

$\eta\pi$ phase shift determination
from $\eta' \rightarrow \eta\pi\pi$ experimental
data

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In collaboration with:

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Outline

- 1) Motivation
- 2) Formalism
- 3) Dispersive relations
- 4) Numerical strategy
- 5) Inputs
- 6) Perspectives and conclusions

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Motivation



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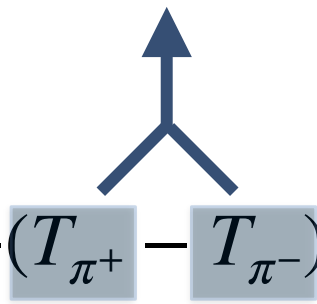
$$Y = \frac{m_\eta + 2m_\pi}{m_\pi} \frac{T_\eta}{Q} - 1$$

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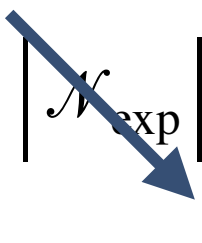
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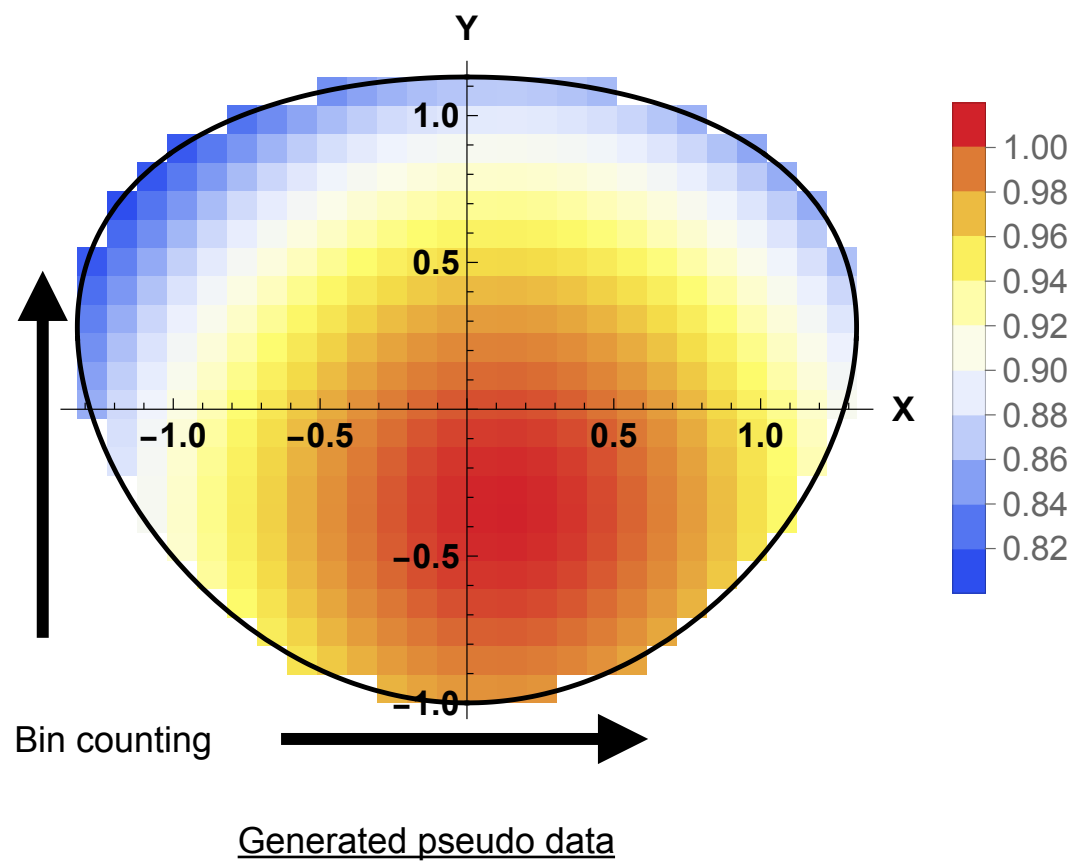
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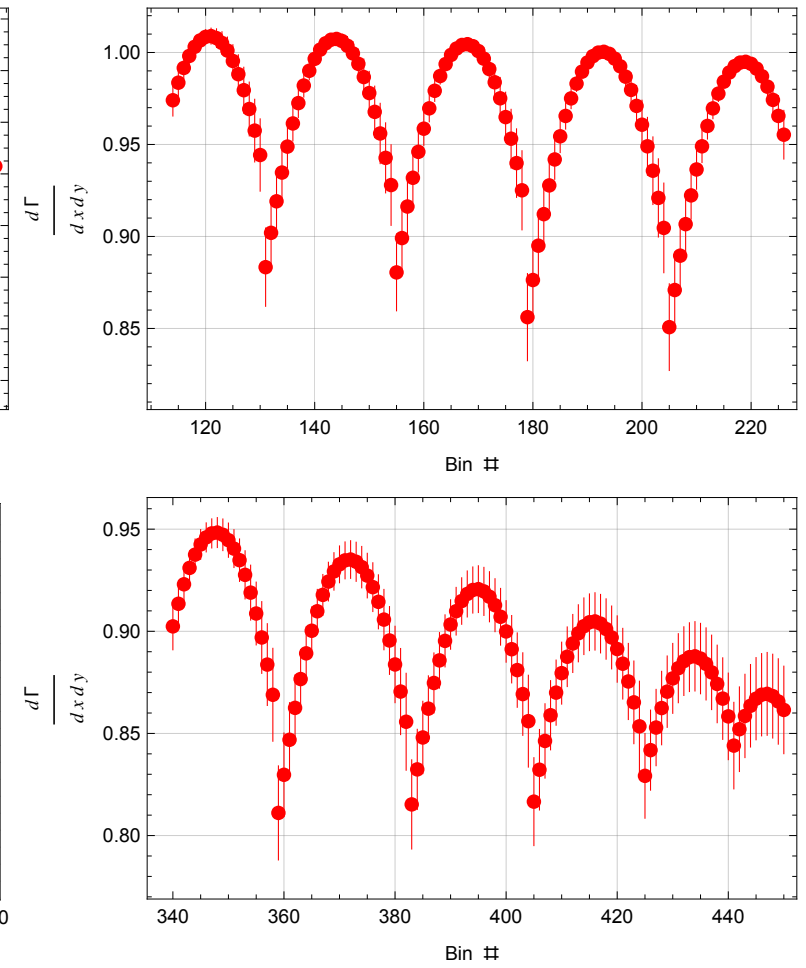
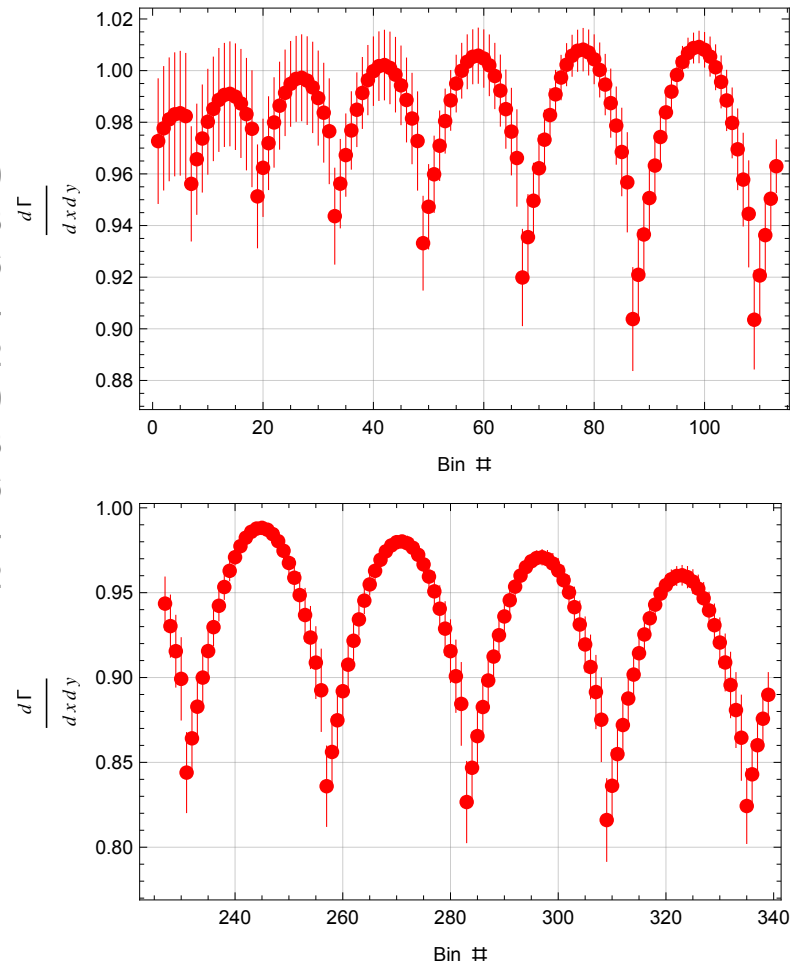
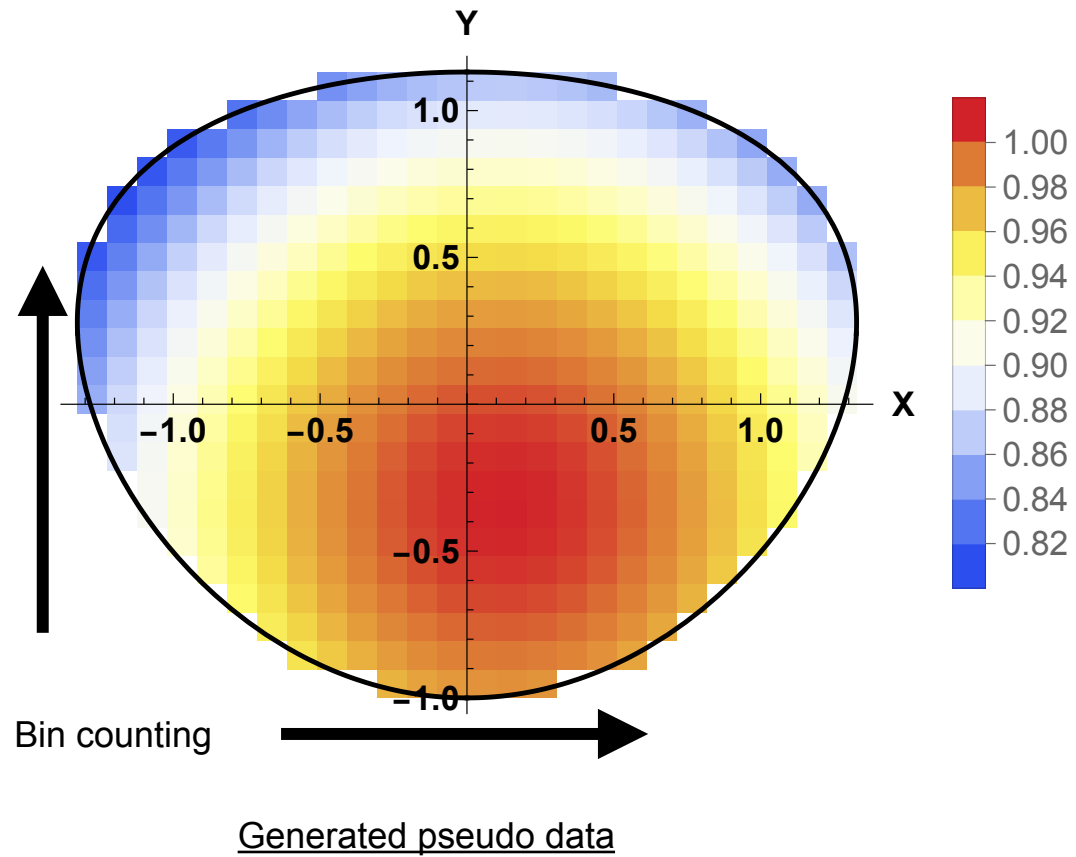

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Tree body final state interactions (FSI)

Incorporated using inelastic inputs in KT
elastic formalism

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The only possibilities are:

$$I_{\pi\pi} = 0, I_{\eta\pi} = 1$$





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$$\begin{aligned} \text{Disc} \left(\mathcal{M}(s, t, u) \right) = & \frac{i}{32\pi^2} \frac{\rho_{\pi\pi}(s)}{2} T_{I=0 \pi\pi}^*(s, \theta_s) \int d\Omega_{s'} \mathcal{M}(s, t', u') \\ & + \frac{i}{32\pi^2} \rho_{\pi\eta}(t) T_{\eta\pi}^*(t, \theta_t) \int d\Omega_{t'} \mathcal{M}(s', t, u') \\ & + \frac{i}{32\pi^2} \rho_{\pi\eta}(u) T_{\eta\pi}^*(u, \theta_u) \int d\Omega_{u'} \mathcal{M}(s', t', u) \end{aligned}$$

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$$\rho_{ij}(s) = \frac{\lambda^{1/2} \left(s, m_i^2, m_j^2 \right)}{s}$$





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$\hat{\mathcal{M}}_0(s)$

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1) Homogeneous Omnès problem



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$$\text{disc } (f(s)) = 2i \lim_{\epsilon \rightarrow 0^+} f(s + i\epsilon) \sin \delta(s) \exp[-i\delta(s)]$$

Dispersive relations

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$f(s)$ has a cut in the region $s \in [s_0, \infty)$  $\text{disc}(f(s)) = 2i \lim_{\epsilon \rightarrow 0^+} f(s + i\epsilon) \sin \delta(s) \exp[-i\delta(s)]$

Dispersive relations

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$f(s)$ has a cut in the region $s \in [s_0, \infty)$  $\text{disc}(f(s)) = 2i \lim_{\epsilon \rightarrow 0^+} f(s + i\epsilon) \sin \delta(s) \exp[-i\delta(s)]$

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Related with the subtractions

Outline

- 1) Motivation
- 2) Formalism
- 3) Dispersive relations
- 4) Numerical strategy
- 5) Inputs
- 6) Perspectives and conclusions

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Numerical strategy



Numerical strategy

1) Basis functions



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$$\begin{array}{ccc} \mathcal{M}_0(s) = s \Omega_0^{\pi\pi^0}(s) & \longrightarrow & \mathcal{M}_0^{\beta}(s) \\ \mathcal{M}_0^{\eta\pi}(s) = 0 & & \mathcal{M}_0^{\eta\pi\beta}(s) \end{array}$$





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$$\hat{\mathcal{M}}_0(s) = \frac{1}{2} \int_{-1}^1 dz_{s'} \left(\mathcal{M}_0^{\pi\eta}(t') + \mathcal{M}_0^{\pi\eta}(u') \right) = \frac{2}{\kappa_{\pi\pi}(s)} \int_{t_-(s)}^{t_+(s)} dt' \mathcal{M}_0^{\pi\eta}(t')$$

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D. Stamen, PhD thesis, (2024).

$$\kappa_{\eta\pi}(s) = \frac{\lambda^{1/2}(s, m_\eta^2, m_\pi^2) \lambda^{1/2}(s, m_{\eta'}^2, m_\pi^2)}{s}$$

2) Inhomogeneities

$\pi\pi$

$$\hat{\mathcal{M}}_0(s) = \frac{1}{2} \int_{-1}^1 dz_{s'} \left(\mathcal{M}_0^{\pi\eta}(t') + \mathcal{M}_0^{\pi\eta}(u') \right) = \frac{2}{\kappa_{\pi\pi}(s)} \int_{t_-(s)}^{t_+(s)} dt' \mathcal{M}_0^{\pi\eta}(t')$$

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$$\hat{\mathcal{M}}_0^{\eta\pi}(t) = \int_{-1}^1 dz_{t'} \left(\mathcal{M}_0^{\pi\eta}(u') - \mathcal{M}_0(s') \right) = \frac{1}{\kappa_{\eta\pi}(t)} \left(\int_{u_-(t)}^{u_+(t)} du' \mathcal{M}_0^{\pi\eta}(u') - \int_{s_-(t)}^{s_+(t)} ds' \mathcal{M}_0(s') \right)$$

S. Schneider, PhD thesis, (2012). First attempt to describe $\eta' \rightarrow \eta\pi\pi$, not successful

T. Isken, PhD thesis, (2021). Inclusion of high energy inelastic effects in the $\pi\pi$ and $\eta\pi$ phase shifts

M. Akdag, PhD thesis, (2023). Incorporation of CP-violating terms

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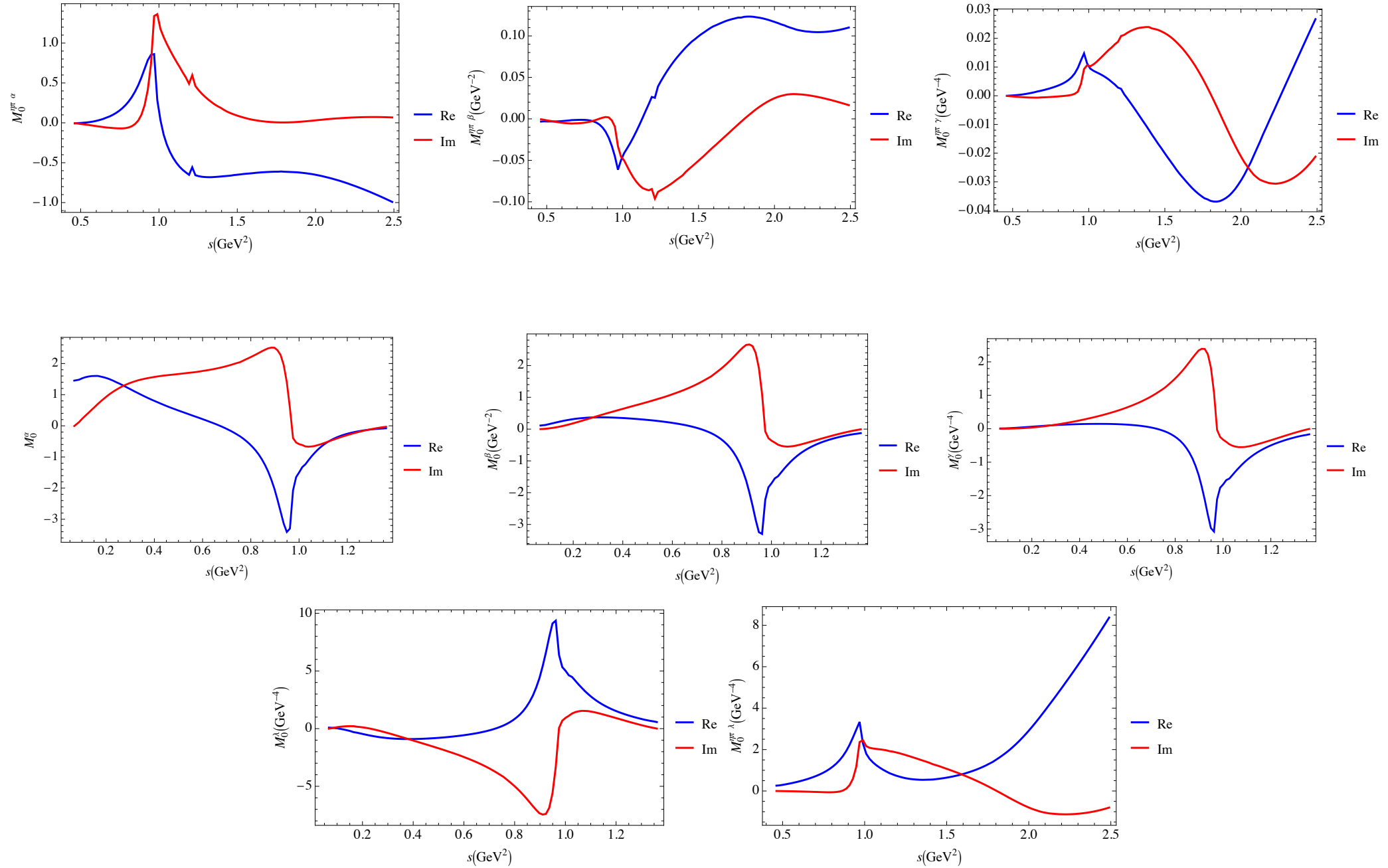
D. Winney, IterateKT, github repository. KT code with 3π in the final states

$$\kappa_{\eta\pi}(s) = \frac{\lambda^{1/2}(s, m_\eta^2, m_\pi^2) \lambda^{1/2}(s, m_{\eta'}^2, m_\pi^2)}{s}$$

Code crosscheck with Akdag2023 reported basis functions.



Code crosscheck with Akdag2023 reported basis functions.



Outline

- 1) Motivation
- 2) Formalism
- 3) Dispersive relations
- 4) Numerical strategy
- 5) Inputs
- 6) Perspectives and conclusions

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Inputs



Inputs

$$\text{Im} \left(\mathcal{T}_I^{ab\ 0}(s) \right) = \rho_{ab}(s) \mathcal{T}_I^{ab\ 0\ *} (s) \mathcal{T}_I^{ab\ 0}(s)$$



Inputs

$$\text{Im} \left(\mathcal{T}_I^{ab\ 0}(s) \right) = \rho_{ab}(s) \mathcal{T}_I^{ab\ 0\ *} (s) \mathcal{T}_I^{ab\ 0}(s)$$

$$\text{Im} \left(\mathcal{T}_I^{ab\ 0}(s)^{-1} \right) = - \rho_{ab}(s) \theta(s - (m_a + m_b)^2)$$

Inputs

$$\text{Im} \left(\mathcal{T}_I^{ab\ 0}(s) \right) = \rho_{ab}(s) \mathcal{T}_I^{ab\ 0*}(s) \mathcal{T}_I^{ab\ 0}(s)$$

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$$\mathcal{T}_I^{ab\ 0}(s)^{-1} = \mathcal{T}_I^{ab\ 0}(\tilde{s}_M)^{-1} + \frac{s - \tilde{s}_M}{\pi} \int_L \frac{ds'}{s' - \tilde{s}_M} \frac{\text{Im} \left(\mathcal{T}_I^{ab\ 0}(s')^{-1} \right)}{s' - s} + \frac{s - \tilde{s}_M}{\pi} \int_{s_0}^{\infty} \frac{ds'}{s' - \tilde{s}_M} \frac{-\rho_{ab}(s')}{s' - s}$$

Inputs

$$\text{Im} \left(\mathcal{T}_I^{ab\ 0}(s) \right) = \rho_{ab}(s) \mathcal{T}_I^{ab\ 0*}(s) \mathcal{T}_I^{ab\ 0}(s)$$

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Subtraction

Inputs

$$\text{Im} \left(\mathcal{T}_I^{ab\ 0}(s) \right) = \rho_{ab}(s) \mathcal{T}_I^{ab\ 0*}(s) \mathcal{T}_I^{ab\ 0}(s)$$

$$\text{Im} \left(\mathcal{T}_I^{ab\ 0}(s)^{-1} \right) = -\rho_{ab}(s) \theta(s - (m_a + m_b)^2)$$

$$\mathcal{T}_I^{ab\ 0}(s)^{-1} = \underbrace{\mathcal{T}_I^{ab\ 0}(\tilde{s}_M)^{-1}}_{\text{Subtraction}} + \underbrace{\frac{s - \tilde{s}_M}{\pi} \int_L \frac{ds'}{s' - \tilde{s}_M} \frac{\text{Im} \left(\mathcal{T}_I^{ab\ 0}(s')^{-1} \right)}{s' - s}}_{\text{Left-Hand-Cut (LHC)}} + \frac{s - \tilde{s}_M}{\pi} \int_{s_0}^{\infty} \frac{ds'}{s' - \tilde{s}_M} \frac{-\rho_{ab}(s')}{s' - s}$$

Inputs

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Inputs

$$\text{Im} \left(\mathcal{T}_I^{ab\ 0}(s) \right) = \rho_{ab}(s) \mathcal{T}_I^{ab\ 0*}(s) \mathcal{T}_I^{ab\ 0}(s)$$

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Analytic in the s complex plane except in the LHC

$$\mathcal{T}_I^{ab\ 0}(s)^{-1} = \underbrace{\mathcal{T}_I^{ab\ 0}(\tilde{s}_M)^{-1}}_{\text{Subtraction}} + \underbrace{\frac{s - \tilde{s}_M}{\pi} \int_L \frac{ds'}{s' - \tilde{s}_M} \frac{\text{Im} \left(\mathcal{T}_I^{ab\ 0}(s')^{-1} \right)}{s' - s}}_{\text{Left-Hand-Cut (LHC)}} + \frac{s - \tilde{s}_M}{\pi} \int_{s_0}^{\infty} \frac{ds'}{s' - \tilde{s}_M} \frac{-\rho_{ab}(s')}{s' - s}$$

Right-Hand-Cut (RHC)

Inputs

$$\text{Im} \left(\mathcal{T}_I^{ab\ 0}(s) \right) = \rho_{ab}(s) \mathcal{T}_I^{ab\ 0*}(s) \mathcal{T}_I^{ab\ 0}(s)$$

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Analytic in the s
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$$\mathcal{T}_I^{ab\ 0}(s)^{-1} = \underbrace{\mathcal{T}_I^{ab\ 0}(\tilde{s}_M)^{-1}}_{\text{Subtraction}} + \underbrace{\frac{s - \tilde{s}_M}{\pi} \int_L \frac{ds'}{s' - \tilde{s}_M} \frac{\text{Im} \left(\mathcal{T}_I^{ab\ 0}(s')^{-1} \right)}{s' - s}}_{\text{Left-Hand-Cut (LHC)}} + \frac{s - \tilde{s}_M}{\pi} \int_{s_0}^{\infty} \frac{ds'}{s' - \tilde{s}_M} \frac{-\rho_{ab}(s')}{s' - s}$$

Right-Hand-Cut (RHC)

$$\mathcal{T}_I^{ab\ 0}(s)^{-1} \approx \sum_{n=0} C_n \omega^n(s) + \frac{s - \tilde{s}_M}{\pi} \int_{s_0}^{\infty} \frac{ds'}{s' - \tilde{s}_M} \frac{-\rho_{ab}(s')}{s' - s}$$

Inputs

$$\text{Im} \left(\mathcal{T}_I^{ab\ 0}(s) \right) = \rho_{ab}(s) \mathcal{T}_I^{ab\ 0*}(s) \mathcal{T}_I^{ab\ 0}(s)$$

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Analytic in the s complex plane except in the LHC

$$\mathcal{T}_I^{ab\ 0}(s)^{-1} = \underbrace{\mathcal{T}_I^{ab\ 0}(\tilde{s}_M)^{-1}}_{\text{Subtraction}} + \underbrace{\frac{s - \tilde{s}_M}{\pi} \int_L \frac{ds'}{s' - \tilde{s}_M} \frac{\text{Im} \left(\mathcal{T}_I^{ab\ 0}(s')^{-1} \right)}{s' - s}}_{\text{Left-Hand-Cut (LHC)}} + \underbrace{\frac{s - \tilde{s}_M}{\pi} \int_{s_0}^{\infty} \frac{ds'}{s' - \tilde{s}_M} \frac{-\rho_{ab}(s')}{s' - s}}_{\text{Right-Hand-Cut (RHC)}}$$

$$\mathcal{T}_I^{ab\ 0}(s)^{-1} \approx \sum_{n=0} C_n \omega^n(s) + \frac{s - \tilde{s}_M}{\pi} \int_{s_0}^{\infty} \frac{ds'}{s' - \tilde{s}_M} \frac{-\rho_{ab}(s')}{s' - s}$$

$\omega(s)$ is a conformal mapping that transforms the LHC into the boundary circle $|\omega| = 1$ in the ω complex plane. The rest of the s complex plane transforms into the interior of the circle. $\omega(s)$ is real if s belongs to $[s_0, \infty)$.

Inputs

$$\text{Im} \left(\mathcal{T}_I^{ab\ 0}(s) \right) = \rho_{ab}(s) \mathcal{T}_I^{ab\ 0*}(s) \mathcal{T}_I^{ab\ 0}(s)$$

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I. Danilkin et al., Phys. Rev. D 107, 074021 (2023).

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Perspectives and conclusions

- We have generated pseudo data that reproduces the experimental fit of BESIII
- We have coded a KT code to compute the basis functions. It converges after 1 or 2 iterations in the majority of the cases. This code have been crosschecked with the basis functions reported in Akdag2023.
- The next step is to fit the pseudo data with the basis functions that generates the dispersive inverse amplitude to fix the conformal mapping parameters.
- The most important challenge is to constrain well enough the dispersive inverse amplitude with χPT inputs.

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