# Scotogenic mechanism from a 3221 symmetry JHEP 10 (2025) 129 [2507.21223]

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From ancient greek

Σκότος: Darkness

Γένος: Kin, generation

• Take the SM and add a  $Z_2$  symmetry

$$SU(3)_C \times SU(2)_L \times U(1)_Y \times Z_2$$



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$$\eta \sim (2, 1/2, -) \quad N^{\sigma} \sim (1, 0, -)$$

• The  $\eta_0$  VEV is zero, so that  $Z_2$  is exact

$$\langle \eta_0 \rangle = 0 \rightarrow Z_2$$
 does not break

 Very minimal model with interesting phenomenology

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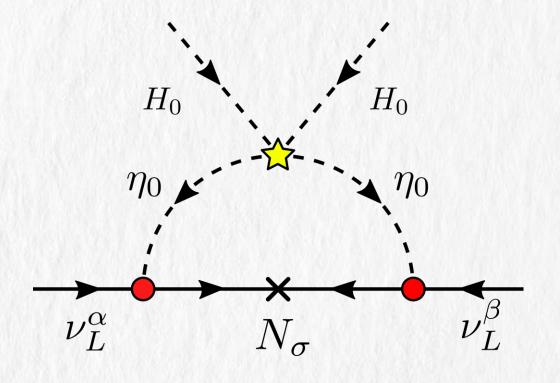
• **DM candidates** (bosonic or fermionic), stable thanks to an exact  $Z_2$ 

**Bosonic:** Lightest  $Re(\eta_0)$  or  $Im(\eta_0)$ 

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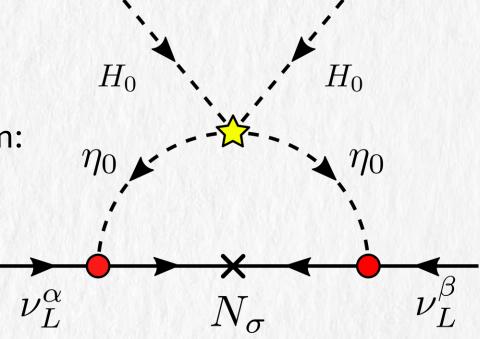


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Radiative Majorana neutrino masses at 1-loop LO

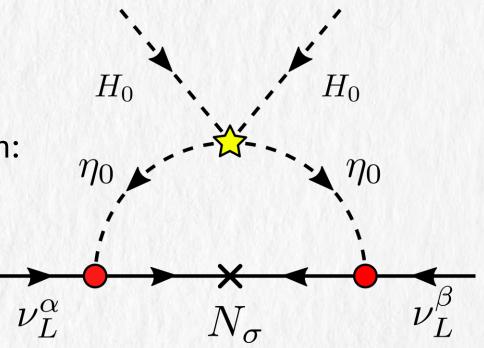
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A Yukawa interaction:

$$Y^{\alpha\sigma}(\tilde{\eta}^{\dagger}L^{\alpha}\overline{N}^{\sigma}) + h.c.$$



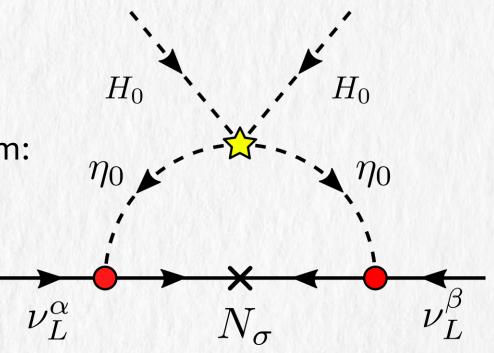
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• A quartic scalar interaction:

$$\frac{1}{2}\lambda_5\left[\left(H^{\dagger}\eta\right)^2+h.c.\right]$$



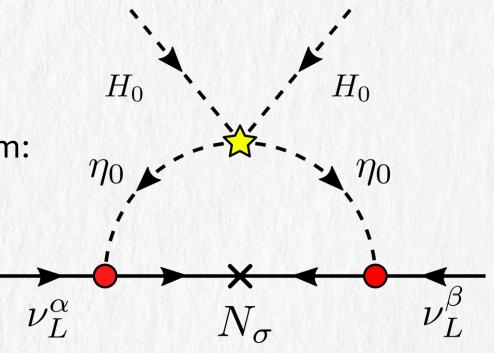
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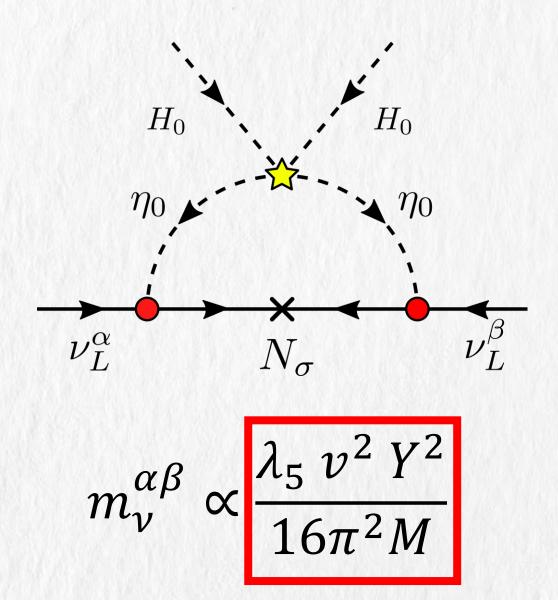
$$\frac{1}{2}\lambda_5\left[\left(H^{\dagger}\eta\right)^2+h.c.\right]$$



$$m_{\nu}^{\alpha\beta} \propto \frac{\lambda_5 \, v^2 \, Y^2}{16\pi^2 M}$$

• We can get **neutrino masses**  $\sim$  **eV** with **loop masses**  $M \sim 1$  **TeV** if:

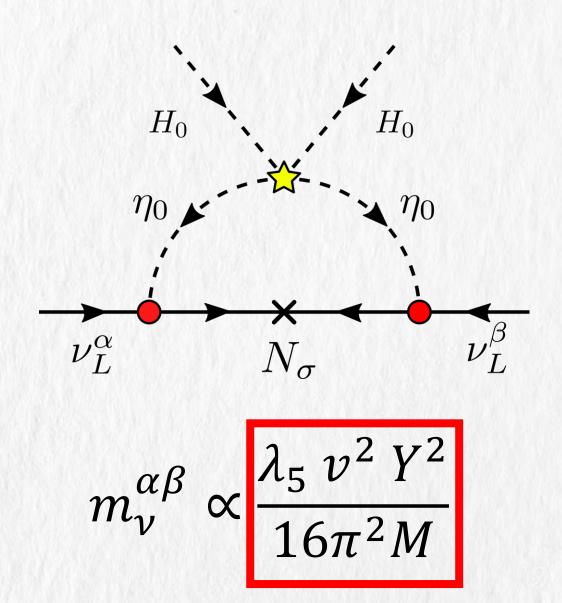
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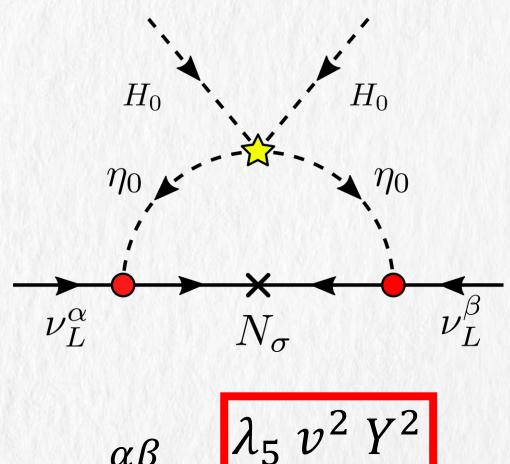


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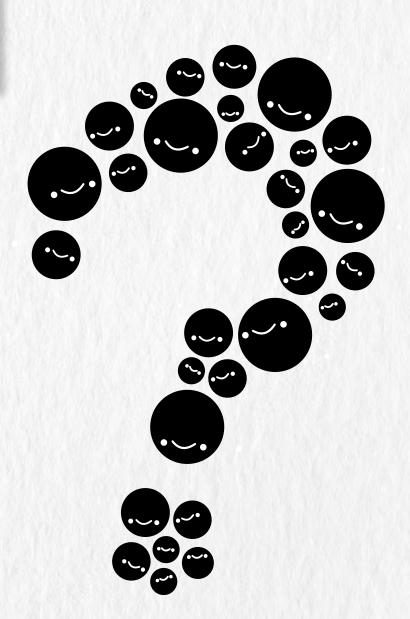
## $\triangle$ But why's $\lambda_5$ small?



$$m_{\nu}^{\alpha\beta} \propto \frac{\lambda_5 \, v^2 \, Y^2}{16\pi^2 M}$$

## There is always a bigger fish

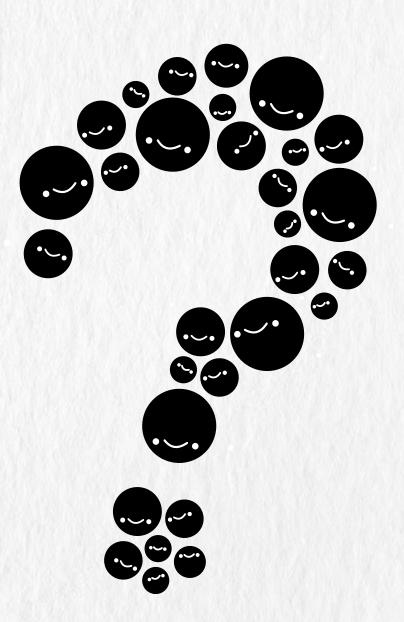
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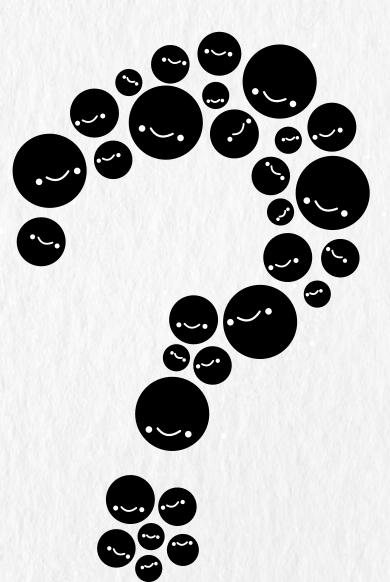


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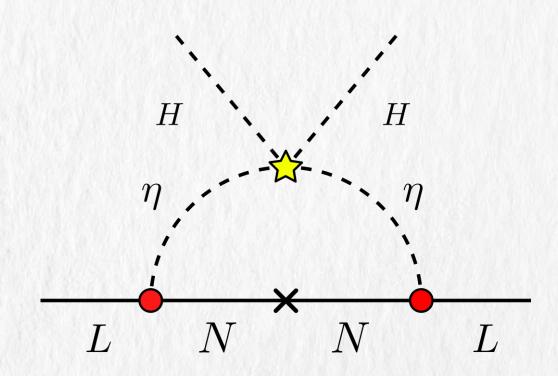
 We want to build a SM extension that includes the Scotogenic at low energies...

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We will also generate Z<sub>2</sub> as an accidental symmetry



$$SU(3)_C \times SU(2)_1 \times SU(2)_2 \times U(1)_Y$$



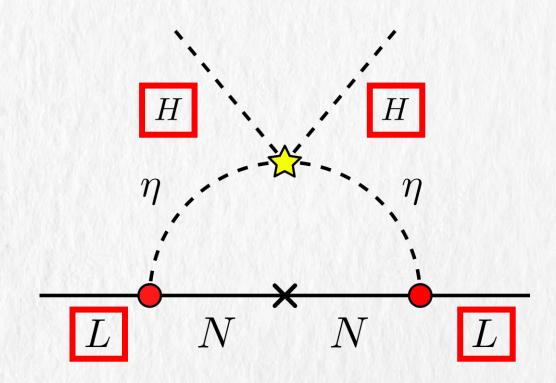
The usual SM representations

$$SU(3)_C \times SU(2)_1 \times SU(2)_2 \times U(1)_Y$$

-1/2

H

**1** +1/2



The scotogenic scalar doublet

$$SU(3)_C \times SU(2)_1 \times SU(2)_2 \times U(1)_Y$$

L

1

2

1

-1/2

H

1

2

V.

+1/2

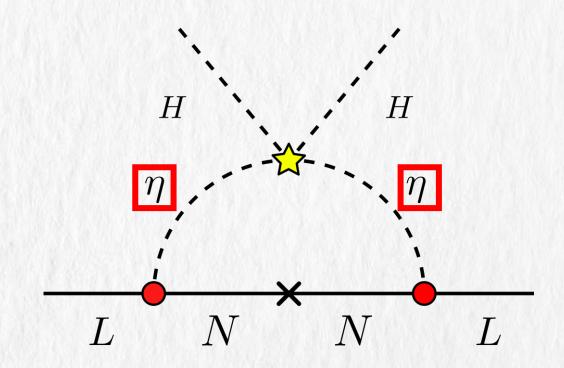
η

1

1

2

+1/2



This vertex is now messed up

$$SU(3)_C \times SU(2)_1 \times SU(2)_2 \times U(1)_Y$$

L

2

1

-1/2

H

1

2

W

+1/2

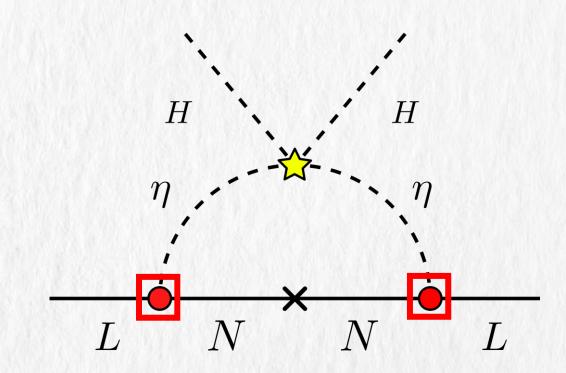
η

1

1

2

+1/2



$$(2,1) \times (1,2) \times (1,1) \not \supset (1,1)$$

Change the singlet for a bidoublet

$$SU(3)_C \times SU(2)_1 \times SU(2)_2 \times U(1)_Y$$

L 1

2

1

-1/2

H :

2

1

+1/2

η

1

1

2

+1/2

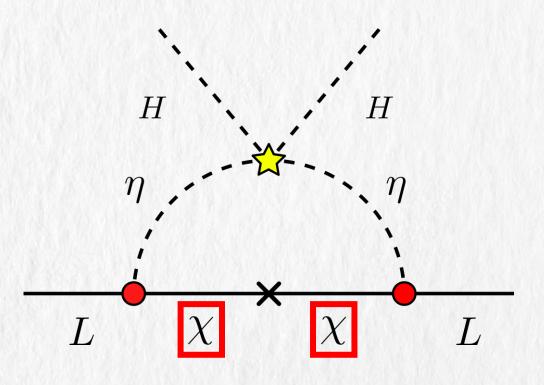
X

1

2

**2** 

0



$$(2,1) \times (1,2) \times (2,2) \supset (1,1)$$

This vertex is messed up too

$$SU(3)_C \times SU(2)_1 \times SU(2)_2 \times U(1)_Y$$

L 1

2

1

-1/2

H

2

1

+1/2

η

1

1

2

+1/2

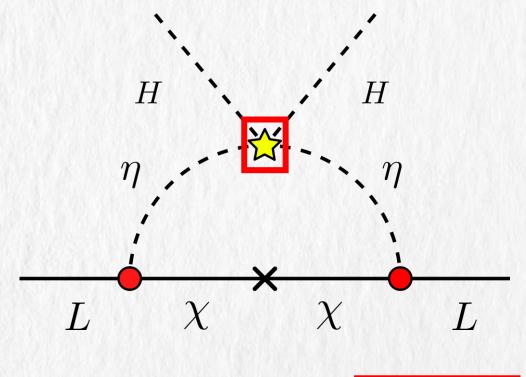
X

1

2

**2** 

0



$$(2,1)^2 \times (1,2)^2 \not\supset (1,1)$$

$$SU(3)_C \times SU(2)_1 \times SU(2)_2 \times U(1)_Y$$

L 1

2

1

-1/2

H

1

2

1

+1/2

η

1

1

2

+1/2

X

1

2

2

0

Δ

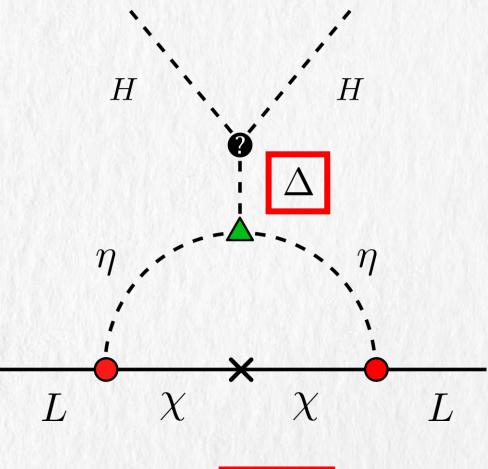
1

1

3

+1

Just add a scalar triplet, all good



$$(1,2)^2 \times (1,3) \supset (1,1)$$

$$SU(3)_C \times SU(2)_1 \times SU(2)_2 \times U(1)_Y$$

L 1

2

1

-1/2

Н

1

2

1

+1/2

η

1

1

2

+1/2

X

1

2

7

0

Λ

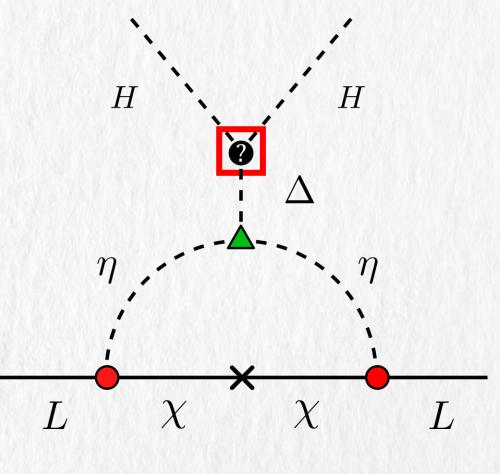
1

1

3

+1

And this last vertex...



$$(2,1)^2 \times (1,3) \not\supset (1,1)$$

$$SU(3)_C \times SU(2)_1 \times SU(2)_2 \times U(1)_Y$$

 L
 1
 2
 1
 -1/2 

 H
 1
 2
 1
 +1/2 

  $\eta$  1
 1
 2
 +1/2 

  $\chi$  1
 2
  $\overline{2}$  0

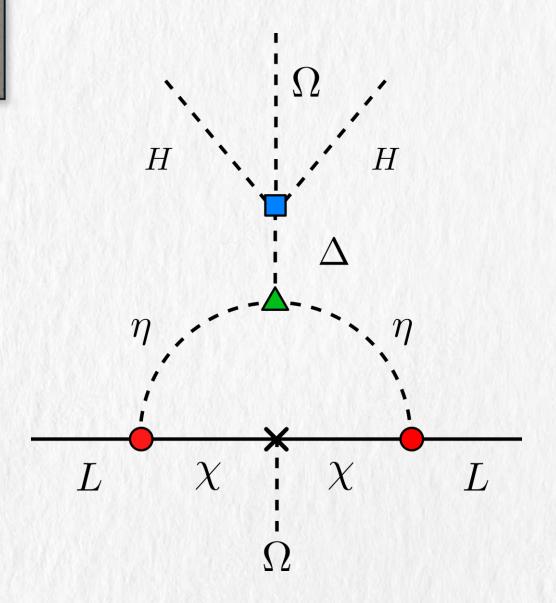
+1

 $\Omega$  1 3  $\overline{3}$  0

...is fixed by a scalar bitriplet  $(2,1)^2 \times (1,3) \times (3,3) \supset (1,1)$ 

## Going down the ladder

 How do we return to the scotogenic's diagram?



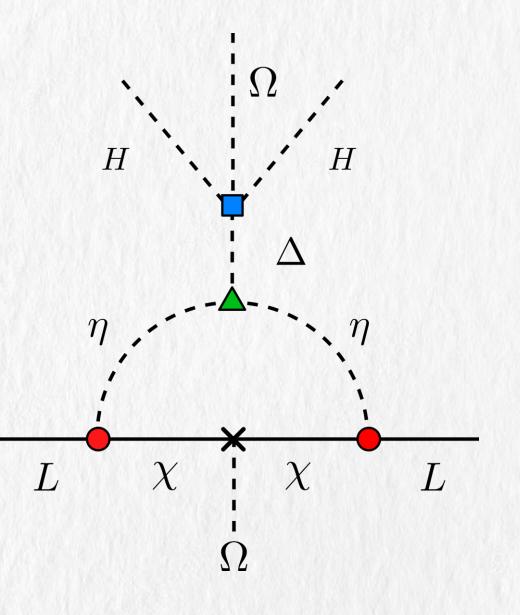
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Consider the following VEVs and hierarchy

$$\langle \Omega \rangle \propto \begin{pmatrix} v_{\Omega} & 0 & 0 \\ 0 & v_{\xi} & 0 \\ 0 & 0 & v_{\Omega} \end{pmatrix} \qquad \langle \Delta \rangle \propto \begin{pmatrix} 0 \\ 0 \\ v_{\Delta} \end{pmatrix}$$

$$\langle H \rangle \propto \begin{pmatrix} 0 \\ v_H \end{pmatrix}$$
  $v_{\Omega}, v_{\xi} \gg v_H \gg v_{\Delta}$ 



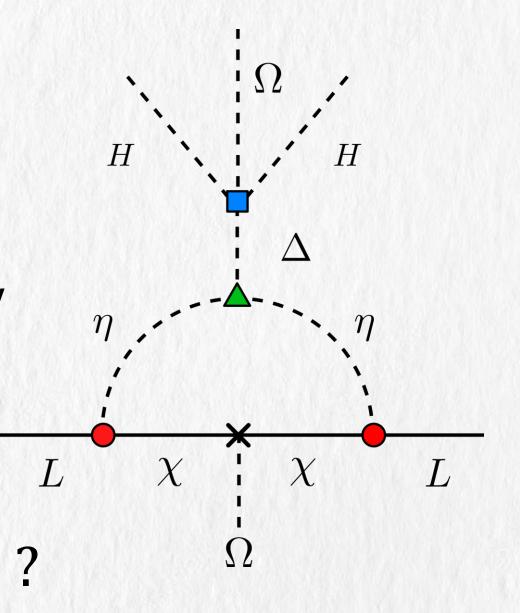
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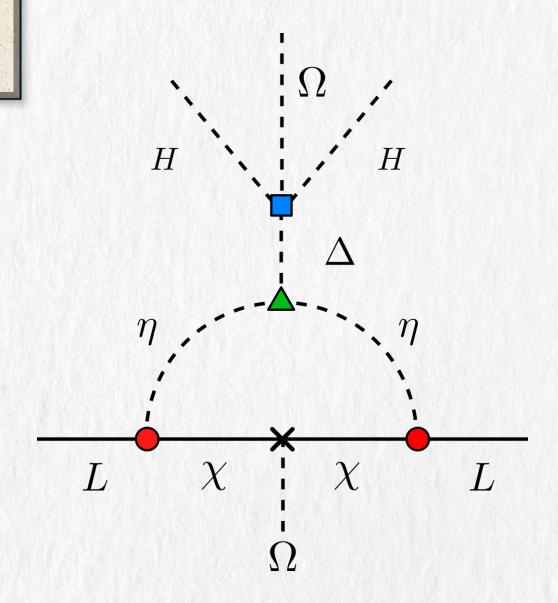
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?



## Why is $v_{\Delta}$ the smallest?

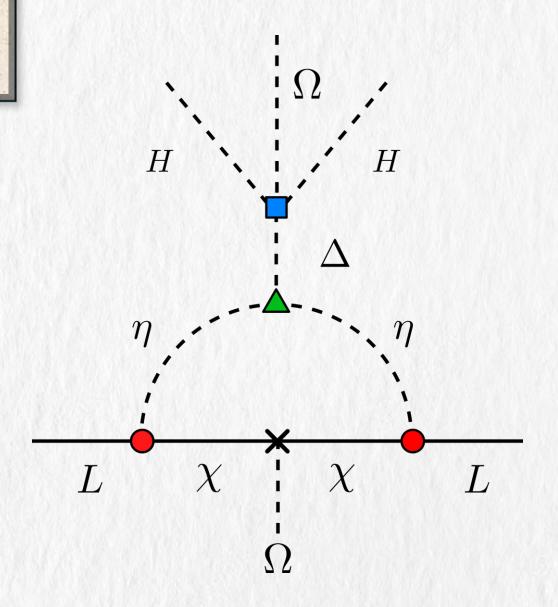
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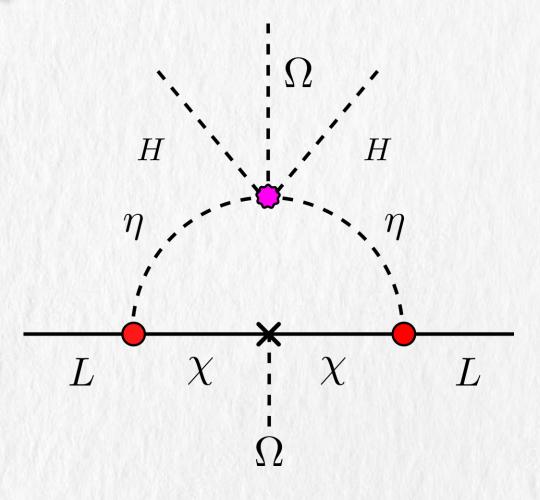
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- The smallness of  $\lambda_5$  comes from it being an effective operator after  $\Delta$  is integrated out.
- Therefore we need  $m_{\it \Delta} \gg$  any other mass scale
- From the tadpole equations, this forces small  $v_A$

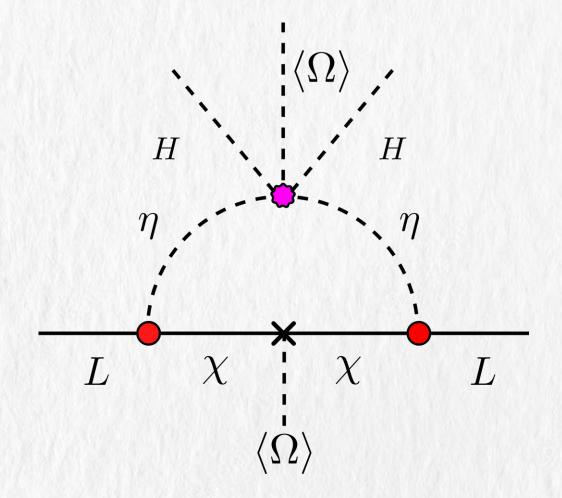
# Effective theory!



## Symmetry breaking

• The first symmetry breaking step comes from  $v_\Omega$  and  $v_\xi$ 

$$v_\Omega, v_\xi \gg v_H \gg v_\Delta$$

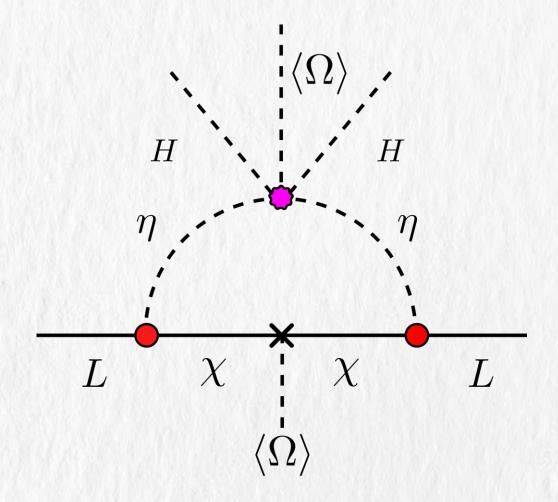


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$$v_\Omega$$
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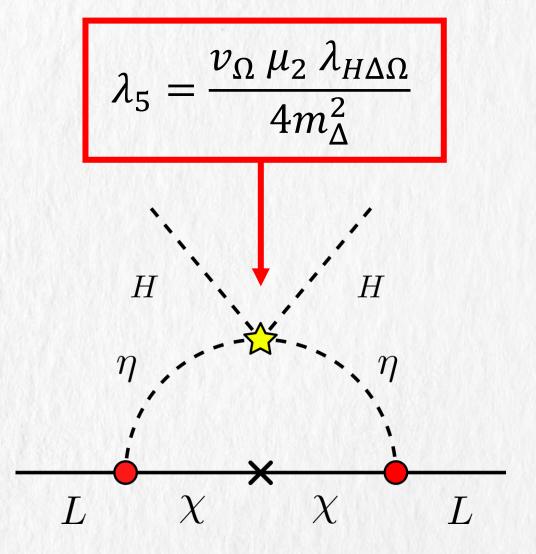
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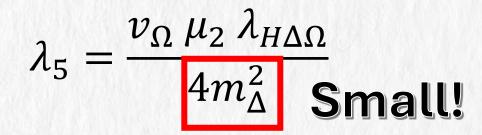
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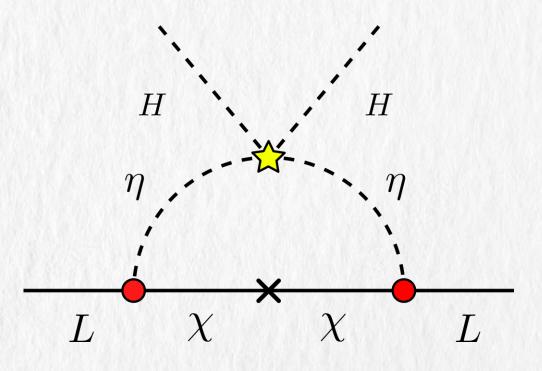
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• Our model also induces an accidental  $Z_2$  symmetry



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$$\eta \sim (\mathbf{1}, \mathbf{2}) \longrightarrow \mathbf{2}$$

$$\chi \sim (\mathbf{2}, \overline{\mathbf{2}}) \longrightarrow \mathbf{3} \oplus \mathbf{1}$$

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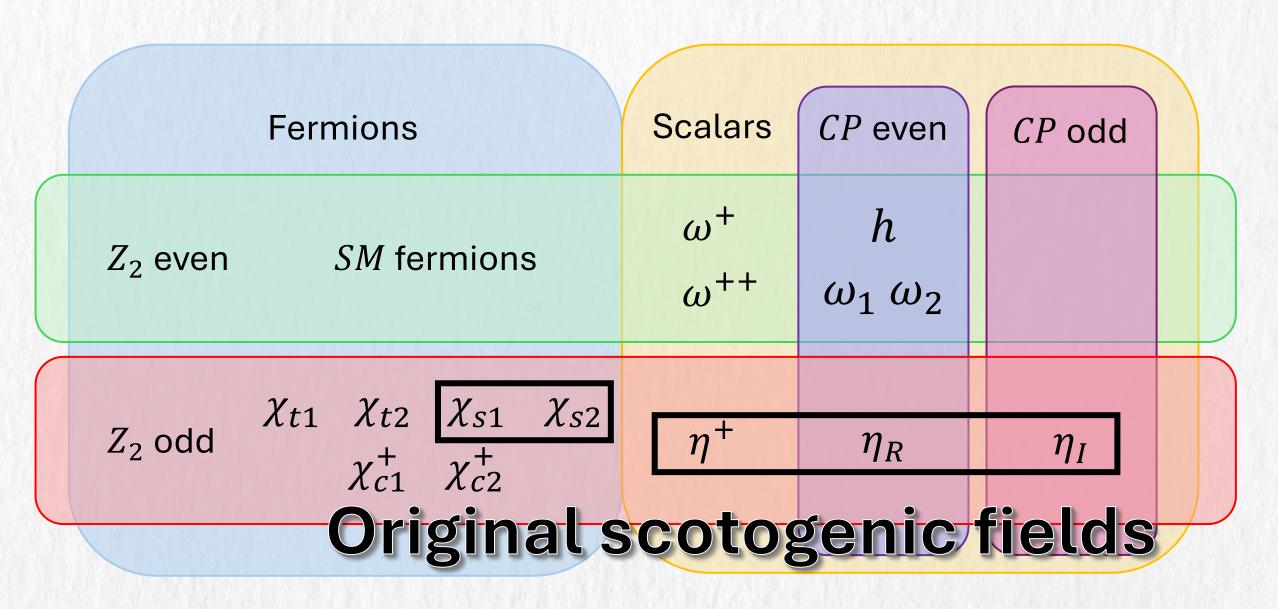
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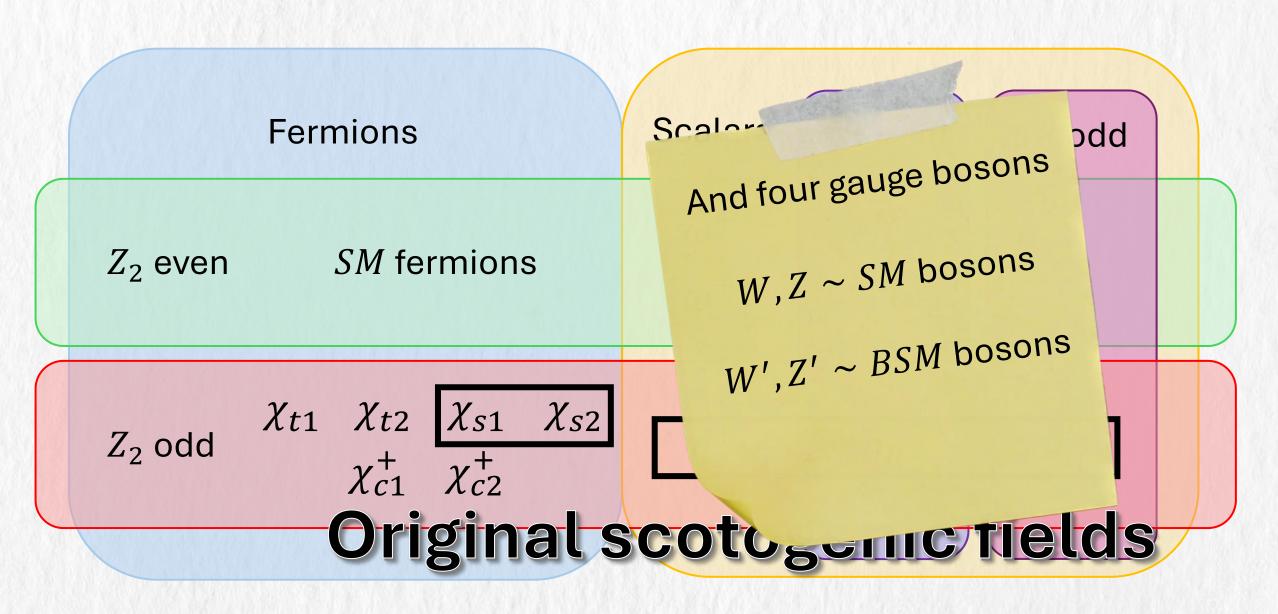


Fermions		Scalars	CP even	(CP odd	
$Z_2$ even	SM fermions	$\omega^+$ $\omega^{++}$	$h$ $\omega_1  \omega_2$		
$Z_2$ odd	$\chi_{t1}$ $\chi_{t2}$ $\chi_{s1}$ $\chi_{s2}$ $\chi_{c1}^+$ $\chi_{c2}^+$	$\eta^+$	$\eta_R$	$\eta_I$	

#### All BSM masses have contributions from $v_\Omega$



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Very important!

Scotogenic mechanism from an extended \$SU(2)\_1 \times SU(2)\_2 \times U(1)\_Y\$ electroweak symmetry

 $\equiv$ 

We propose an extension of the electroweak sector of the Standard Model in which the gauge group \$SU(2)\_L\$ is promoted to \$SU(2)\_1 \times SU(2)\_2\$. This framework naturally includes a viable dark matter candidate and generates neutrino masses radiatively à la Scotogenic. Our scenario can be viewed as an ultravioles expension et the Scotogenic.

(it says phenomenology somewhere in here)



Presentador Javier Perez-Soler

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- Assuming  $g_1 = g_2$  for simplicity, the **EW** precision parameter...

$$\rho = 1.00031 \pm 0.00019$$

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Scotogenic mechanism from an extend \$SU(2)

These limits also suppress gauge boson mixing

gauge

 $\theta_c, \theta_n \le 10^{-3}$ 

New physics scales are multi-TeV...
 cannot see anything

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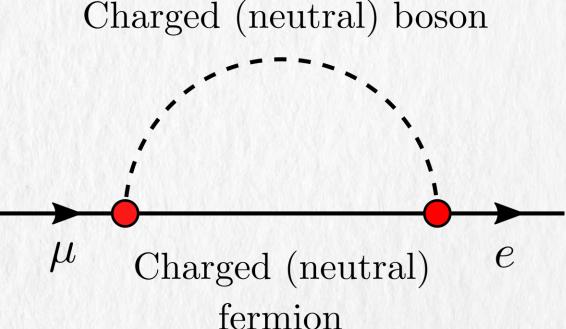
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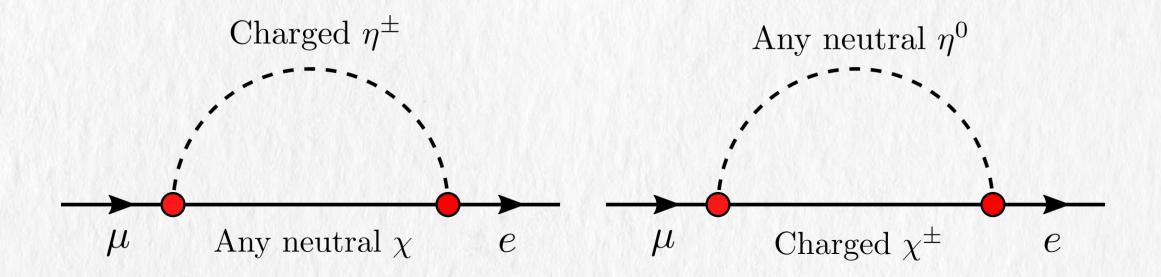
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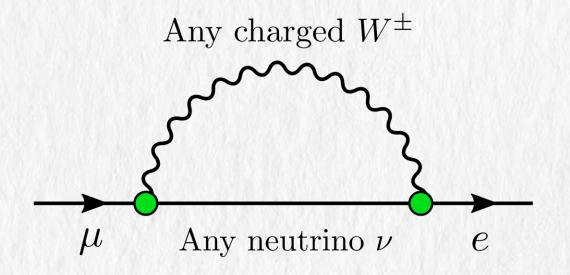
•  $\mu \rightarrow e + \gamma$  is a classic choice

• BR(
$$\mu \to e + \gamma$$
) < 1.5 · 10<sup>-13</sup> (90% CL, MEG II 2021-2022)

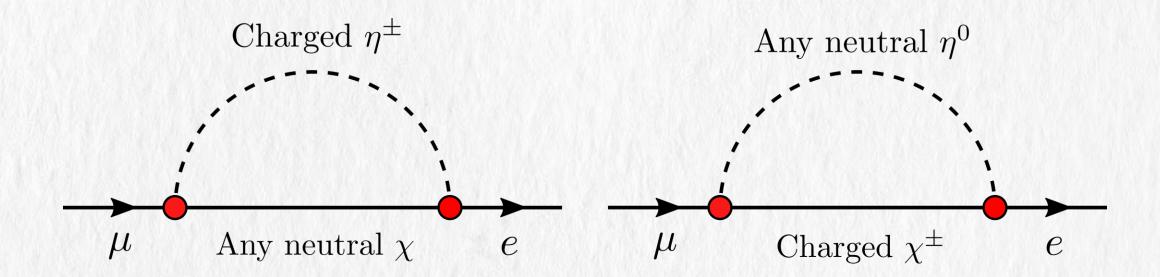


(All particles mass eigenstates)



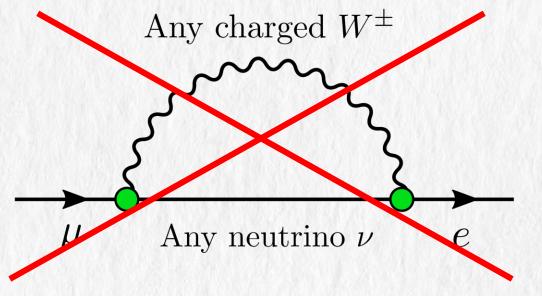


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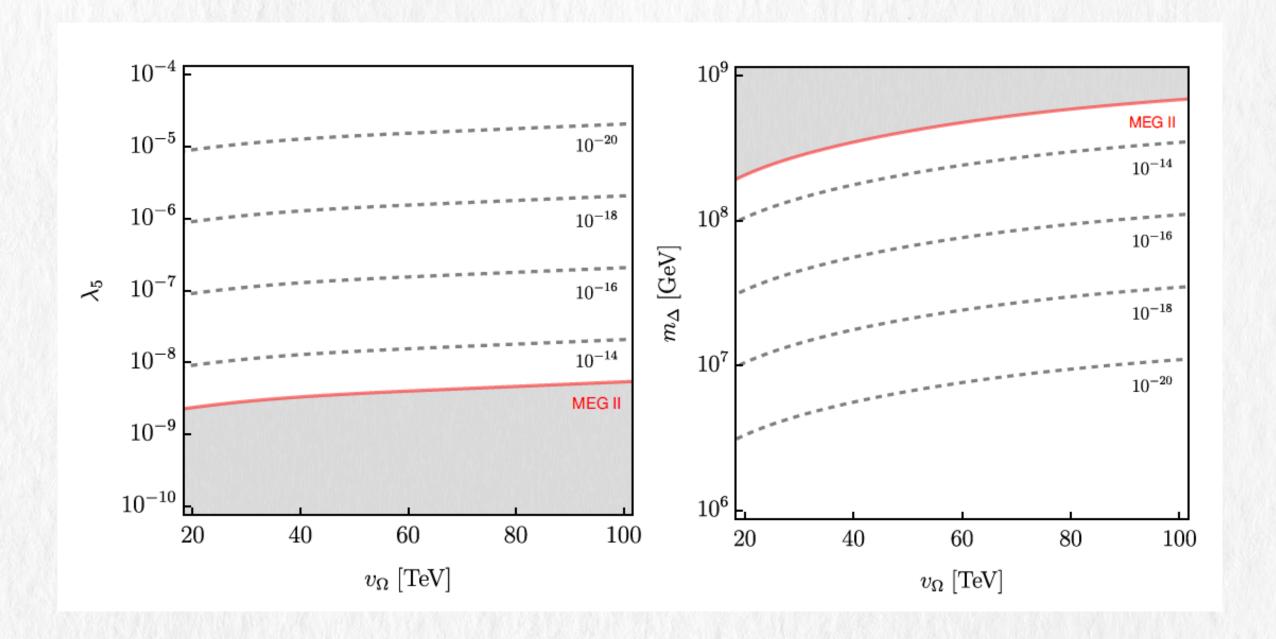


#### Mega suppressed

$$\sim \left(\frac{m_{\nu}}{m_W}\right)^4$$

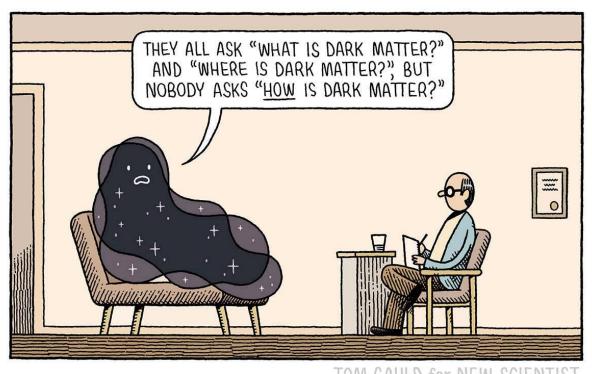


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### Final point: Dark matter

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TOM GAULD for NEW SCIENTIST

Actual picture displayed on the door to my office

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- For scalars, it's the scotogenic guys  $\eta_R$  and  $\eta_I$
- For fermions, we have two,  $\chi_{t1}$  and  $\chi_{s1}$ ; coming from the neutral components of





TOM GAULD for NEW SCIENTIS

Actual picture displayed on the door to my office

### Summary of the summary

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Can get BSM phenomenology at the TeV scale



## What to do if i have theories?

Question

I contacted a college and they ignored me







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