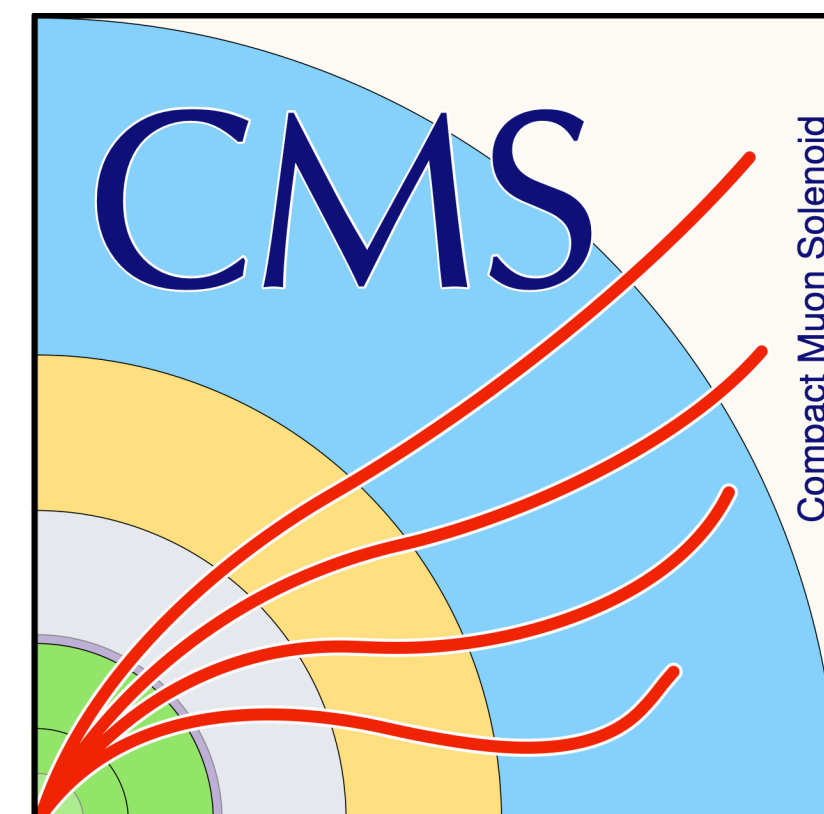




Alexander von
HUMBOLDT
STIFTUNG



Top quark pairs in the threshold at the LHC

(Based on 2412.16685, in collaboration with Garzelli, Moch, Steinhauser, Zenaiev)

Giovanni Limatola
DESY

Workshop on Top Quark Mass Measurement
Valencia
10th November 2025

Introduction



- The LHC has provided astonishing results in testing the Standard Model
- No clear signal of Beyond the Standard Model effects (yet!)
- The upcoming High-Luminosity regime at the LHC is expected to improve both the statistical accuracy and the coverage of phase space
- This poses a significant challenge to the theory community
- Emphasis on high multiplicity processes and/or massive particles production

The Top Quark



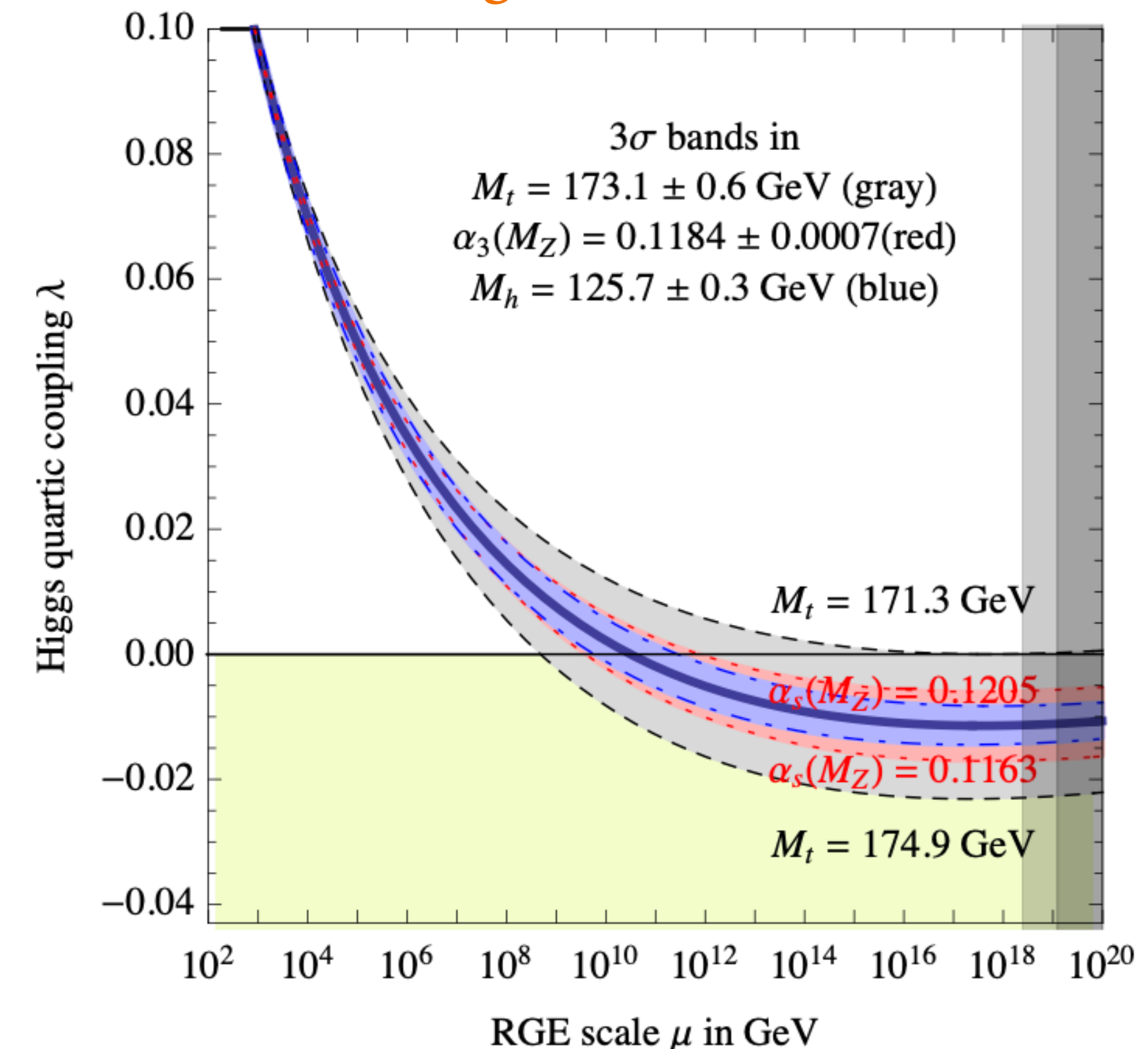
[PDG 2018]

- The heaviest known fundamental particle $m_t = 172.9 \pm 0.4 \text{ GeV}$
- Large Yukawa coupling to the Higgs Boson λ_t
- m_t plays a key role in the stability of the Universe
- $\Gamma_t(t \rightarrow Wb) \simeq \frac{G_F m_t^3}{8\pi\sqrt{2}} \sim 1.6 \text{ GeV} \gg \Lambda_{\text{QCD}}$

Life-time too short for hadron formation

Top quark properties can be studied directly!

Degrassi et al. [1205.6497]

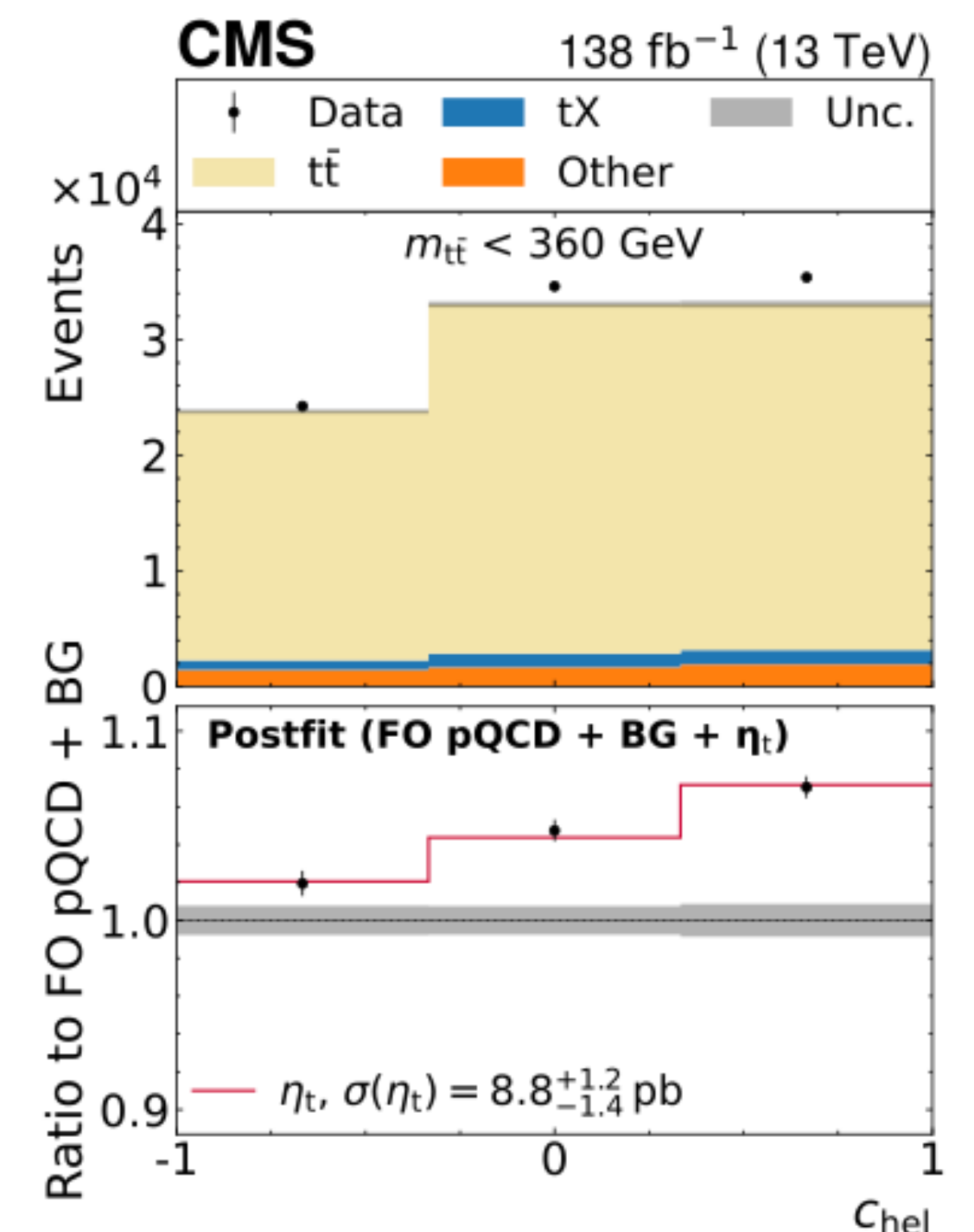




Results from CMS and ATLAS

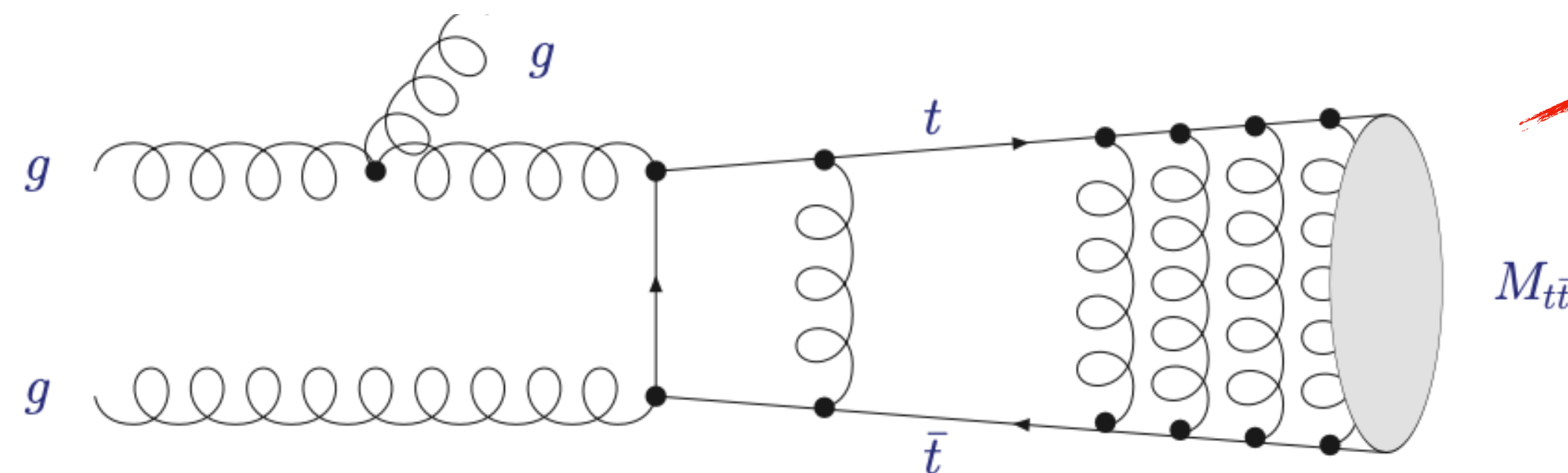
- The spin information of the top quark is entirely transferred to its leptonic decay product
- Spin correlation measurements performed introducing two angular observables $\hat{c}_{\text{hel}}$, $\hat{c}_{\text{han}}$
- Observation of a significant excess of events in the threshold region $M_{t\bar{t}} \simeq 2m_t$ w.r.t. to pQCD predictions
- $\sigma_{\text{CMS}} = 8.8^{+1.2}_{-1.4} \text{ pb}$ [CMS, 2503.22382]
 $\sigma_{\text{ATLAS}} = 9.0 \pm 1.3 \text{ pb}$ [ATLAS-CONF-2025-008]
- BSM pseudoscalar effects?
 Quasi-bound state effects?

Threshold enhanced effects require dedicated studies!



$t\bar{t}$ production at the threshold

- In the threshold region the top quark velocity $\beta_t = \sqrt{1 - 4m_t^2/M_{t\bar{t}}^2} \sim \alpha_s \sim 0.1$ in the $t\bar{t}$ CM frame
- Significant QCD Coulomb corrections scaling as $(\alpha_s/\beta_t)^k$ due to quasi-bound state interactions
- Quasi-bound state effects can be modelled as multiple soft gluon exchanges between top quarks

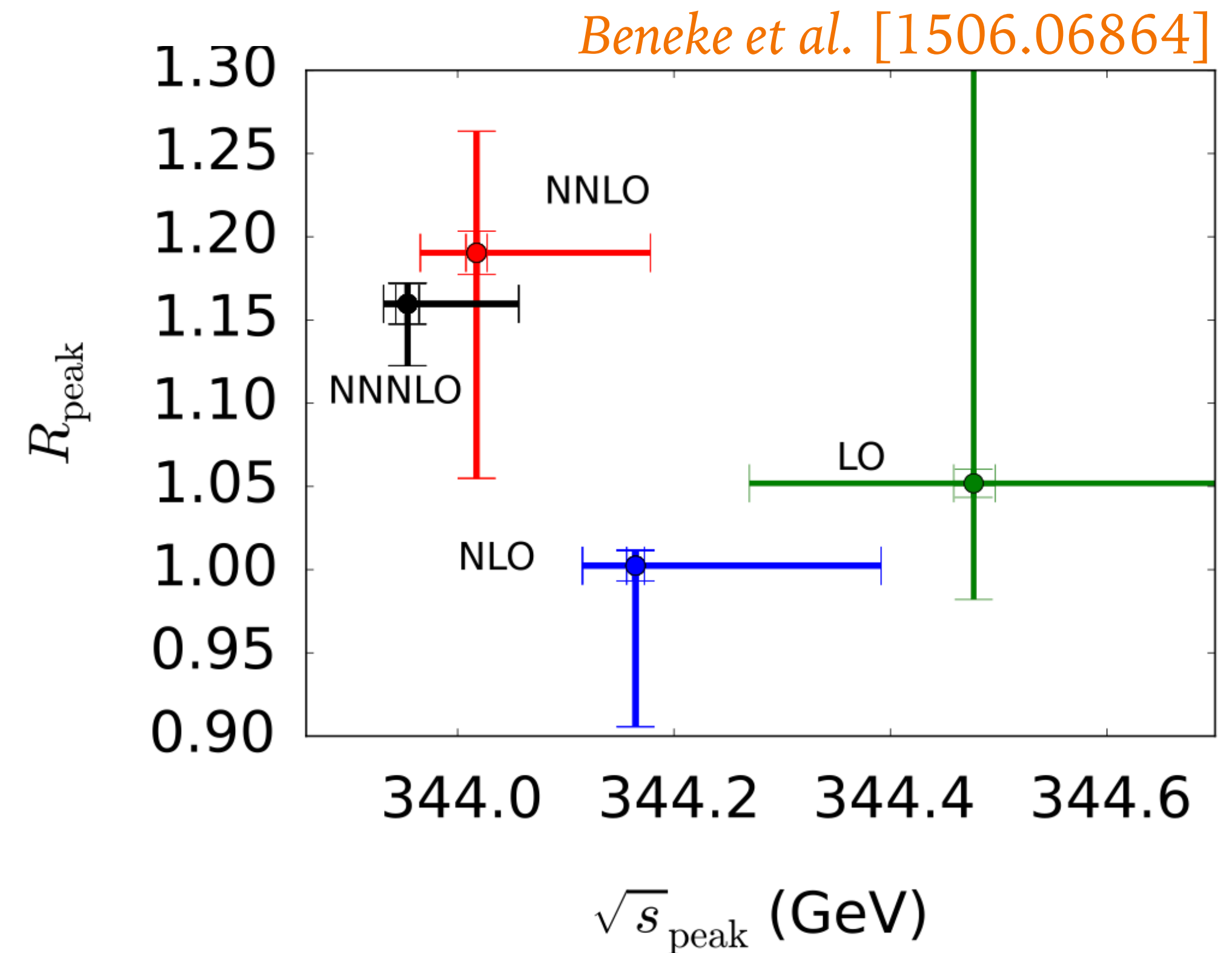
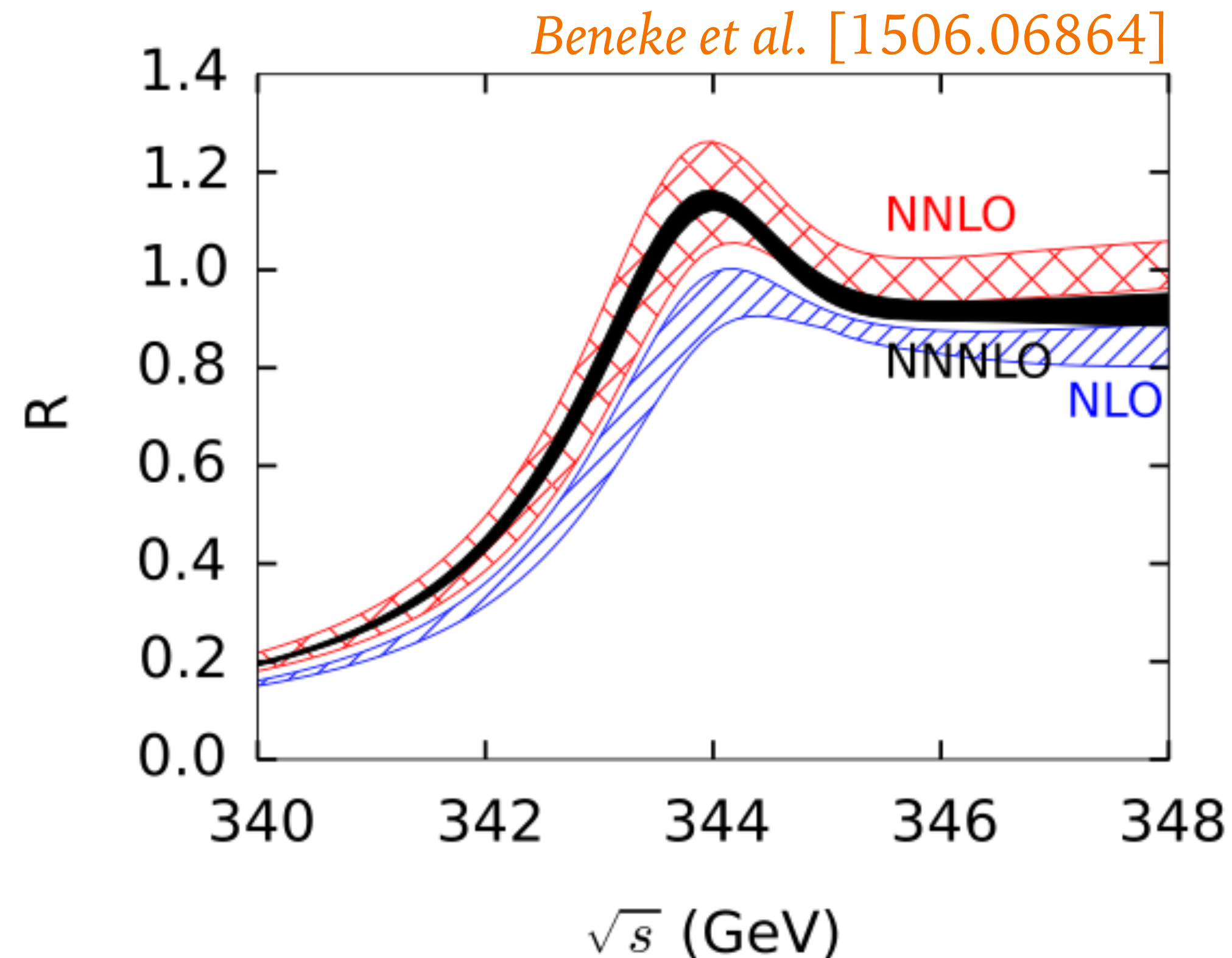


EFT Non-Relativistic QCD
to factorise $(\alpha_s/\beta_t)^k$ at all-orders

$t\bar{t}$ production at the threshold



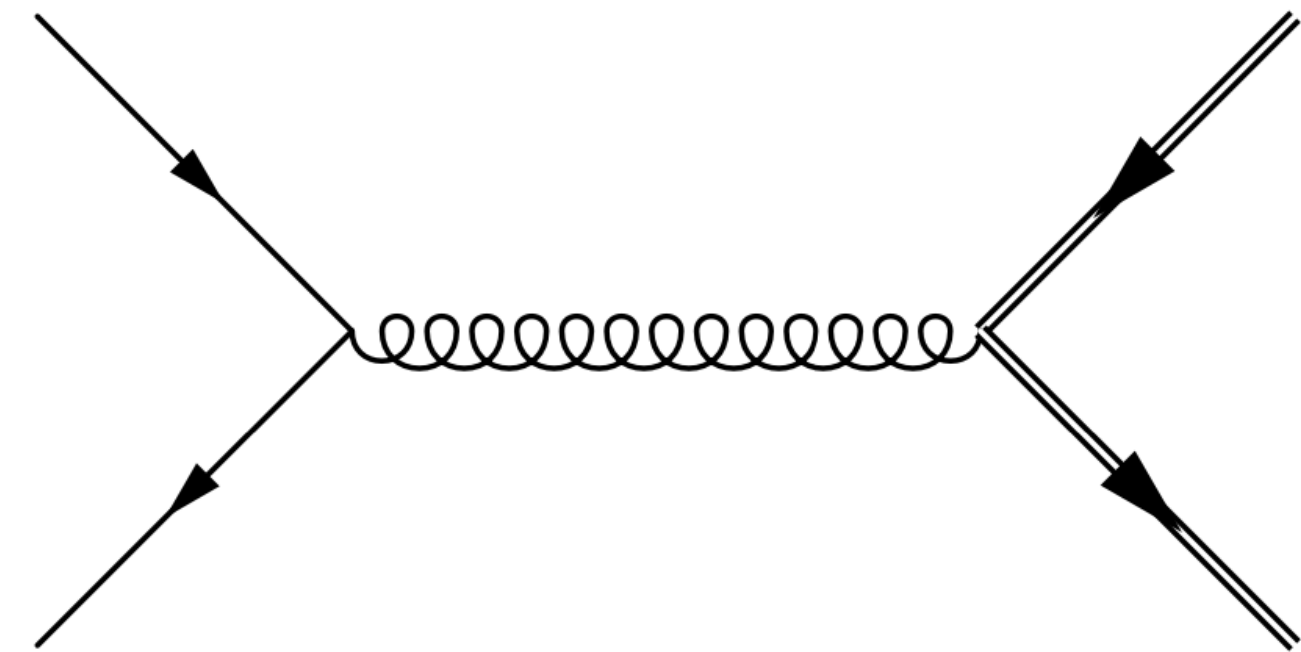
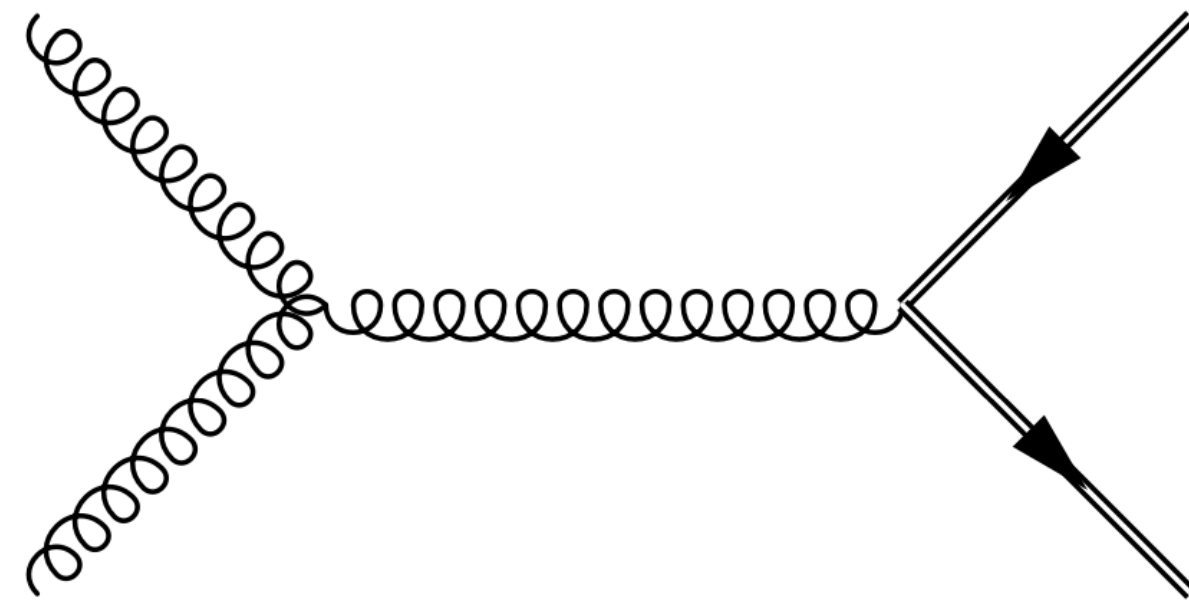
$$R = \frac{\sigma_{e^+e^- \rightarrow t\bar{t}}}{\sigma_{e^+e^- \rightarrow \mu^+\mu^-}} = \beta \sum_k \left(\frac{\alpha_s}{\beta} \right)^k \{1(\text{LO}); \alpha_s, \beta(\text{NLO}); \alpha_s^2, \alpha_s\beta, \beta^2(\text{NNLO}); \dots\}$$



$t\bar{t}$ production at the threshold @ LHC



- $t\bar{t}$ pairs predominantly produced at the threshold
- Threshold scan of the invariant mass distribution $M_{t\bar{t}}$ for top mass extraction
- NRQCD predictions for colour singlet/octet state $t\bar{t}$ formation
- Radiative corrections from initial state region need to be considered
[Hagiwara, Sumino, Yokoya '08], [Kiyo, Kühn, Moch, Steinhauser '08], [Sumino, Yokoya '10]
[Garzelli, GL, Moch, Steinhauser, Zenaiev '24]



$t\bar{t}$ production at the threshold @ LHC



- $t\bar{t}$ in the quasi-bound state in terms of the Spin S , the orbital L , the total J angular momentum and the colour state c of the $t\bar{t}$ pair

$$T = 2S+1 \quad L_J^{[c]}$$

- Zero orbital angular momentum in the threshold region $J = S$
- $S = 0, 1$ states are possible
- $c = 1, 8$ for the colour state



$t\bar{t}$ production at the threshold @ LHC

- Invariant mass distribution for $h_1(P_1) h_2(P_2) \rightarrow T$

$$M_{t\bar{t}} \frac{d\sigma_{P_1 P_2 \rightarrow T}}{dM_{t\bar{t}}}(S, M_{t\bar{t}}^2) = \sum_{ij} \int_{\rho}^1 d\tau \left[\frac{d\mathcal{L}_{ij}}{d\tau} \right] M_{t\bar{t}} \frac{d\hat{\sigma}_{ij \rightarrow T}}{dM_{t\bar{t}}}(\hat{s} = \tau S, M_{t\bar{t}}^2)$$

$$\left[\frac{d\mathcal{L}_{ij}}{d\tau} \right] = \int_0^1 dx_1 \int_0^1 dx_2 f_{i/P_1}(x_1, \mu_f^2) f_{j/P_2}(x_2, \mu_f^2) \delta(\tau - x_1 x_2)$$

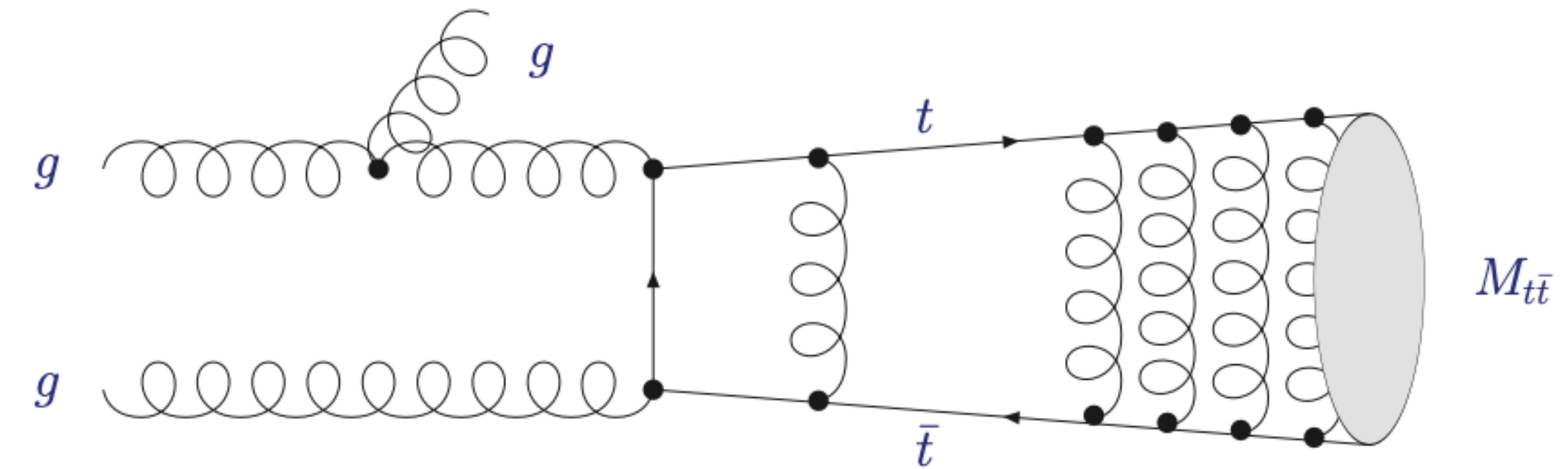
Parton Luminosity

$$\frac{d\hat{\sigma}_{ij \rightarrow T}}{dM_{t\bar{t}}}(\hat{s}, M_{t\bar{t}}^2)$$

Partonic Cross Section

Quasi-bound state in NRQCD

- Multiscale problem
- Hard production at the scale $\sim m_t$
- Quasi-bound state formation at the scale $\sim m_t \beta_t < m_t$
- NRQCD allows the factorisation



$$\frac{d\hat{\sigma}_{ij \rightarrow T}}{dM_{t\bar{t}}}(\hat{s}, M_{t\bar{t}}^2, \mu_f^2) = \underbrace{F_{ij \rightarrow T}(\hat{s}, M_{t\bar{t}}^2, \mu_f^2)}_{\text{Hard Production}} \frac{1}{m_t^2} \underbrace{\text{Im} G^{[1,8]}(M + i\Gamma_t)}_{\text{Quasi-bound state effects}}$$

Separately independent from μ_r



F_{ij} contribution

- Hard production cross section for a state T with invariant mass $M_{t\bar{t}}$
- Evaluated at NLO within pQCD

$$F_{ij \rightarrow T}(\hat{s}, M_{t\bar{t}}^2, \mu_f^2) = \mathcal{N}_{ij} \frac{\pi^2 \alpha_s^2(\mu_r)}{3\hat{s}} \left(1 + \frac{\alpha_s(\mu_r)}{\pi} \mathcal{C}_h \right)$$

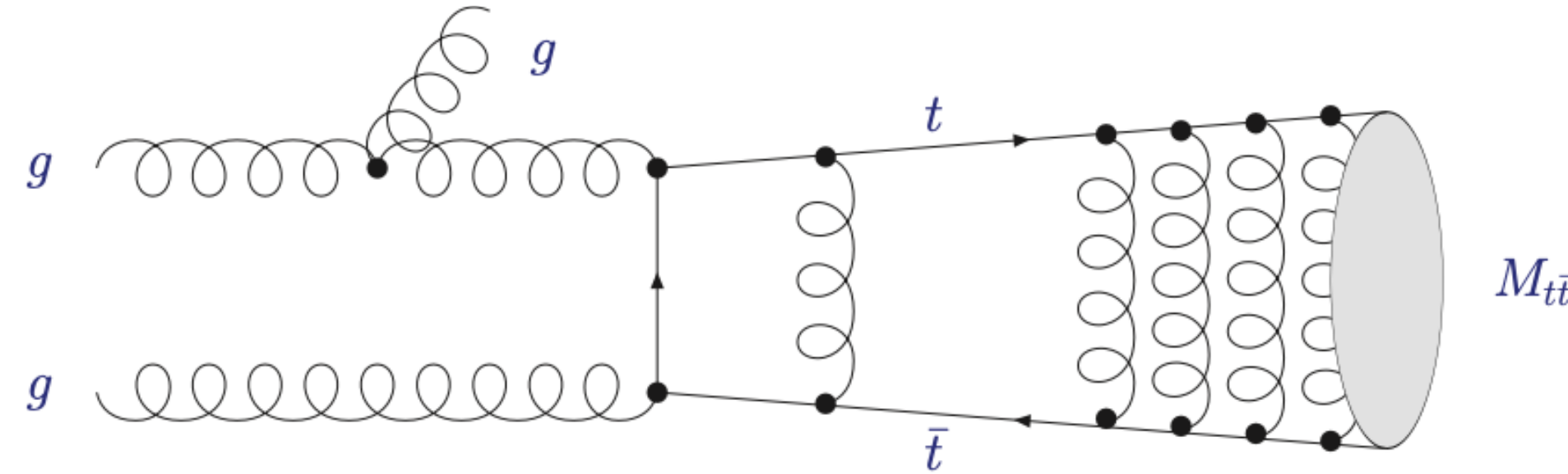
Hard Correction

$$\times \left[\delta_{ij \rightarrow T} \delta(1 - z) + \frac{\alpha_s(\mu_r)}{\pi} (\mathcal{A}_c + \mathcal{A}_{nc}) \right]$$

$$gg \rightarrow {}^1 S_0^{[1,8]}, \quad q\bar{q} \rightarrow {}^3 S_1^{[8]}$$

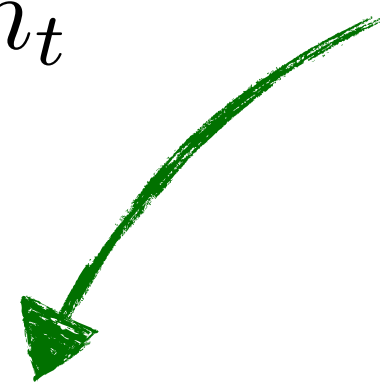
Real Correction from (non) collinear emission

$G(M + i\Gamma_t)$ contribution



- Quasi-bound state due to long-distance effects at the scale $m_t\beta_t$
- Multiple soft gluon exchange modelled with the introduction of a zero-distance Green Function

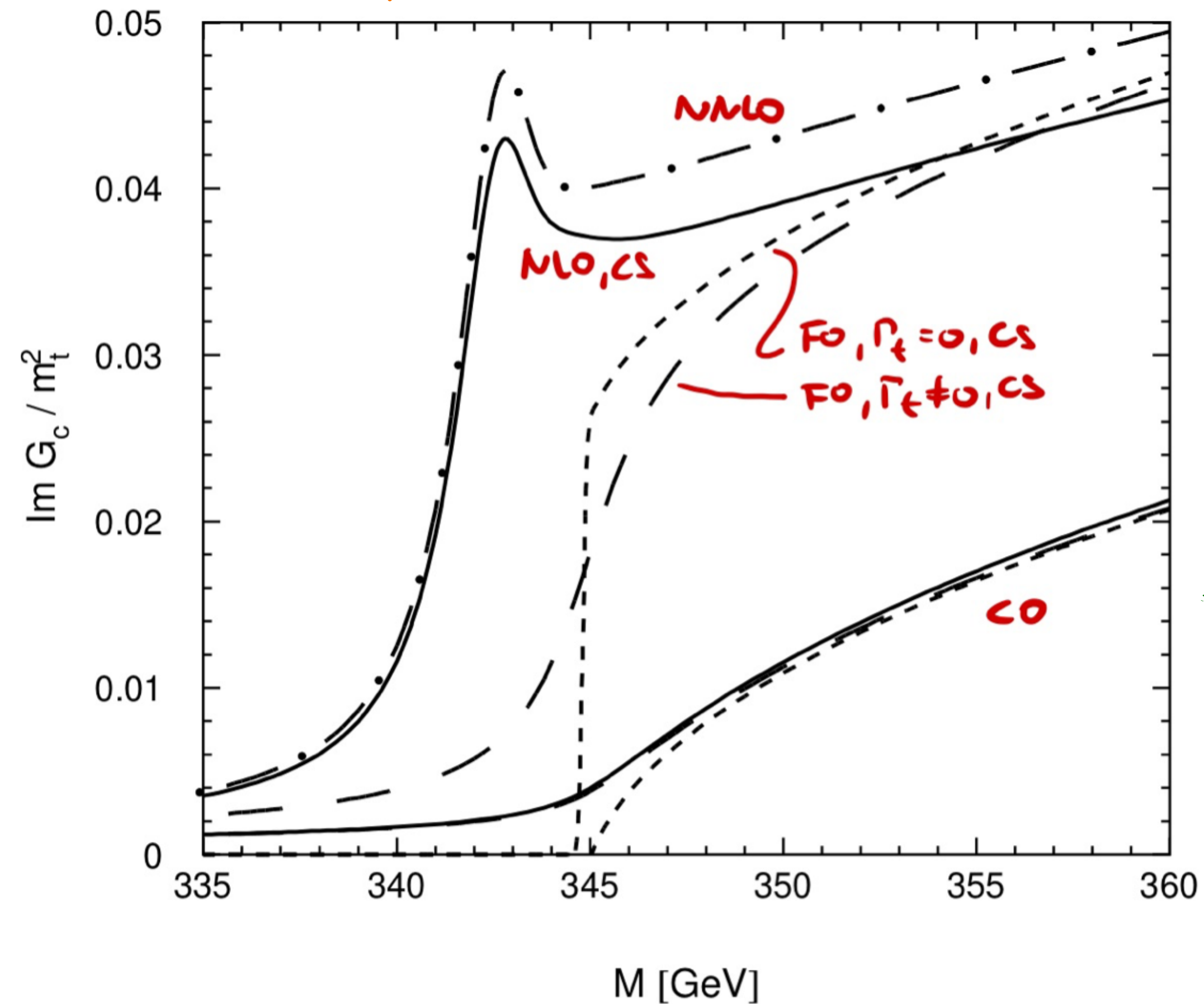
$$\left\{ 2m_t + \left[\frac{(-i\nabla)^2}{m_t} + V_C^{[1,8]}(\vec{r}) \right] - (M + i\Gamma_t) \right\} G^{[1,8]}(\vec{r}; M + i\Gamma_t) = \delta^{(3)}(\vec{r})$$



$$V_C^{[1,8]} = -\frac{4\pi\alpha_s(\mu)C^{[1,8]}}{\vec{q}^2} \left[1 + \frac{\alpha_s(\mu)}{4\pi} \left(\beta_0 \ln \frac{\mu^2}{\vec{q}^2} + a_1 \right) \right] \quad \text{Coulomb-like QCD potential at NLO}$$

$G(M + i\Gamma_t)$ contribution

[Kiyo, Kühn, Moch, Steinhauser, '08]



$$V_C^{[1,8]} = -\frac{4\pi\alpha_s(\mu)C^{[1,8]}}{\vec{q}^2} \left[1 + \frac{\alpha_s(\mu)}{4\pi} \left(\beta_0 \ln \frac{\mu^2}{\vec{q}^2} + a_1 \right) \right]$$

$$C^{[1]} = C_F = 4/3$$

$$C^{[8]} = C_F - C_A/2 = -1/6$$

Thanks to Matthias Steinhauser



F_{ij} : soft gluon resummation

- Logarithmic enhanced terms from $\mathcal{A}_c$ and $\mathcal{A}_{nc}$ due to soft gluon emissions
- Impact over $\mathcal{L} \otimes F$ as long as $z = \rho/\tau \rightarrow 1$
- Working in the Mellin space

$$F_{ij \rightarrow T}(\hat{s}, M_{t\bar{t}}^2, \mu_f^2) = \mathcal{N}_{ij} \frac{\pi^2 \alpha_s^2(\mu_r)}{3\hat{s}} \left(1 + \frac{\alpha_s(\mu_r)}{\pi} \mathcal{C}_h \right) \times \left[\delta_{ij \rightarrow T} \delta(1-z) + \frac{\alpha_s(\mu_r)}{\pi} (\mathcal{A}_c + \mathcal{A}_{nc}) \right]$$

$$F_{ij}^N = \int_0^1 dz z^{N-1} F_{ij}$$

- NLL soft-gluons threshold resummation

$$F_{ij}^{\text{res}} = \int \frac{dN}{2\pi i} x^{-N} \left(F_{ij}^N - F_{ij}^N|_{\text{NLO}} \right) + F_{ij}^{\text{NLO}}$$



F_{ij} : soft gluon resummation

- Run2@LHC

$$\sqrt{S} = 13 \text{ TeV}$$

$$m_t = 172.5 \text{ GeV}$$

$$\Gamma_t = 1.4 \text{ GeV}$$

$$\text{NNPDF_3.1_NNLO_as_0118}$$

$$\alpha_s(M_Z) = 0.1180$$

$$\mu_r = \mu_f \in \{m_t, 2m_t, 4m_t\}$$

	$\mu = m_t \quad 2m_t \quad 4m_t$					
$M_{t\bar{t}} = 2m_t$	NLO			resummed		
$gg \rightarrow {}^1S_0^{[1]}$	18.2	18.7	18.3	19.4	20.5	21.1
$gg \rightarrow {}^1S_0^{[8]}$	55.8	55.2	52.8	60.0	61.5	62.0
$q\bar{q} \rightarrow {}^3S_1^{[8]}$	21.7	22.3	22.0	22.4	22.4	22.0

[Garzelli, GL, Moch, Steinhauser, Zenaiev '24]



Up to 10% from soft gluon resummation
for the leading channels

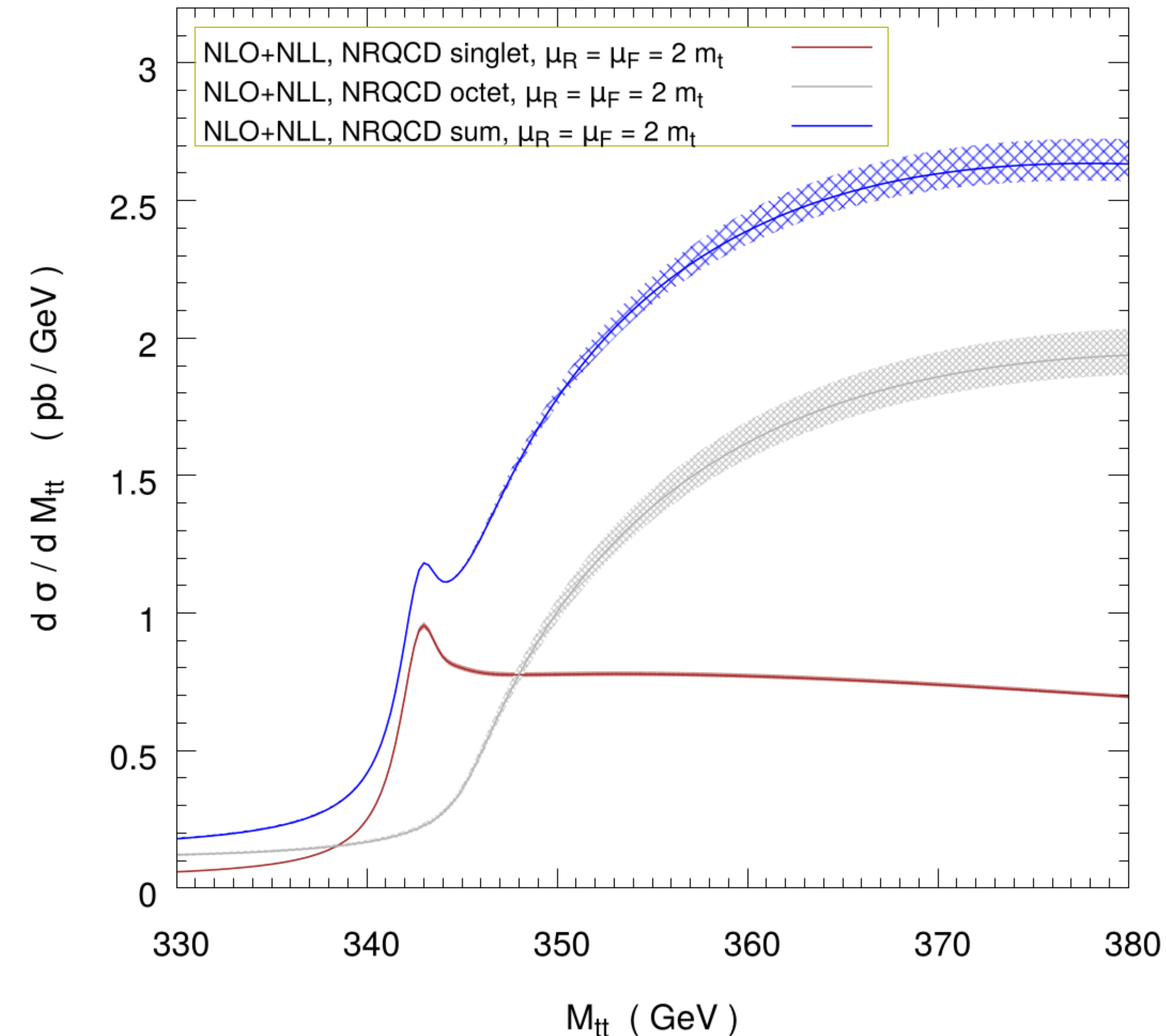
No large threshold logarithms from the S -wave channels

$$gg \rightarrow {}^3S_1^{[1,8]} + q\bar{q} \rightarrow {}^1S_0^{[1,8]} + gq \rightarrow {}^1S_0^{[1,8]} + gq \rightarrow {}^3S_1^{[8]} = \mathcal{O}(5\%)$$

NRQCD Predictions



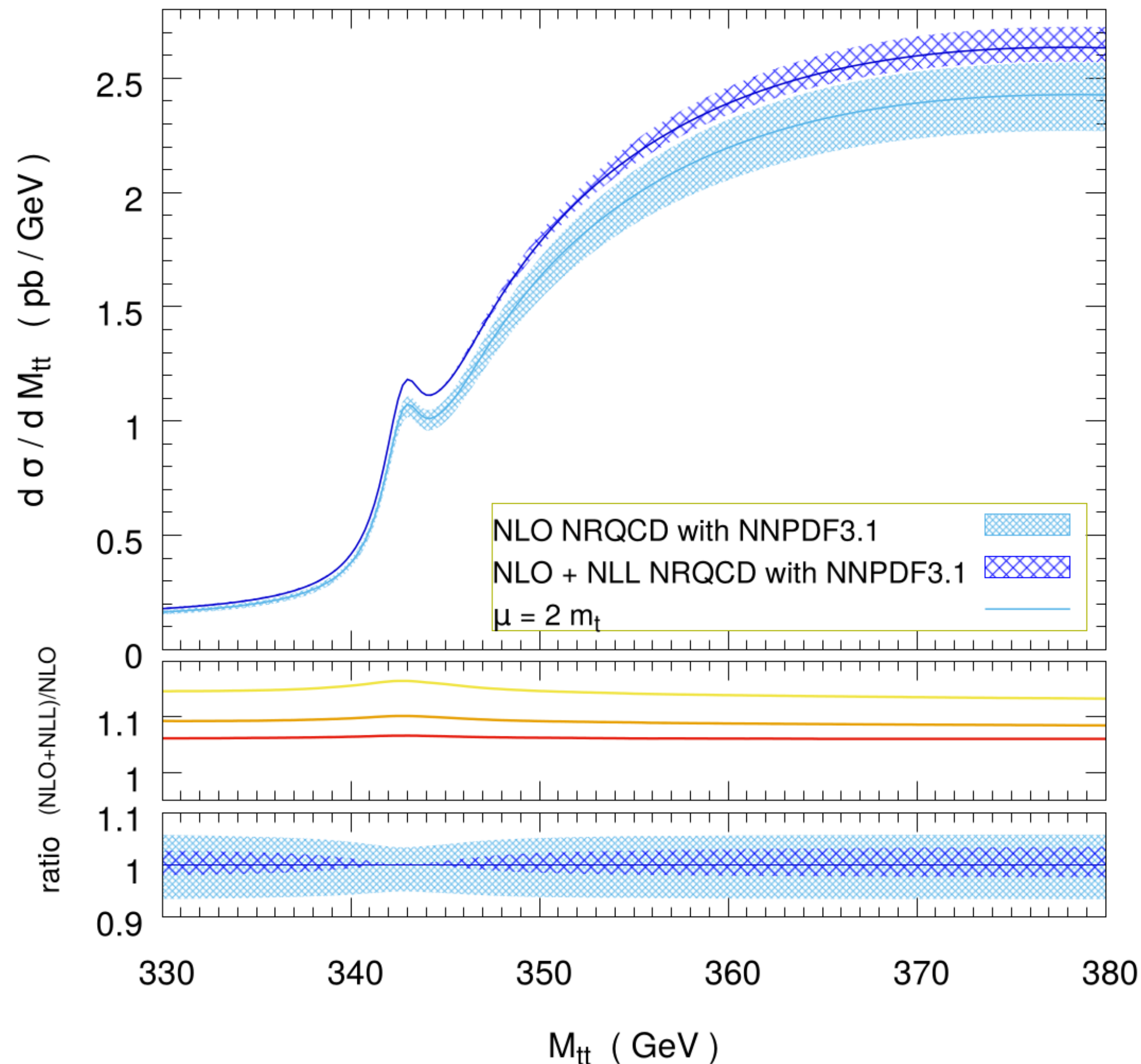
[Garzelli, GL, Moch, Steinhauser, Zenaiev '24]



- $gg \rightarrow {}^1S_0^{[1]}$ shows a resonance peak in the kinematic region $M_{t\bar{t}} \simeq 2m_t$
- Depletion in the $gg \rightarrow {}^1S_0^{[8]} + q\bar{q} \rightarrow {}^3S_1^{[8]}$ due to the repulsive nature of $V_C^{[8]}$
- Impact of Γ_t in the region below the threshold due to $G^{[8]}(M_{t\bar{t}} + i\Gamma_t)$ falling off slowly in this region

NRQCD Predictions

[Garzelli, GL, Moch, Steinhauser, Zenaiev '24]

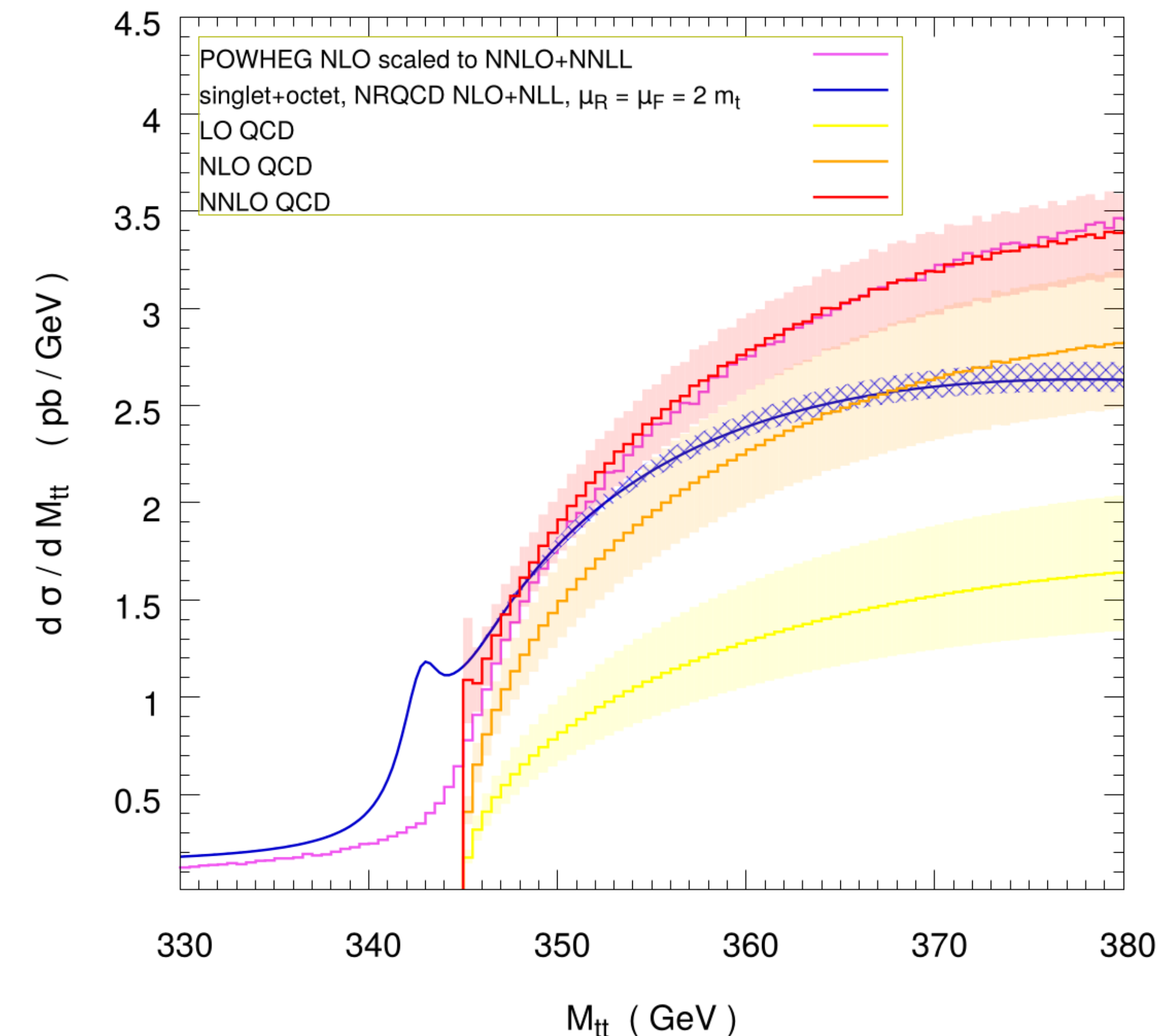


- $\mu = m_t, 2m_t, 4m_t$
- Scale stability improved in the threshold region $M_{t\bar{t}} \simeq 2m_t$
- $+\mathcal{O}(10\%)$ from resummation at the central scale choice

$M_{t\bar{t}} = 2m_t$	NLO			resummed		
$gg \rightarrow {}^1S_0^{[1]}$	18.2	18.7	18.3	19.4	20.5	21.1
$gg \rightarrow {}^1S_0^{[8]}$	55.8	55.2	52.8	60.0	61.5	62.0
$q\bar{q} \rightarrow {}^3S_1^{[8]}$	21.7	22.3	22.0	22.4	22.4	22.0

Threshold and continuum

[Garzelli, GL, Moch, Steinhauser, Zenaiev '24]

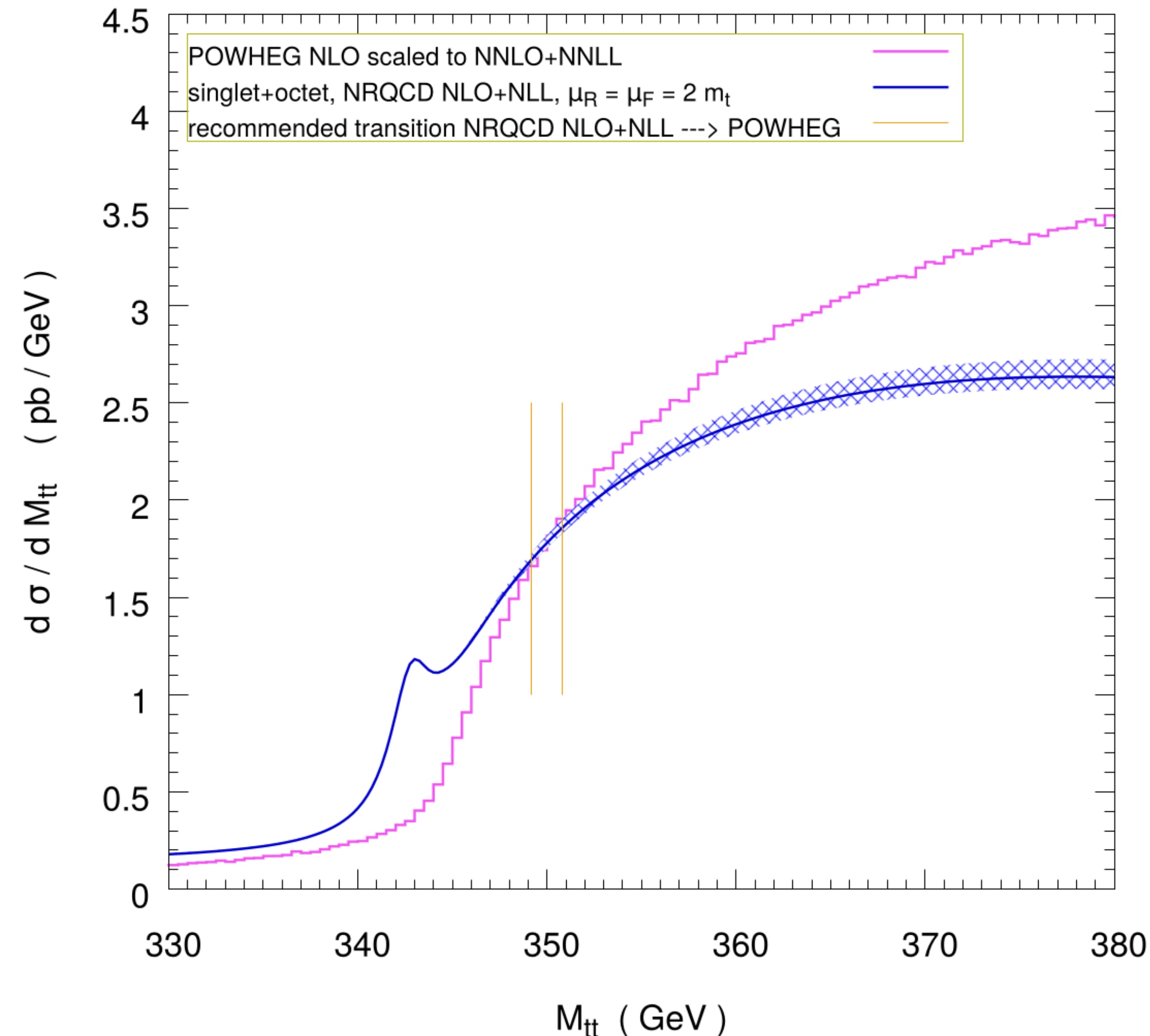


- **CMS-PAS-HIG-22-013** based on NLO pQCD
POWHEG predictions rescaled to NNLO+NNLL
also accounting for NLO EW effects
- **NRQCD** predictions valid in the region

$$|M_{t\bar{t}} - 2m_t| \lesssim 5 \text{ GeV}$$
- Fixed-order pQCD at **LO**, **NLO**, **NNLO**
- The transition from the threshold to the continuum needs accurate study

Threshold and continuum

[Garzelli, GL, Moch, Steinhauser, Zenaiev '24]



- NRQCD predictions cease to be valid for

$$|M_{t\bar{t}} - 2m_t| \gtrsim 5 \text{ GeV}$$

- $\sigma_{\text{NRQCD}}(340 \text{ GeV} \lesssim M_{t\bar{t}} \lesssim 350 \text{ GeV}) = 12 \text{ pb}$
 $\sigma_{\text{POWHEG}}(340 \text{ GeV} \lesssim M_{t\bar{t}} \lesssim 350 \text{ GeV}) = 8.4 \text{ pb}$
- $\sigma_{\text{CMS}} = 8.8^{+1.2}_{-1.4} \text{ pb}$
 $\sigma_{\text{ATLAS}} = 9.0 \pm 1.3 \text{ pb}$

Improvements in NRQCD estimate are needed

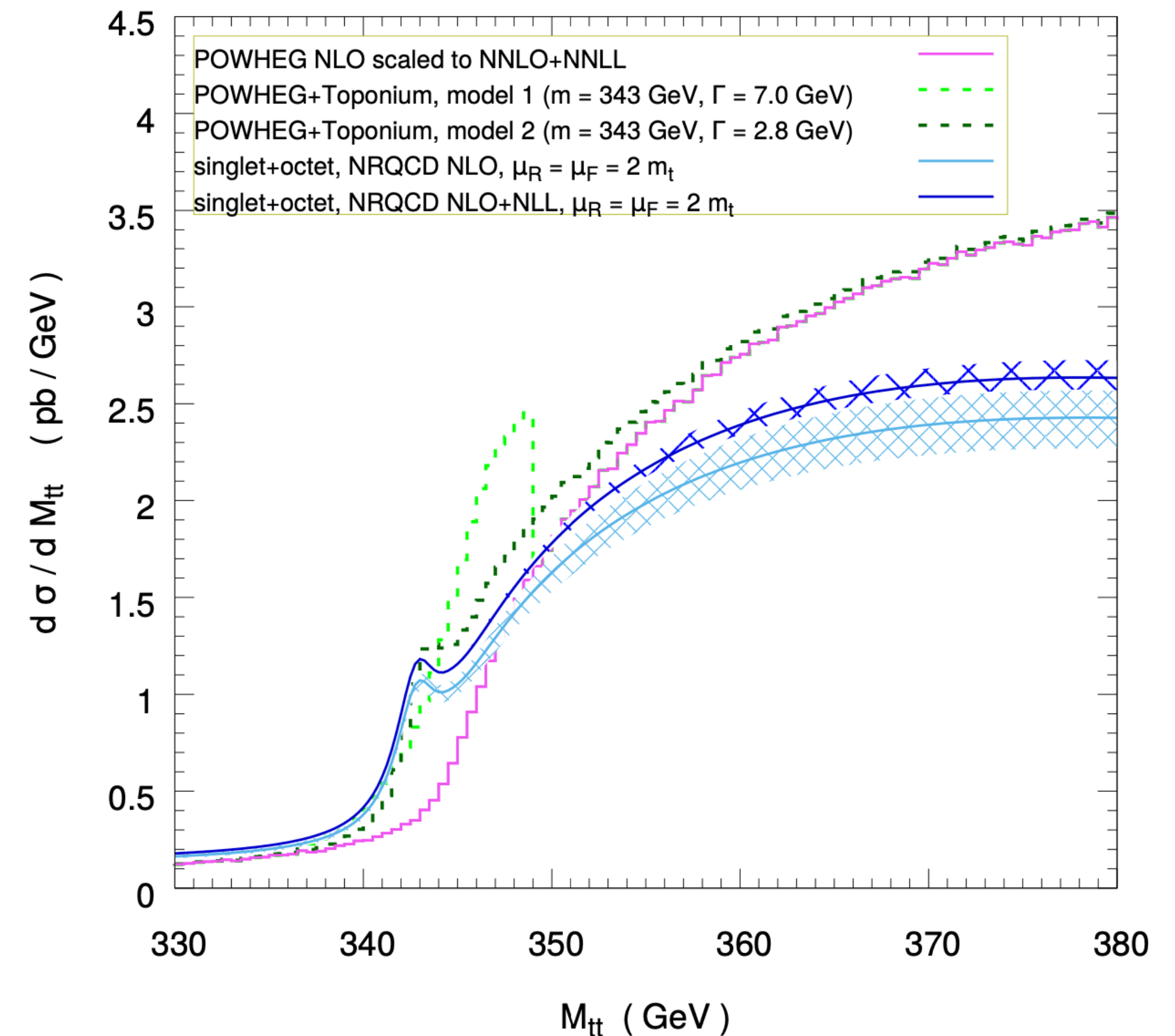
Modelling $t\bar{t}$ resonances



- Quasi-bound state due to a pseudoscalar particle η_t
Fuks, Hagiwara, Ma, Zheng [2102.11281], Maltoni et al. [2401.08751]
- LO MEs reweighted with the LO zero-distance Green Function
Fuks, Hagiwara, Ma, Zheng [2411.18962]
- Resonance effect added on the FO pQCD predictions
[CMS-TOP-24-007]
- The model only accounts for colour-singlet final states
- Colour repulsive interactions not included

$$\mathcal{L}_{\eta_t} = \frac{1}{2} \partial_\mu \phi_{\eta_t} \partial^\mu \phi_{\eta_t} - \frac{1}{2} m_{\eta_t} \phi_{\eta_t}^2 - \frac{1}{4} g_{gg\eta_t} \phi_{\eta_t} G_{\mu\nu}^a \tilde{G}^{a\mu\nu} - i g_{t\bar{t}\eta_t} \phi_{\eta_t} \bar{t} \gamma_5 t$$

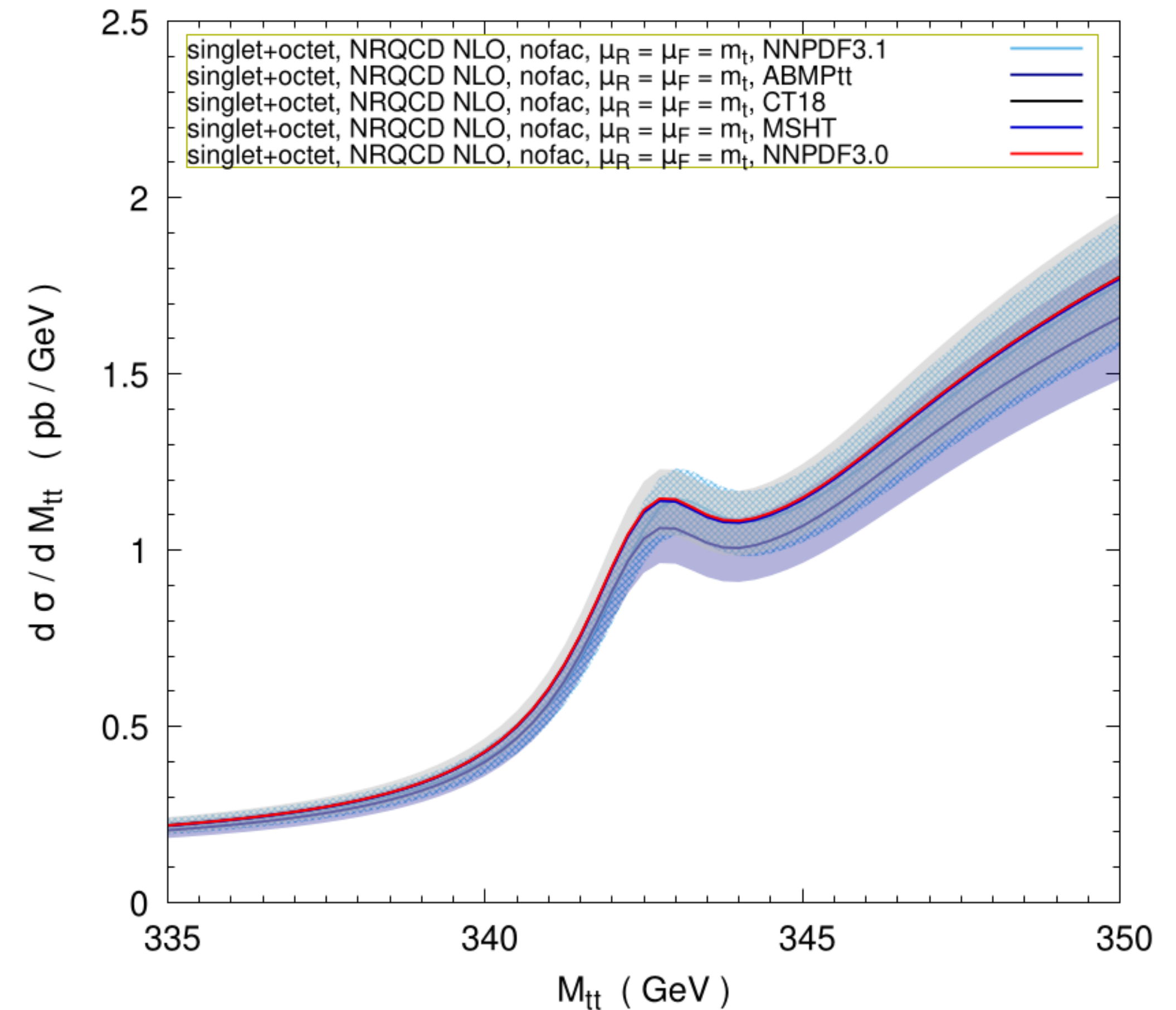
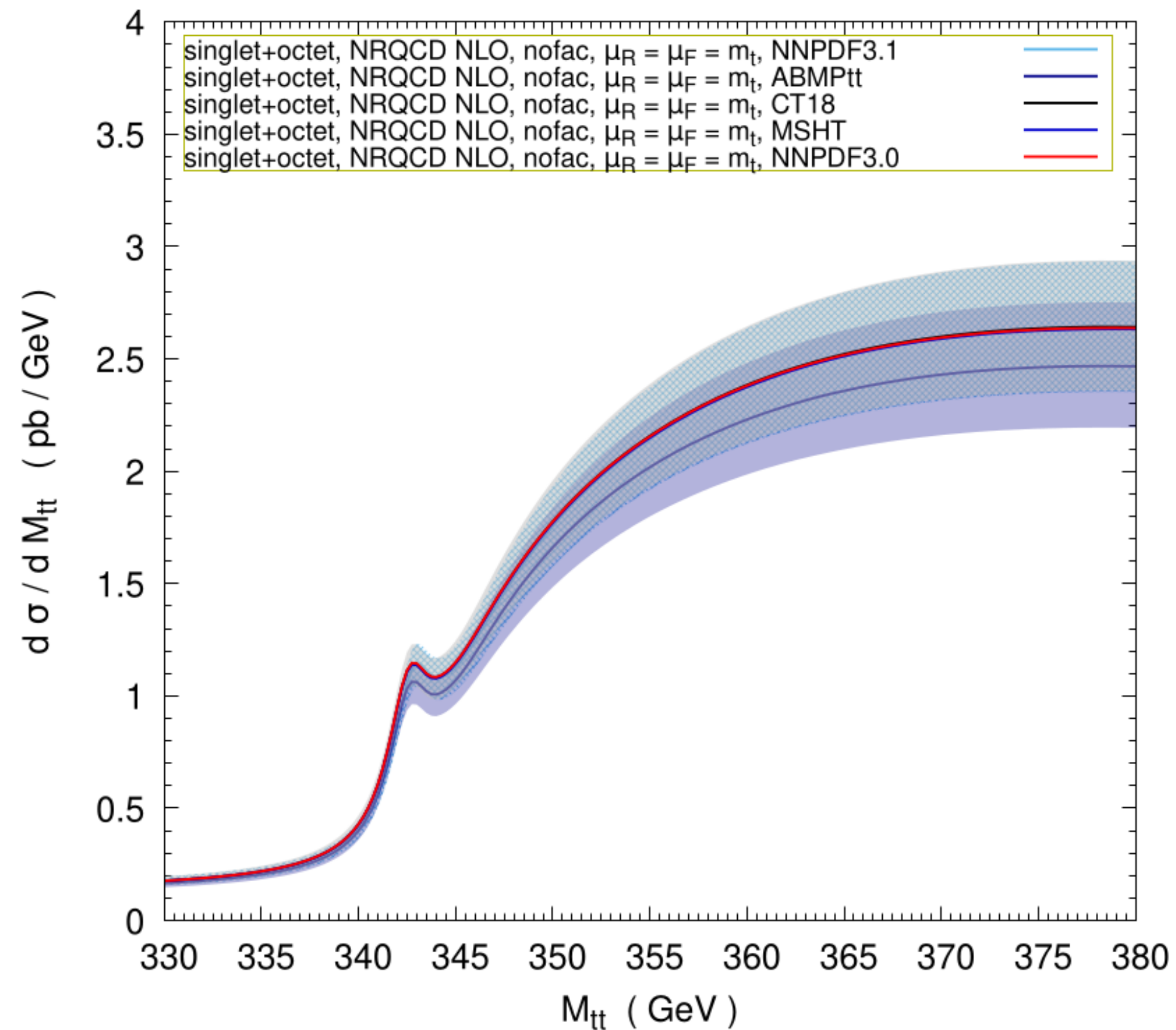
Thanks to Anuar, Garzelli, Grohsjean, Jeppe, Schwanenberger



NRQCD Predictions:PDF variations



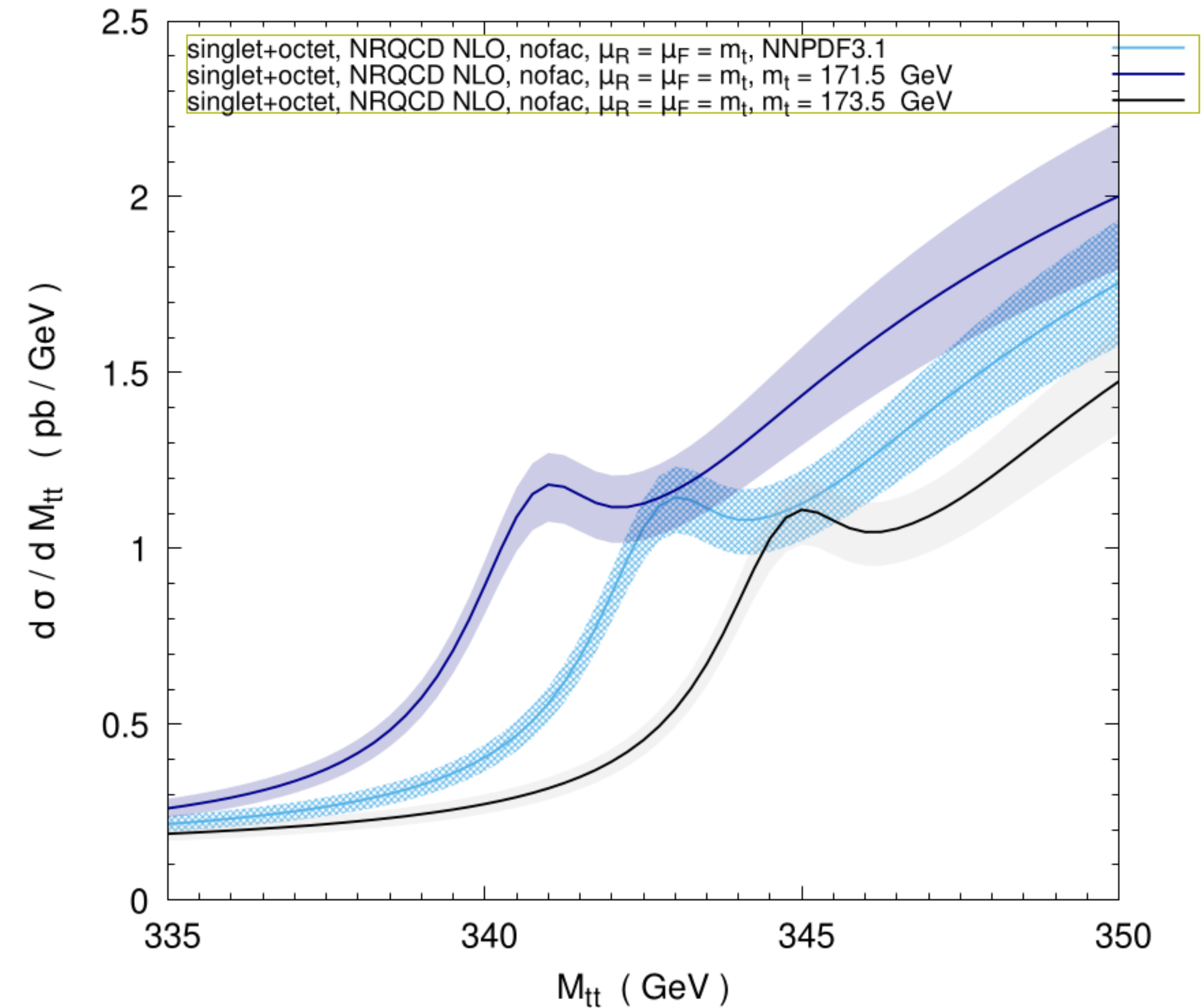
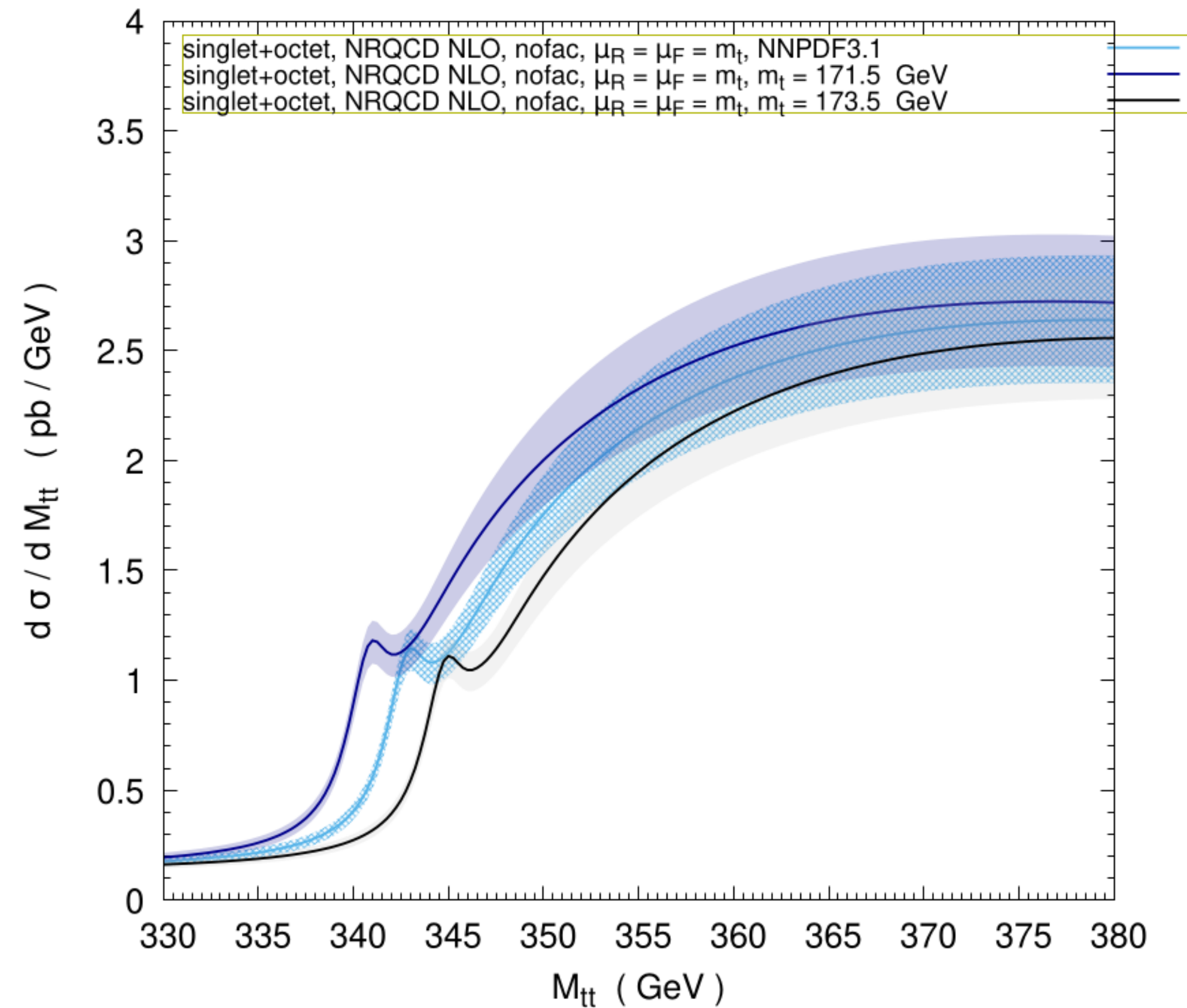
[Garzelli, GL, Moch, Steinhauser, Zenaiev to appear soon]



NRQCD Predictions: m_t variations

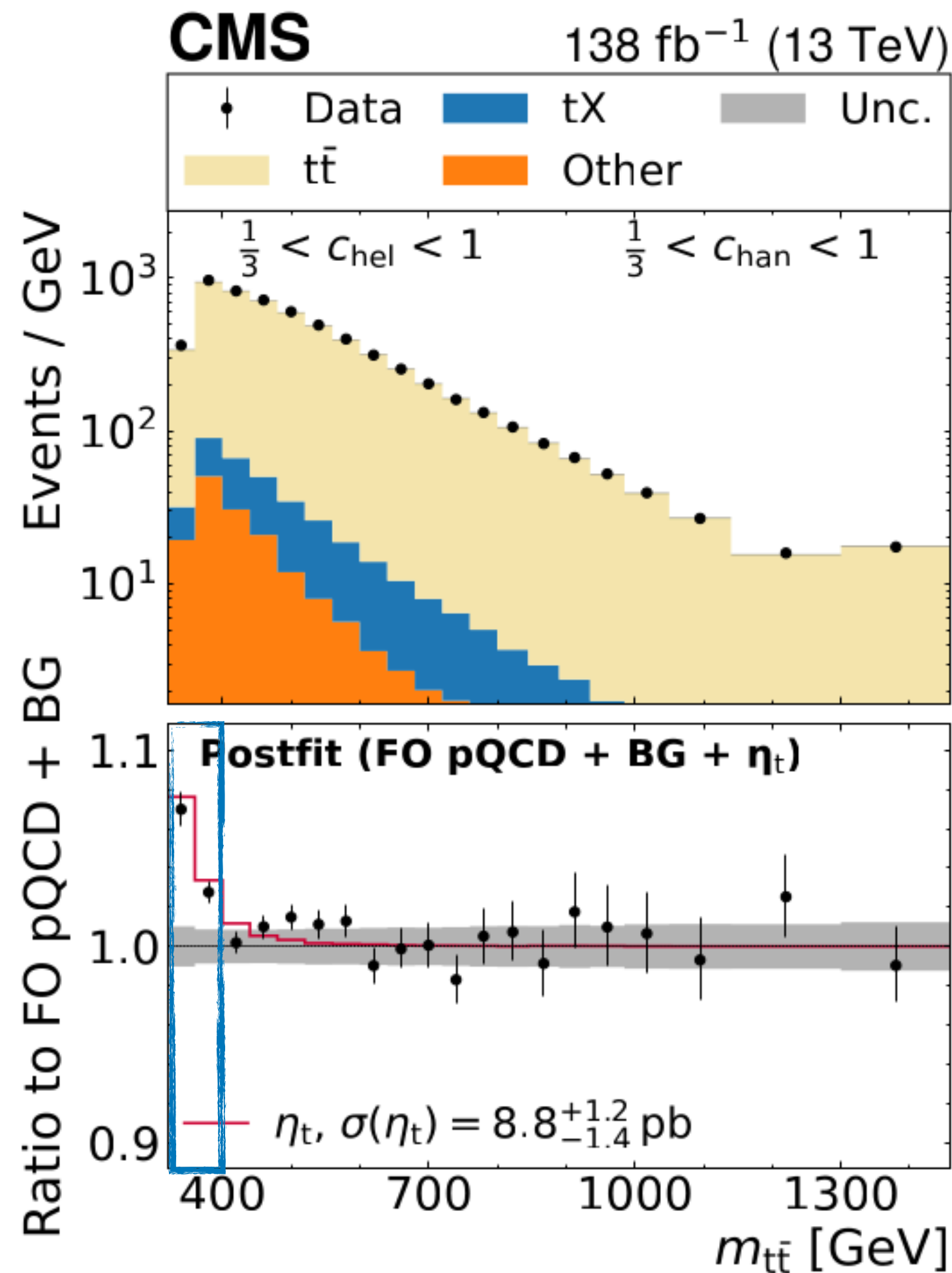


[Garzelli, GL, Moch, Steinhauser, Zenaiev to appear soon]



Other Approaches

[Nason, Re, Rottoli '25]



- Large experimental resolution
- Excess observed in the first bin $M_{t\bar{t}} \lesssim 380$ GeV
- Fixed-order pQCD can be applied
- Quasi-bound state at N3LO $\mathcal{O}(\alpha_s^3)$

Is the excess of event only related to quasi-bound state interactions?

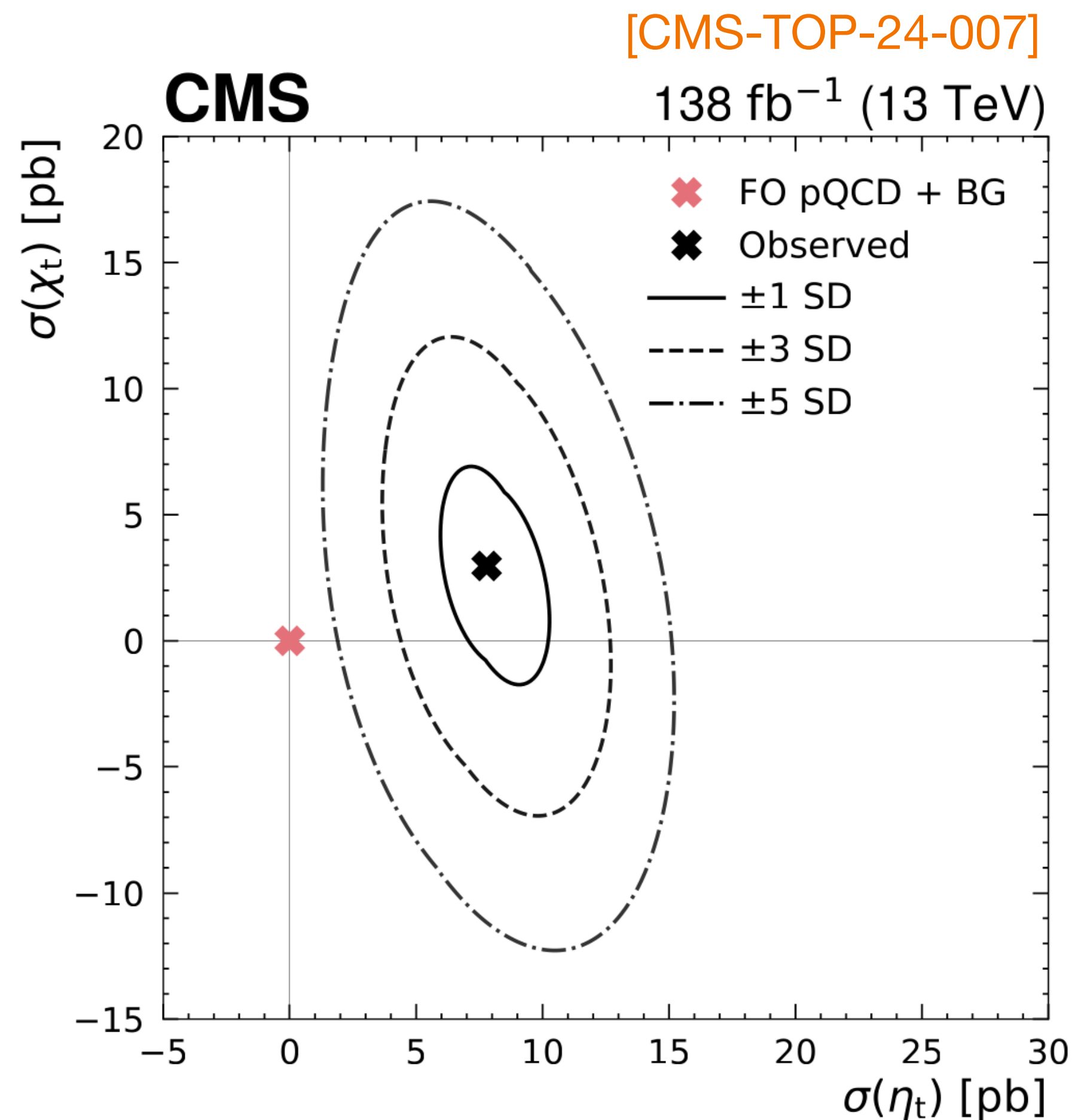


Conclusions and Outlooks

- Excess of events in $M_{t\bar{t}}$ distribution from CMS and ATLAS
- NRQCD to study threshold effects in $t\bar{t}$ pair production at the LHC
- Uncertainties due to missing higher orders effects in the NRQCD expansion
- Well formulated matching prescription is demanded
- NRQCD predictions for other observables than $M_{t\bar{t}}$

BACKUP

Conclusions and Outlooks



- Simultaneous fit including scalar χ_t and pseudoscalar η_t contributions

$\sigma(\chi_t)$ compatible with zero within one SD