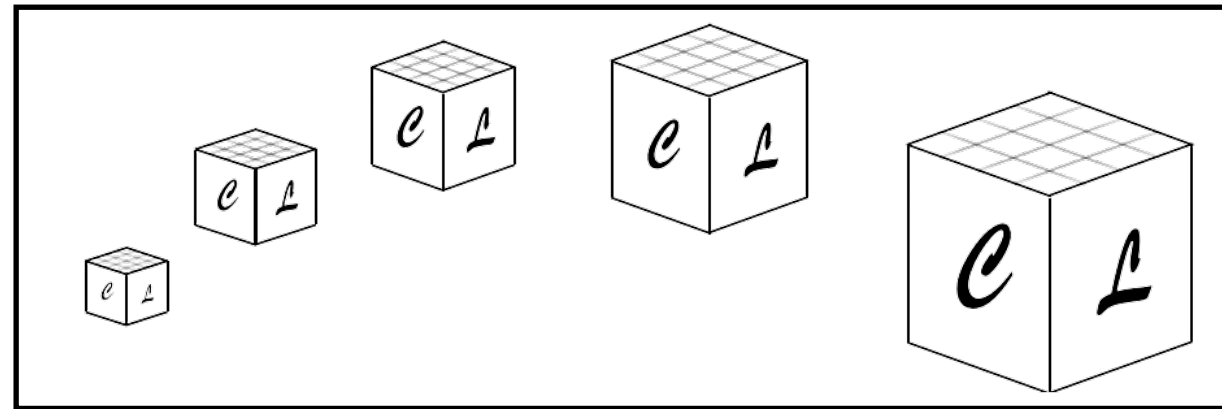
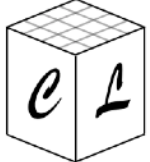




Cosmo $\mathcal{L}$ attice school:

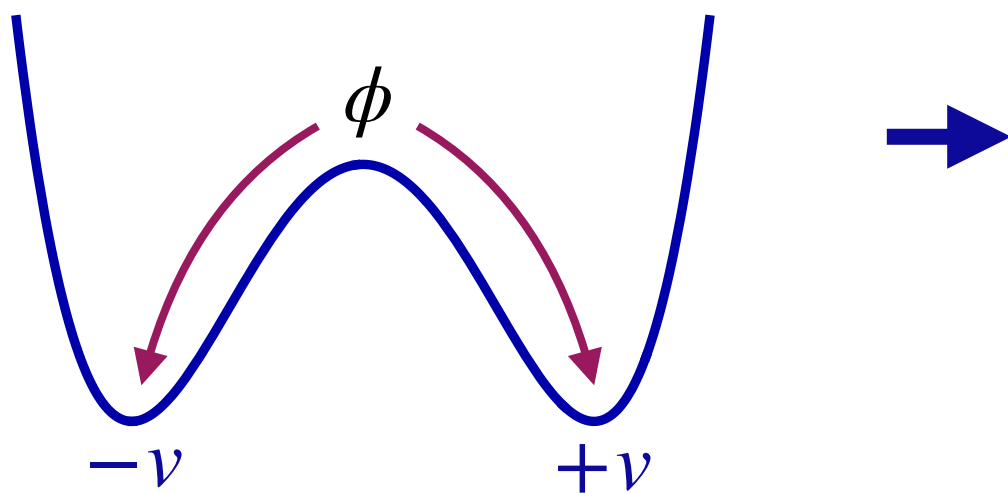
Lecture 11: Lattice simulations of domain walls

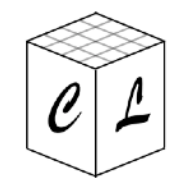
Francisco Torrentí
ICCUB & U. Barcelona



Domain walls

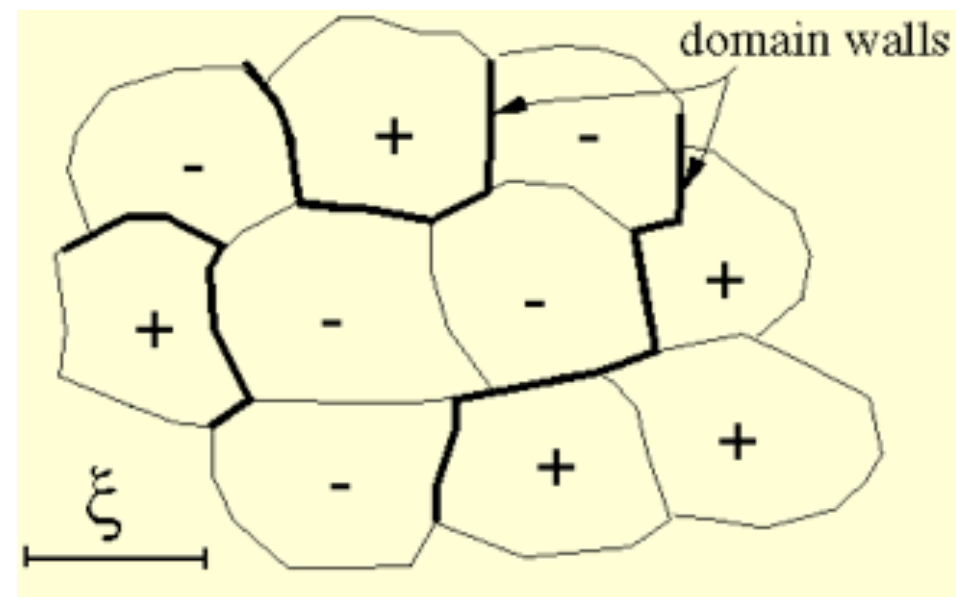
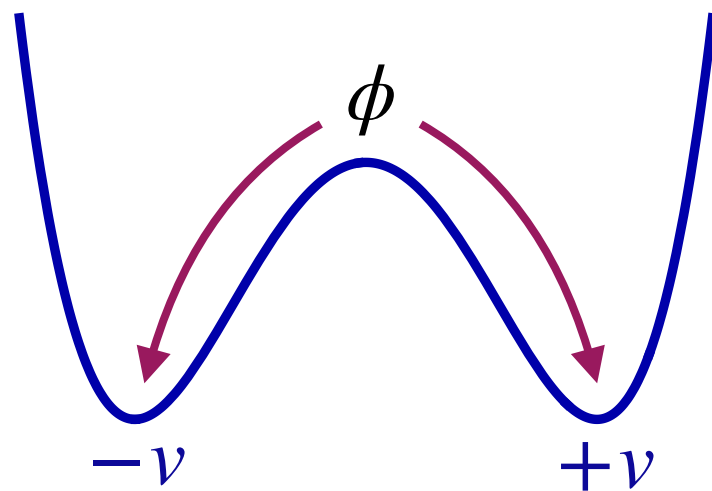
$$V(\phi) = \frac{\lambda}{4}(\phi^2 - v^2)^2$$

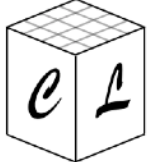




Domain walls

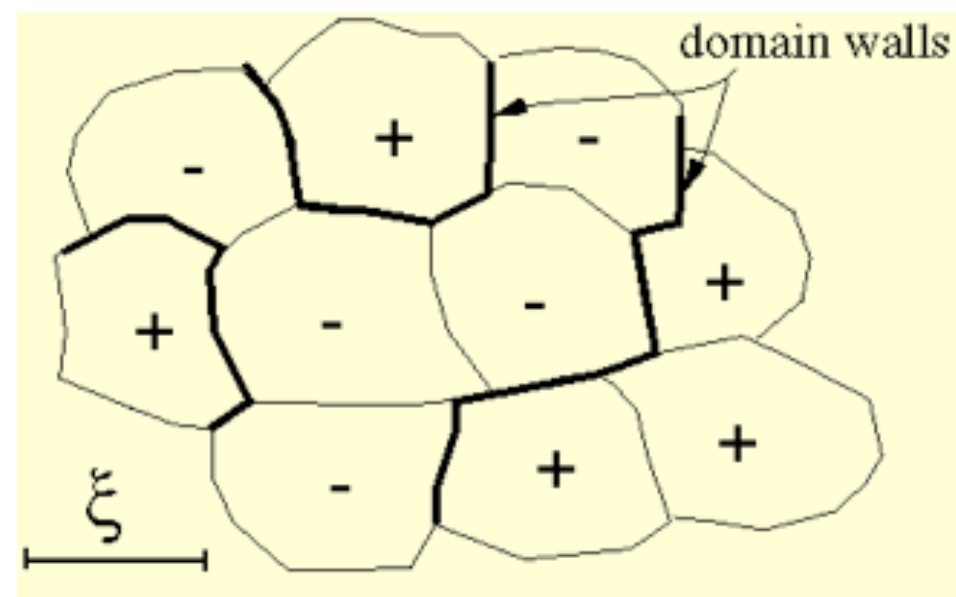
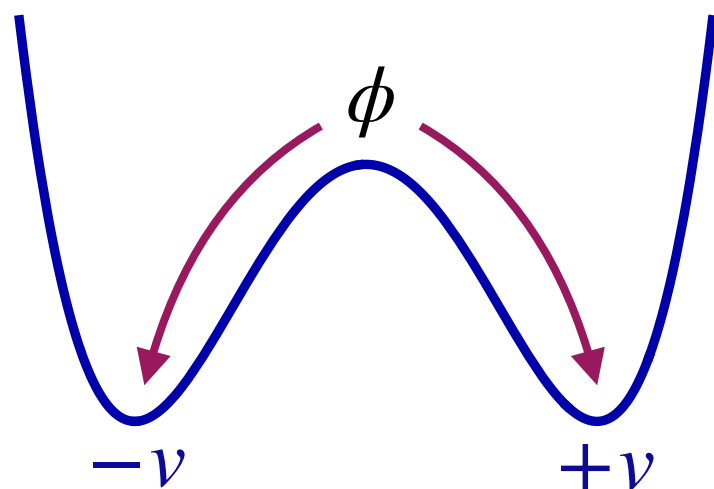
$$V(\phi) = \frac{\lambda}{4}(\phi^2 - v^2)^2$$





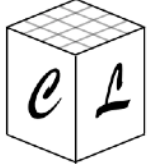
Domain walls

$$V(\phi) = \frac{\lambda}{4}(\phi^2 - v^2)^2$$



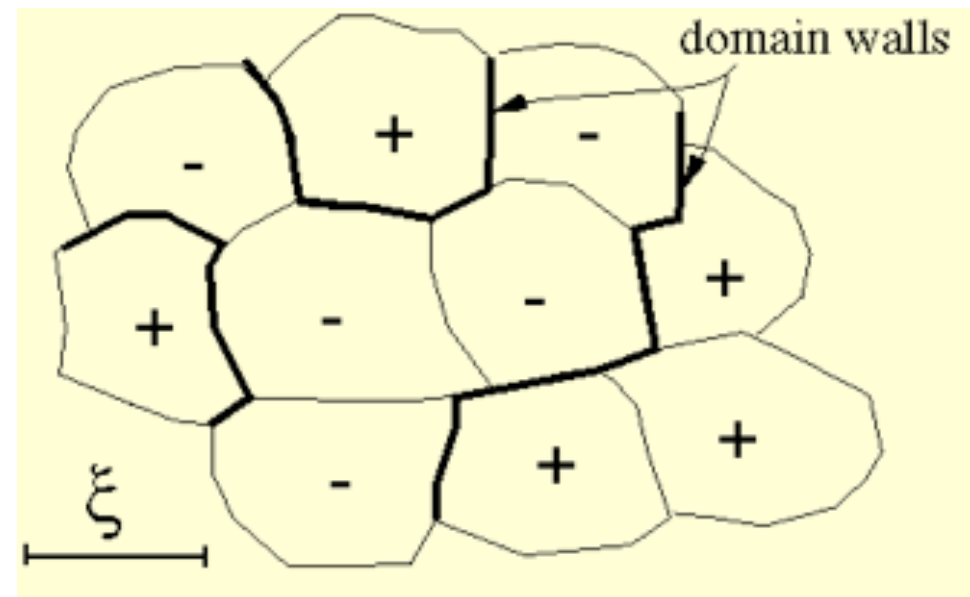
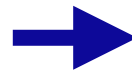
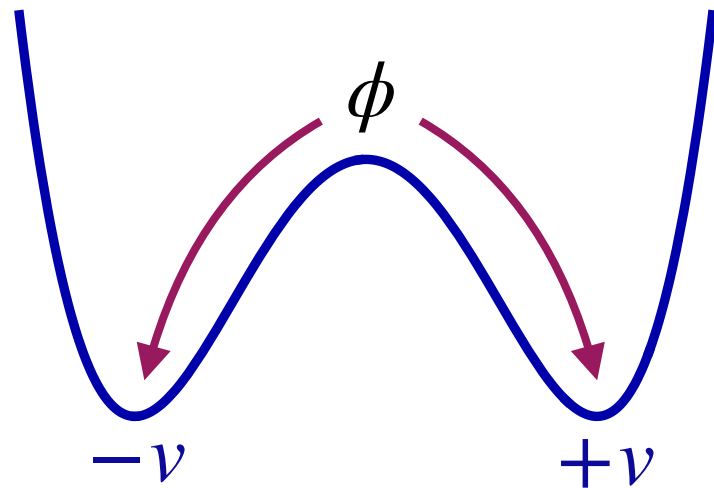
► 1-D solution:

$$\frac{\partial^2 \phi}{\partial z^2} = \frac{dV}{d\phi}$$



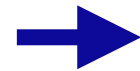
Domain walls

$$V(\phi) = \frac{\lambda}{4}(\phi^2 - v^2)^2$$



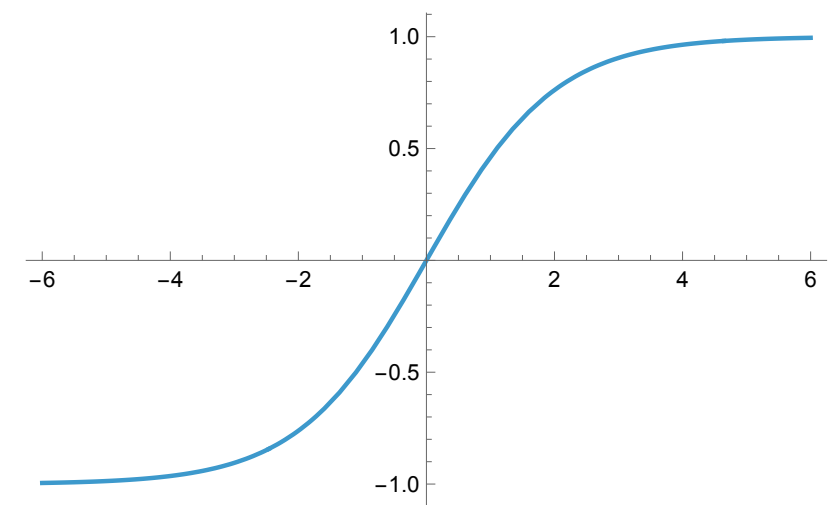
► 1-D solution:

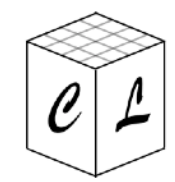
$$\frac{\partial^2 \phi}{\partial z^2} = \frac{dV}{d\phi}$$



$$\phi(z) = v \tanh \left(\frac{m}{2}(z - z_0) \right)$$

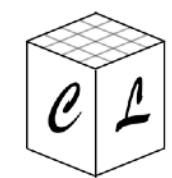
$$m \equiv \sqrt{2\lambda}v$$





Domain walls: quartic potential

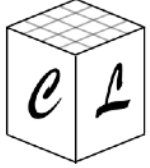
$$V(\phi) = \frac{\lambda}{4}(\phi^2 - v^2)^2$$



Domain walls: quartic potential

$$V(\phi) = \frac{\lambda}{4}(\phi^2 - v^2)^2$$

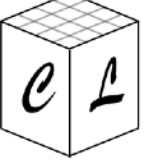
- The DW network soon achieves a “scaling regime” (~one DW per Hubble volume):



Domain walls: quartic potential

$$V(\phi) = \frac{\lambda}{4}(\phi^2 - v^2)^2$$

- The DW network soon achieves a “scaling regime” (~one DW per Hubble volume):
 - Width: $\delta_w = m \equiv \sqrt{2\lambda}v$ (physical width)



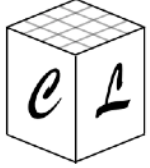
Domain walls: quartic potential

$$V(\phi) = \frac{\lambda}{4}(\phi^2 - v^2)^2$$

► The DW network soon achieves a “scaling regime” (~one DW per Hubble volume):

• Width: $\delta_w = m \equiv \sqrt{2\lambda}v$ (physical width)

• Tension: $\sigma \equiv \int_{-\infty}^{\infty} dz \left[\frac{1}{2}(\phi')^2 + V(\phi) \right] = \int_{-v}^{+v} \sqrt{2V(\phi)} d\phi = \frac{2}{3}\sqrt{2\lambda}v^3$
(energy/area)



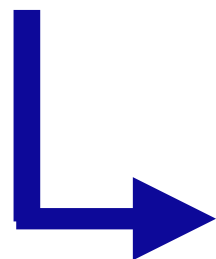
Domain walls: quartic potential

$$V(\phi) = \frac{\lambda}{4}(\phi^2 - v^2)^2$$

► The DW network soon achieves a “scaling regime” (~one DW per Hubble volume):

• Width: $\delta_w = m \equiv \sqrt{2\lambda}v$ (physical width)

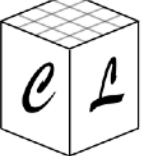
• Tension: $\sigma \equiv \int_{-\infty}^{\infty} dz \left[\frac{1}{2}(\phi')^2 + V(\phi) \right] = \int_{-v}^{+v} \sqrt{2V(\phi)} d\phi = \frac{2}{3}\sqrt{2\lambda}v^3$
(energy/area)



Energy density: $\rho_{\text{dw}} = \frac{\sigma A a^2}{V a^3} = 2\mathcal{A}\sigma H$

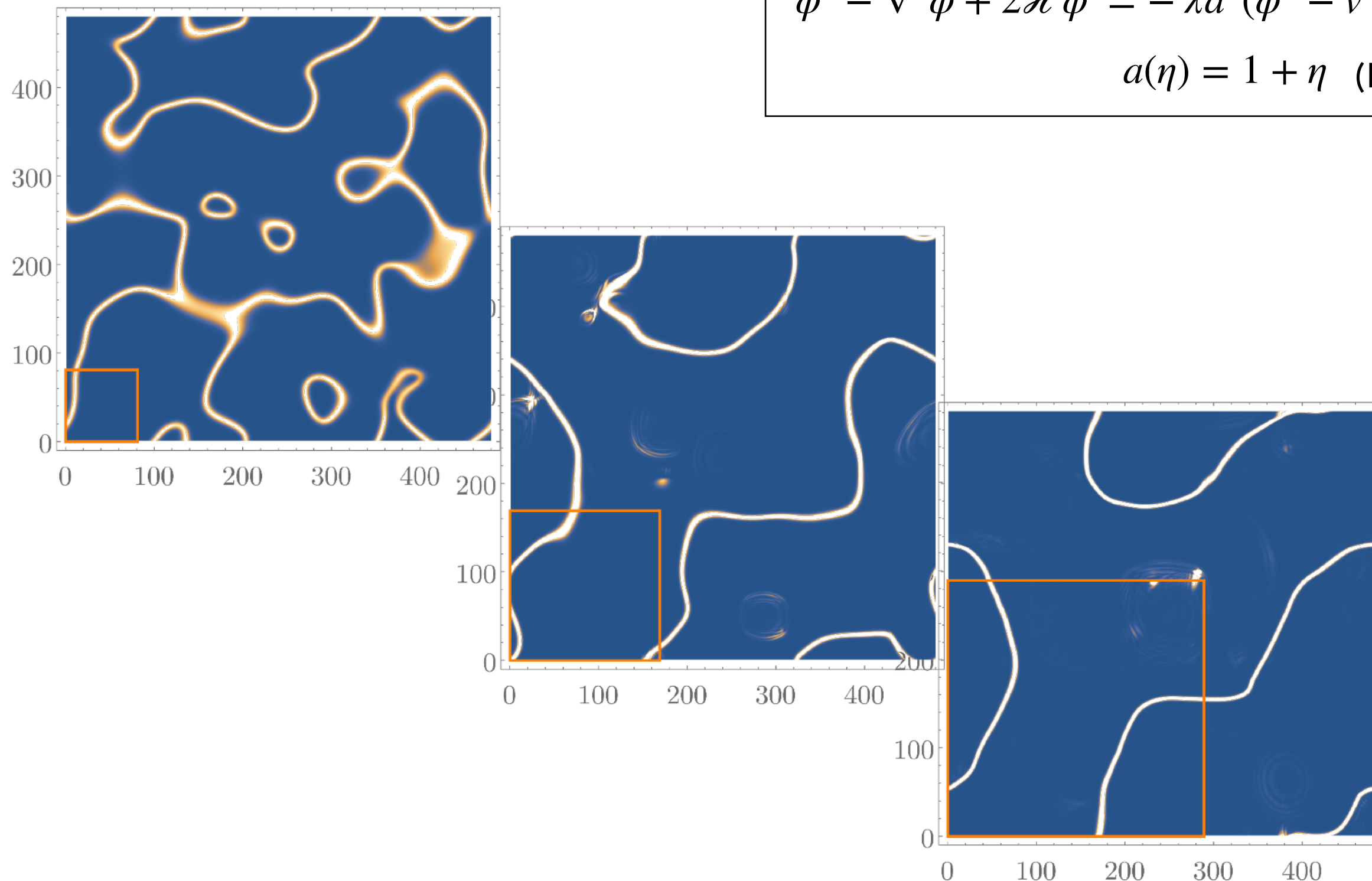
$$\mathcal{A} \equiv \frac{A}{V} \frac{1}{2aH} \simeq \text{const}$$

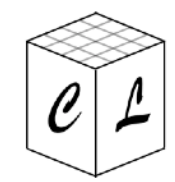
Area parameter



Domain walls: lattice simulations

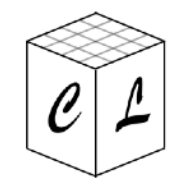
$$\phi'' - \nabla^2 \phi + 2\mathcal{H}\phi' = -\lambda a^2(\phi^2 - v^2)\phi$$
$$a(\eta) = 1 + \eta \quad (\text{RD})$$





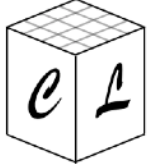
Domain walls: lattice simulations

- Comoving width decreases with time: $\delta_w/a \propto \eta^{-1}$



Domain walls: lattice simulations

- Comoving width decreases with time: $\delta_w/a \propto \eta^{-1}$
- Number of Hubble patches contained in the lattice: $N_H = \left(\frac{a(\eta)L}{H^{-1}(\eta)} \right)^3 \propto \eta^{-3}$

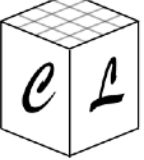


Domain walls: lattice simulations

- Comoving width decreases with time: $\delta_w/a \propto \eta^{-1}$
- Number of Hubble patches contained in the lattice: $N_H = \left(\frac{a(\eta)L}{H^{-1}(\eta)} \right)^3 \propto \eta^{-3}$

Separation of comoving momenta scales:

$$k_w \equiv \frac{2\pi a(\eta)}{\delta_w} \propto \eta; \quad k_h \equiv \frac{2\pi a(\eta)}{H^{-1}(\eta)} \propto \eta^{-1} \quad \rightarrow \quad \text{We require: } k_{\text{IR}} < k_h < k_w < k_{\text{UV}}$$



Domain walls: lattice simulations

- Comoving width decreases with time: $\delta_w/a \propto \eta^{-1}$
- Number of Hubble patches contained in the lattice: $N_H = \left(\frac{a(\eta)L}{H^{-1}(\eta)} \right)^3 \propto \eta^{-3}$

Separation of comoving momenta scales:

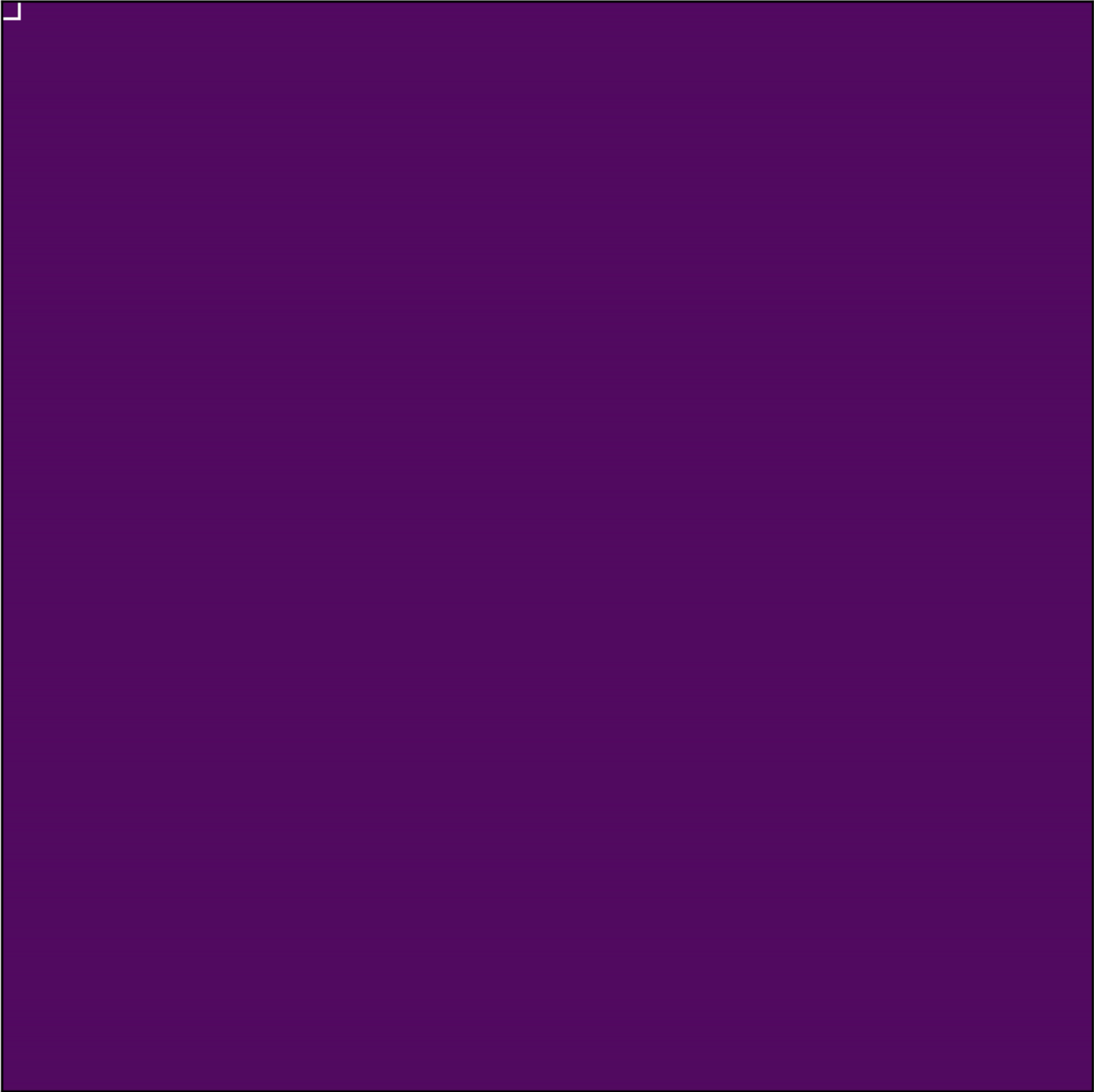
$$k_w \equiv \frac{2\pi a(\eta)}{\delta_w} \propto \eta; \quad k_h \equiv \frac{2\pi a(\eta)}{H^{-1}(\eta)} \propto \eta^{-1} \quad \rightarrow \quad \text{We require: } k_{\text{IR}} < k_h < k_w < k_{\text{UV}}$$

- Given a lattice of N lattice points and comoving length L , the maximum time for which DWs can be resolved are:

$$\eta < \eta_{\text{max}} = \left(\frac{\sqrt{3}}{2} \frac{N}{H_i L} - 1 \right) H_i^{-1}$$

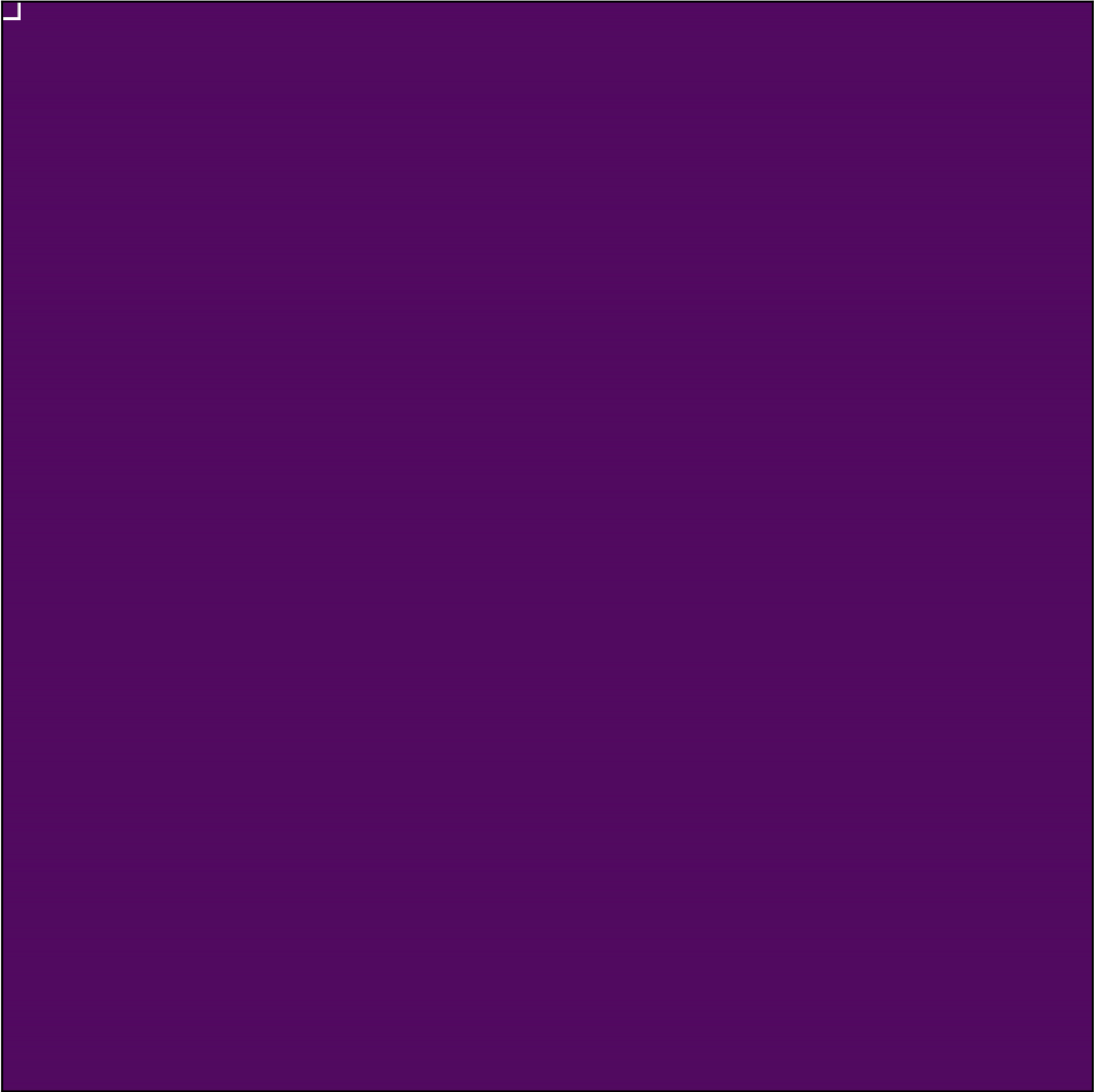
$N = 3060$

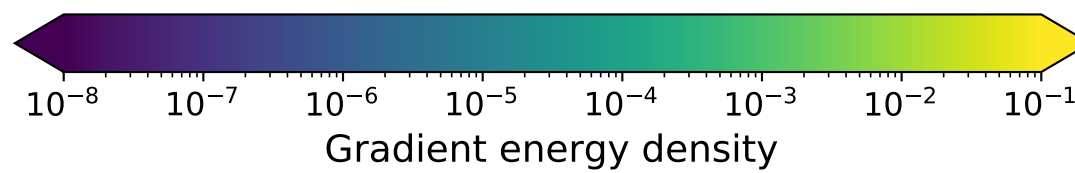
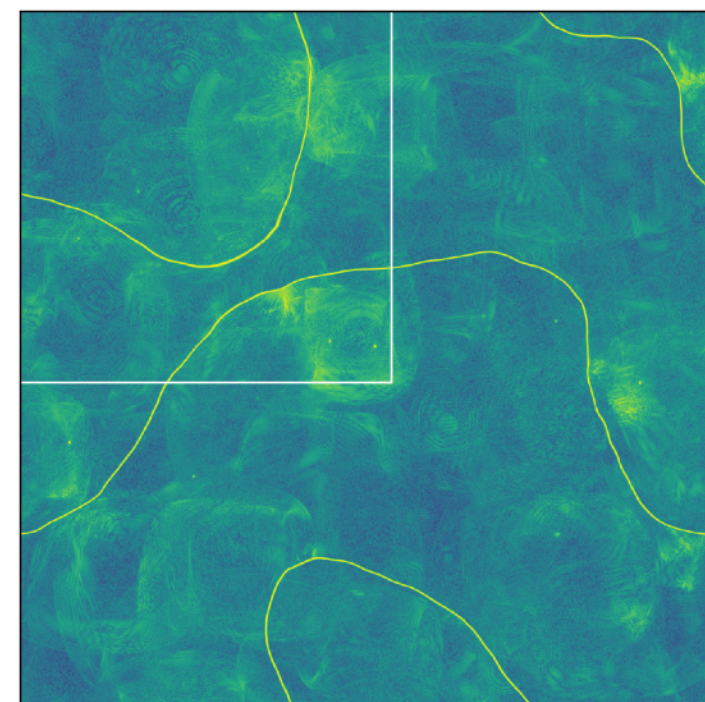
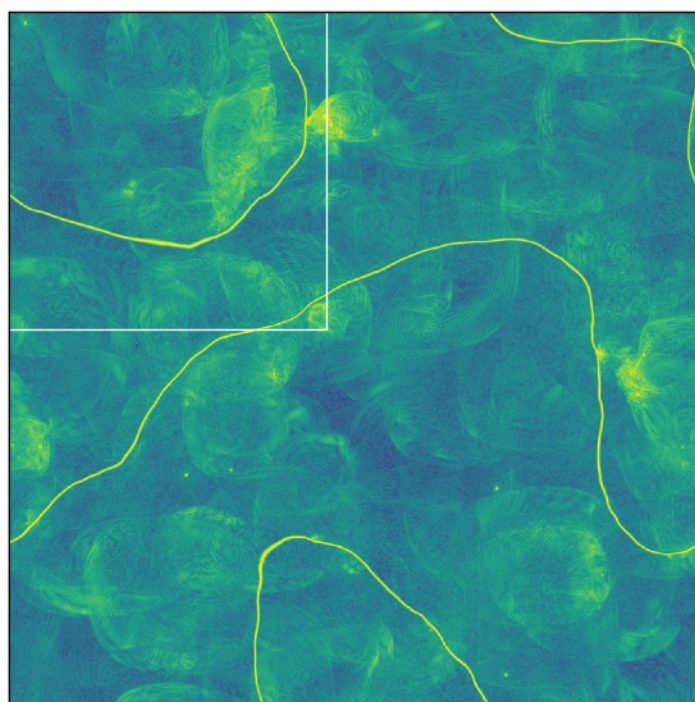
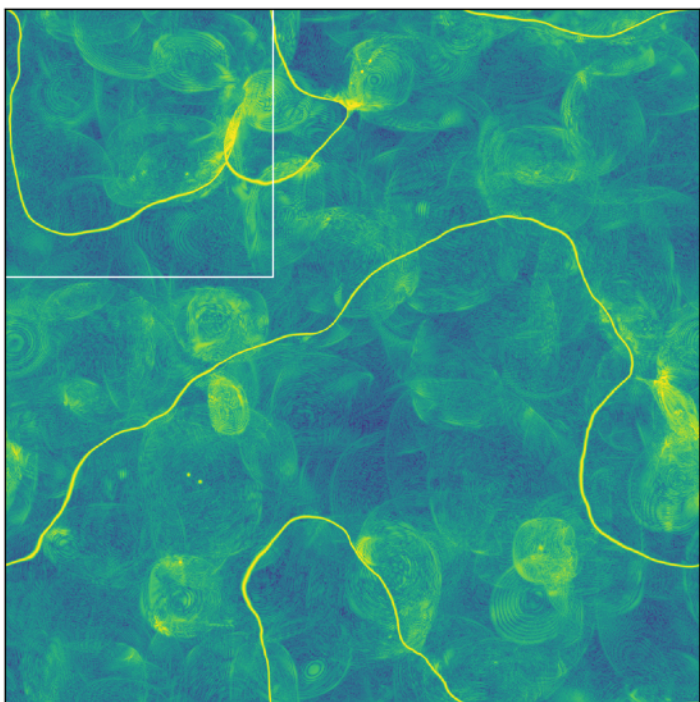
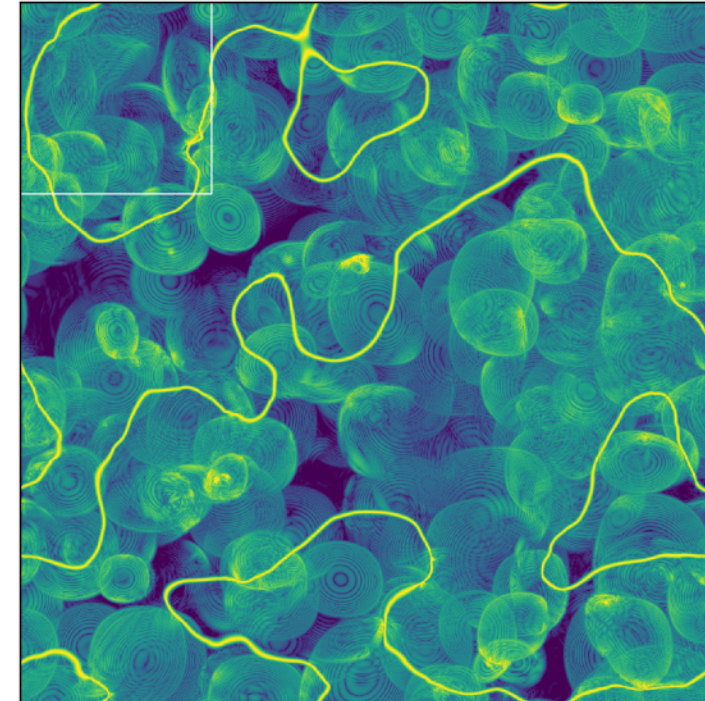
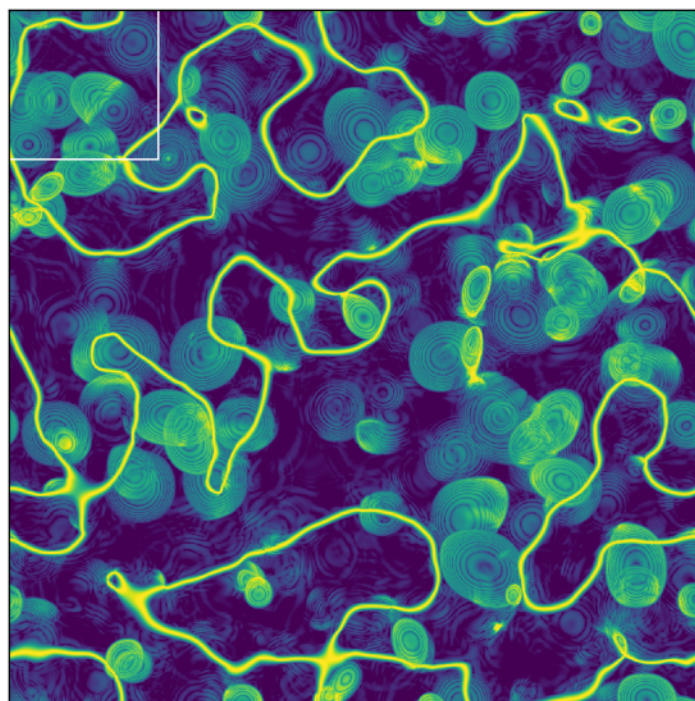
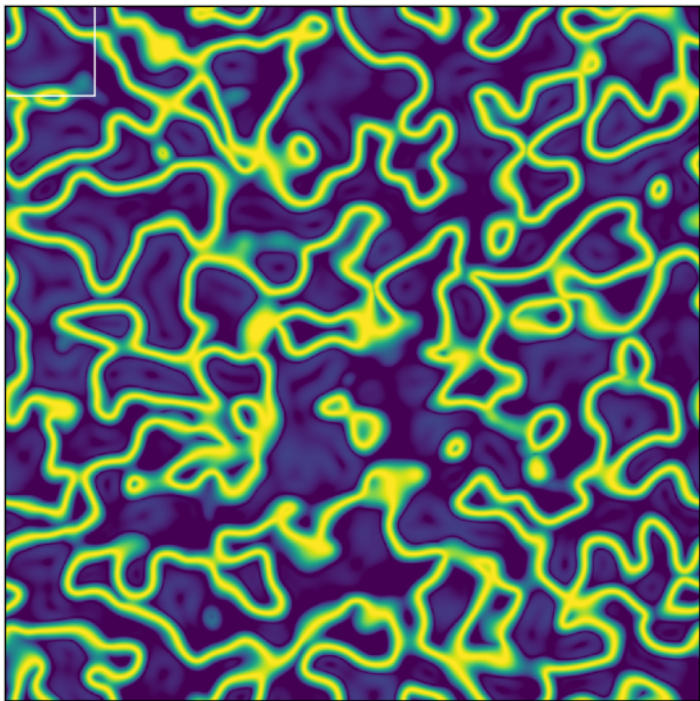
$L = 70H_i^{-1}$

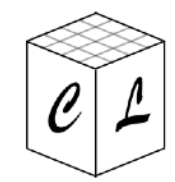


$N = 3060$

$L = 70H_i^{-1}$

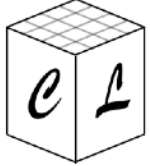






Domain wall evolution: scaling

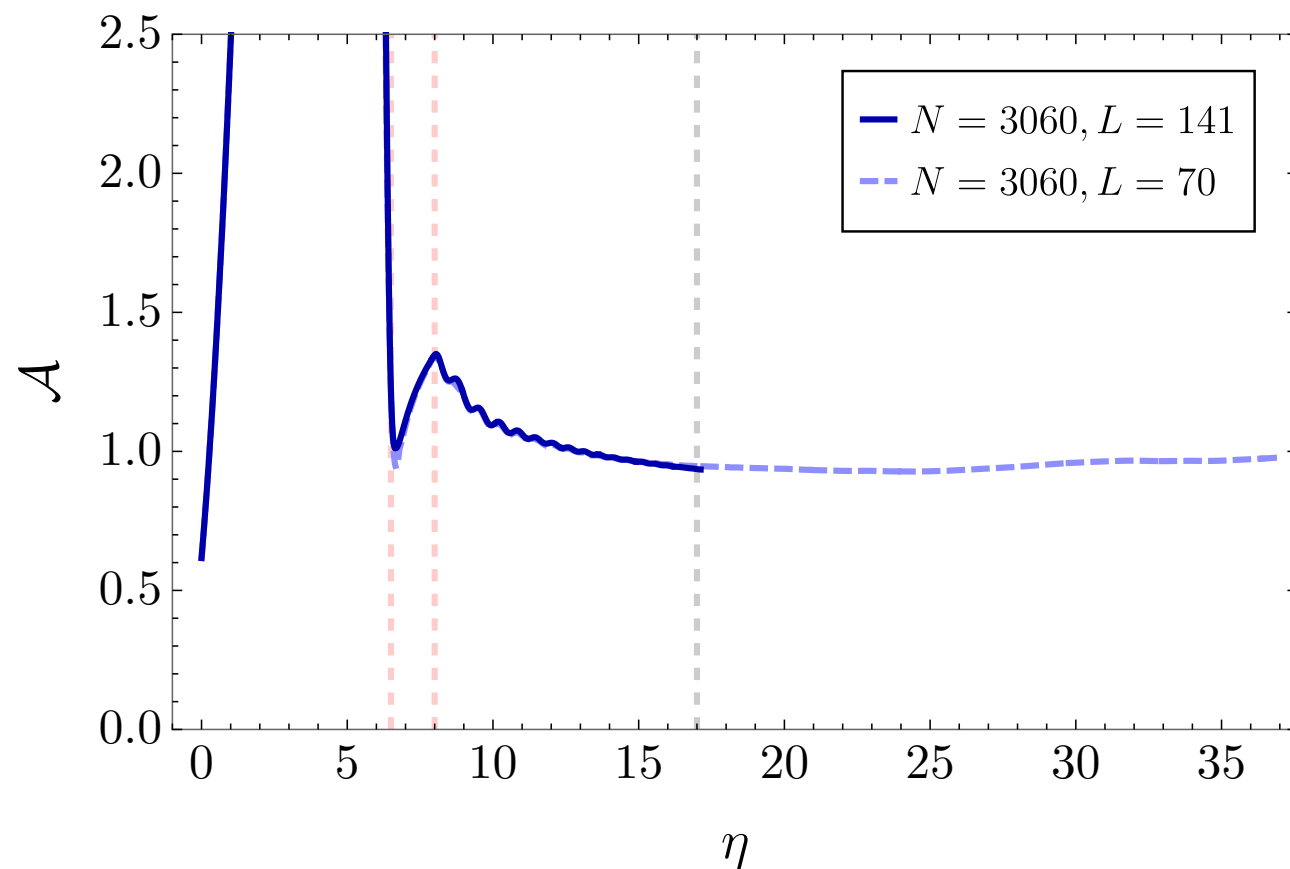
$$[H_i = m = 1]$$



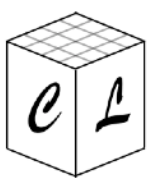
Domain wall evolution: scaling

$$[H_i = m = 1]$$

Area parameter



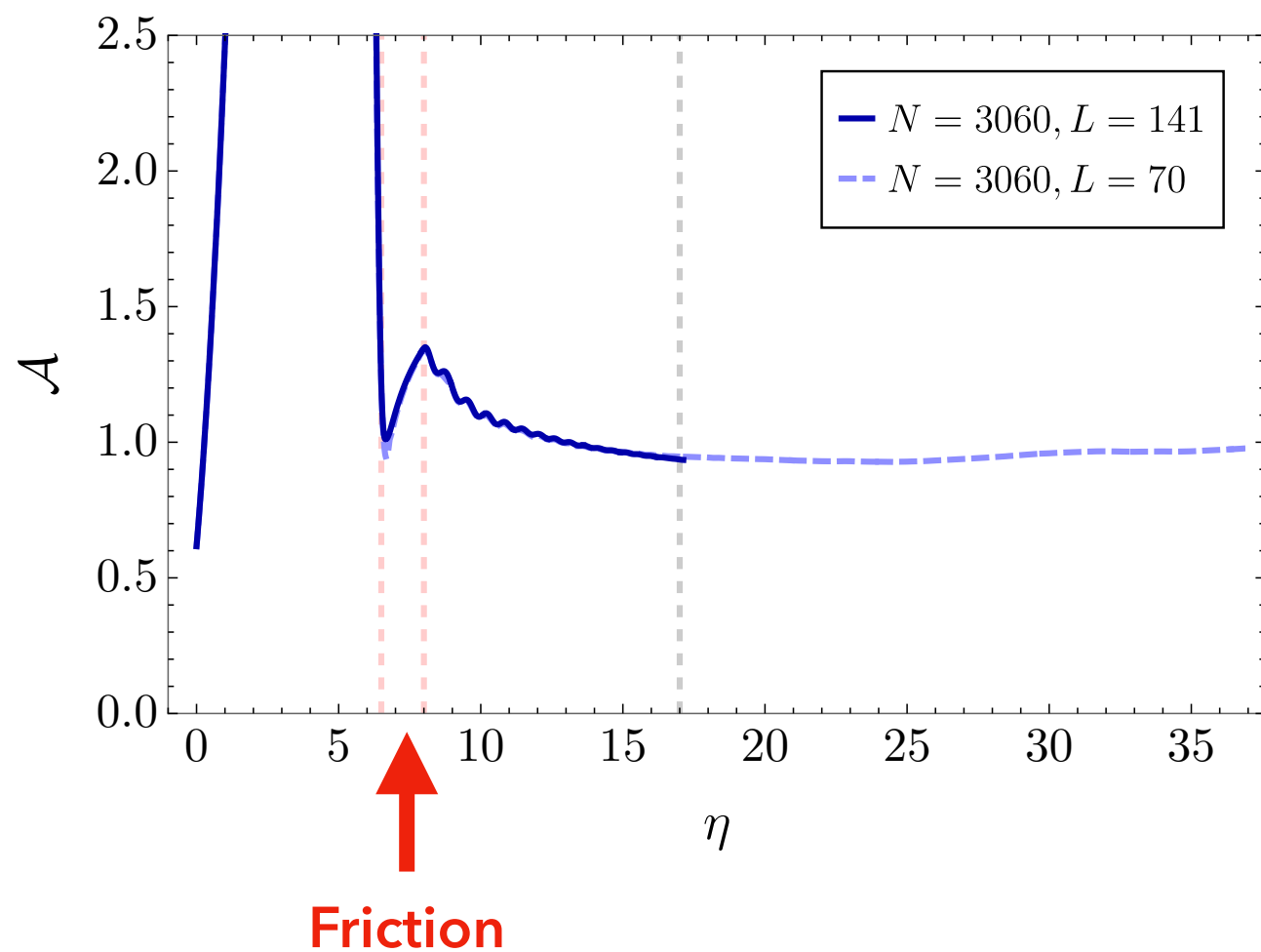
Notari, Rompineve & F.T (JCAP, 2025)



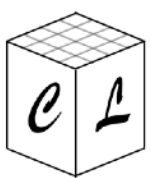
Domain wall evolution: scaling

$$[H_i = m = 1]$$

Area parameter



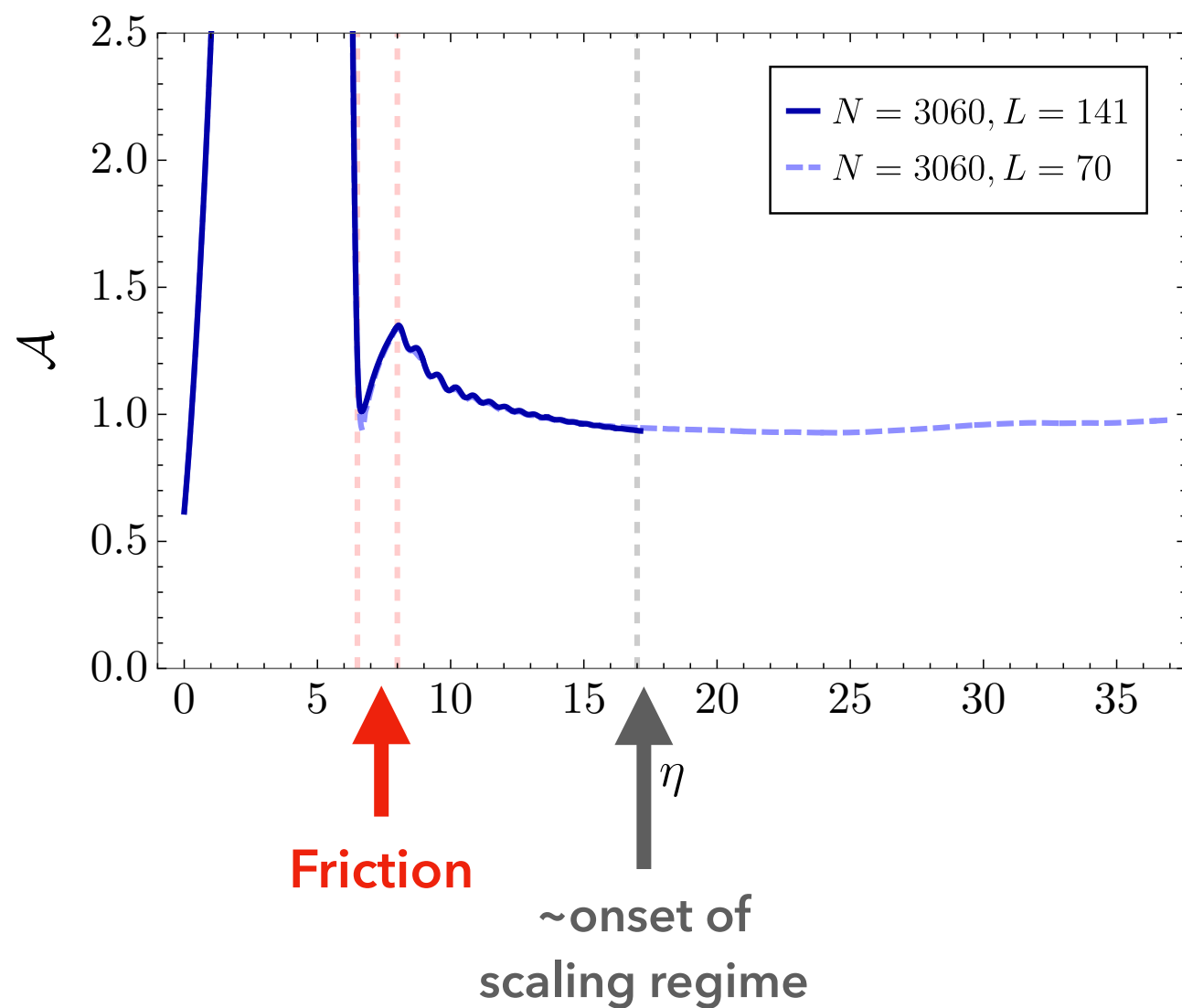
Notari, Rompineve & F.T (JCAP, 2025)



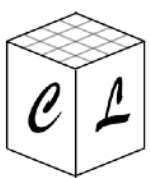
Domain wall evolution: scaling

$$[H_i = m = 1]$$

Area parameter



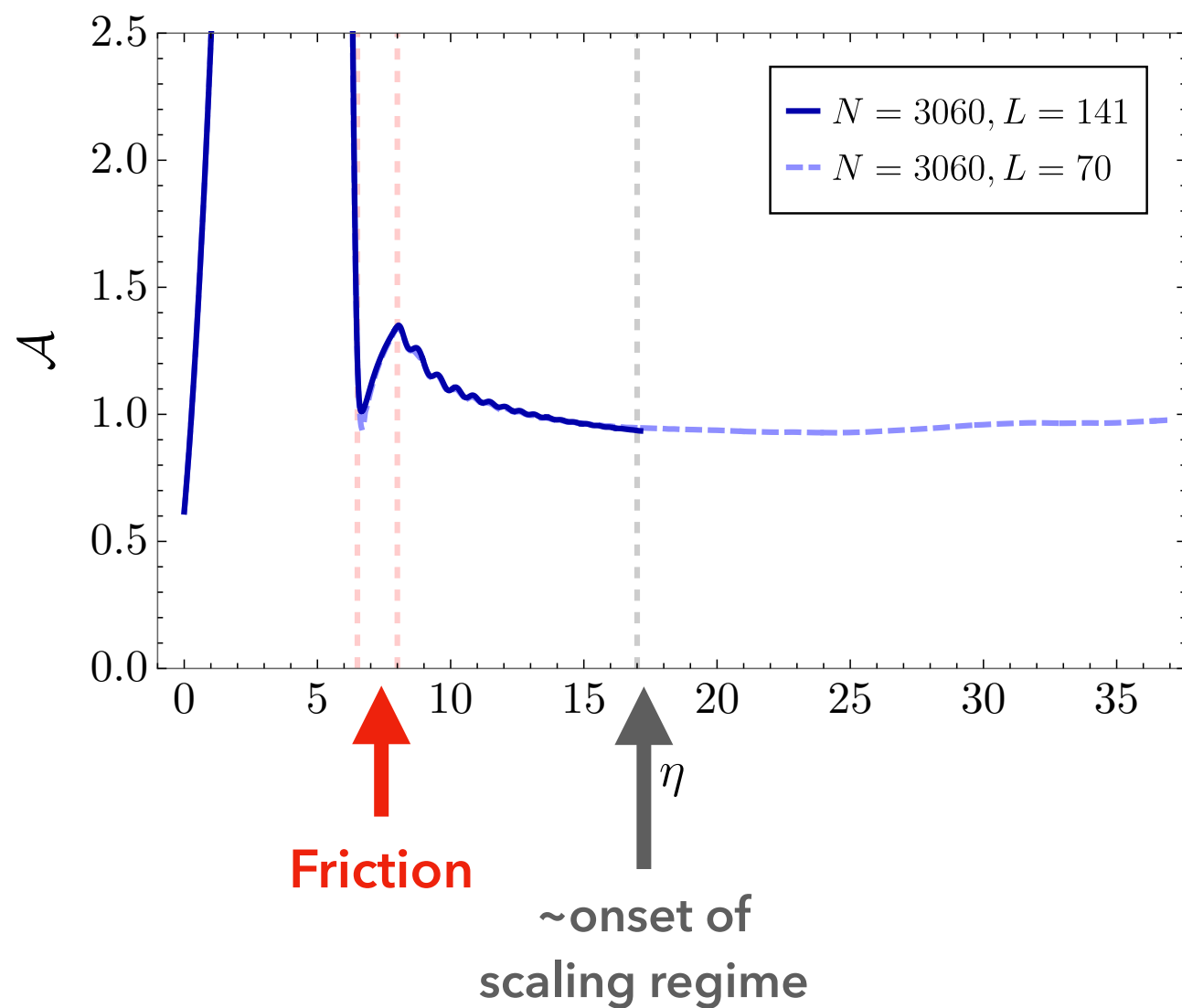
Notari, Rompineve & F.T (JCAP, 2025)



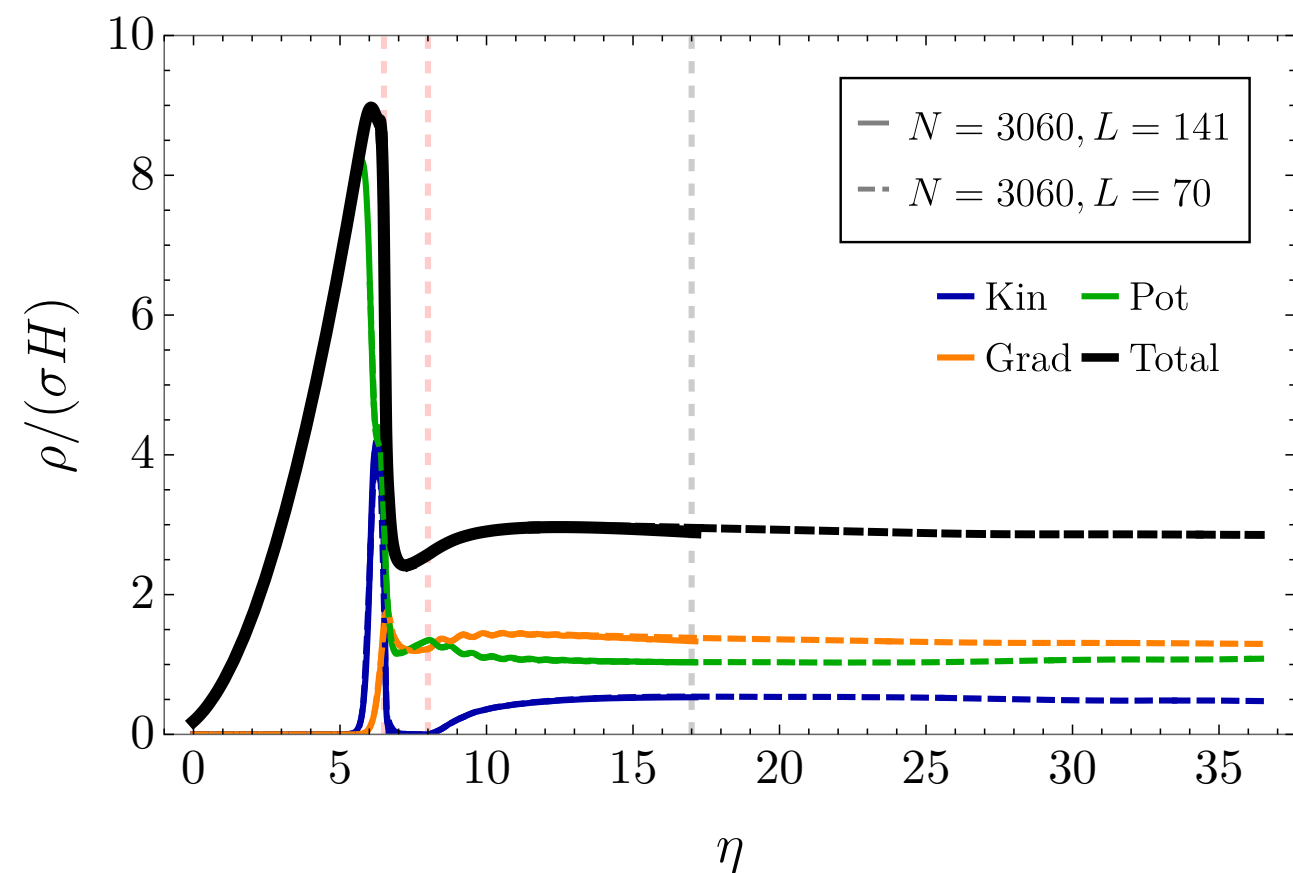
Domain wall evolution: scaling

$$[H_i = m = 1]$$

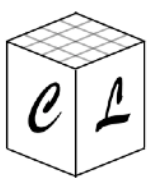
Area parameter



Energy density



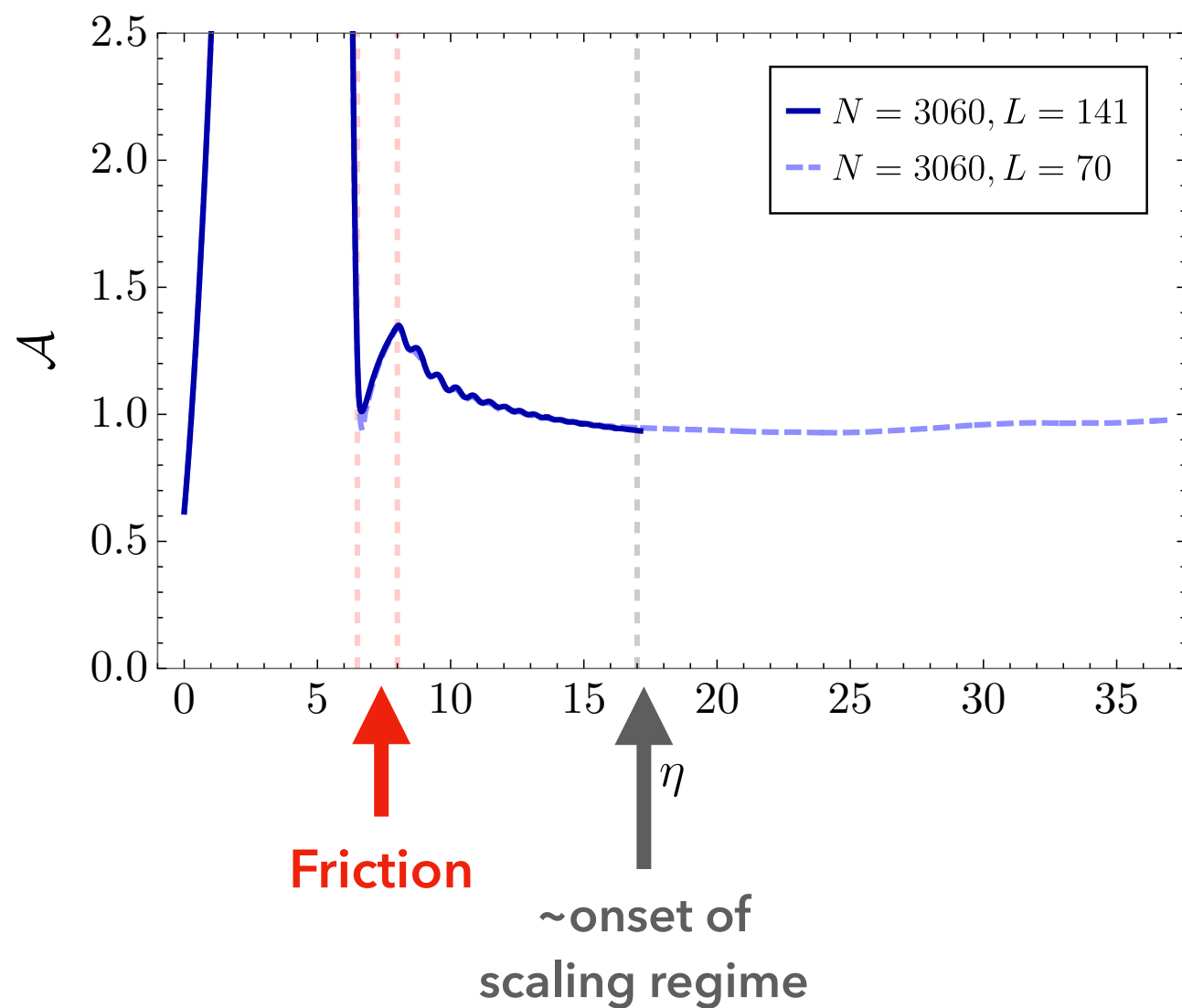
Notari, Rompineve & F.T (JCAP, 2025)



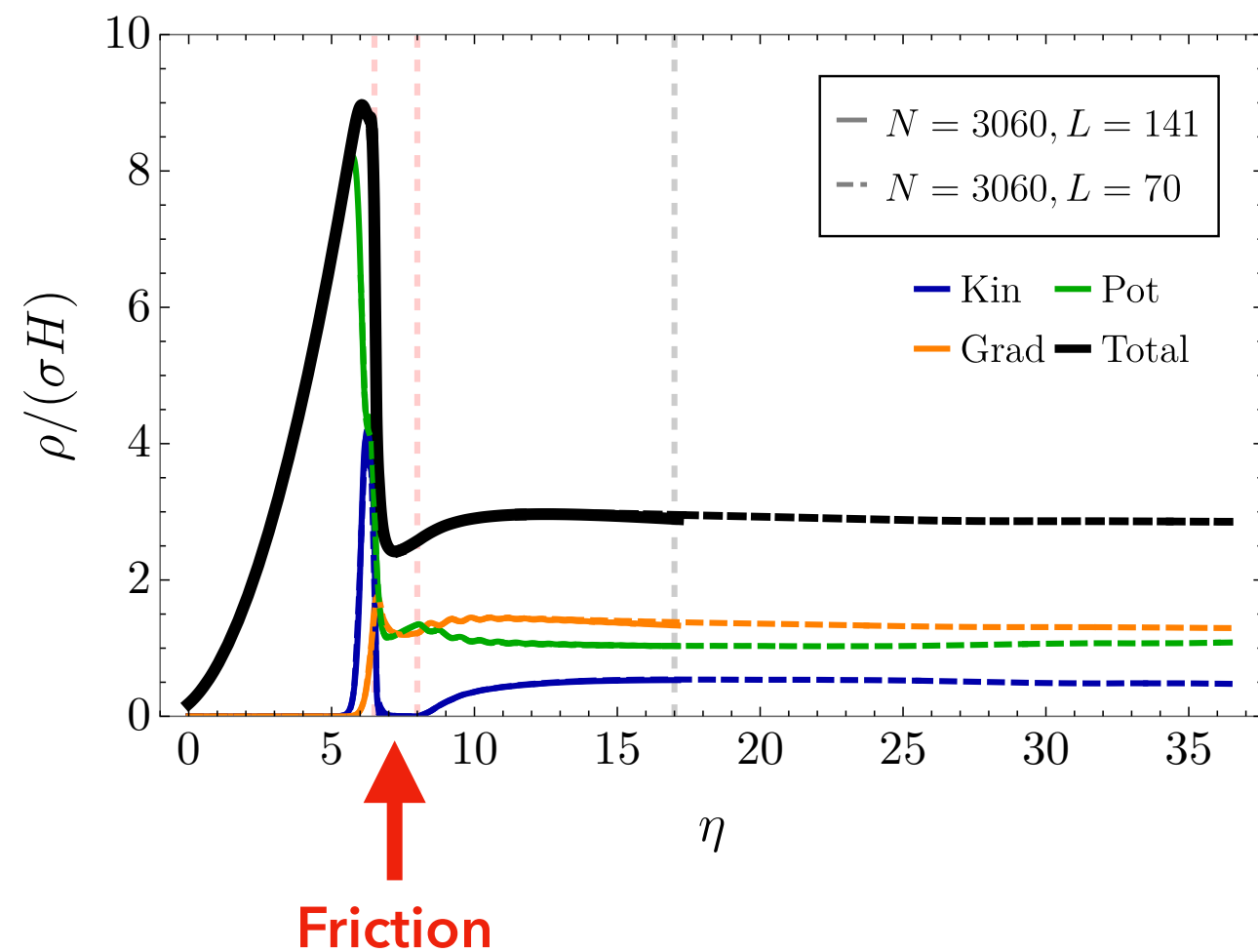
Domain wall evolution: scaling

$$[H_i = m = 1]$$

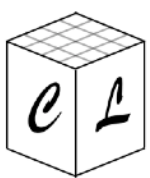
Area parameter



Energy density



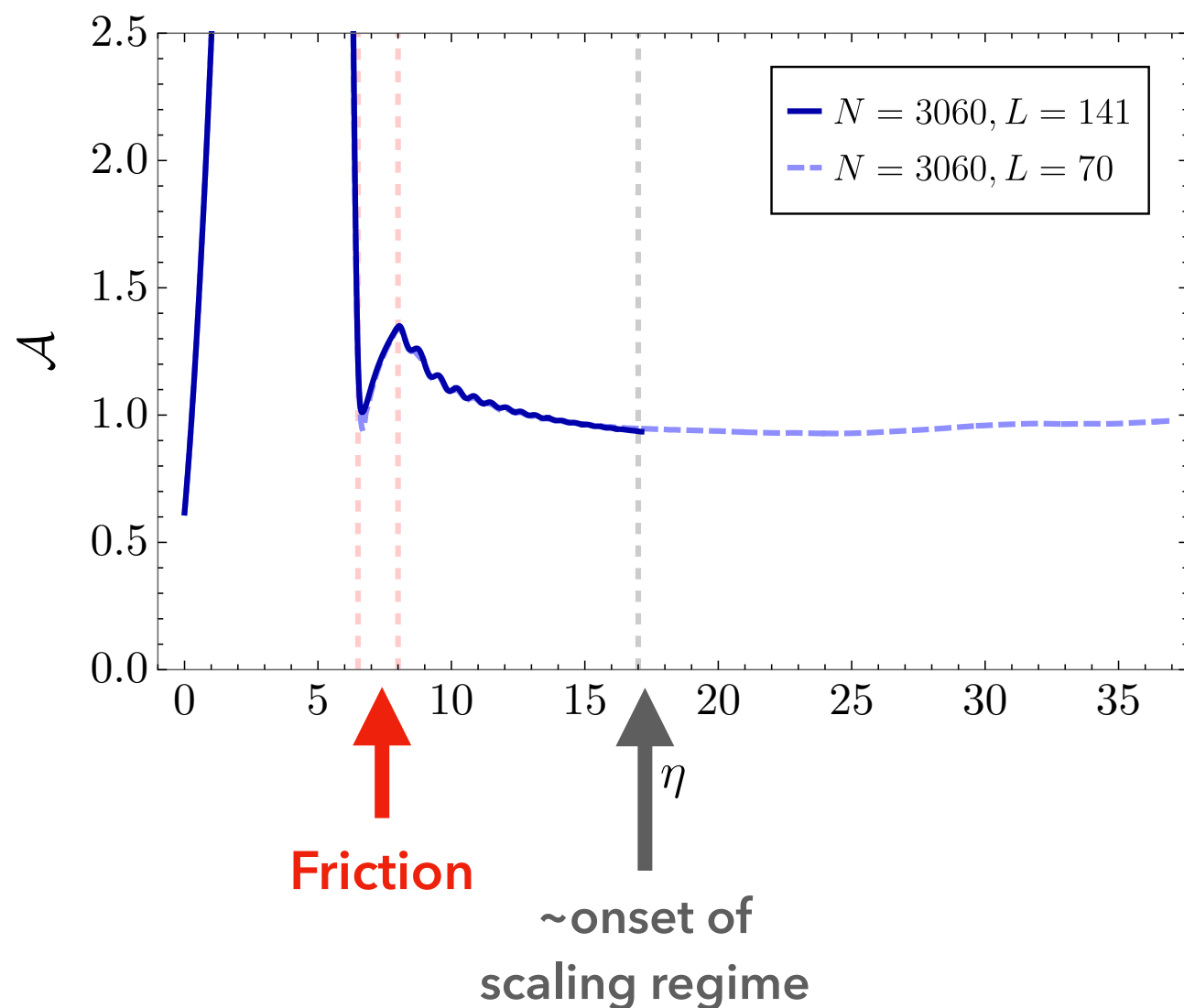
Notari, Rompineve & F.T (JCAP, 2025)



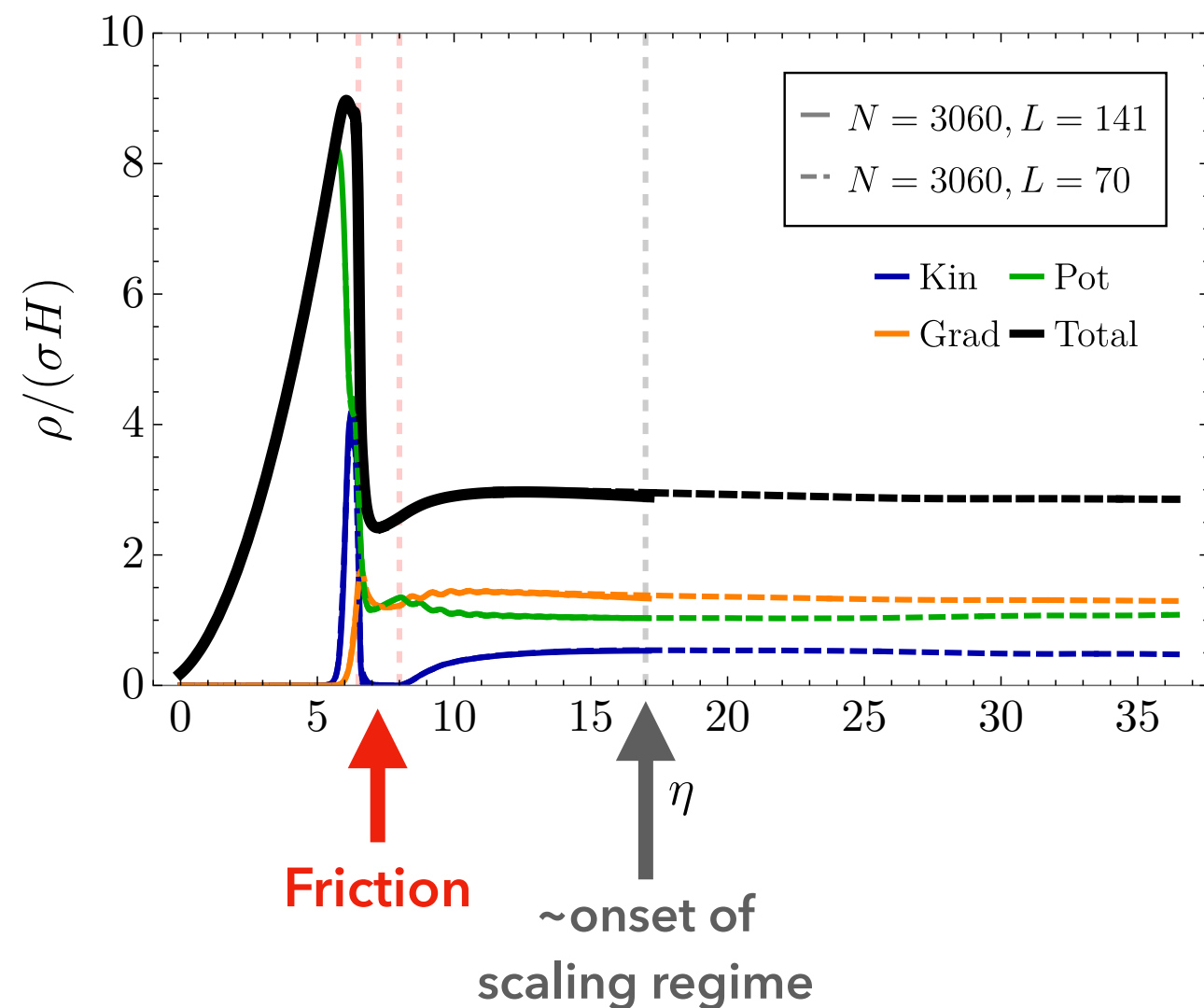
Domain wall evolution: scaling

$$[H_i = m = 1]$$

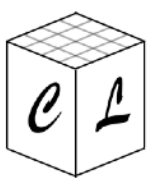
Area parameter



Energy density



Notari, Rompineve & F.T (JCAP, 2025)

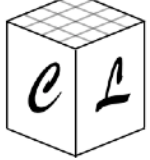


Domain wall evolution: scaling

$$\phi'' - \nabla^2 \phi + (2\mathcal{H} + \Gamma(a))\phi' = -\lambda a^2(\phi^2 - v^2)\phi$$

$$\Gamma(a) \equiv \frac{\Gamma_0/m}{(1 + e^{\gamma(a(\eta_1) - a)})(1 + e^{\gamma(a - a(\eta_2))})}$$

Friction term

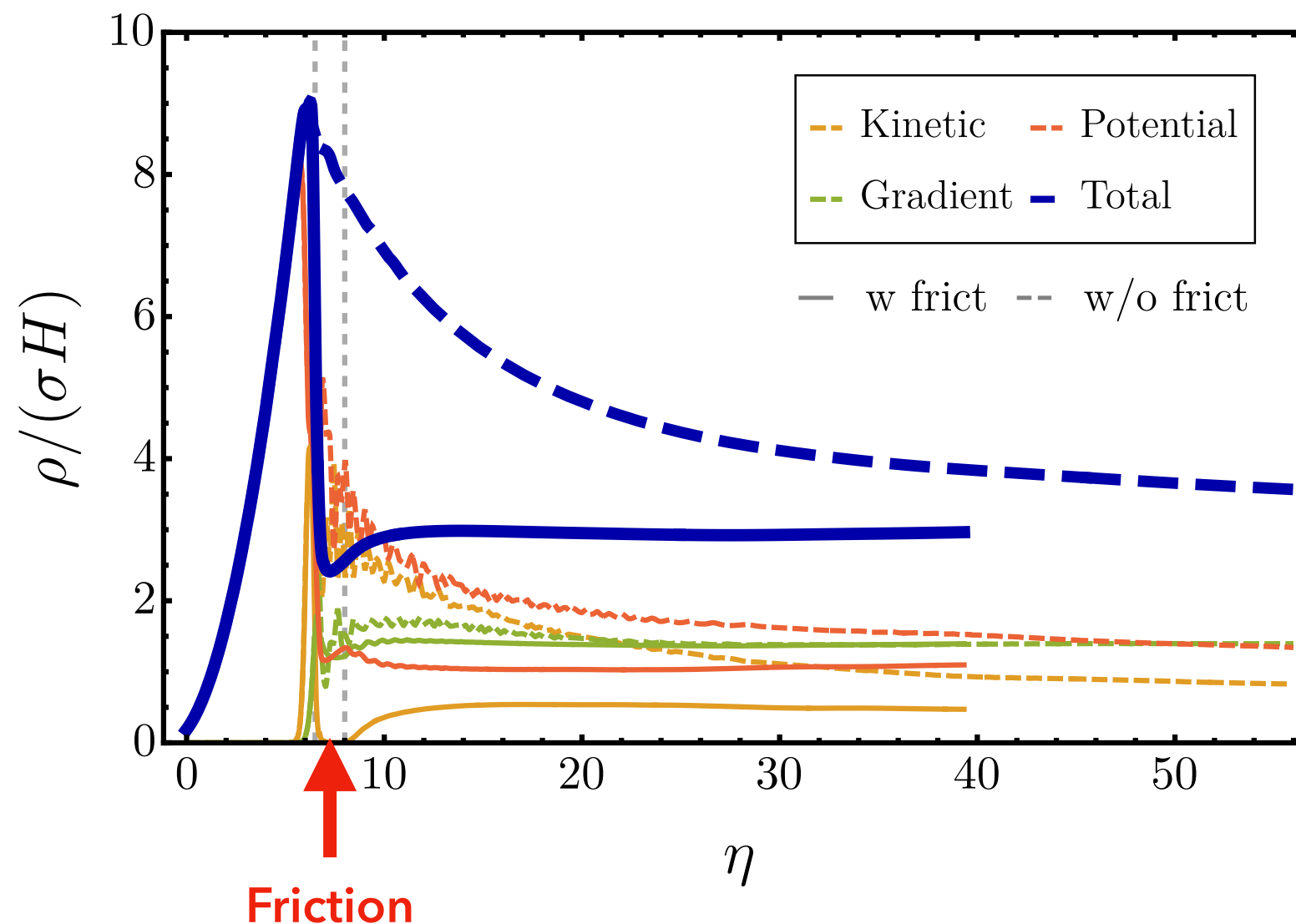


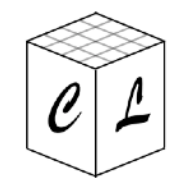
Domain wall evolution: scaling

$$\phi'' - \nabla^2 \phi + (2\mathcal{H} + \Gamma(a))\phi' = -\lambda a^2(\phi^2 - v^2)\phi$$

$$\Gamma(a) \equiv \frac{\Gamma_0/m}{(1 + e^{\gamma(a(\eta_1)-a)})(1 + e^{\gamma(a-a(\eta_2))})}$$

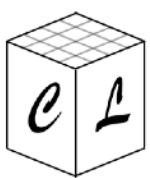
Friction term





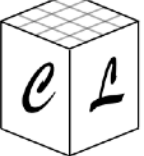
GWs from domain walls (scaling)

$$h''_{ij} - \nabla^2 h_{ij} + 2\mathcal{H}h'_{ij} = \frac{2}{m_p^2}(\partial_i\phi\partial_j\phi)^{\text{TT}}$$



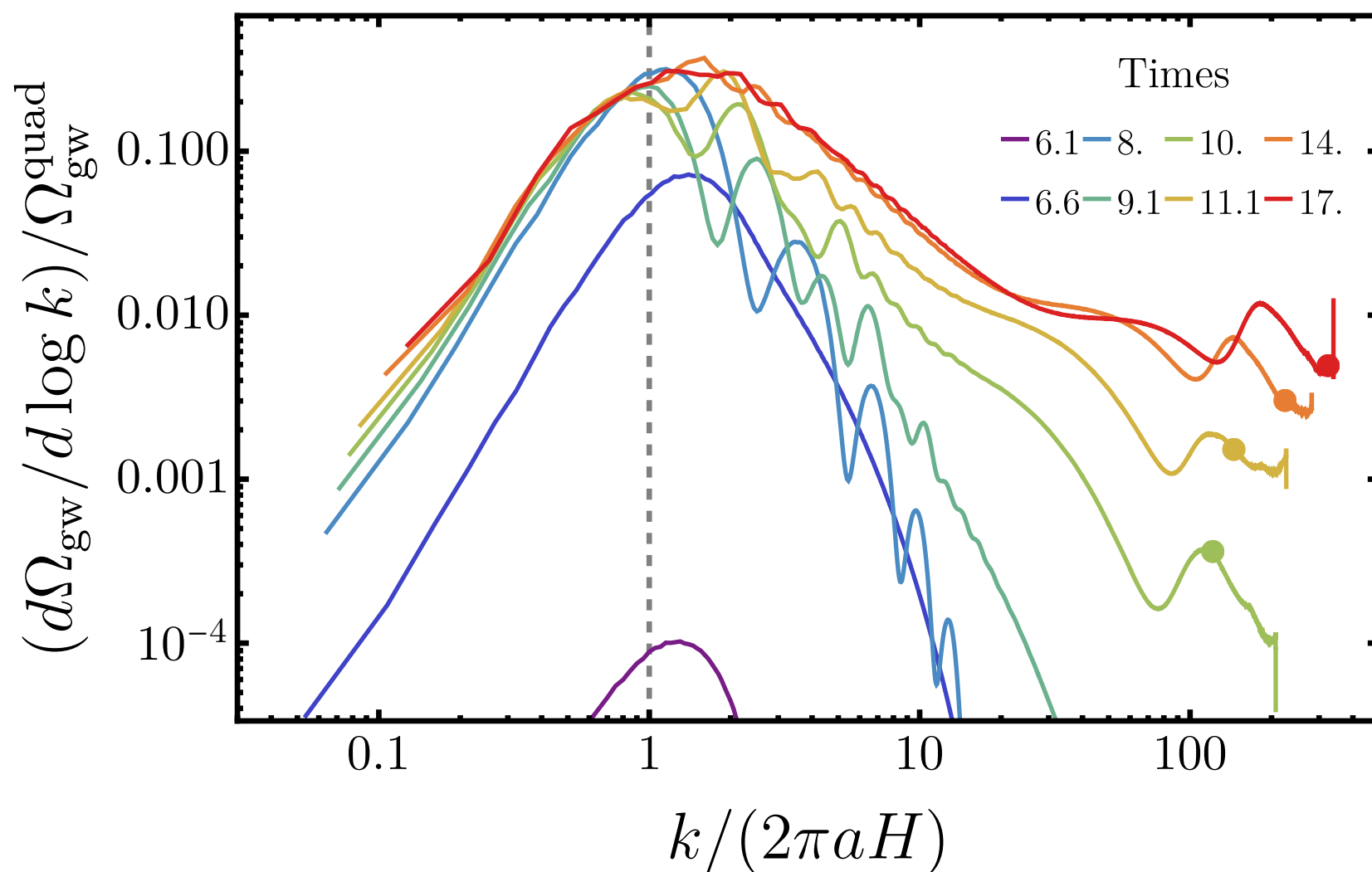
GWs from domain walls (scaling)

$$h''_{ij} - \nabla^2 h_{ij} + 2\mathcal{H}h'_{ij} = \frac{2}{m_p^2}(\partial_i\phi\partial_j\phi)^{\text{TT}} \quad \longrightarrow \quad \rho_{\text{gw}} \simeq \frac{m_p^2}{4a^2}\langle h'_{ij}h'_{ij}\rangle_V; \quad \boxed{\frac{d\rho_{\text{gw}}}{d\log k} = \frac{m_p^2 k^3}{8\pi^2 a^2 V} \int \frac{d\Omega_k}{4\pi} h'_{ij}(\mathbf{k}, \eta) h'^*_{ij}(\mathbf{k}, \eta)}$$



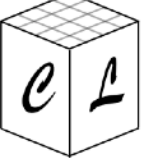
GWs from domain walls (scaling)

$$h''_{ij} - \nabla^2 h_{ij} + 2\mathcal{H}h'_{ij} = \frac{2}{m_p^2}(\partial_i\phi\partial_j\phi)^{\text{TT}} \quad \longrightarrow \quad \rho_{\text{gw}} \simeq \frac{m_p^2}{4a^2}\langle h'_{ij}h'_{ij} \rangle_V; \quad \boxed{\frac{d\rho_{\text{gw}}}{d\log k} = \frac{m_p^2 k^3}{8\pi^2 a^2 V} \int \frac{d\Omega_k}{4\pi} h'_{ij}(\mathbf{k}, \eta) h'^*_{ij}(\mathbf{k}, \eta)}$$



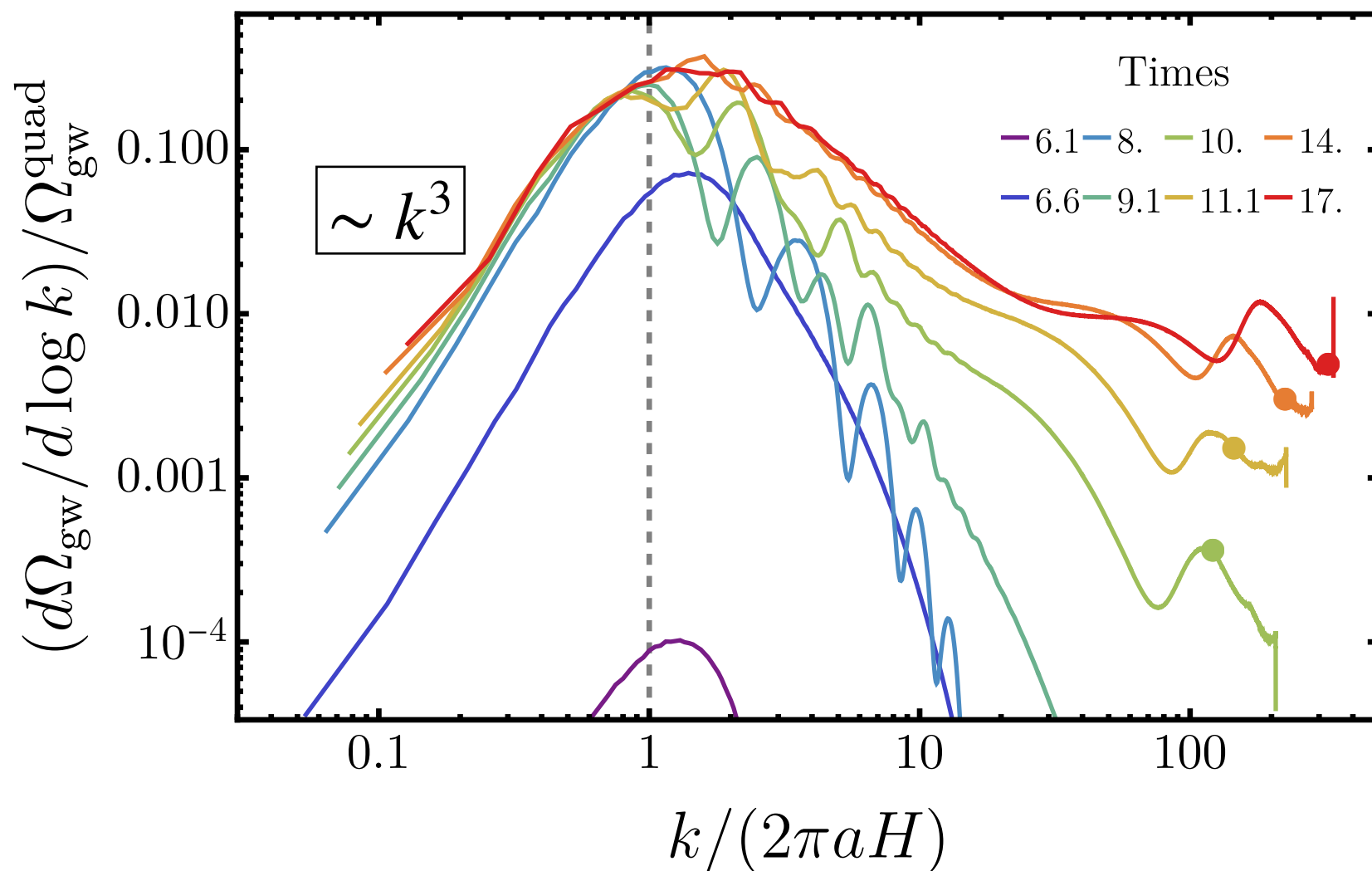
Notari, Rompineve & F.T (JCAP, 2025)

$$\Omega_{\text{gw}}^{\text{quad}} \equiv 3/(32\pi)[2\sigma H/\rho_c^2]$$



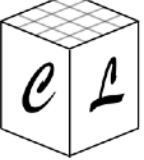
GWs from domain walls (scaling)

$$h''_{ij} - \nabla^2 h_{ij} + 2\mathcal{H}h'_{ij} = \frac{2}{m_p^2}(\partial_i\phi\partial_j\phi)^{\text{TT}} \quad \longrightarrow \quad \rho_{\text{gw}} \simeq \frac{m_p^2}{4a^2}\langle h'_{ij}h'_{ij} \rangle_V; \quad \boxed{\frac{d\rho_{\text{gw}}}{d\log k} = \frac{m_p^2 k^3}{8\pi^2 a^2 V} \int \frac{d\Omega_k}{4\pi} h'_{ij}(\mathbf{k}, \eta) h'^*_{ij}(\mathbf{k}, \eta)}$$



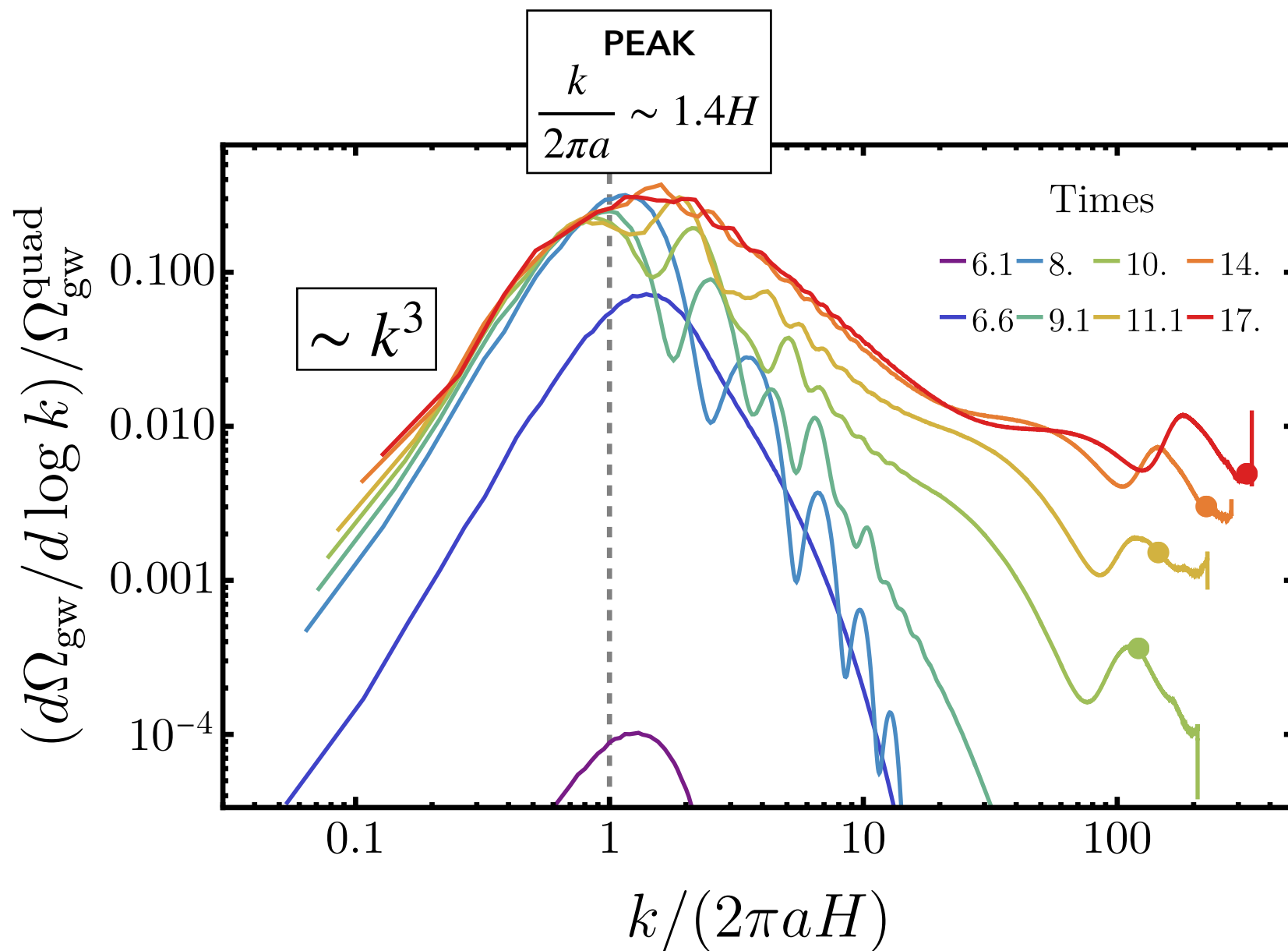
Notari, Rompineve & F.T (JCAP, 2025)

$$\Omega_{\text{gw}}^{\text{quad}} \equiv 3/(32\pi)[2\sigma H/\rho_c^2]$$



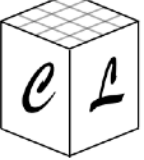
GWs from domain walls (scaling)

$$h''_{ij} - \nabla^2 h_{ij} + 2\mathcal{H}h'_{ij} = \frac{2}{m_p^2}(\partial_i\phi\partial_j\phi)^{\text{TT}} \quad \longrightarrow \quad \rho_{\text{gw}} \simeq \frac{m_p^2}{4a^2}\langle h'_{ij}h'_{ij} \rangle_V; \quad \boxed{\frac{d\rho_{\text{gw}}}{d\log k} = \frac{m_p^2 k^3}{8\pi^2 a^2 V} \int \frac{d\Omega_k}{4\pi} h'_{ij}(\mathbf{k}, \eta) h'^*_{ij}(\mathbf{k}, \eta)}$$



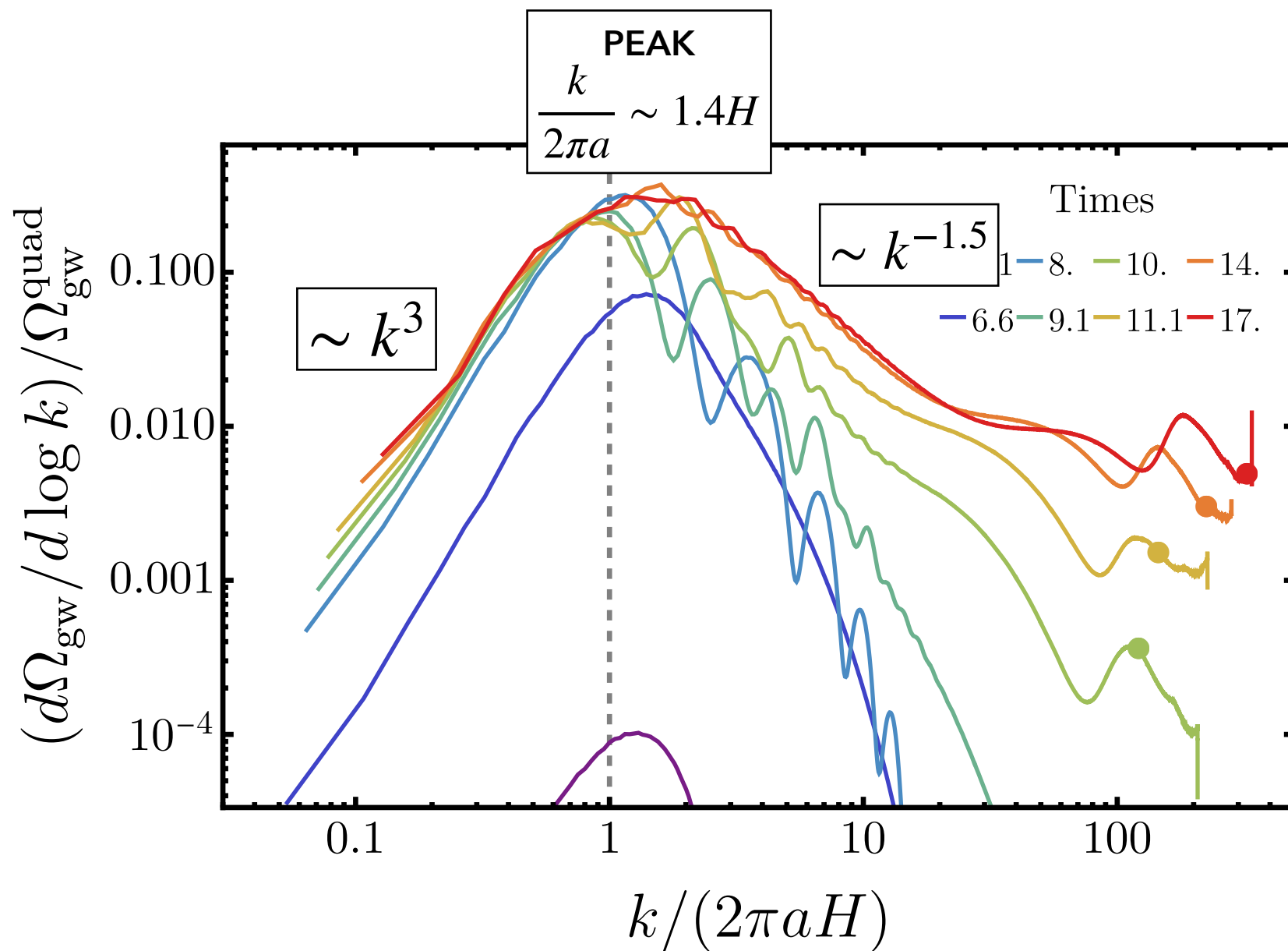
Notari, Rompineve & F.T (JCAP, 2025)

$$\Omega_{\text{gw}}^{\text{quad}} \equiv 3/(32\pi)[2\sigma H/\rho_c^2]$$



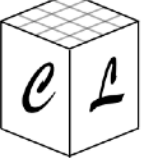
GWs from domain walls (scaling)

$$h''_{ij} - \nabla^2 h_{ij} + 2\mathcal{H}h'_{ij} = \frac{2}{m_p^2}(\partial_i\phi\partial_j\phi)^{\text{TT}} \quad \longrightarrow \quad \rho_{\text{gw}} \simeq \frac{m_p^2}{4a^2}\langle h'_{ij}h'_{ij} \rangle_V; \quad \boxed{\frac{d\rho_{\text{gw}}}{d\log k} = \frac{m_p^2 k^3}{8\pi^2 a^2 V} \int \frac{d\Omega_k}{4\pi} h'_{ij}(\mathbf{k}, \eta) h'^*_{ij}(\mathbf{k}, \eta)}$$



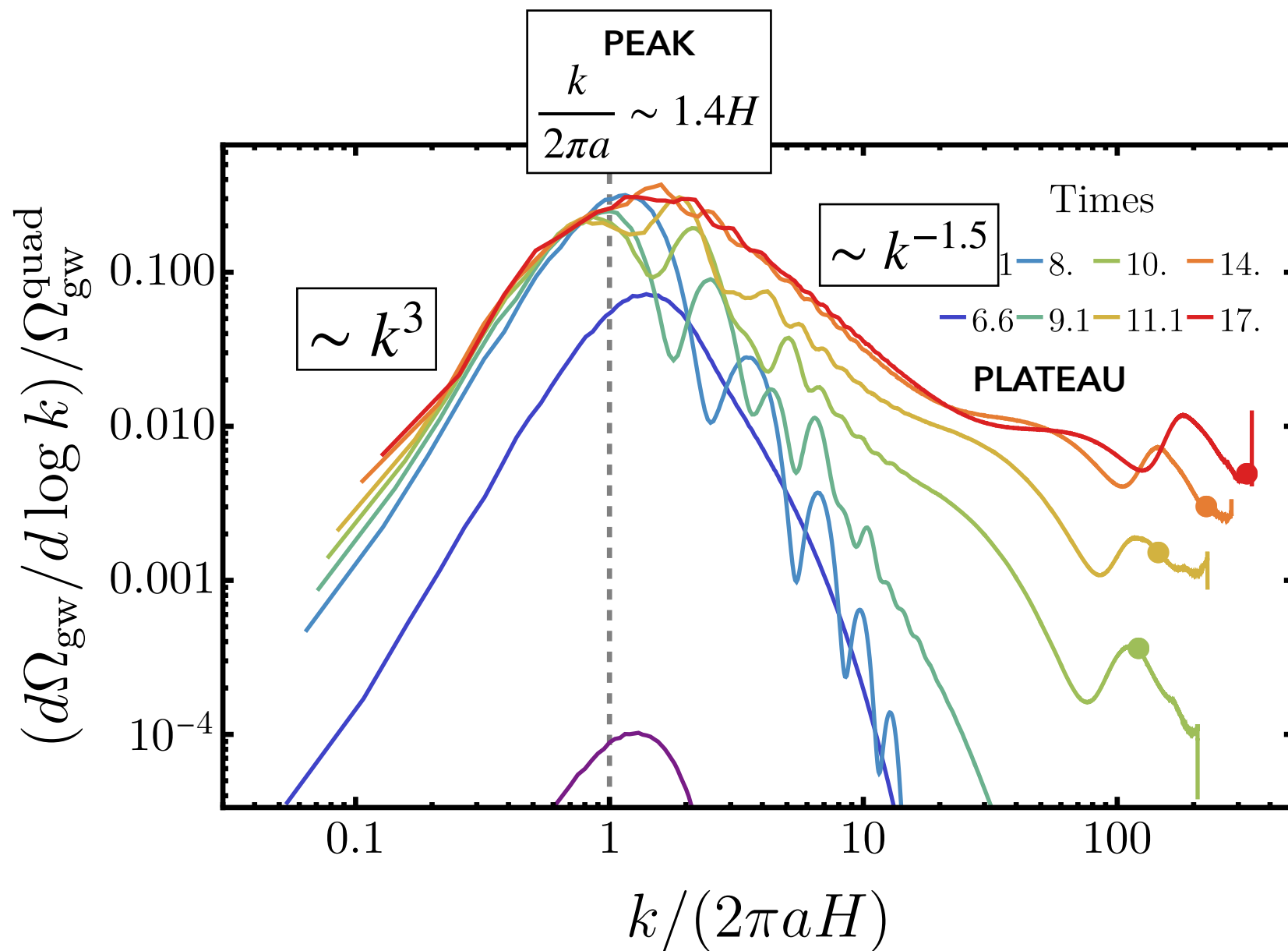
Notari, Rompineve & F.T (JCAP, 2025)

$$\Omega_{\text{gw}}^{\text{quad}} \equiv 3/(32\pi)[2\sigma H/\rho_c^2]$$



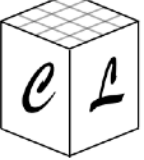
GWs from domain walls (scaling)

$$h''_{ij} - \nabla^2 h_{ij} + 2\mathcal{H}h'_{ij} = \frac{2}{m_p^2}(\partial_i\phi\partial_j\phi)^{\text{TT}} \quad \longrightarrow \quad \rho_{\text{gw}} \simeq \frac{m_p^2}{4a^2}\langle h'_{ij}h'_{ij} \rangle_V; \quad \boxed{\frac{d\rho_{\text{gw}}}{d\log k} = \frac{m_p^2 k^3}{8\pi^2 a^2 V} \int \frac{d\Omega_k}{4\pi} h'_{ij}(\mathbf{k}, \eta) h'^*_{ij}(\mathbf{k}, \eta)}$$



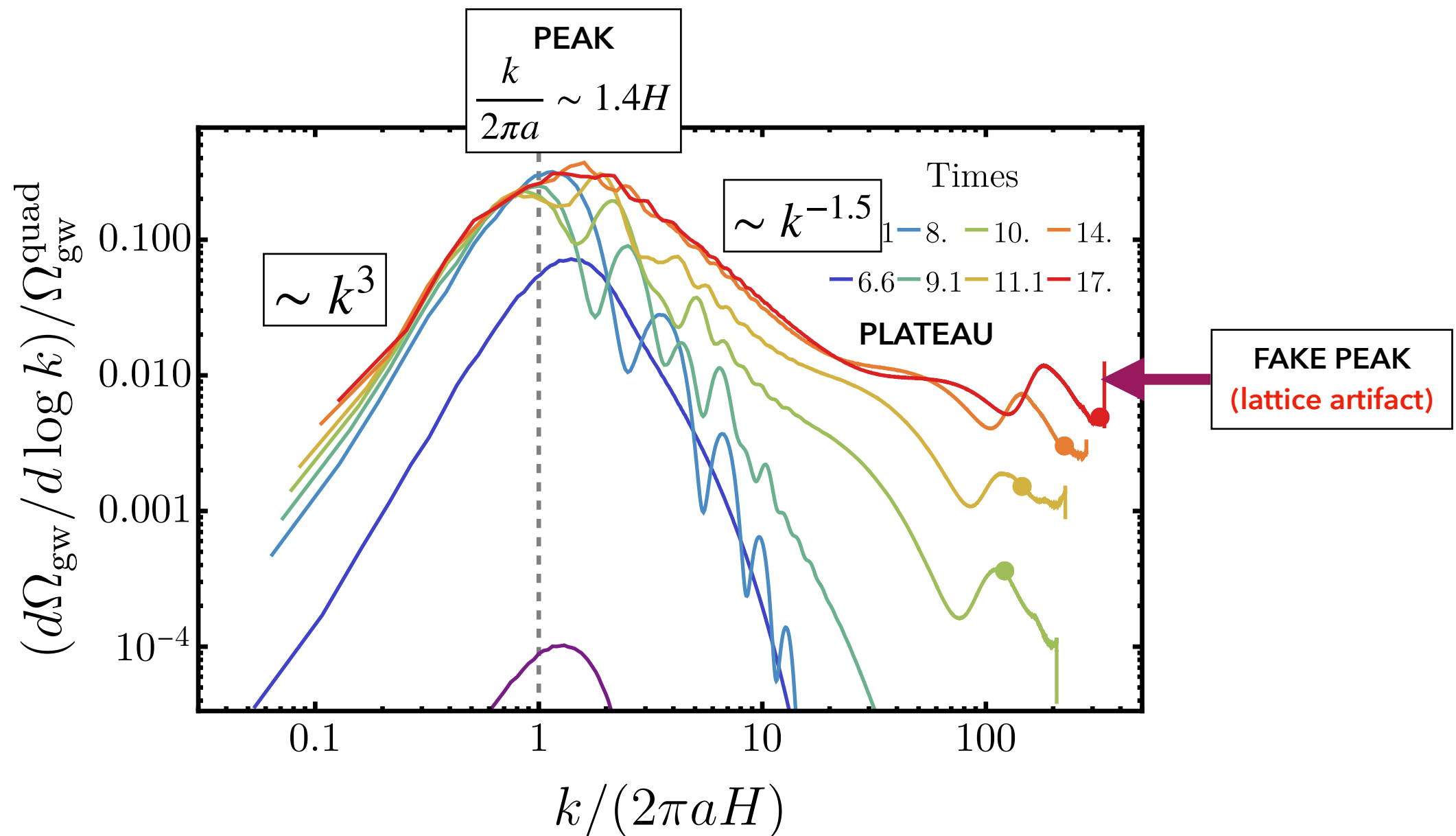
Notari, Rompineve & F.T (JCAP, 2025)

$$\Omega_{\text{gw}}^{\text{quad}} \equiv 3/(32\pi)[2\sigma H/\rho_c^2]$$



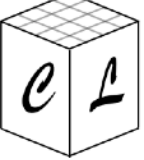
GWs from domain walls (scaling)

$$h''_{ij} - \nabla^2 h_{ij} + 2\mathcal{H}h'_{ij} = \frac{2}{m_p^2}(\partial_i\phi\partial_j\phi)^{\text{TT}} \quad \longrightarrow \quad \rho_{\text{gw}} \simeq \frac{m_p^2}{4a^2}\langle h'_{ij}h'_{ij} \rangle_V; \quad \boxed{\frac{d\rho_{\text{gw}}}{d\log k} = \frac{m_p^2 k^3}{8\pi^2 a^2 V} \int \frac{d\Omega_k}{4\pi} h'_{ij}(\mathbf{k}, \eta) h'^*_{ij}(\mathbf{k}, \eta)}$$



Notari, Rompineve & F.T (JCAP, 2025)

$$\Omega_{\text{gw}}^{\text{quad}} \equiv 3/(32\pi)[2\sigma H/\rho_c^2]$$



GWs from domain walls (scaling)

$$[\nabla_i f]^{(2)} \equiv \frac{f_{\mathbf{n}+\hat{i}} - f_{\mathbf{n}-\hat{i}}}{2 \Delta x} \longrightarrow \partial_i f(\mathbf{x}) + \mathcal{O}(\Delta x^2)$$

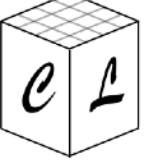
$$[\nabla_i^2 f]^{(4)} \equiv \frac{f_{\mathbf{n}+\hat{i}} - 2f_{\mathbf{n}} + f_{\mathbf{n}-\hat{i}}}{2 \Delta x^2} \longrightarrow \partial_i^2 f(\mathbf{x}) + \mathcal{O}(\Delta x^2)$$

ORDER 2

$$[\nabla_i f]^{(4)} \equiv \frac{f_{\mathbf{n}-2\hat{i}} - 8f_{\mathbf{n}-\hat{i}} + 8f_{\mathbf{n}+\hat{i}} - f_{\mathbf{n}+2\hat{i}}}{12 \Delta x} \longrightarrow \partial_i f(\mathbf{x}) + \mathcal{O}(\Delta x^4)$$

$$[\nabla_i^2 f]^{(4)} \equiv \frac{-f_{\mathbf{n}+2\hat{i}} + 16f_{\mathbf{n}+\hat{i}} - 30f_{\mathbf{n}} + 16f_{\mathbf{n}-\hat{i}} - f_{\mathbf{n}-2\hat{i}}}{12 \Delta x^2} \longrightarrow \partial_i^2 f(\mathbf{x}) + \mathcal{O}(\Delta x^4)$$

ORDER 4



GWs from domain walls (scaling)

$$[\nabla_i f]^{(2)} \equiv \frac{f_{\mathbf{n}+\hat{i}} - f_{\mathbf{n}-\hat{i}}}{2 \Delta x} \longrightarrow \partial_i f(\mathbf{x}) + \mathcal{O}(\Delta x^2)$$

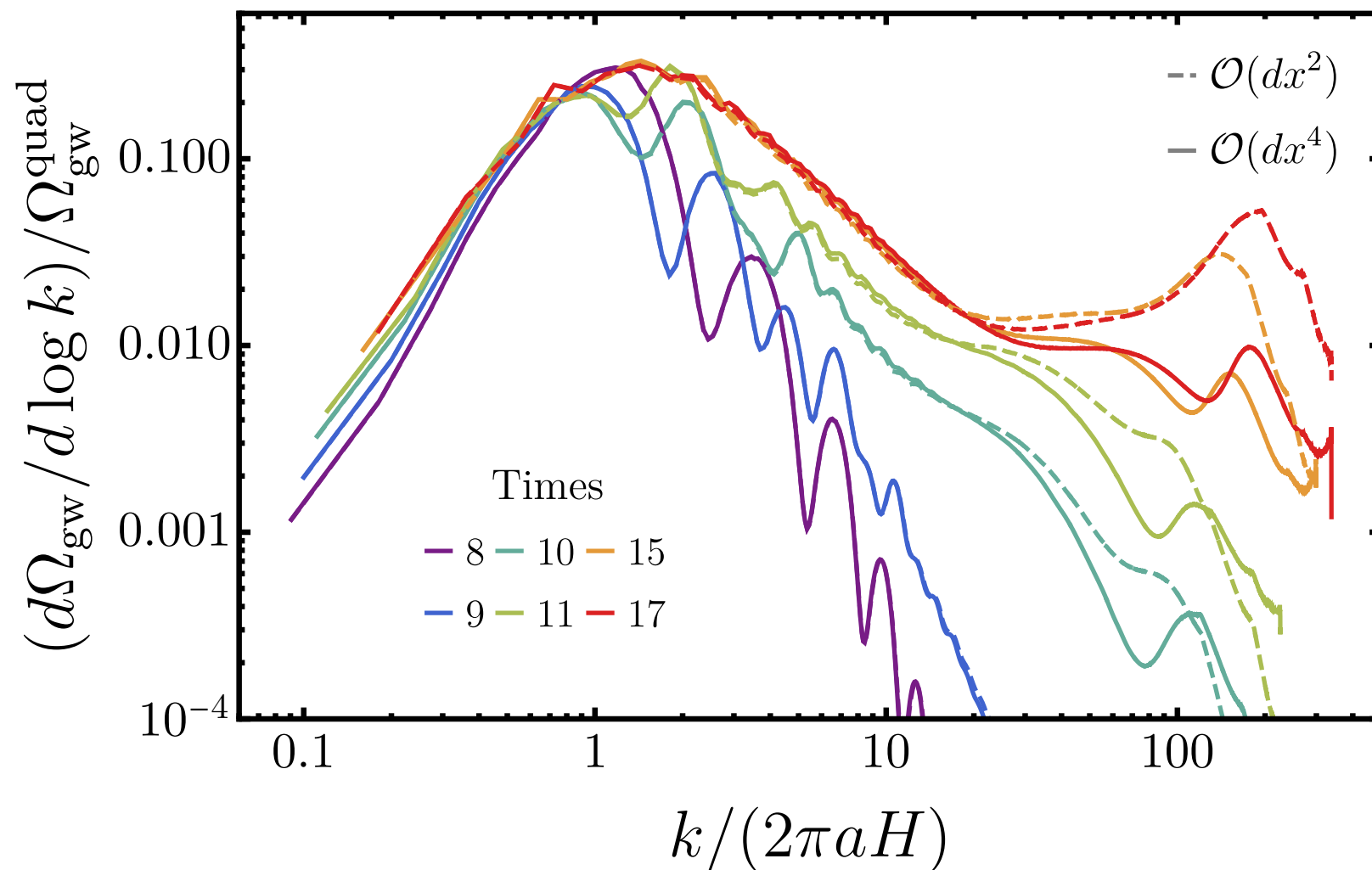
$$[\nabla_i^2 f]^{(4)} \equiv \frac{f_{\mathbf{n}+\hat{i}} - 2f_{\mathbf{n}} + f_{\mathbf{n}-\hat{i}}}{2 \Delta x^2} \longrightarrow \partial_i^2 f(\mathbf{x}) + \mathcal{O}(\Delta x^2)$$

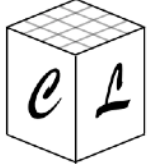
ORDER 2

$$[\nabla_i f]^{(4)} \equiv \frac{f_{\mathbf{n}-2\hat{i}} - 8f_{\mathbf{n}-\hat{i}} + 8f_{\mathbf{n}+\hat{i}} - f_{\mathbf{n}+2\hat{i}}}{12 \Delta x} \longrightarrow \partial_i f(\mathbf{x}) + \mathcal{O}(\Delta x^4)$$

$$[\nabla_i^2 f]^{(4)} \equiv \frac{-f_{\mathbf{n}+2\hat{i}} + 16f_{\mathbf{n}+\hat{i}} - 30f_{\mathbf{n}} + 16f_{\mathbf{n}-\hat{i}} - f_{\mathbf{n}-2\hat{i}}}{12 \Delta x^2} \longrightarrow \partial_i^2 f(\mathbf{x}) + \mathcal{O}(\Delta x^4)$$

ORDER 4

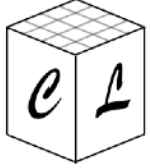




Annihilating domain walls

- DWs eventually dominate the energy density of the universe if $\sigma \gtrsim (\text{MeV})^3$:

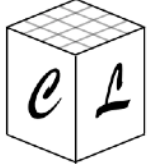
$$\rho_{\text{dw}} \sim \sigma H \sim a^{-2}$$



Annihilating domain walls

- DWs eventually dominate the energy density of the universe if $\sigma \gtrsim (\text{MeV})^3$:

$$\rho_{\text{dw}} \sim \sigma H \sim a^{-2} \rightarrow \text{DW problem!}$$



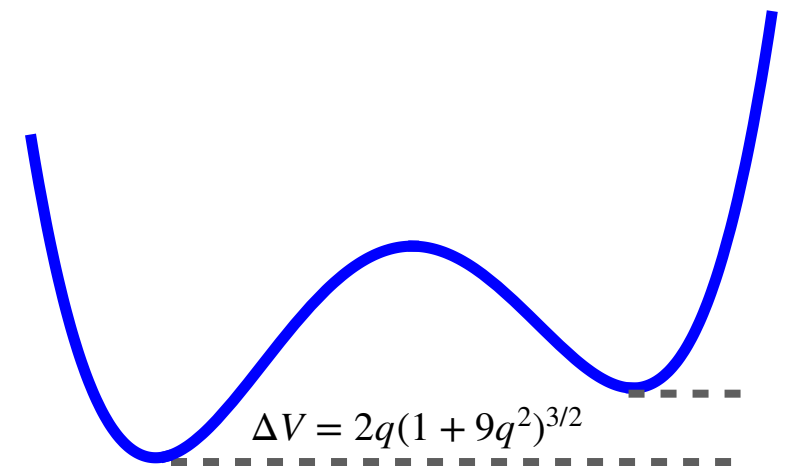
Annihilating domain walls

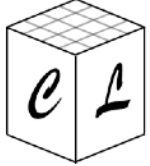
- DWs eventually dominate the energy density of the universe if $\sigma \gtrsim (\text{MeV})^3$:

$$\rho_{\text{dw}} \sim \sigma H \sim a^{-2} \quad \rightarrow \quad \text{DW problem!}$$

- **Biased potential:** ($q \ll 1$)

$$V = V_{\mathbb{Z}_2} + V_{\text{bias}} = \frac{\lambda}{4} (\phi^2 - v^2)^4 + qv\phi^3$$





Annihilating domain walls

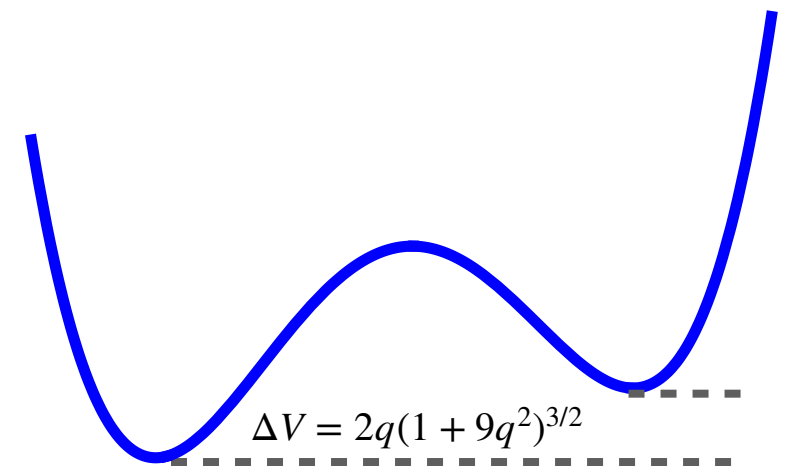
- DWs eventually dominate the energy density of the universe if $\sigma \gtrsim (\text{MeV})^3$:

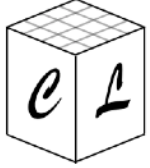
$$\rho_{\text{dw}} \sim \sigma H \sim a^{-2} \quad \rightarrow \quad \text{DW problem!}$$

- **Biased potential:** ($q \ll 1$)

$$V = V_{\mathbb{Z}_2} + V_{\text{bias}} = \frac{\lambda}{4} (\phi^2 - v^2)^4 + qv\phi^3$$

Annihilation starts when: $\sigma H(\eta_{\Delta V}) = \Delta V$





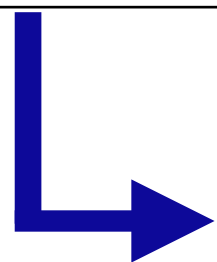
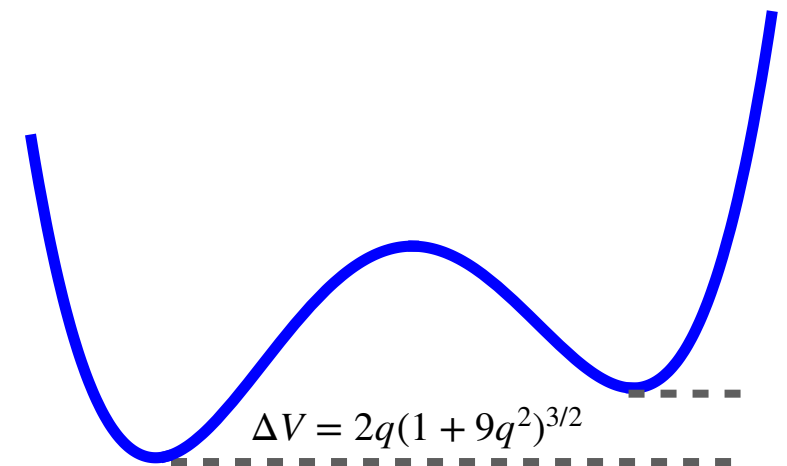
Annihilating domain walls

- DWs eventually dominate the energy density of the universe if $\sigma \gtrsim (\text{MeV})^3$:

$$\rho_{\text{dw}} \sim \sigma H \sim a^{-2} \rightarrow \text{DW problem!}$$

- **Biased potential:** ($q \ll 1$)

$$V = V_{\mathbb{Z}_2} + V_{\text{bias}} = \frac{\lambda}{4} (\phi^2 - v^2)^4 + qv\phi^3$$



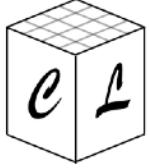
Annihilation starts when: $\sigma H(\eta_{\Delta V}) = \Delta V$

- Recent papers studying GW production from annihilating DWs with simulations:

Cyr, Cotterill, Battye (2025)

Notari, Rompineve & F.T (2025)

Babichev, Dankovsky, Gorbunov, Ramazanov, Vikman (2025)



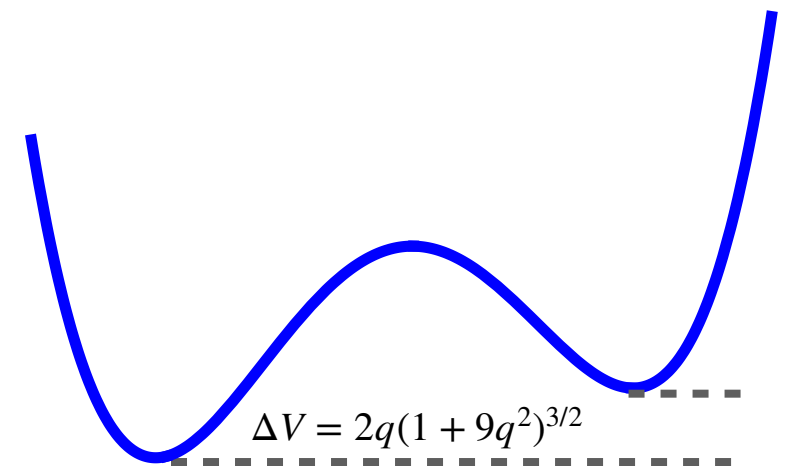
Annihilating domain walls

- DWs eventually dominate the energy density of the universe if $\sigma \gtrsim (\text{MeV})^3$:

$$\rho_{\text{dw}} \sim \sigma H \sim a^{-2} \rightarrow \text{DW problem!}$$

- **Biased potential:** ($q \ll 1$)

$$V = V_{\mathbb{Z}_2} + V_{\text{bias}} = \frac{\lambda}{4} (\phi^2 - v^2)^4 + qv\phi^3$$



Annihilation starts when: $\sigma H(\eta_{\Delta V}) = \Delta V$

- Recent papers studying GW production from annihilating DWs with simulations:

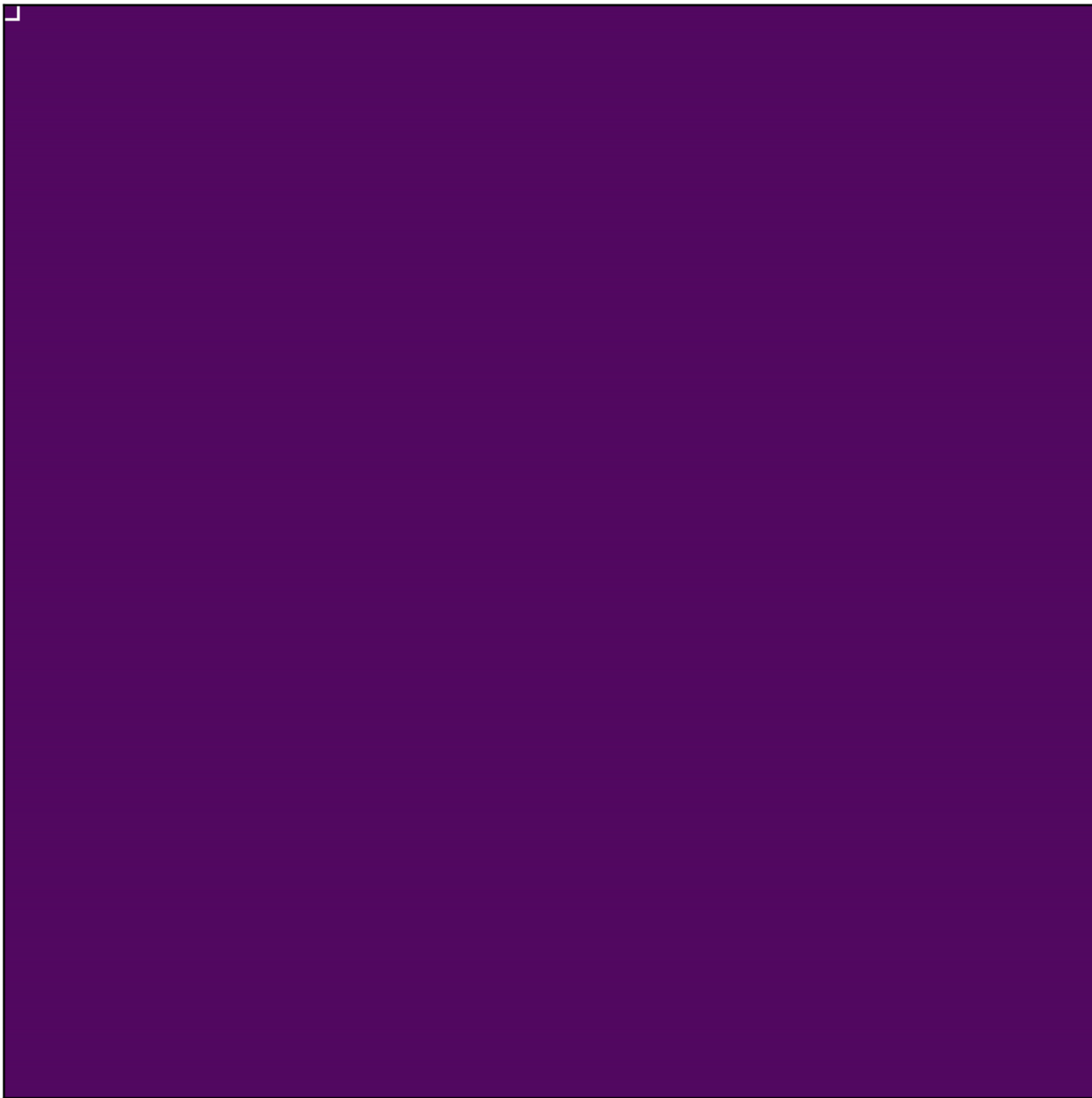
Cyr, Cotterill, Battye (2025)

Notari, Rompineve & F.T (2025)

Babichev, Dankovsky, Gorbunov, Ramazanov, Vikman (2025)

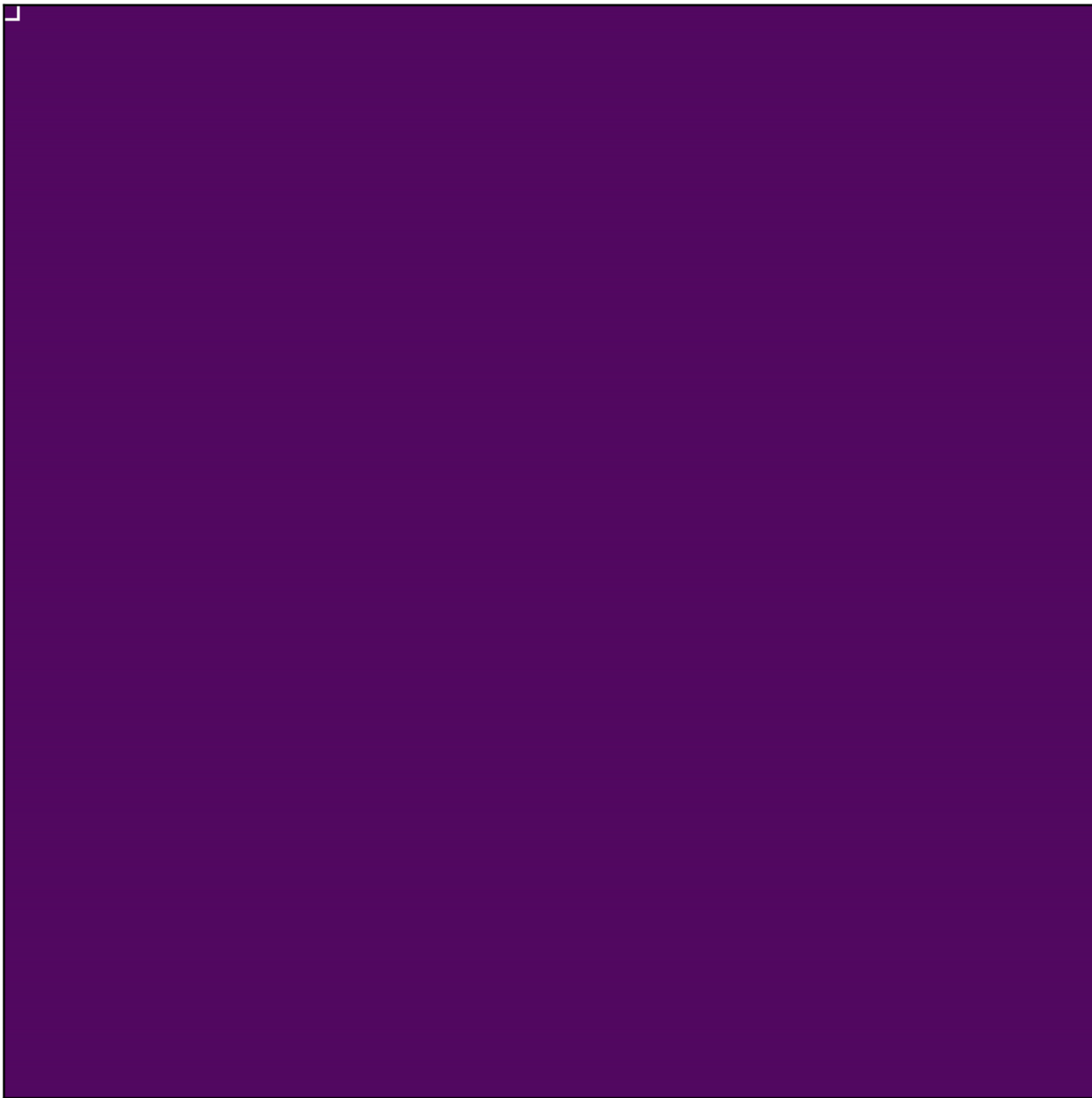
$$N = 3060$$

$$L = 80H_i^{-1}$$

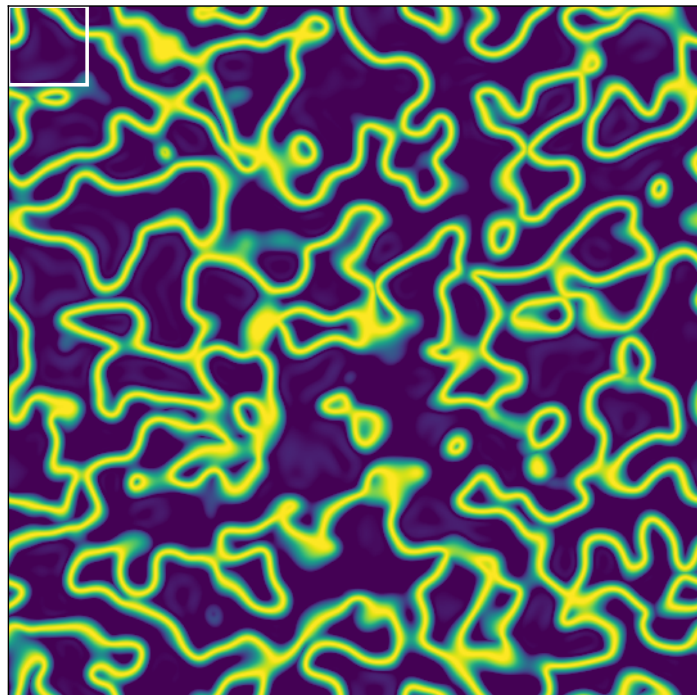


$$N = 3060$$

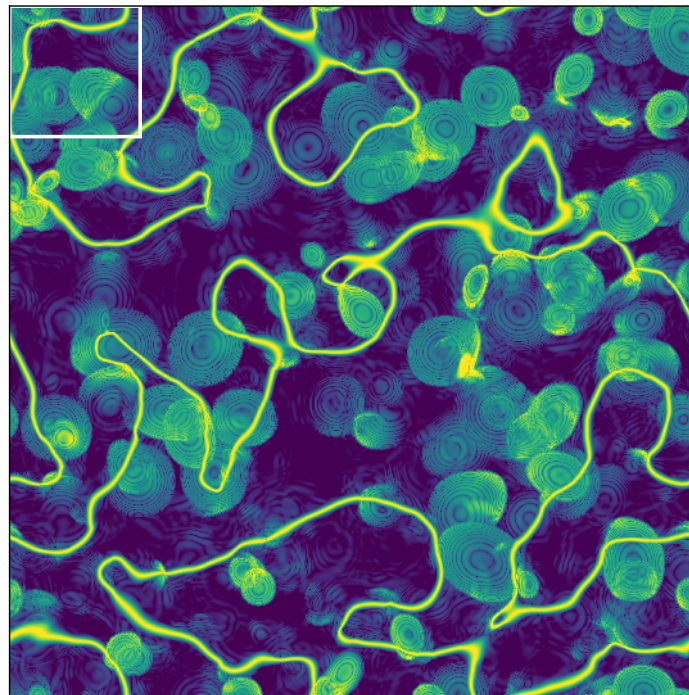
$$L = 80H_i^{-1}$$



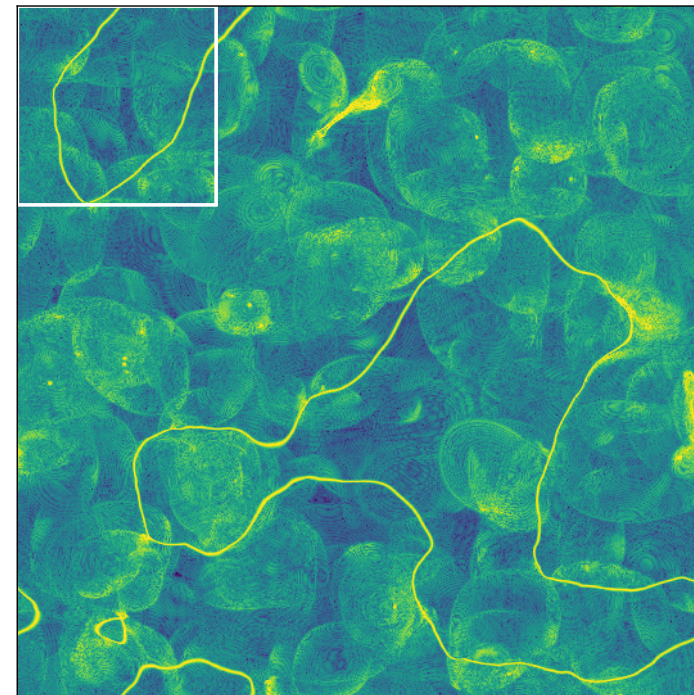
$\eta = 8.1$



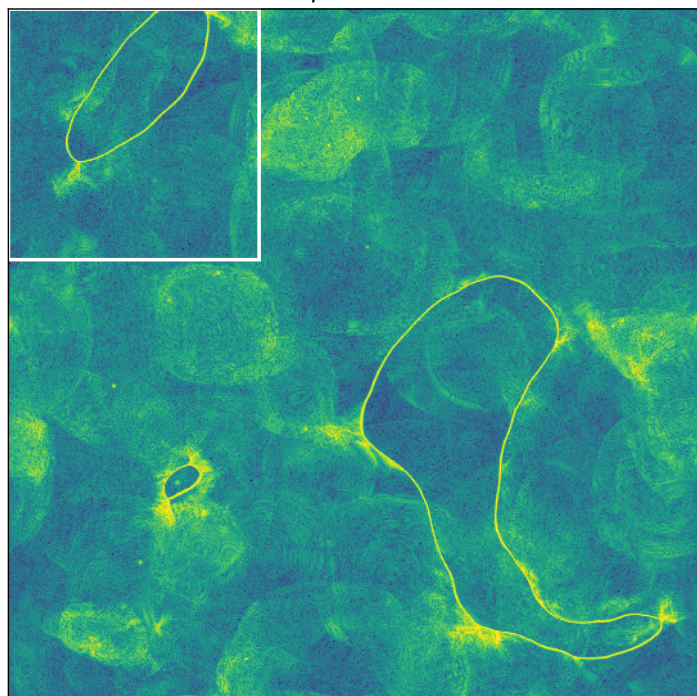
$\eta = 14.1$



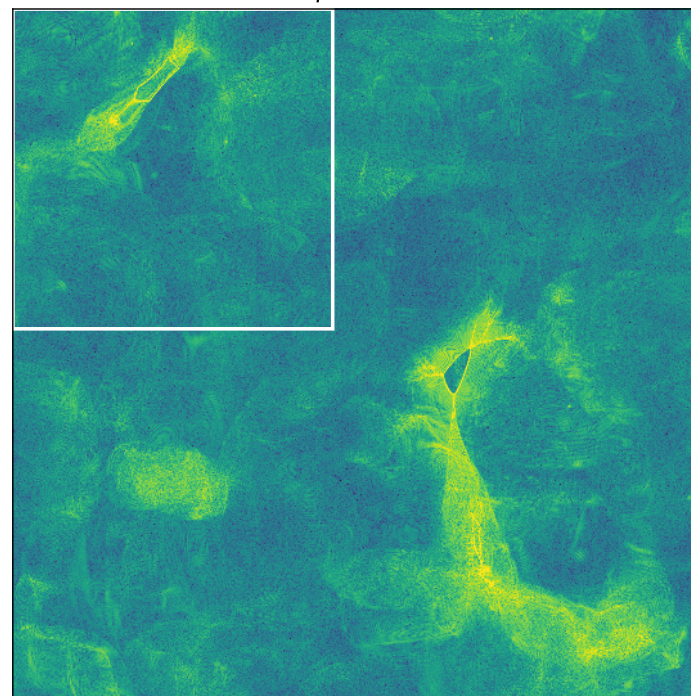
$\eta = 22.0$



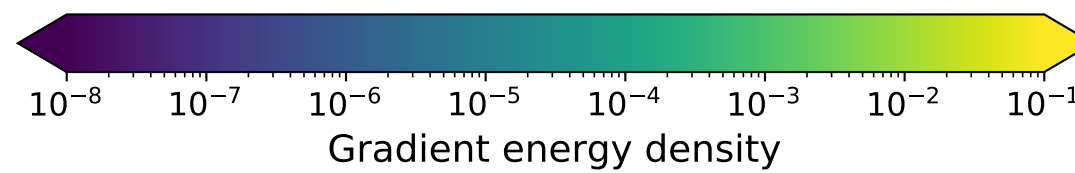
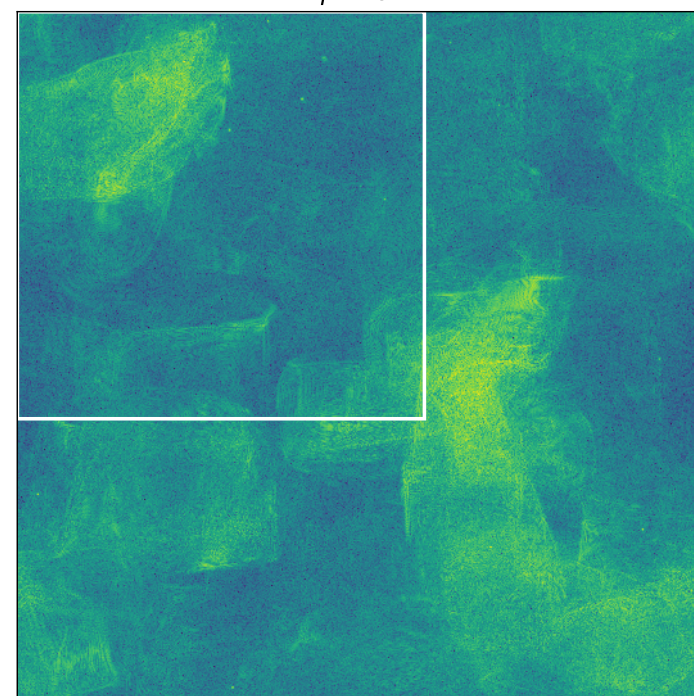
$\eta = 28.0$

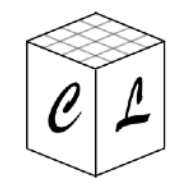


$\eta = 36.0$



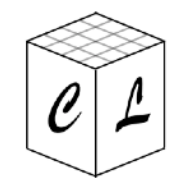
$\eta = 46.1$





Domain wall evolution (biased case)

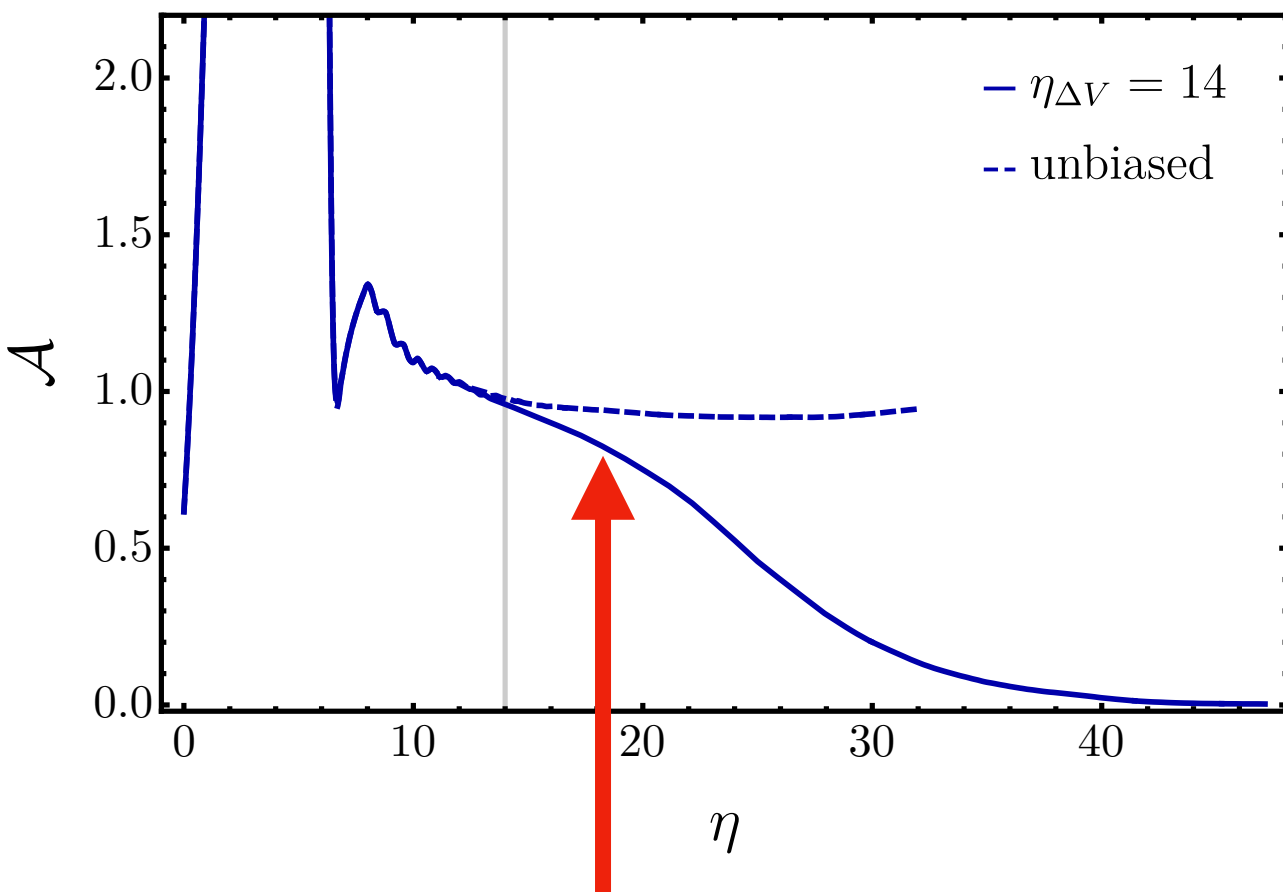
$$\eta_{\Delta V} = 14$$



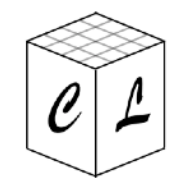
Domain wall evolution (biased case)

$$\eta_{\Delta V} = 14$$

Area parameter



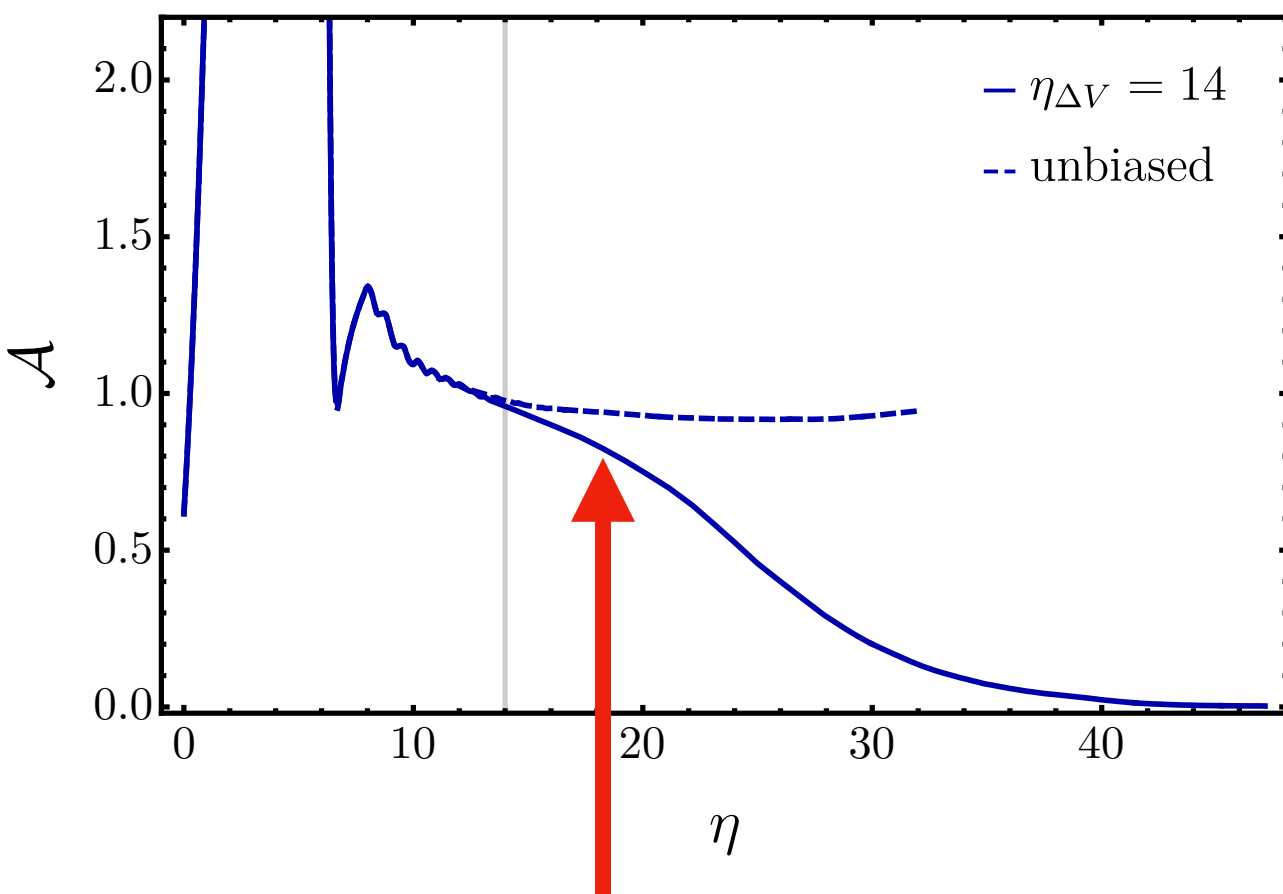
Annihilation



Domain wall evolution (biased case)

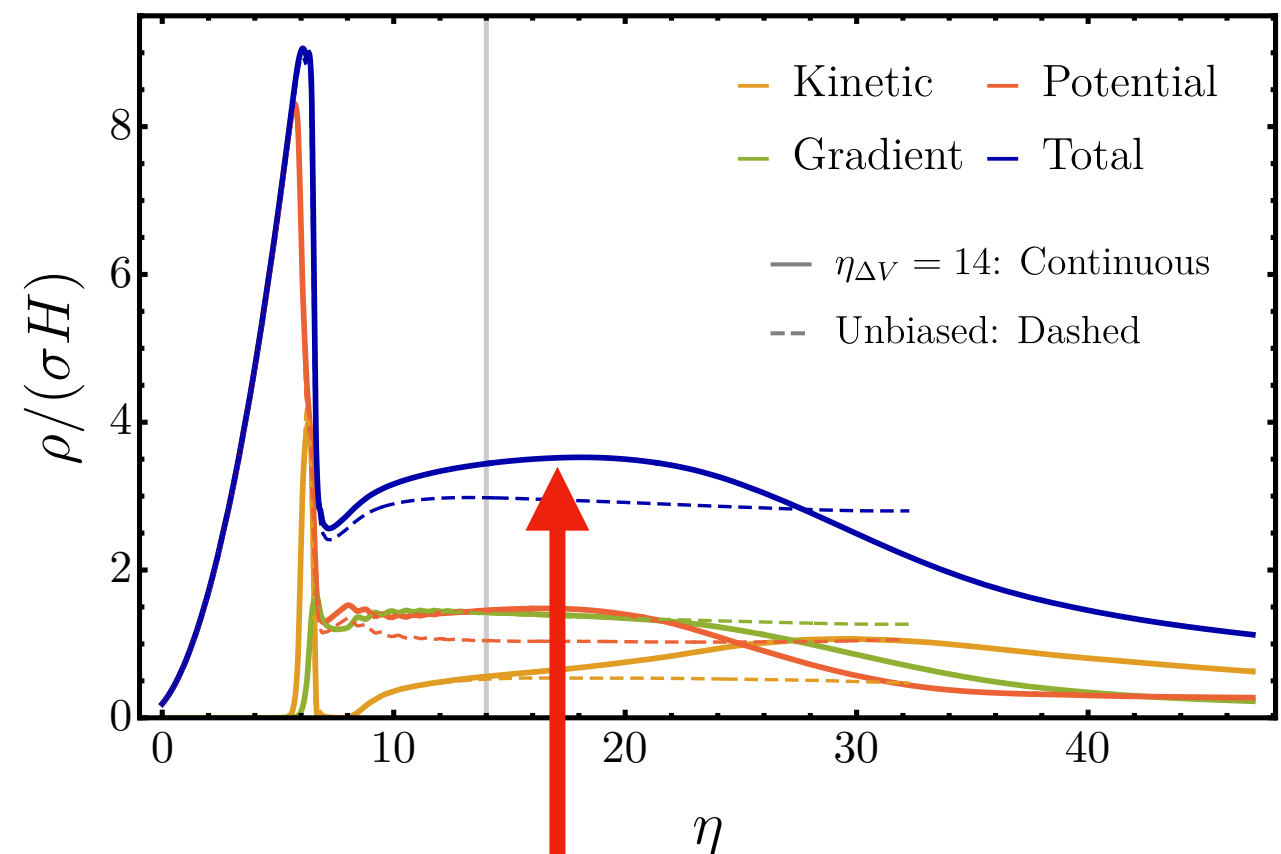
$$\eta_{\Delta V} = 14$$

Area parameter

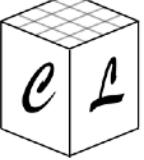


Annihilation

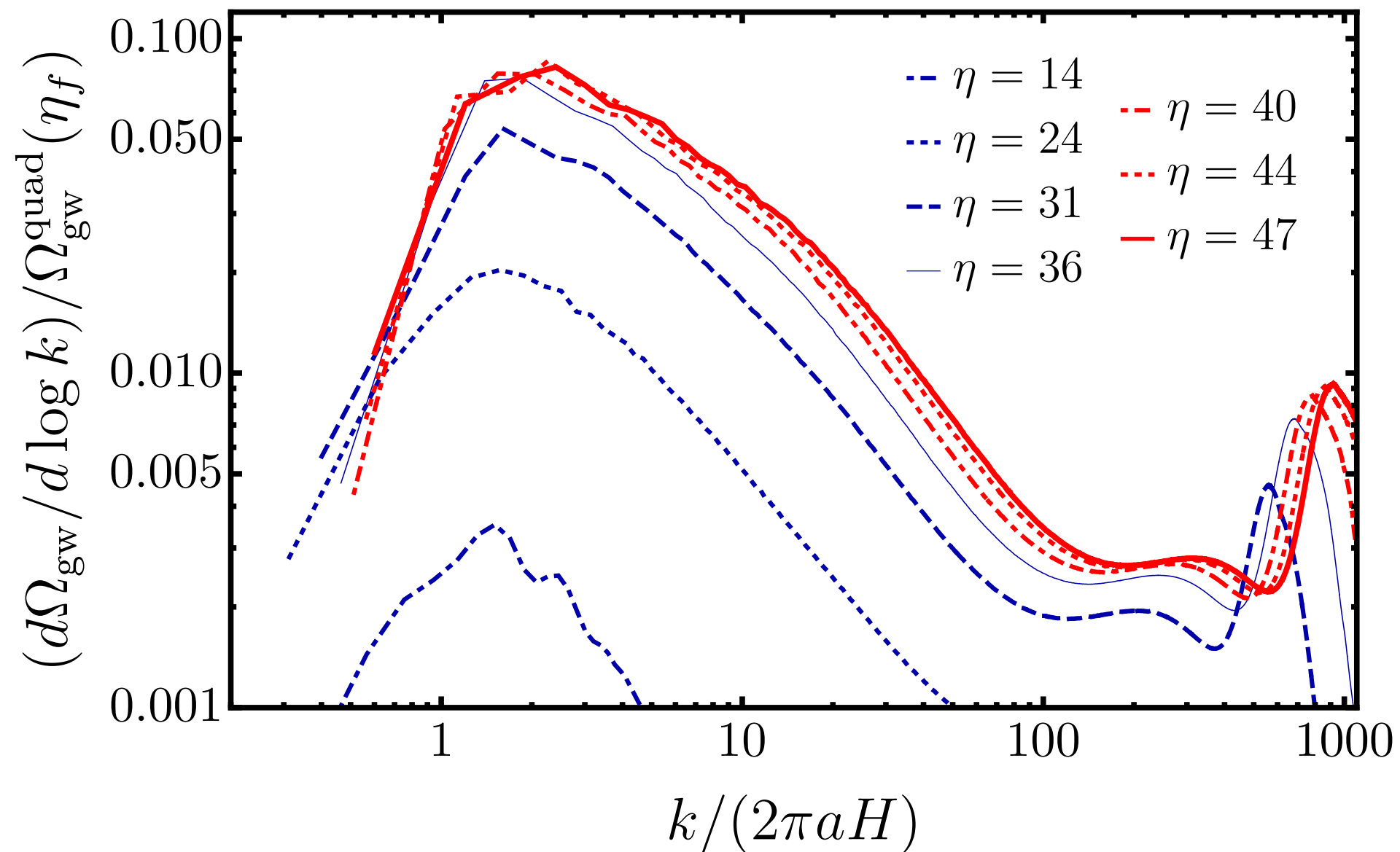
Energy density



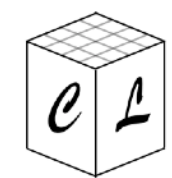
Annihilation



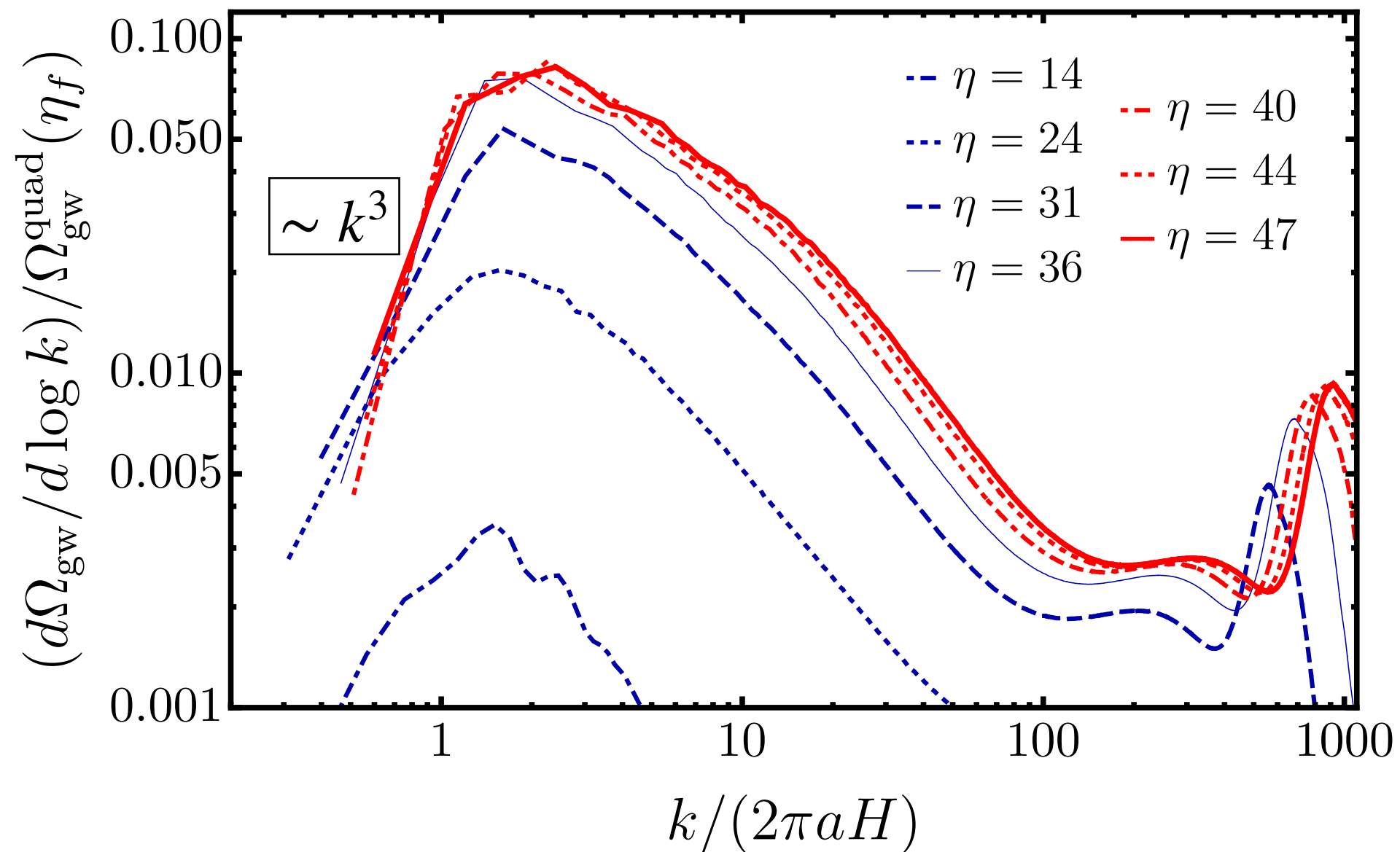
GWs from annihilating DW networks



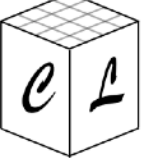
$$\Omega_{\text{gw}}^{\text{quad}} = 3/(32\pi)[2\sigma H/(3H^2 M_p)^2]$$



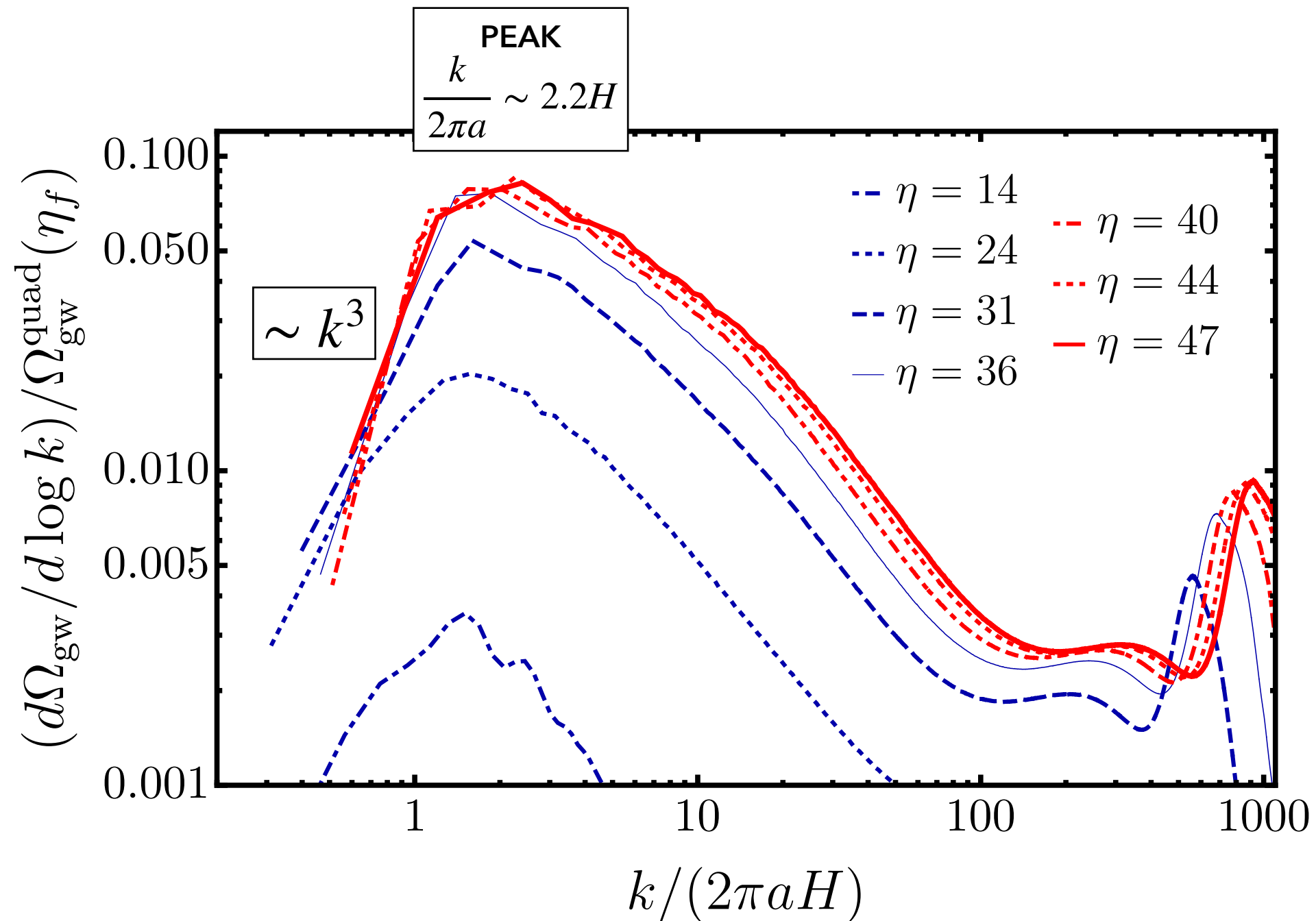
GWs from annihilating DW networks



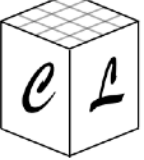
$$\Omega_{\text{gw}}^{\text{quad}} = 3/(32\pi)[2\sigma H/(3H^2 M_p)^2]$$



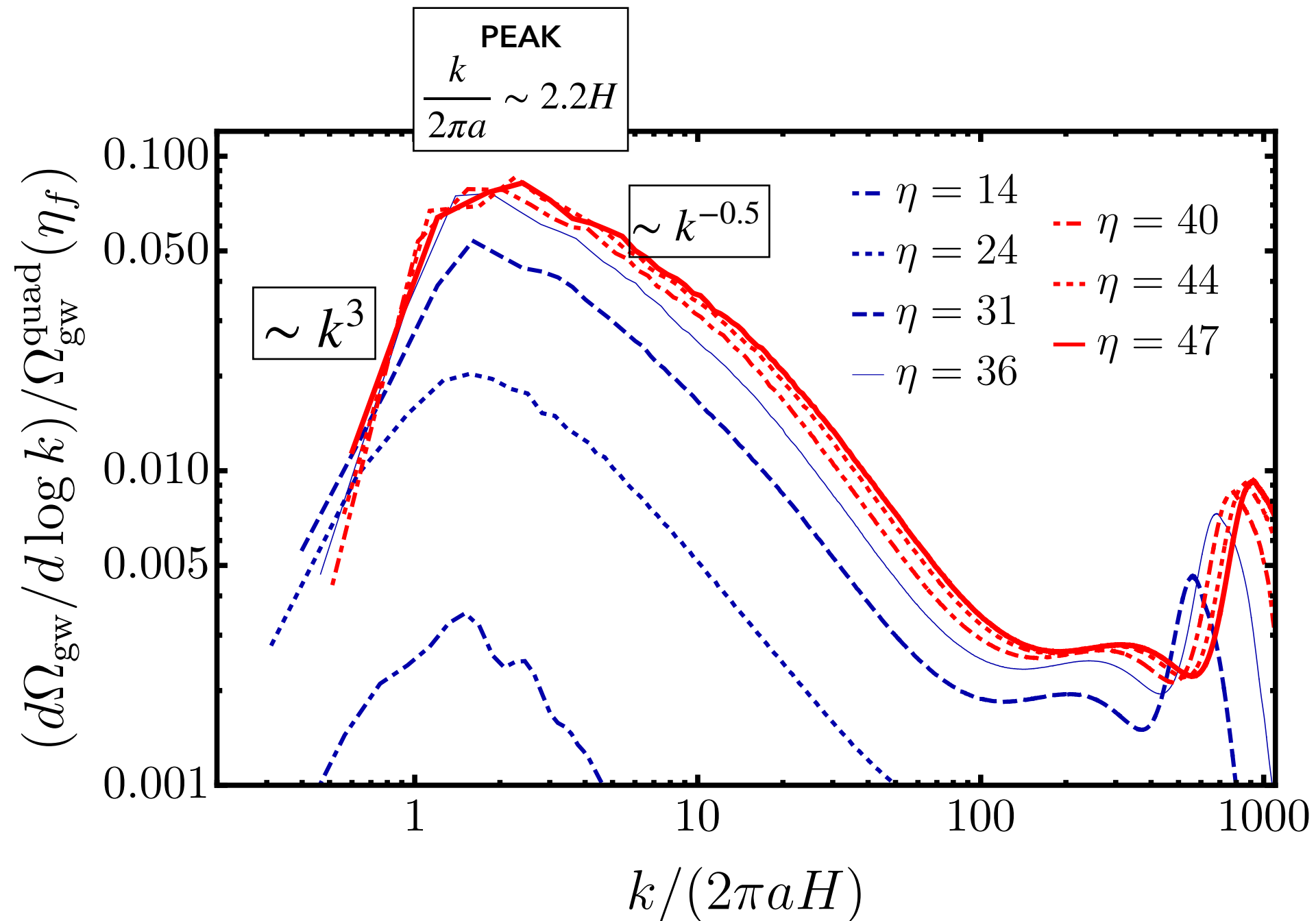
GWs from annihilating DW networks



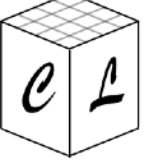
$$\Omega_{\text{gw}}^{\text{quad}} = 3/(32\pi)[2\sigma H/(3H^2 M_p)^2]$$



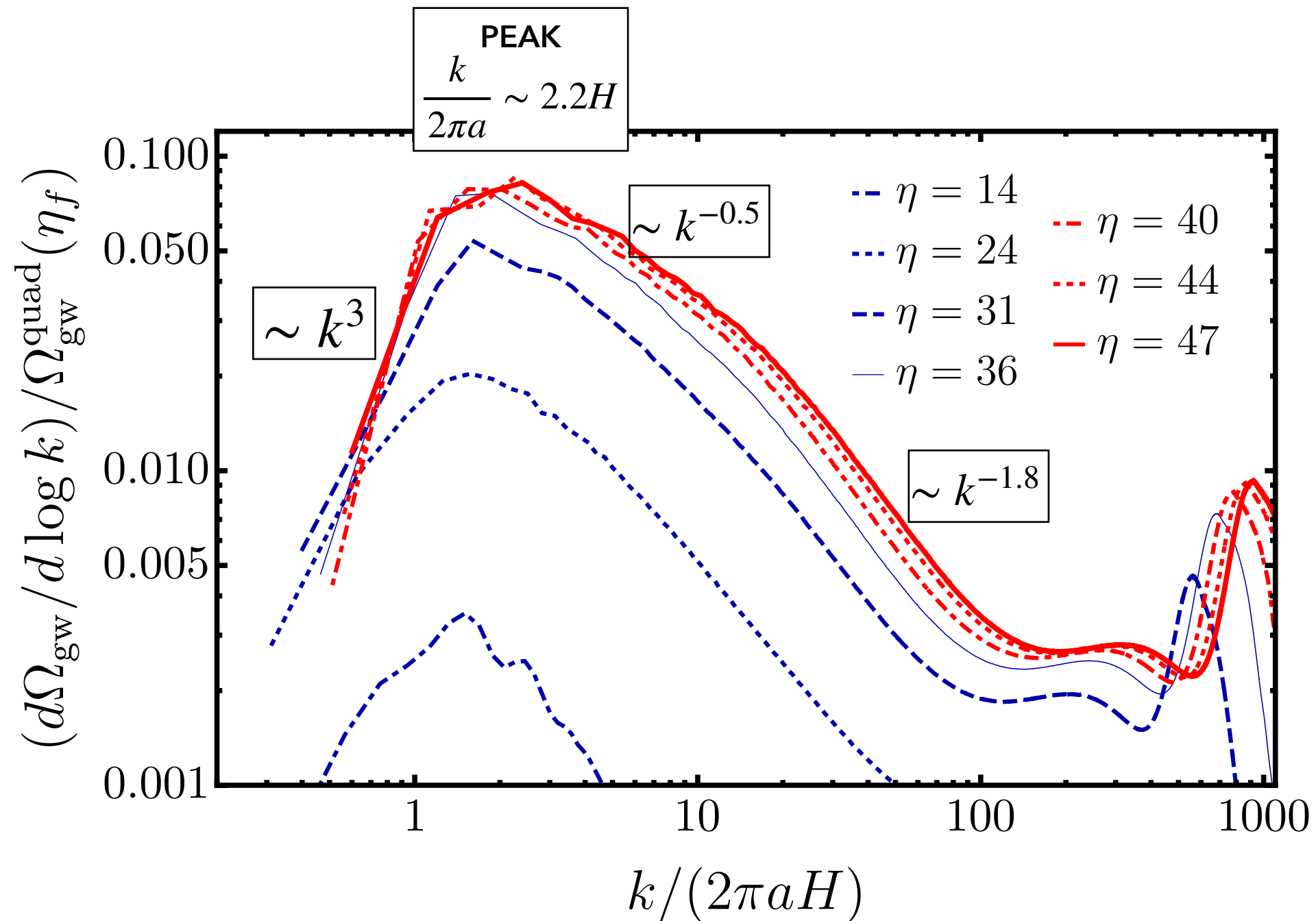
GWs from annihilating DW networks



$$\Omega_{\text{gw}}^{\text{quad}} = 3/(32\pi)[2\sigma H/(3H^2 M_p)^2]$$

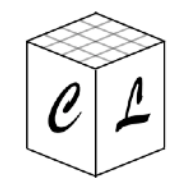


GWs from annihilating DW networks

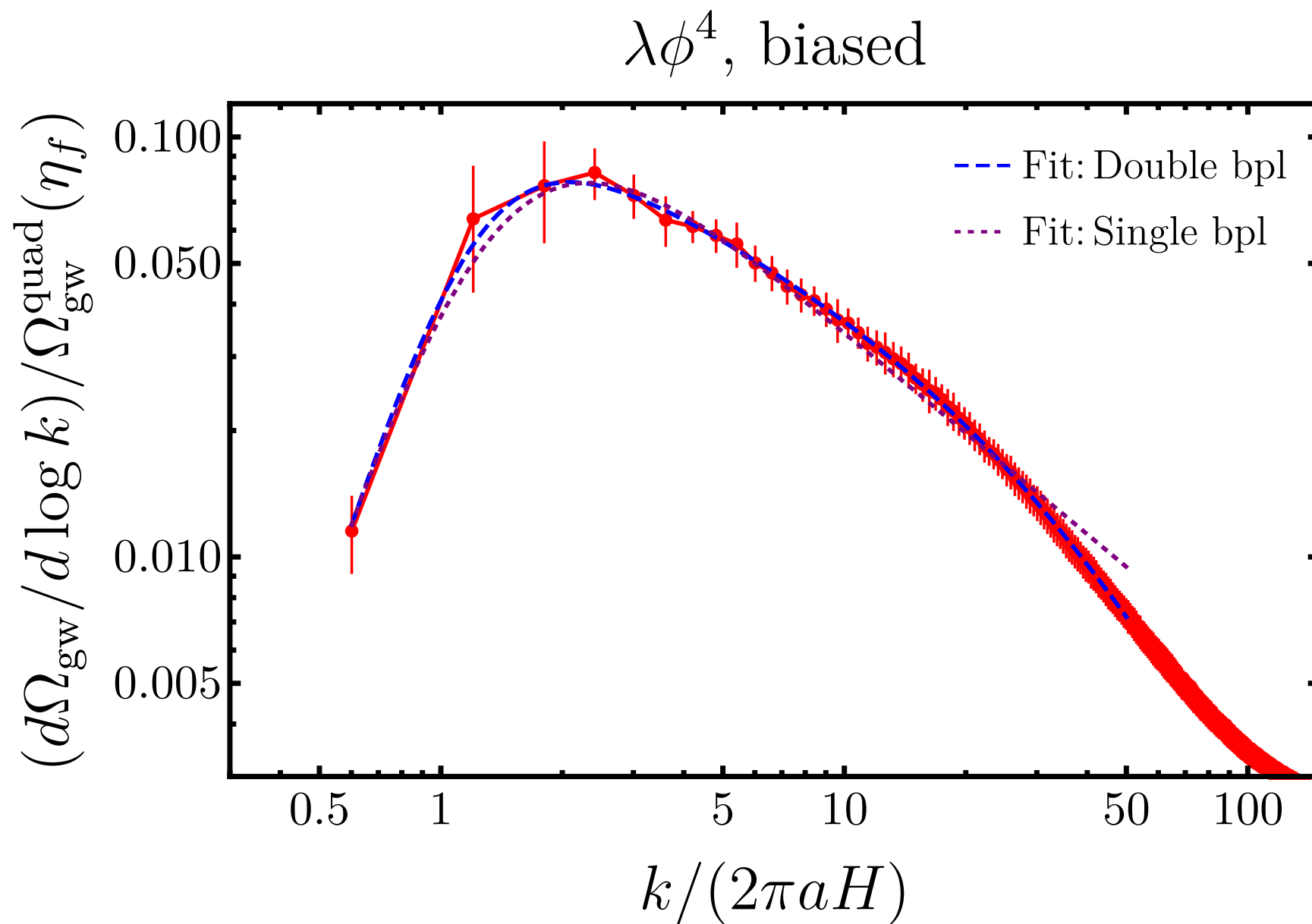


$$\Omega_{\text{gw}}^{\text{quad}} = 3/(32\pi)[2\sigma H/(3H^2 M_p)^2]$$

THANK YOU!



GWs from annihilating DW networks



Notari, Rompineve & F.T (JCAP, 2025)