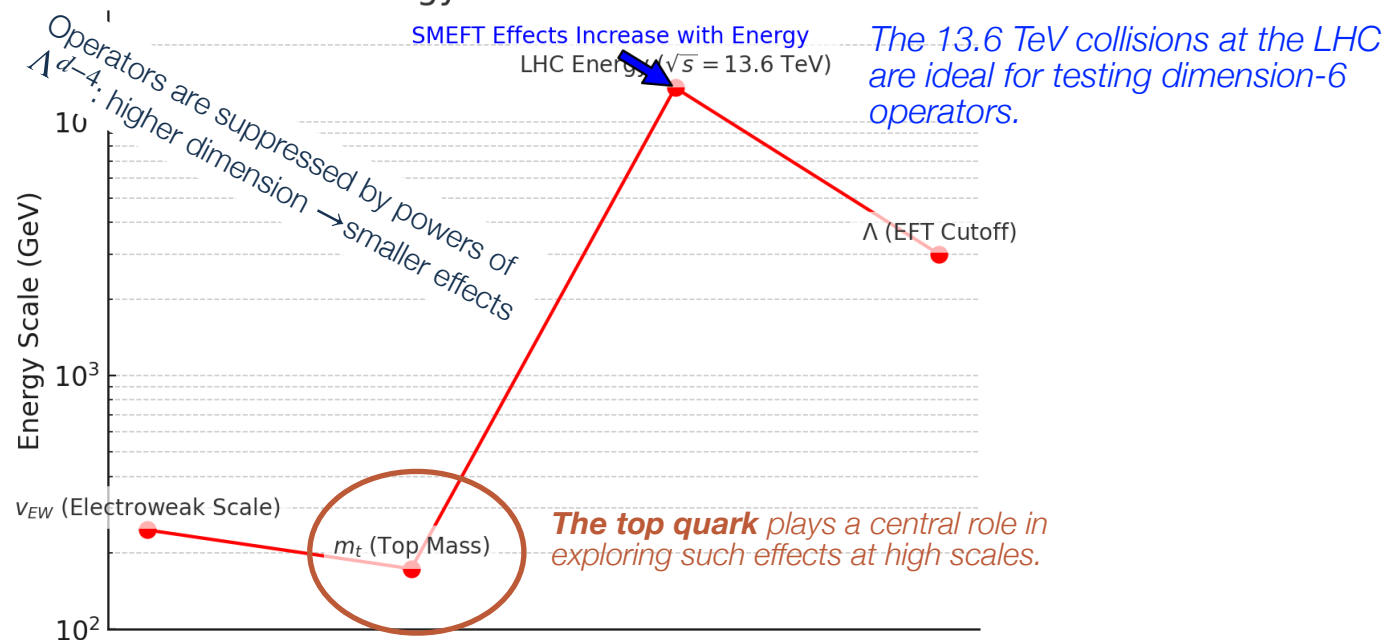
$$\mathcal{L} = \mathcal{L}_{SM} + \sum_i \frac{C_i}{\Lambda^2} \mathcal{O}_i^{(6)} + \sum_j \frac{\tilde{C}_j}{\Lambda^4} \mathcal{O}_j^{(8)} + \dots$$

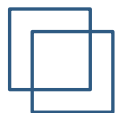
$\mathcal{O}_i^{(d)}$ — gauge-invariant operators

C_i — Wilson coefficients

The energy handle. Dimension-6 effects scale as $\delta\sigma \propto (C_i/\Lambda^2) E^2$: high-energy top-quark interactions amplify new physics. Tied to the electroweak scale, the top quark is a prime window on these operators.

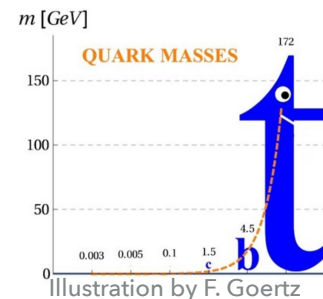
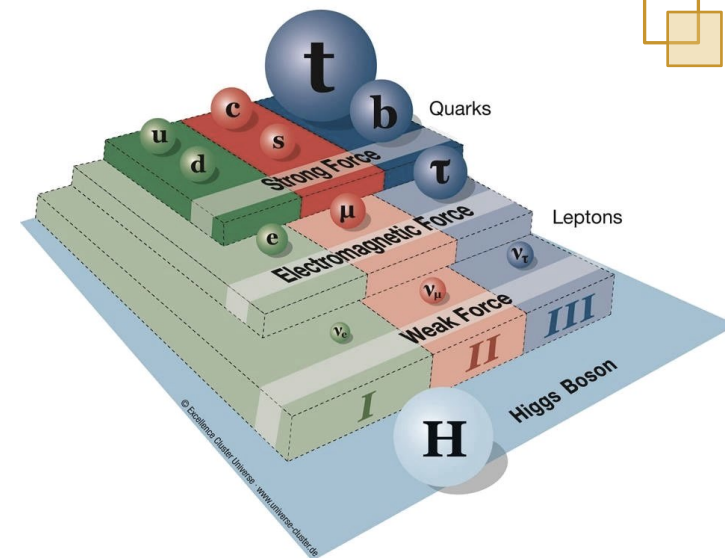
SM vs. SMEFT Energy Scales and Motivation for LHC





The Top Quark: A Unique Probe

The heaviest elementary particle — the only quark that decays before it hadronizes.



1

Heaviest of them all

$m_t \approx 172.5 \text{ GeV}$ — Its Yukawa coupling $y_t \approx 1$ ties it directly to the Higgs and electroweak symmetry breaking.

2

It decays as a free quark

$\tau_t \approx 5 \times 10^{-25} \text{ s}$, shorter than the hadronization time — so its spin is transferred, undiluted, to its decay products.

3

A clean decay chain

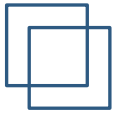
$t \rightarrow Wb$ with $|V_{tb}| \approx 0.999$ (>99%), then $W \rightarrow \ell\nu$. Leptons make an especially clean, fully-reconstructable final state.

4

Copiously produced at the LHC

$\sigma(t\bar{t}) \approx 830 \text{ pb}$ and single-top (t-channel dominant) at 13 TeV — abundant events for differential studies.

Why it matters: the top's spin polarization survives into the **angles**. Those angular distributions are precision observables — and exactly where new physics leaves its fingerprint.

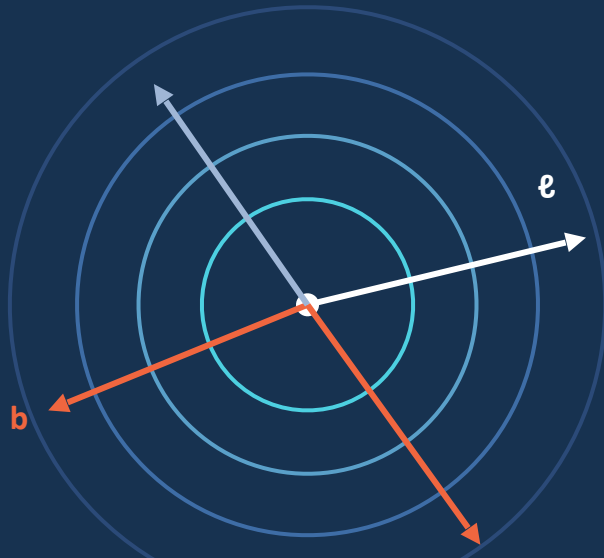


Angular Measurements: The Sharpest Tool



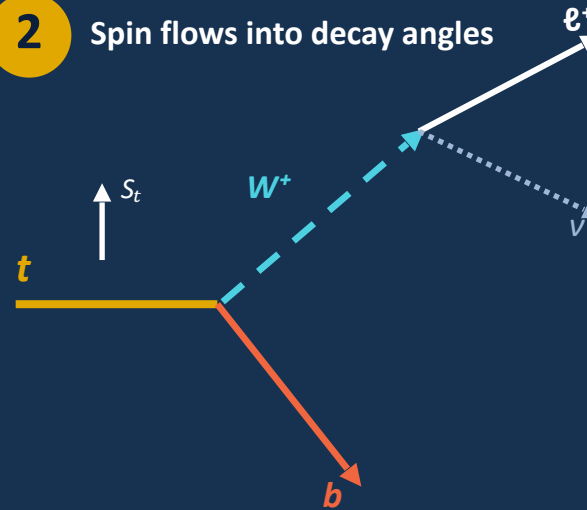
Top spin survives, so decay-product angles are precision observables — where SMEFT shows up most cleanly.

1 ATLAS reconstructs the products



Charged leptons, b-jets and missing energy are measured layer-by-layer, giving each decay product a precise momentum vector.

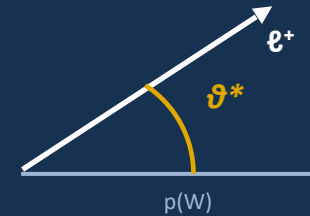
2 Spin flows into decay angles



$$t \rightarrow W^+ b, \quad W^+ \rightarrow \ell^+ \nu$$

The top decays before it can hadronize ($\tau \approx 5 \times 10^{-25}$ s), so its spin is transferred intact to the lepton — whose emission angle carries the polarization.

3 Angular observables test the SM

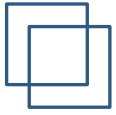


W helicity	SM (NNLO)	ATLAS
F_0 (long.)	0.687	0.709 ± 0.019
F_L (left)	0.311	0.299 ± 0.015
F_R (right)	0.0017	-0.008 ± 0.014

Phys. Rev. D 81 (2010) 111503(R) Eur. Phys. J. C 77 (2017) 264

Spin correlation & W-helicity fractions agree with the Standard Model to percent-level precision.

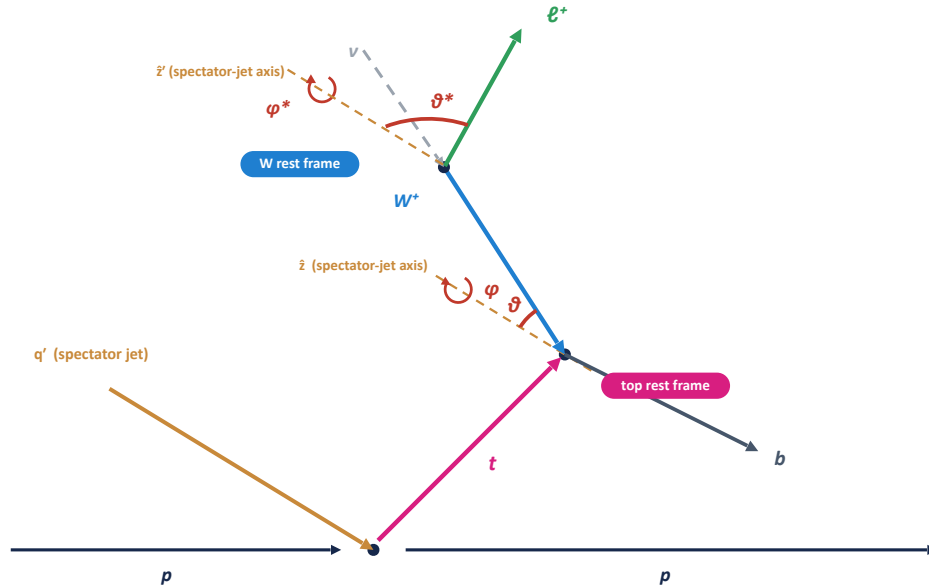
The through-line: ATLAS turns detector hits into decay-product 4-vectors → the **angular distributions** of those products encode top-quark polarization, W-helicity and $t\bar{t}$ spin correlation — precision tests of the SM.



The Challenge: a Fully-Differential Decay Rate



The $t \rightarrow Wb \rightarrow \ell\nu b$ decay is described by four angles — binning a 4-D distribution is hopeless.



Four angles fully specify the decay: (θ, ϕ) orient the W in the top rest frame; (θ^*, ϕ^*) orient the lepton in the W rest frame.

ATLAS · arXiv:2510.23372, Fig. 2

A finite-basis framework

The fully-differential rate of $t \rightarrow Wb \rightarrow \ell\nu b$ is expanded in **orthogonal basis functions**. A finite set of coefficients captures the full multidimensional shape and normalisation — extracted directly from data by projection, then used to fit multiple EFT Wilson coefficients simultaneously. **140 fb⁻¹ · ATLAS · 13 TeV**.

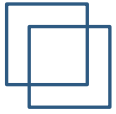
Why not just bin a histogram?

Curse of dimensionality. Even 5 bins per angle $\Rightarrow 5^4 = 625$ — and $5^{10} \approx 10^7$ for the $t\bar{t}$ decay chain, far more than the available events. Projecting to 1-D discards correlations and needs assumptions about the unseen variables.

Most top measurements are projections of this space

W-helicity uses $\cos \theta^*$ alone; **spin correlation** uses azimuthal differences. Each is a 1-D/2-D slice of the same joint density — the M-function method keeps all of it at once.

A ladder runs from 1 angle (W-helicity) $\rightarrow$ 4 angles (this analysis) $\rightarrow$ 10 angles (full $t\bar{t}$). One formalism climbs all of it.



The Idea: Orthogonal-Series Density Estimation



Expand the density in an orthonormal basis; the coefficients are just averages over events.

1

Choose an orthonormal basis

The M-functions are built from Wigner D-functions — a product covering (θ, ϕ) and (θ^*, ϕ^*) . They satisfy $\int M_n M_n'^* d\Omega = \delta_{nn'}$.

2

Expand the density

$\rho(\Omega) = \sum c_n M_n(\Omega)$. Completeness $\Rightarrow$ the series carries all shape information, including correlations between angles.

3

Project to get coefficients

$c_n = \int \rho M_n^* d\Omega = \langle M_n^* \rangle$ — the coefficient is the average of the basis function over the events. No fit, no binning.

4

Errors come for free

The sample covariance $\text{Cov}(c_n, c_m) = [\langle M_n M_m \rangle - \langle M_n^* \rangle \langle M_m^* \rangle] / N$ gives every coefficient a statistical uncertainty and full correlations.

The M-function

$$M^{j_1 j_2}_{m_1 m_2} = (2j_1+1)^{1/2} (2j_2+1)^{1/2} / 4\pi$$

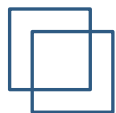
$$\times D^{j_1}_{m_1}(\cos \theta, \varphi) \cdot D^{j_2}_{m_2}(\cos \theta^*, \varphi^*)$$

Wigner D-functions — the natural language of angular momentum and helicity.

Coefficient = average over data

$$c_n = \langle M_n^* \rangle = (\sum w_i M_n^*(\Omega_i)) / (\sum w_i)$$

Unbinned · analytic · every event contributes to every coefficient. A χ^2 distance in coefficient space compares data to model.



Why This Basis? The Physics Lives in 23 Numbers



The basis is chosen from the helicity structure of the decay — so the truncation is physics, not guesswork.

Helicity amplitudes $\Rightarrow$ only four survive

Angular-momentum conservation in $t \rightarrow Wb \rightarrow \ell \nu$ leaves exactly four non-zero reduced amplitudes ($a_{1,1/2}$, $a_{0,1/2}$, $a_{0,-1/2}$, $a_{-1,-1/2}$). This dictates the harmonic content of the decay distribution.

Truncation with a clean cut

At parton level only $j_1 < 1$, $j_2 < 2$ survive $\Rightarrow$ exactly **23 coefficients** completely describe the decay. Detector effects leak power into higher j — which are dropped because the low-order terms carry the EFT physics.

Single top is special: the t-channel produces a **polarised** top quark — giving access to the **imaginary (CP-odd)** parts of the couplings through the ϕ , ϕ^* dependence.

From a distribution to a short vector

$\cos \theta$

ϕ

$\cos \theta^*$

ϕ^*

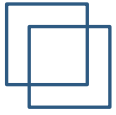
4-D fully-differential decay rate

project onto M-functions

23

M-function coefficients

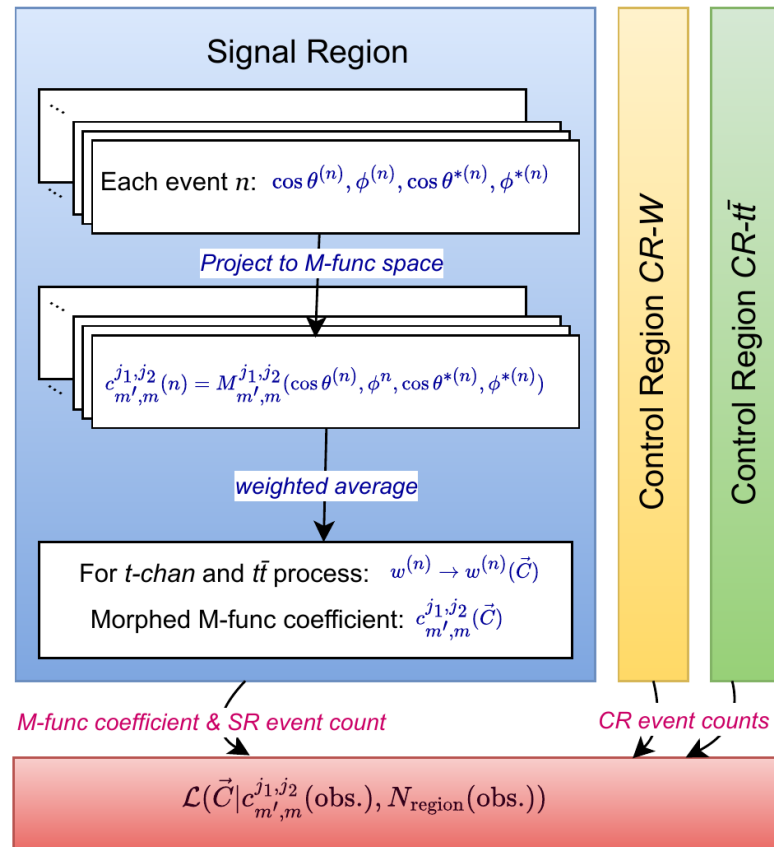
Complete · compact · each with a statistical error and correlations



From Events to Physics: Morphing & the Likelihood

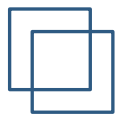


Per event $\rightarrow$ M-functions $\rightarrow$ averaged coefficients $\rightarrow$ a likelihood built directly in coefficient space.



ATLAS · arXiv:2510.23372, Fig. 3

- 1 Reconstruct & project**
Each selected event gives $(\cos \theta, \phi, \cos \theta^*, \phi^*)$; the 23 M-functions are evaluated per event and averaged over the sample (weighted for MC).
- 2 Morph the EFT dependence**
Every coefficient is a polynomial in the Wilson coefficients — a 203-term morphing function (up to 4th order), obtained by reweighting MC. Interference between operators is built in.
- 3 Detector at reco level**
Resolution & acceptance push power into high- j coefficients; restricting to $j_1 < 1, j_2 < 2$ keeps the EFT-sensitive terms and suppresses detector effects.
- 4 One Gaussian likelihood**
Observables = 23×4 coefficients (SR) + event yields (SR + 2 CRs). A χ^2 distance in coefficient space is minimised over the Wilson coefficients.

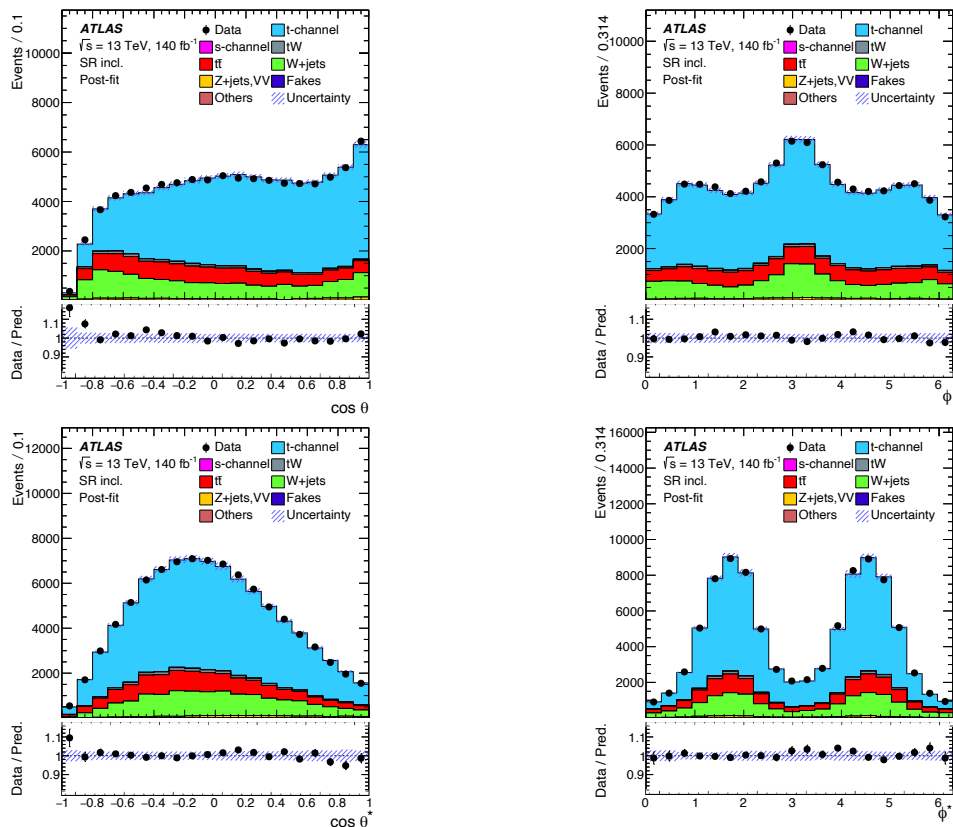


The Observables in ATLAS Data



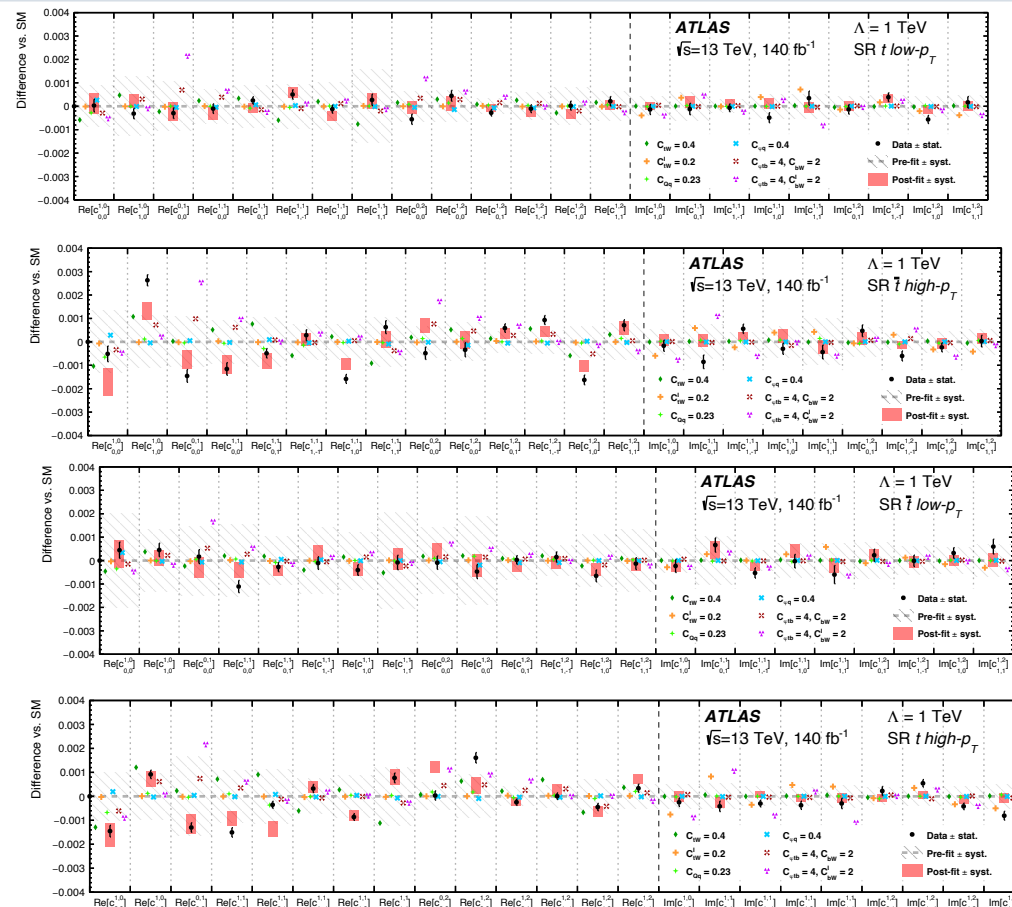
The same information, two views: the four angular shapes — and their compact coefficient encoding.

The four angular distributions



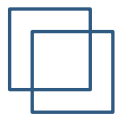
ATLAS · arXiv:2510.23372, Fig. 4

... encoded as M-function coefficients



ATLAS · arXiv:2510.23372, Fig. 6 — deviations of the 23x4 coefficients from the SM

The 4-D distribution on the left is compressed, with no loss, into the coefficient vector on the right — each point an unbinned measurement with its own error, directly sensitive to the Wilson coefficients.



One Formalism, Many Angles: $2 \rightarrow 3 \rightarrow 4 \rightarrow 10$

The same machinery scales from a single helicity angle to the full $t\bar{t}$ production-and-decay chain.



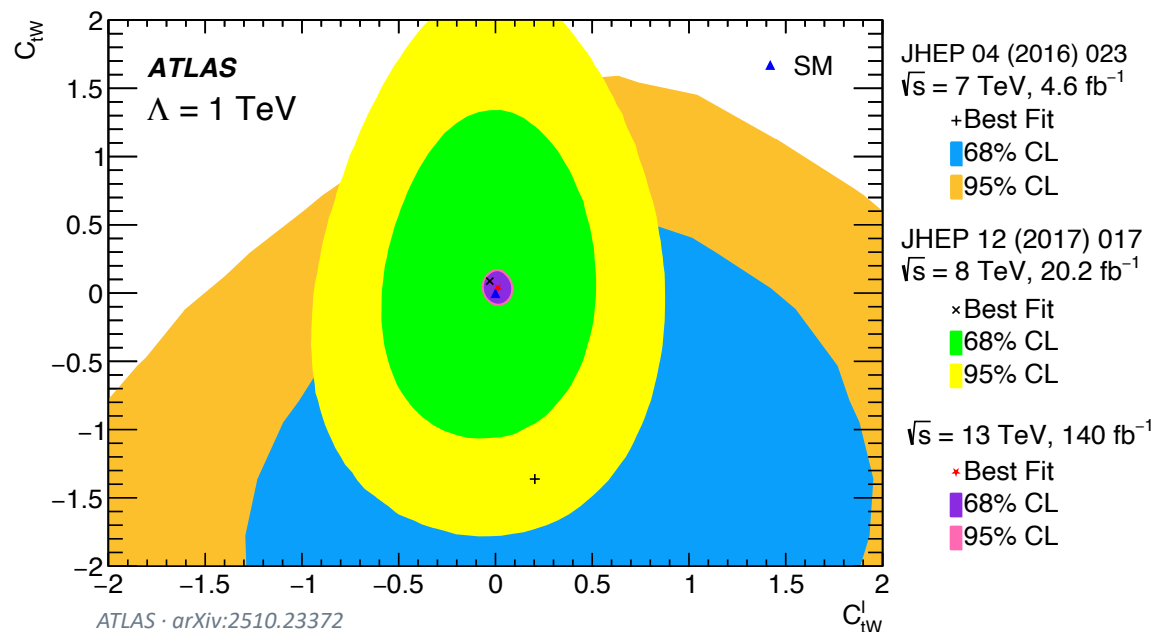
10 **10 angles** — Full $t\bar{t}$: both decays + production angle + rapidity Y . The next frontier.

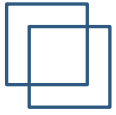
4 **4 angles** — Full $t \rightarrow Wb \rightarrow \ell\nu b$ decay — 23 coefficients. This ATLAS single-top analysis.

3 **3 angles** — Add a production/plane angle — partial spin information and CP-sensitive azimuths.

2 **2 angles** — Top polarisation — first handle on the couplings.

increasing angular dimensionality $\rightarrow$ increasing physics reach

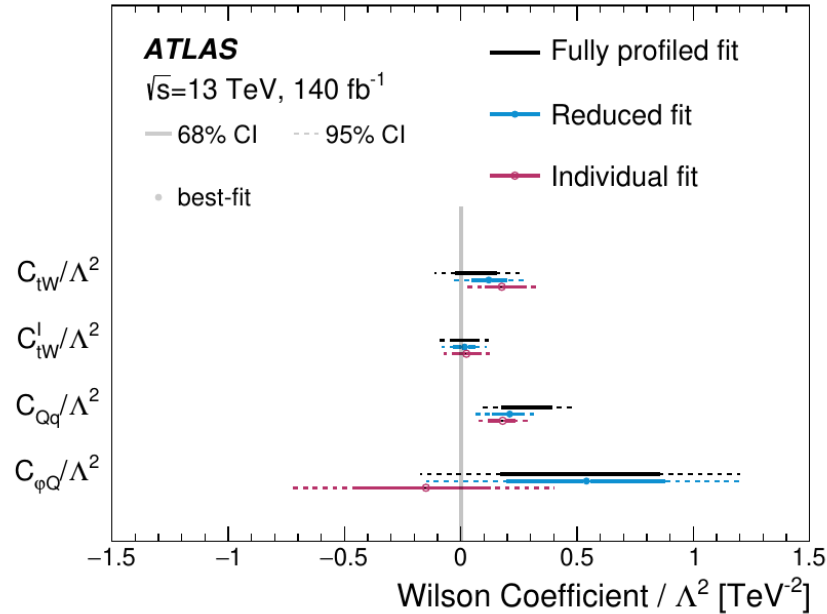




Results — Wilson-Coefficient Constraints



Seven EFT coefficients fitted simultaneously in 140 fb^{-1} — all best-fit points consistent with the Standard Model.



ATLAS · arXiv:2510.23372, Fig. 12(a) — $\text{Re/Im } c_{tW}, C_{Qq}, C_{\phi Q}$ (fine scale)

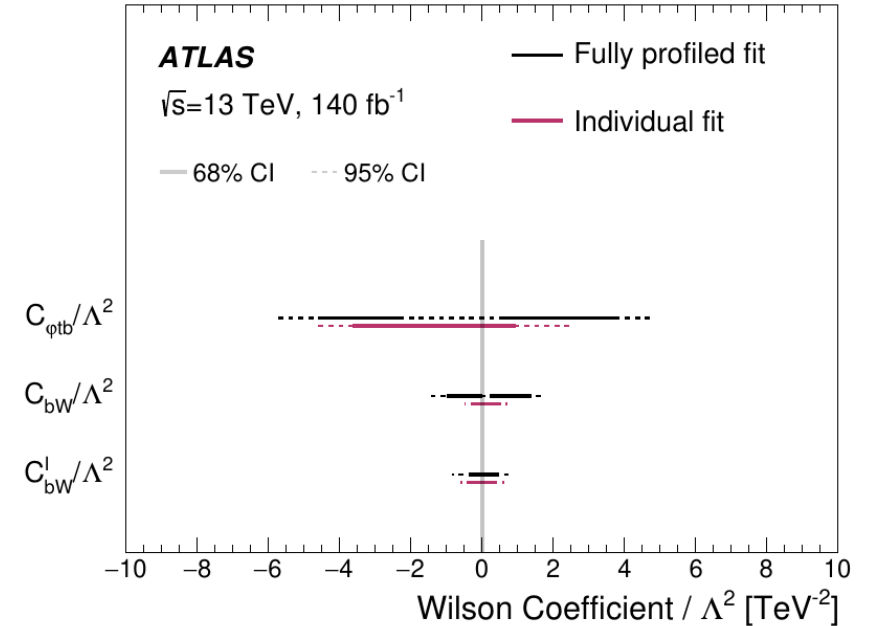


Fig. 12(b) — $c_{\phi tb}, c_{bW}, \text{Im } c_{bW}$ (wide scale; less SM interference)

68% CI (thick) · 95% CI (dotted) · best-fit points · fully-profiled / reduced / individual fit schemes

Tightest on c_{tW}

The most restrictive limits are on Re and Im c_{tW} . This improves the best previous C_{tW} bound (8 TeV W-helicity) by a **factor of 2** — without assuming the other WCs vanish.

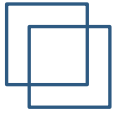
Access to the CP-odd sector

Imaginary parts (C_{tW}^I , and $c_{\phi tb}, c_{bW}$) are constrained by the **imaginary M-function coefficients**; negligible $C_{tW}-C_{tW}^I$ correlation. $c_{bW}, c_{\phi tb}$ also break a direction to which F_R is blind.

Rate + shape together

$C_{\phi Q}$ and C_{Qq} come from event yields in two p_T bins. **C_{Qq} shows a 2.8σ high- p_T deviation** (data vs Powheg) — flagged for NNLO follow-up.

All best-fit points consistent with the SM · seven Wilson coefficients, three fitting schemes · 140 fb^{-1} , $\sqrt{s} = 13 \text{ TeV}$, ATLAS Run 2



The 10-Angle $t\bar{t}$ Analysis: the Full Picture



One likelihood fit to the tenfold-differential production-and-decay rate of $t\bar{t}$ pairs.

Ten kinematic variables

$$\theta_+, \varphi_+$$

Top decay $t \rightarrow W^+ b$
direction of W^+ in the top rest frame

$$\theta_-, \varphi_-$$

Antitop decay $\bar{t} \rightarrow W^- \bar{b}$
direction of W^- in the antitop rest frame

$$\theta_+^*, \varphi_+^*$$

$W^+ \rightarrow \ell^+ \nu$
lepton direction in the W^+ rest frame

$$\theta_-^*, \varphi_-^*$$

$W^- \rightarrow \ell^- \bar{\nu}$
lepton direction in the W^- rest frame

$$\theta, Y$$

Production
top angle in the ZMF · rapidity of the $t\bar{t}$ system

Lepton+jets channel (~44%) — the only channel allowing full angular reconstruction (dilepton hides two neutrinos). b - and c -tagging identify the fermions.

The 10-angle $t\bar{t}$ analysis

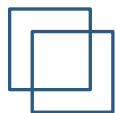
A single likelihood fit to the tenfold differential production + decay rate. The M-function generalises to a product:

$$M = Y(\theta_+, \varphi_+) Y(\theta_-, \varphi_-) Y(\theta_+^*, \varphi_+^*) Y(\theta_-^*, \varphi_-^*) \cdot \tilde{P}_1(\cos \theta) \cdot \tilde{h}_n(Y)$$

spherical harmonics × Legendre (production) × Hermite (rapidity)

Dimensionality explodes: $(N+1)^{10}$ coefficients — 59,409 at $N=2$, $\sim 10^6$ at $N=3$.

Tamed automatically: project onto the SM-tangent subspace (1st + 2nd order), then prune covariance degeneracies → ~27 sensitive discriminants $\tilde{Q}$ carrying spin correlation, both W -helicities, all tWb couplings, even chromo-dipoles.



Taming the Dimensionality



A complete basis is astronomically large — a physics-driven projection keeps only the sensitive directions.

From a complete basis...

$$(N+1)^{10}$$

$\approx 10^6$ (N=3) · 59,409 (N=2)

project onto SM-tangent + 2nd order

35 directions

1st + 2nd-order variations

drop null directions

32 directions

null directions removed

prune covariance degeneracies

27 $\tilde{Q}$

sensitive discriminants

1

Restrict to the SM-tangent subspace

As the parameters vary, the coefficient vector $\vec{c}(a)$ sweeps a submanifold whose dimension is just the number of physics parameters. Keep only the directions tangent to it at the SM point — deviations $\hat{e}^{(i)}$ that form the rows of a transfer matrix H.

2

Add second-order sensitivity

Parameters that don't interfere with the SM have a vanishing linear response. Quadratic and bilinear terms are added so they aren't lost.

3

Drop null directions

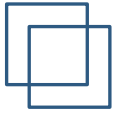
Some $\hat{e}$, $\vec{f}$ vanish identically and carry no information — removed empirically, taking 35 → 32 directions.

4

Prune degeneracies

Diagonalise the covariance $C = A D A^T$ and cut eigenvalues below ~1% of the largest. Symmetries of the cross-section collapse together, leaving 27 independent discriminants $\tilde{Q}$.

~10⁶ possible coefficients → 27 physics-carrying discriminants. A physics-derived analogue of PCA — sensitivity, not variance, sets the axes.



The 10-Angle M-Function: the Mathematics



Same construction as the 4-angle case — a direct product over the ten variables of $t\bar{t}$ production and decay.

Basis function — a product over the ten kinematic variables

$$M_n(\Omega) = Y_{l_1 m_1}(\theta_+, \varphi_+) \cdot Y_{l_2 m_2}(\theta_-, \varphi_-) \cdot Y_{l_3 m_3}(\theta^*_+, \varphi^*_+) \cdot Y_{l_4 m_4}(\theta^*_, \varphi^*_) \cdot \tilde{P}_l(\cos \theta) \cdot \tilde{h}_n(Y)$$

collective index $n = \{ l_1, l_2, l_3, l_4, m_1, m_2, m_3, m_4, l, n \}$ $\tilde{P}_l(x) = \sqrt{(2l+1)/2} P_l(x)$ $h_n(x) = e^{-(x^2/2)} H_n(x) / (\pi^{1/4} \sqrt{2^n n!})$

Orthonormal & complete

$$\int M_n M_n'^* d\tilde{\Omega} = \delta_{nn'}$$

with $d\tilde{\Omega} = d\Omega_+ d\Omega_- d\Omega^*_+ d\Omega^*_- d\theta dY$. Spherical harmonics, Legendre and Hermite are each orthogonal on their own domain — so the product is a complete orthonormal basis.

Density = a series

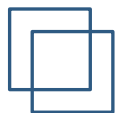
$$\rho(\Omega) = \sum_n c_n M_n(\Omega)$$

Coefficients follow by Fourier projection — which is just a weighted average over events: $c_n = \langle M_n^* \rangle = (\sum_i w_i M_n^*(\Omega_i)) / (\sum_i w_i)$.

Measurement space

$$\chi^2 = (\Delta c)^T (C_d + C_{mc})^{-1} (\Delta c)$$

Covariance $\hat{C}_{nm} = \langle M_n^* M_m^* \rangle - \langle M_n^* \rangle \langle M_m^* \rangle$. The MC coefficients depend on the physics parameters, so $\chi^2(a)$ maps the whole space.



In Preparation: EFT Reach of the 10-Angle $t\bar{t}$ Fit



A simultaneous fit to 11 EFT directions — four-quark production, top decay, and both.

Production — four-quark operators, effective (V–A) basis

C_VV

C_VV'

→ **rate** & $\cos \theta$ production-angle shape

C_AA

C_AA'

→ $t\bar{t}$ **charge (forward–backward) asymmetry** & spin correlation

C_AV

C_VA

C_AV'

→ **top-quark polarisation** (parity-violating)

V / A = vector / axial (light $\times$ top currents) · primes = colour-singlet vs colour-octet. Falkowski & Mahmoudi, JHEP 03 (2011) 125.

Decay — dipole

ctW_Re

ctW_Im

O_{tW} · right-handed tensor coupling

Production + decay — chromo-dipole

ctG_Re

ctG_Im

O_{tG} · chromo-magnetic $\hat{\mu}_t$ & chromo-electric $\hat{d}_t$

Imaginary parts ctW_Im and ctG_Im are CP-odd — a direct handle on CP violation.

11

EFT directions fitted simultaneously

7 production · 2 decay · 2 both

STATUS · analysis in preparation

Common $t\bar{t}$ EFT sample requested for Run 2 (mc20) & Run 3 (mc23)

MadGraph5 + SMEFTsim 3.0 (top, M_W scheme) · Pythia 8 (A14, NNPDF2.3LO) · $\Lambda = 1$ TeV · single + pairwise reweighting grid



Conclusions

Why the M-function angular analysis is so powerful

1

Unbinned & complete

No histograms, no empty bins. Completeness means the coefficients carry every correlation in the full angular density.

2

Compact

A 4-D distribution becomes 23 numbers; the tenfold $t\bar{t}$ rate reduces to ~ 27 discriminants. Maximal information, minimal observables.

3

Errors built in

Each coefficient is a simple event average, so uncertainties and correlations follow analytically from the sample covariance.

4

Physically motivated

The basis follows the helicity structure of the decay, so a clean truncation isolates the EFT-sensitive terms from detector effects.

5

Sensitive to CP

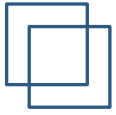
Polarised single tops expose the imaginary parts of the couplings through ϕ , ϕ^* — inaccessible to rate-only measurements.

6

Scalable & unified

One formalism from 1 to 10 angles: shape + rate combined, ready for future differential and EFT interpretations.

Outlook — the M-function coefficients are, in effect, a **tomography of the top-quark spin density matrix**. For the dedicated quantum-tomography measurement, see the companion talk by **Guillem** [slides to be linked].



Results

