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Conselleria de Educació, Cultura,
Universitats y Empleo

PBH PROBES OF HNLS

based on JCAP 04 (2025) 018

in collaboration with Y.F. Perez-Gonzalez & V. De Romeri

Agnese Tolino

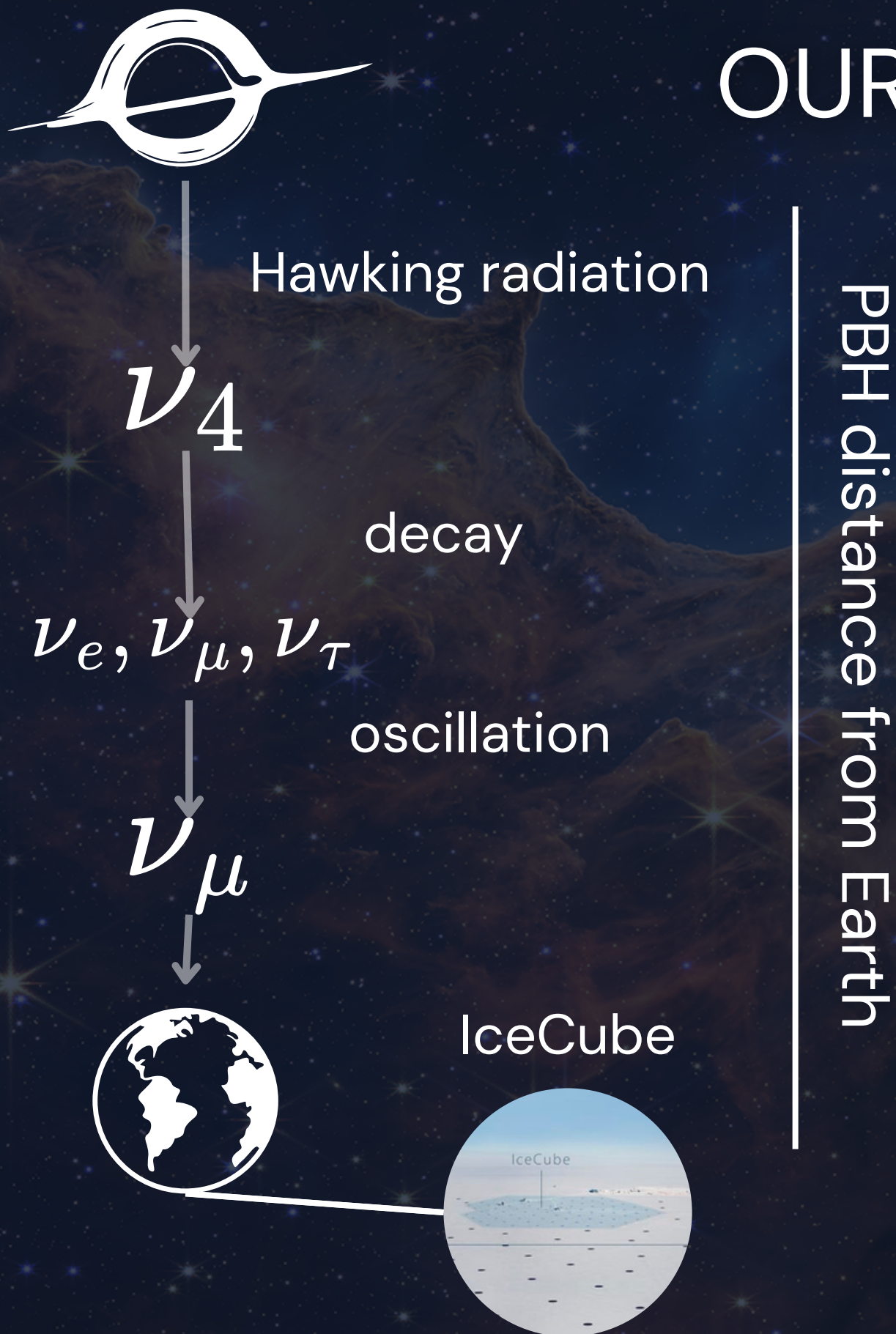
IFIC (CSIC-UV)

November 6th, 2025

TeVPA, Valencia

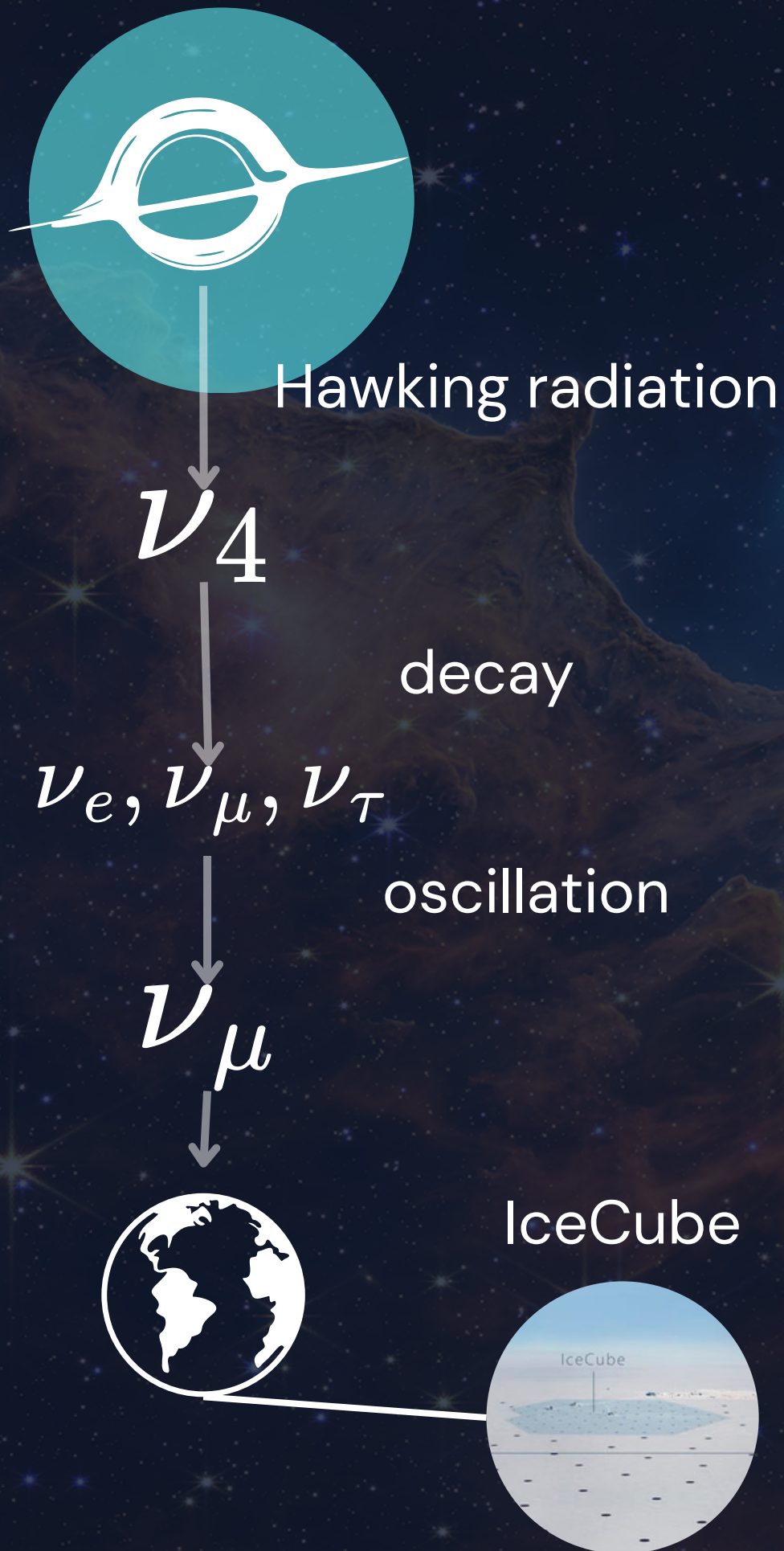


OUR WORK... IN A NUTSHELL



In JCAP 04 (2025) 018,
we estimated the combined **sensitivities**
of **IceCube & HAWK**
to Heavy Neutral Leptons (HNLs) decays
from a 100s
Primordial Black Hole (PBH) **burst**

IDENTIKIT OF PBHS

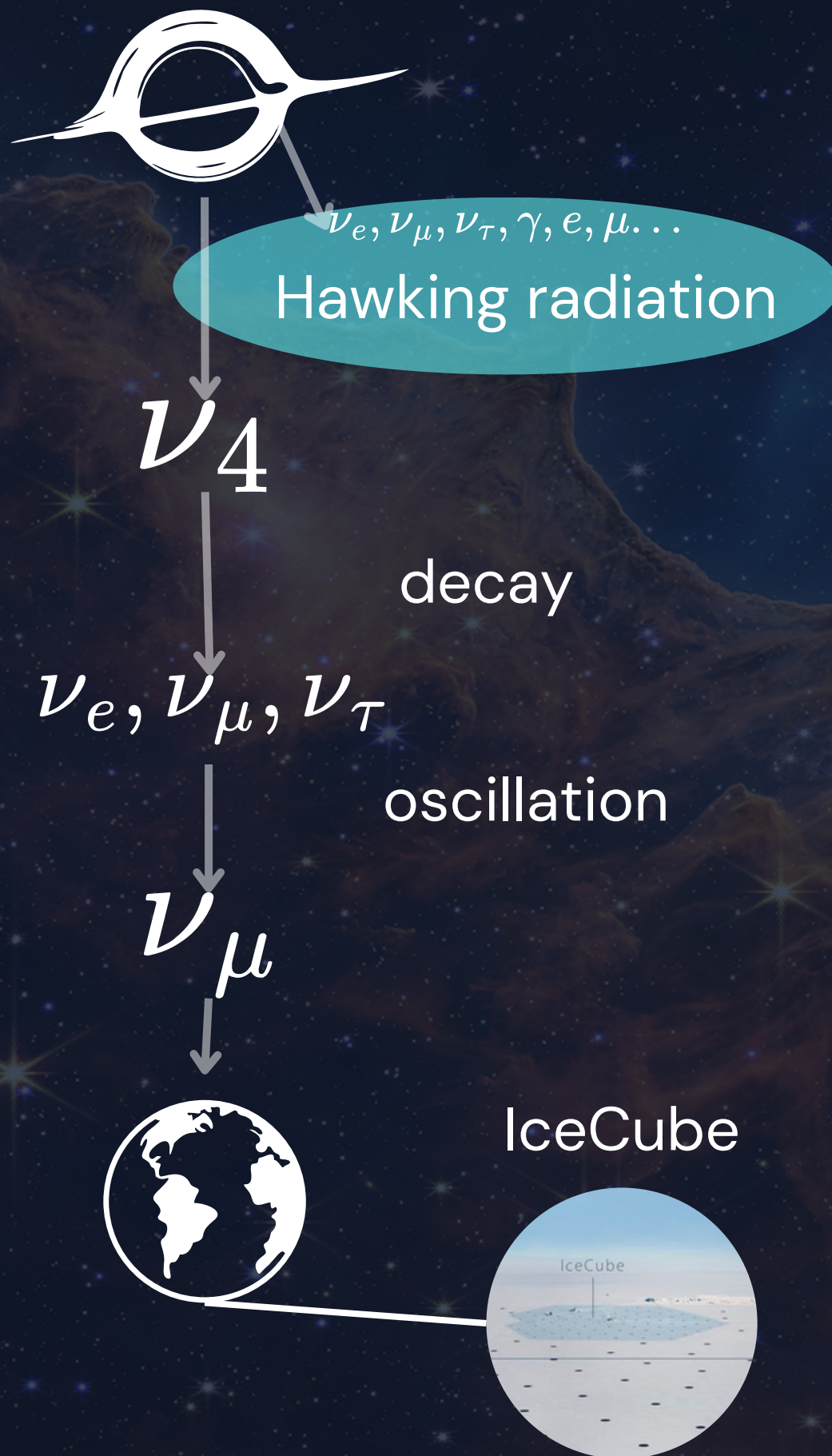


PBH distance from Earth

- Might have formed in the Early Universe by the collapse of **overdensities**
- **Schwarzschild** (i.e. non-rotating & uncharged) PBHs are uniquely described by their **mass**
- Masses span from 10^{-5} g to $10^5 M_\odot$

Hawking, Nature 248 (1974) 30–31
Carr et al., Ann. Rev. Nucl. Part. Sci. 70 (2020)
Carr et al., Rept. Prog. Phys. 84 (2021) 11, 116902

HAWKING RADIATION



- Hawking predicted that PBHs **evaporate** with a temperature

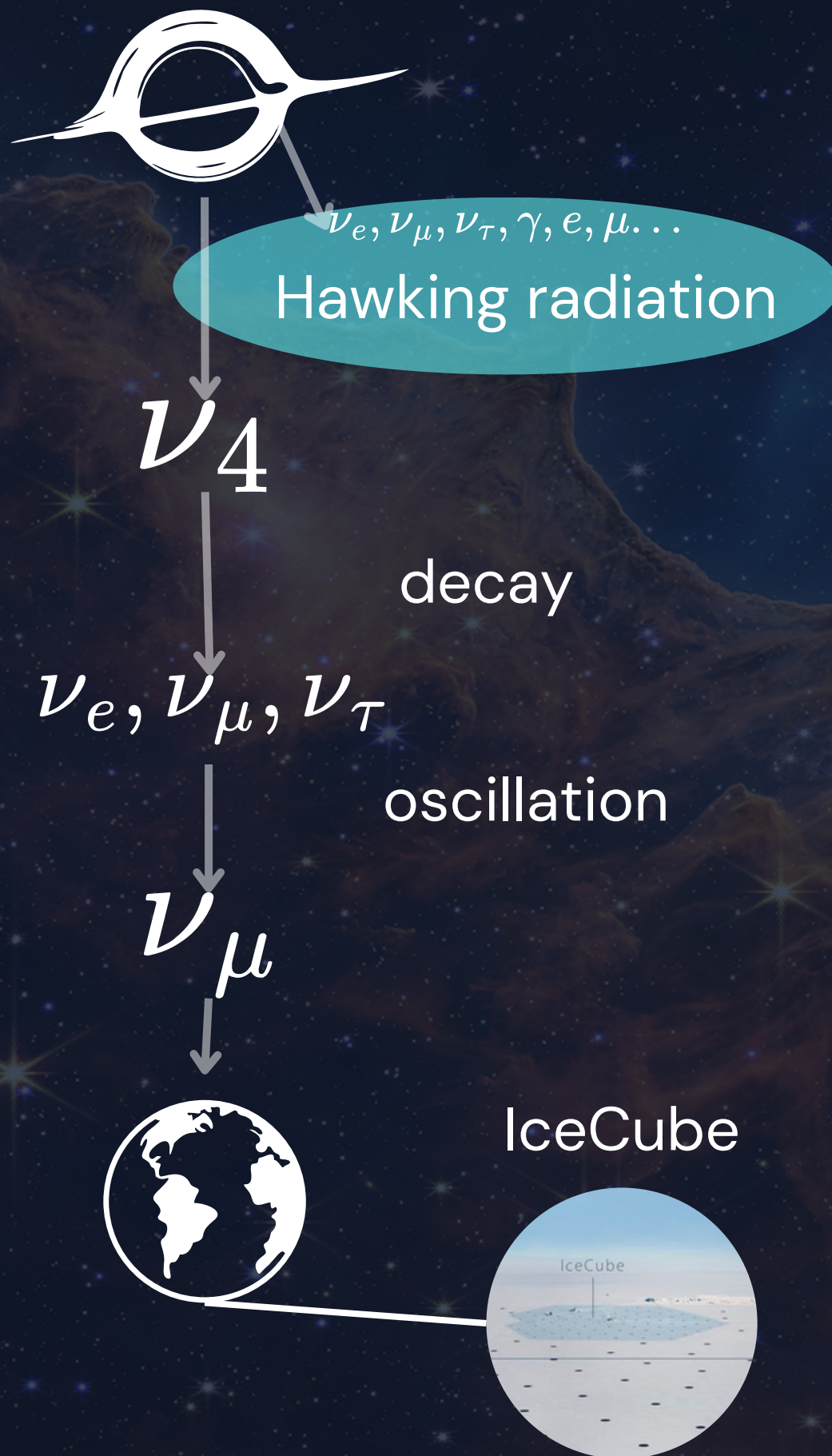
$$T = \frac{1}{8\pi G M_{\text{PBH}}}$$

- Mass loss goes as $\frac{dM}{dt} \sim M_{\text{PBH}}^{-2}$
- The evaporation corresponds to the emission of particles with a semi-thermal spectrum

$$\left. \frac{dN^i}{dEdt} \right|_{\text{prim}} = \frac{g_i \Gamma(M_{\text{PBH}}, E_i)}{2\pi \left(\exp \left\{ \frac{E_i}{T_{\text{PBH}}} \right\} - (-1)^{2s_i} \right)}$$

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 Arbey et al., Eur. Phys. J. C 79 no. 8, (2019) 693

'EXPLODING' PBHS



PBH distance from Earth

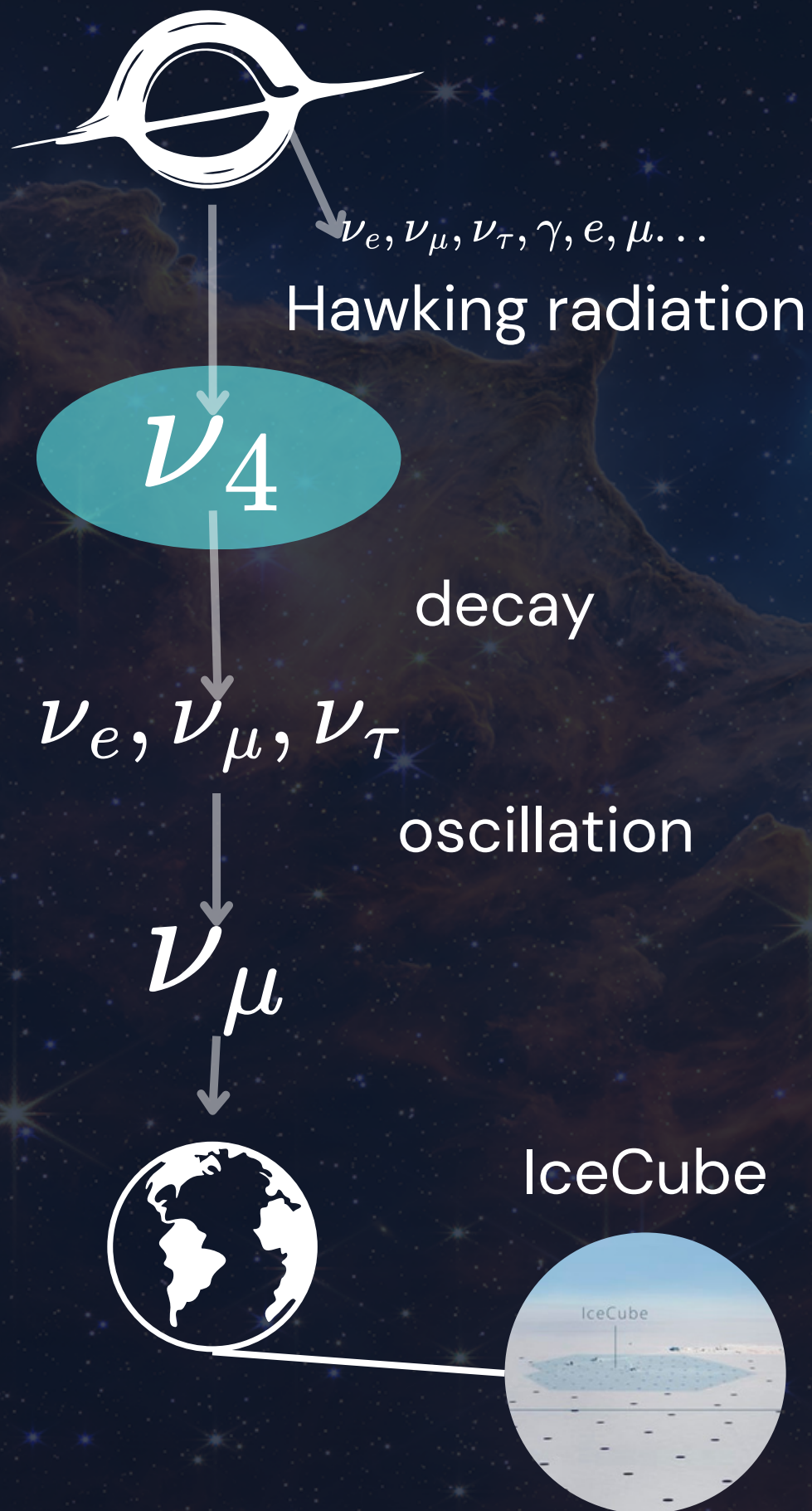
- In the final stage of evaporation, the PBH quickly becomes **hotter** and emits a **burst of particles**
- Particles with a mass up to $m_i \sim T_{\text{PBH}}$ can be emitted, **even BSM particles as HNLs!**

$$\left. \frac{dN^i}{dEdt} \right|_{\text{prim}} = \frac{g_i \Gamma(M_{\text{PBH}}, E_i)}{2\pi \left(\exp \left\{ \frac{E_i}{T_{\text{PBH}}} \right\} - (-1)^{2s_i} \right)}$$

- Our work: **1 PBH** with $M_{\text{PBH}}^{\text{in}} \sim 10^{15} \text{ g}$ exploding in a **100s burst**

Hawking, Nature 248 (1974) 30–31
 Carr et al., Ann. Rev. Nucl. Part. Sci. 70 (2020)
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HNLS



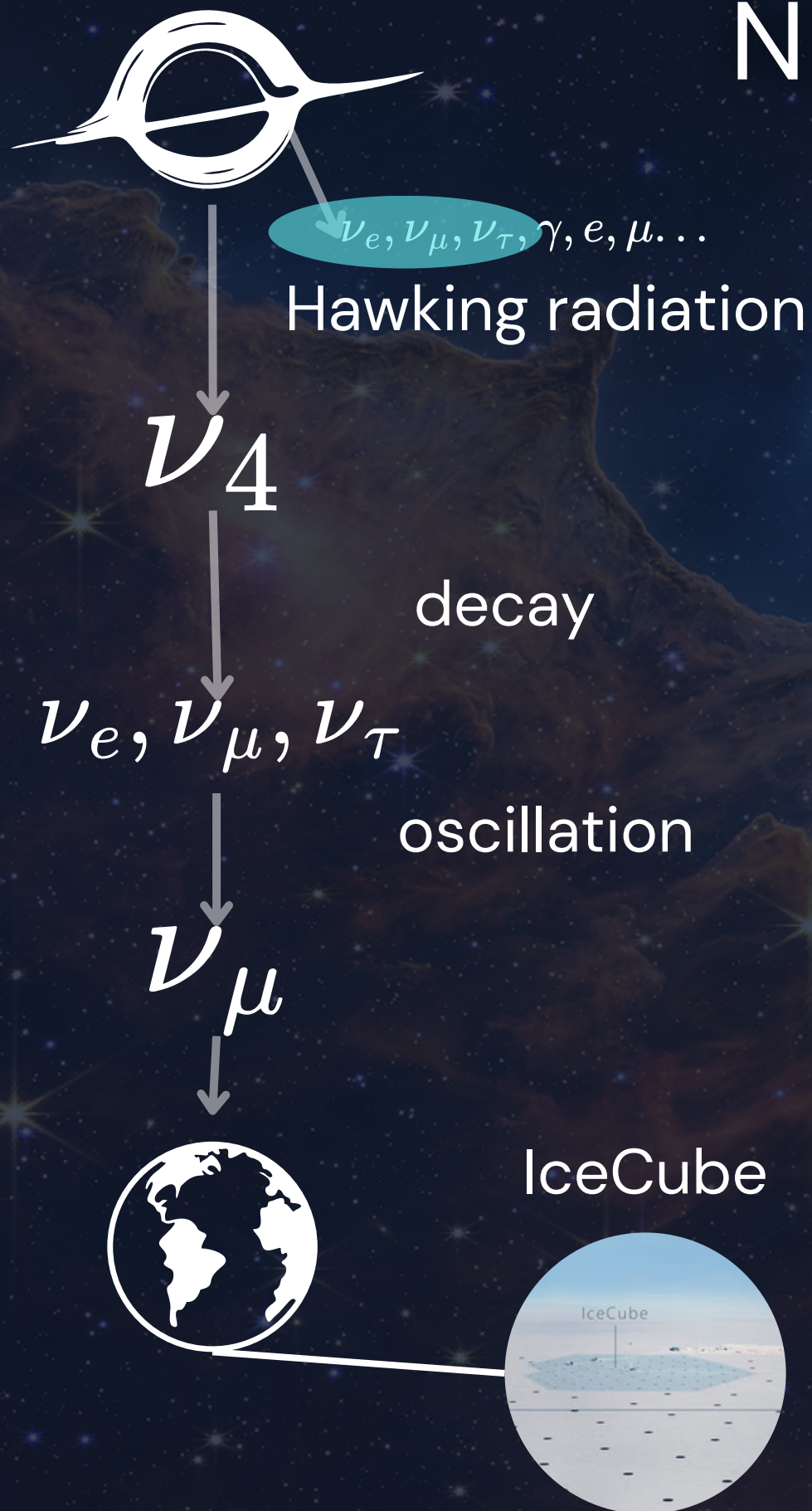
PBH distance from Earth

- Introduced to explain neutrino masses in SM extensions
- Phenomenological study: 1 HNL
- Considered the case in which only 1 active neutrino $\alpha = e, \mu, \tau$ at time mixes with the HNL:

$$|U_{\alpha 4}|^2 \neq 0$$

(So 1:0:0 : only electron neutrino mixing, and so on)

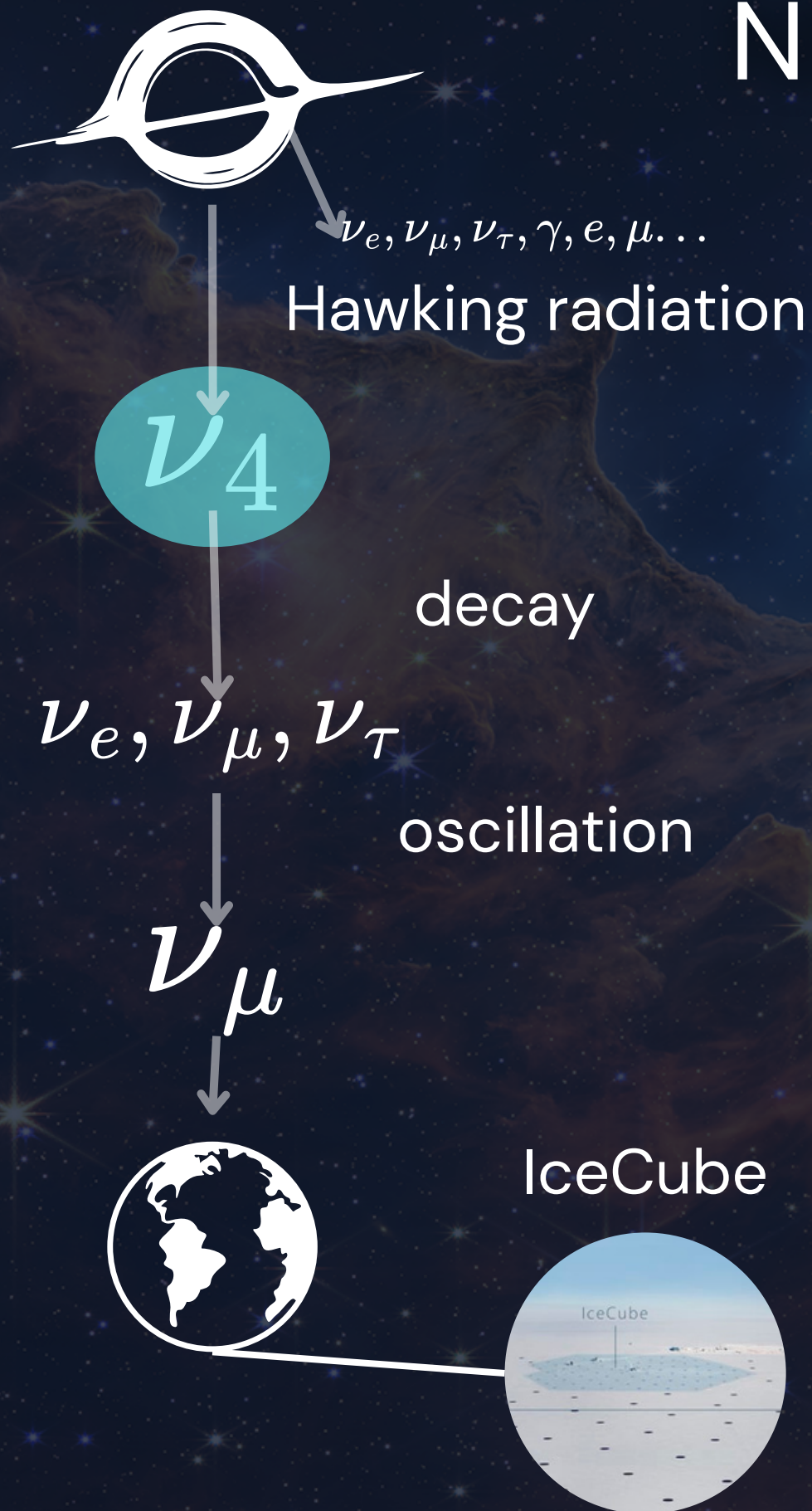
NEUTRINO SIGNALS FROM PBHS



- SM neutrinos are directly emitted by the PBH

Atre et al., JHEP 05 (2009) 030
Mastrototaro et al., JCAP 01 (2020) 010
Coloma et al., Phys. J. C 81 no. 1, (2021) 78
Akita et al., arXiv:2312.1362

NEUTRINO SIGNALS FROM PBHS

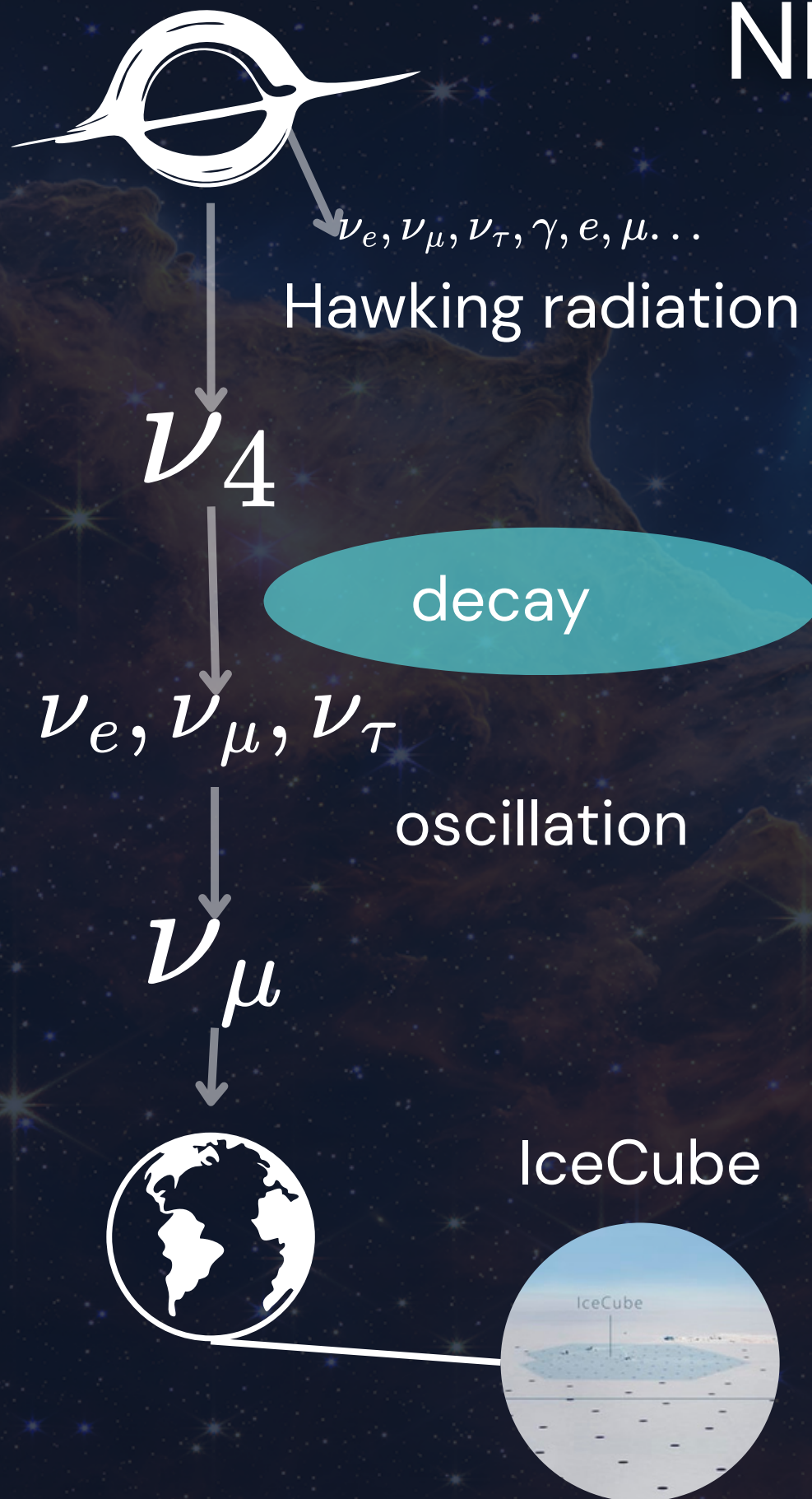


PBH distance from Earth

- SM neutrinos & HNLs are directly emitted by the PBH

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NEUTRINO SIGNALS FROM PBHS



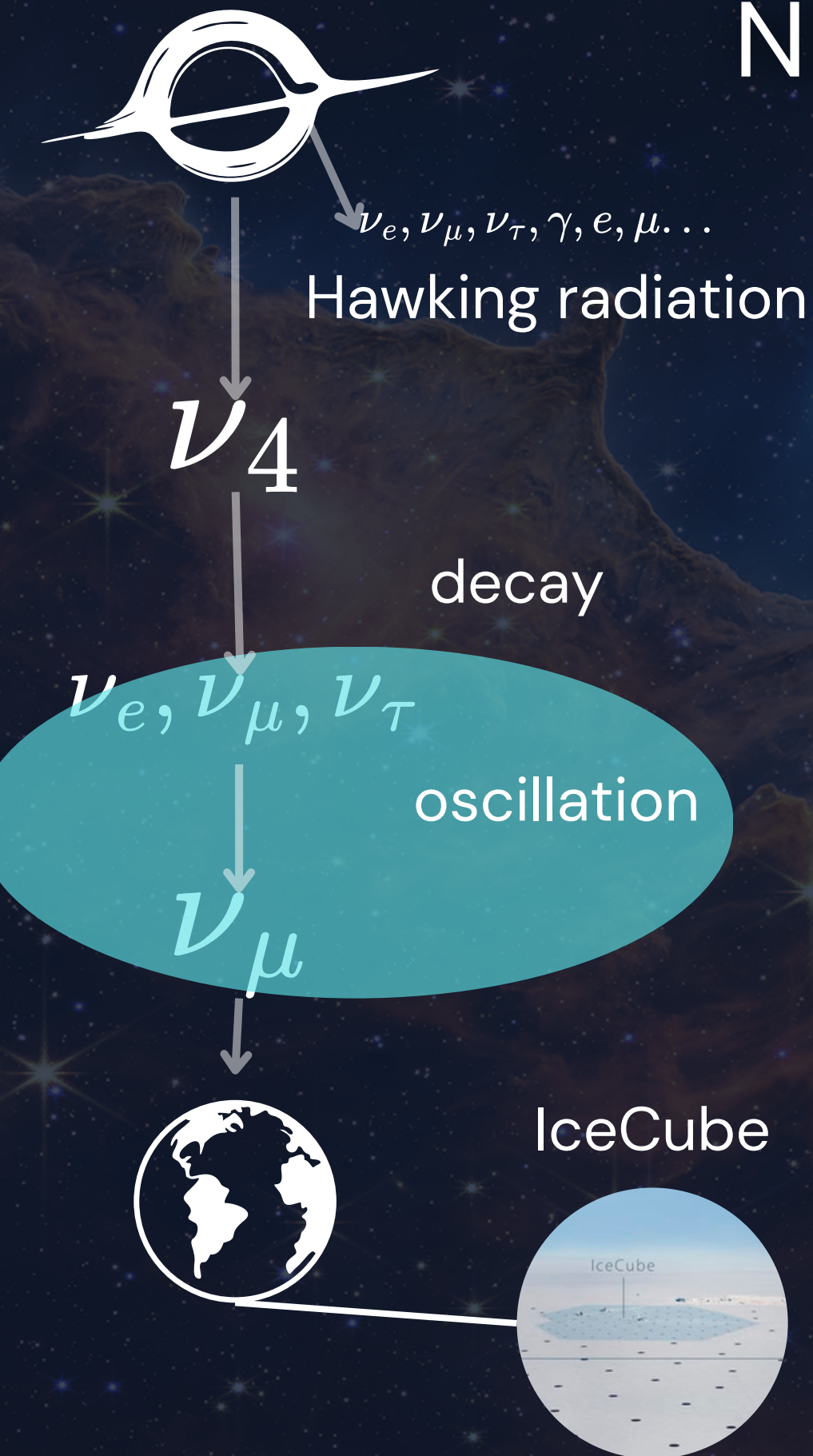
PBH distance from Earth

- SM neutrinos & HNLs are directly emitted by the PBH
- Depending on the mass, HNLs might decay into active neutrinos:

$\nu_4 \rightarrow \nu\nu\nu \ \& \ \nu_4 \rightarrow \nu\pi,$	if $m_4 \in [0.1, 1] \text{ GeV}$	Light mass regime
$\nu_4 \rightarrow H/Z\nu \rightarrow \nu \ \& \ \nu_4 \rightarrow W\mu \rightarrow \nu,$	if $m_4 \in [0.5, 2] \text{ TeV}$	Heavy mass regime

Atre et al., JHEP 05 (2009) 030
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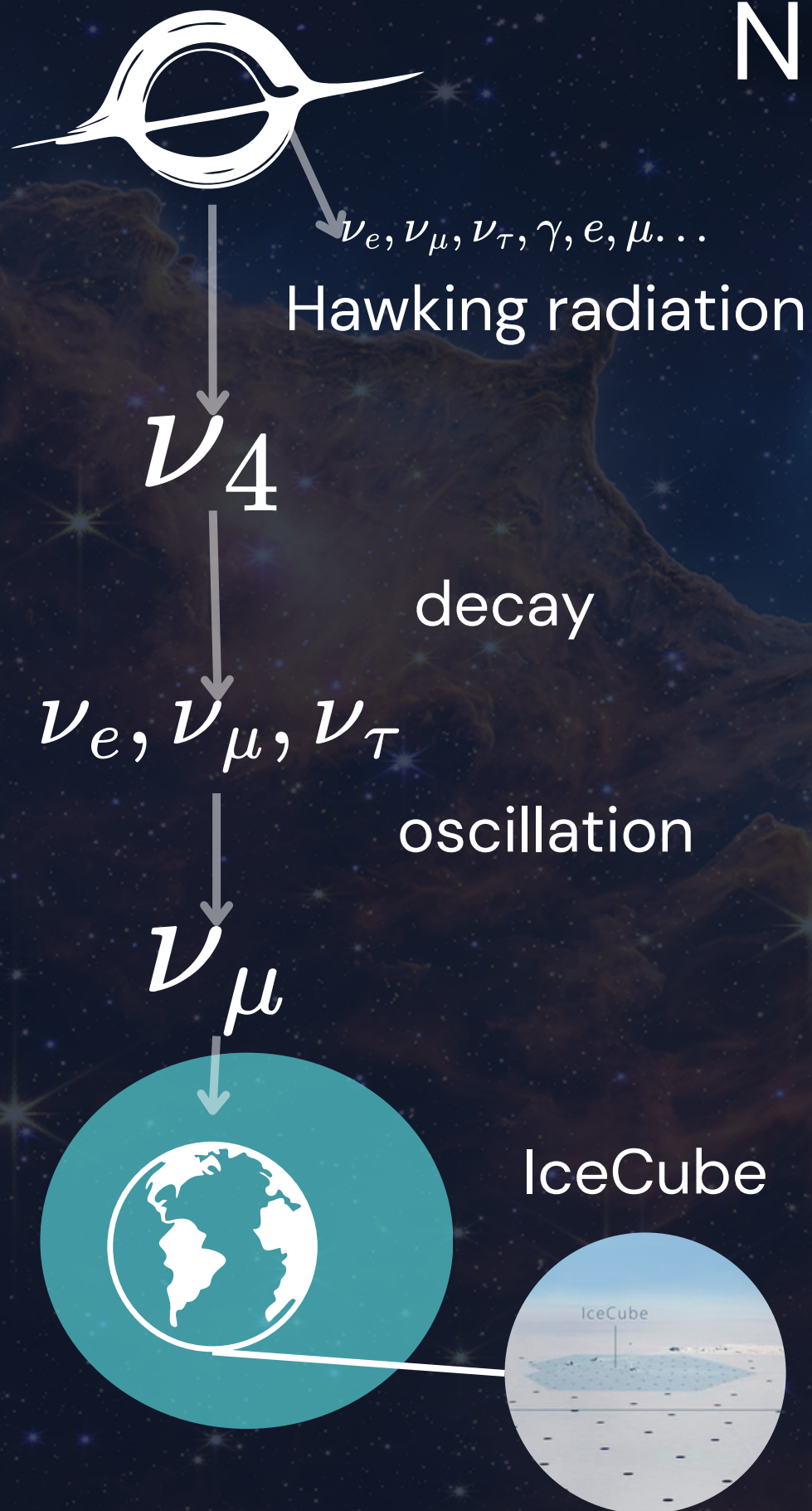
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- The active neutrinos oscillate into muon neutrinos

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NEUTRINO SIGNALS FROM PBHS



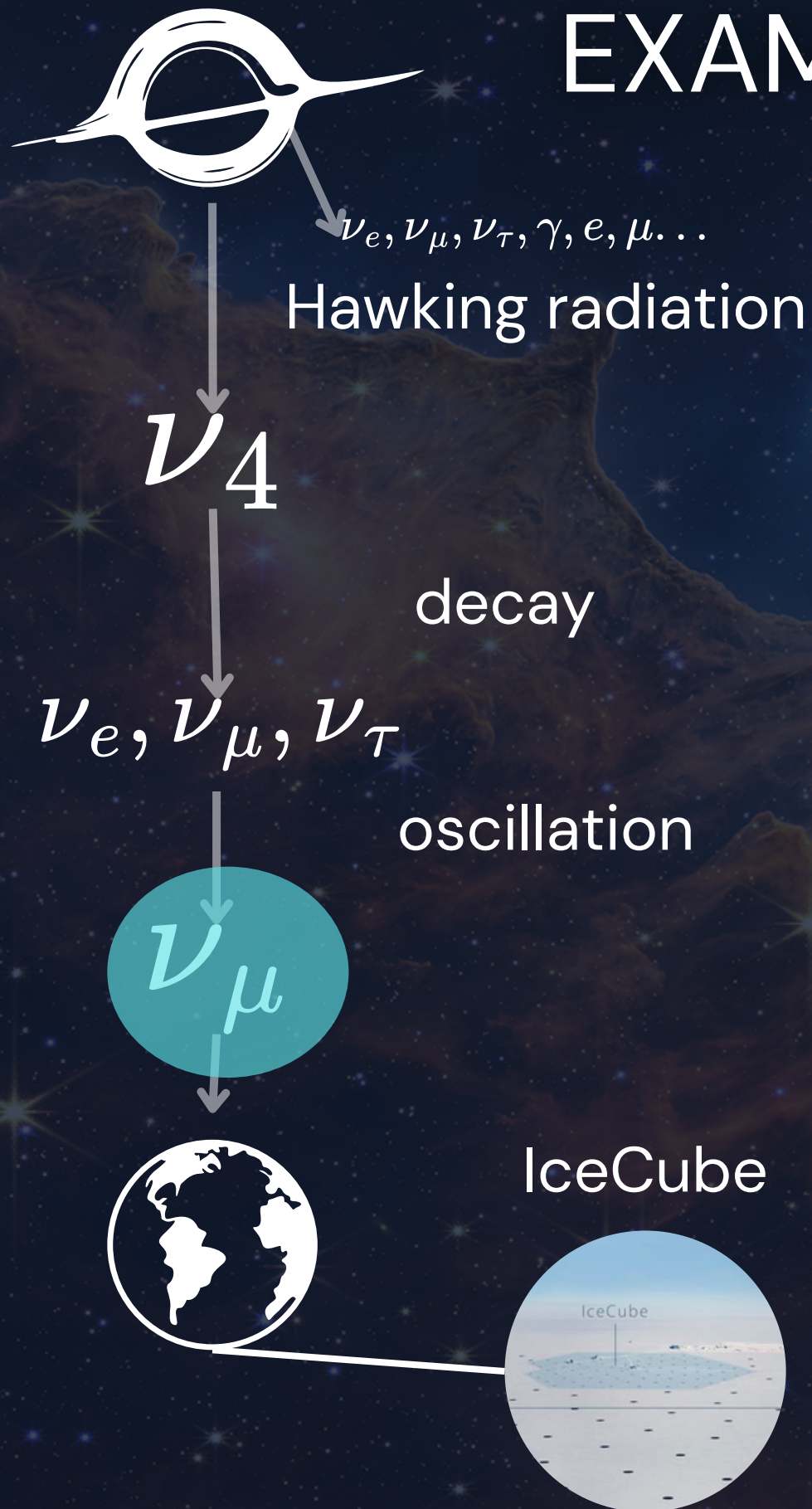
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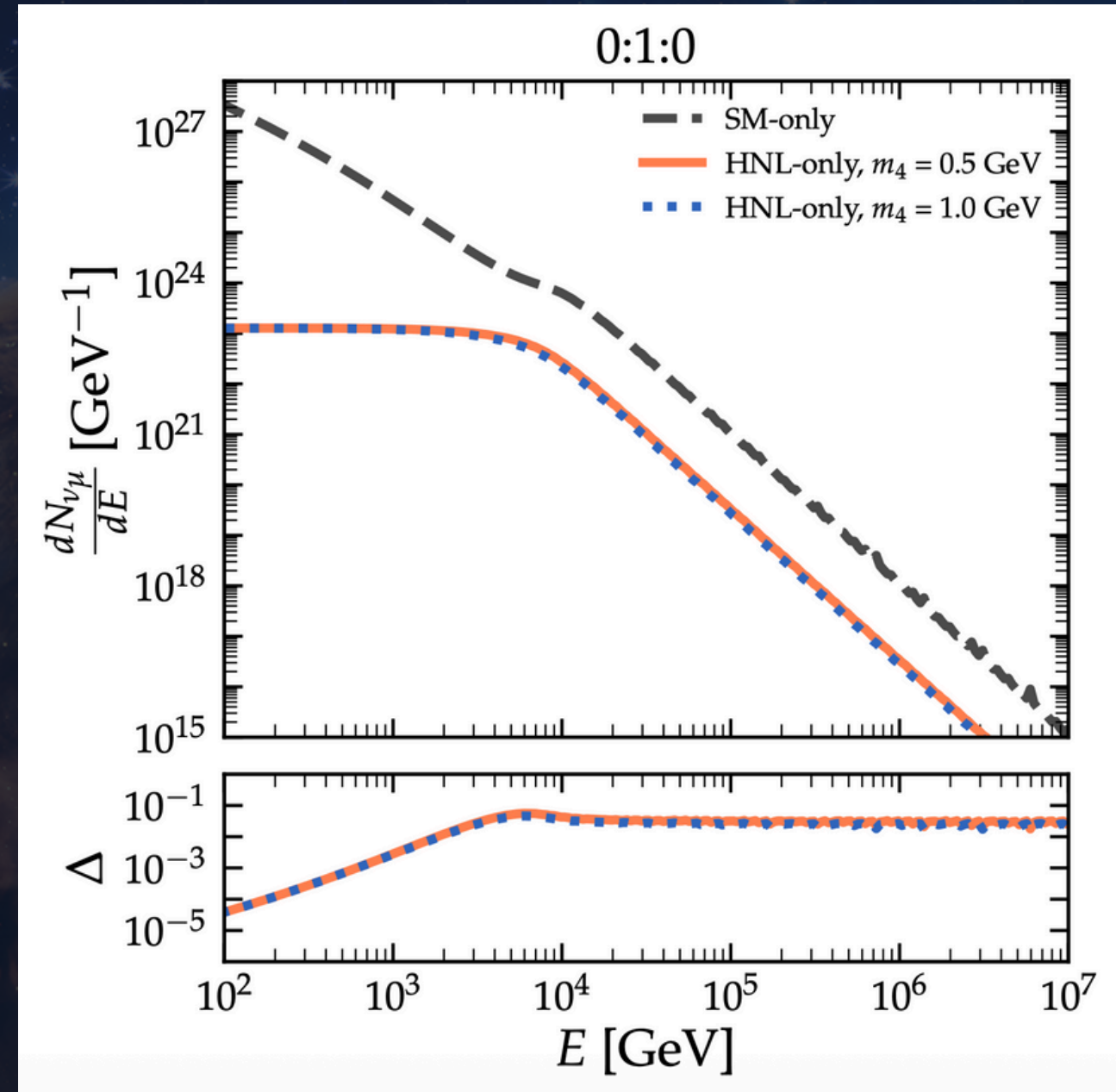
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- The active neutrinos oscillate into muon neutrinos
- Overall neutrino spectrum at Earth from a 100s PBH burst:
SM-only + HNL-decay contributions

Atre et al., JHEP 05 (2009) 030
 Mastrototaro et al., JCAP 01 (2020) 010
 Coloma et al., Phys. J. C 81 no. 1, (2021) 78
 Akita et al., arXiv:2312.1362

EXAMPLE OF SM-ONLY + HNL-ONLY NEUTRINO SPECTRA



PBH distance from Earth



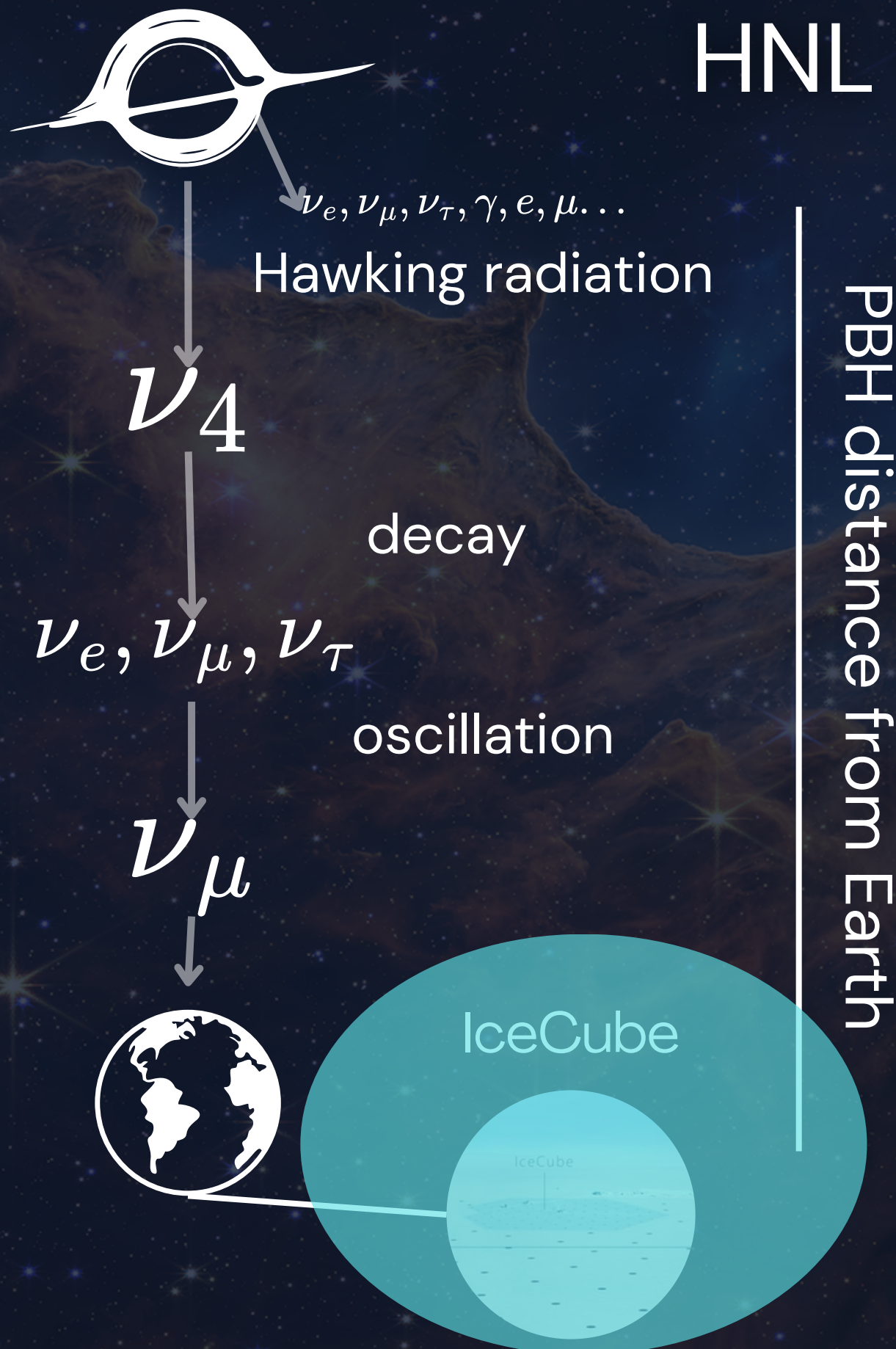
Expected muon neutrino spectrum from HNL-only (color) and SM-only (grey) contributions in a 100s PBH burst

$$\frac{dN_{\nu_\mu}}{dE} \propto B_\alpha(\text{channel})$$

Note that if only one active neutrino mixes with the heavy, the branching ratio, and hence the spectrum, will not depend on the mixing!

The spectra are obtained through BlackHawk + HDMSpectra

HNL SIGNATURES AT ICECUBE

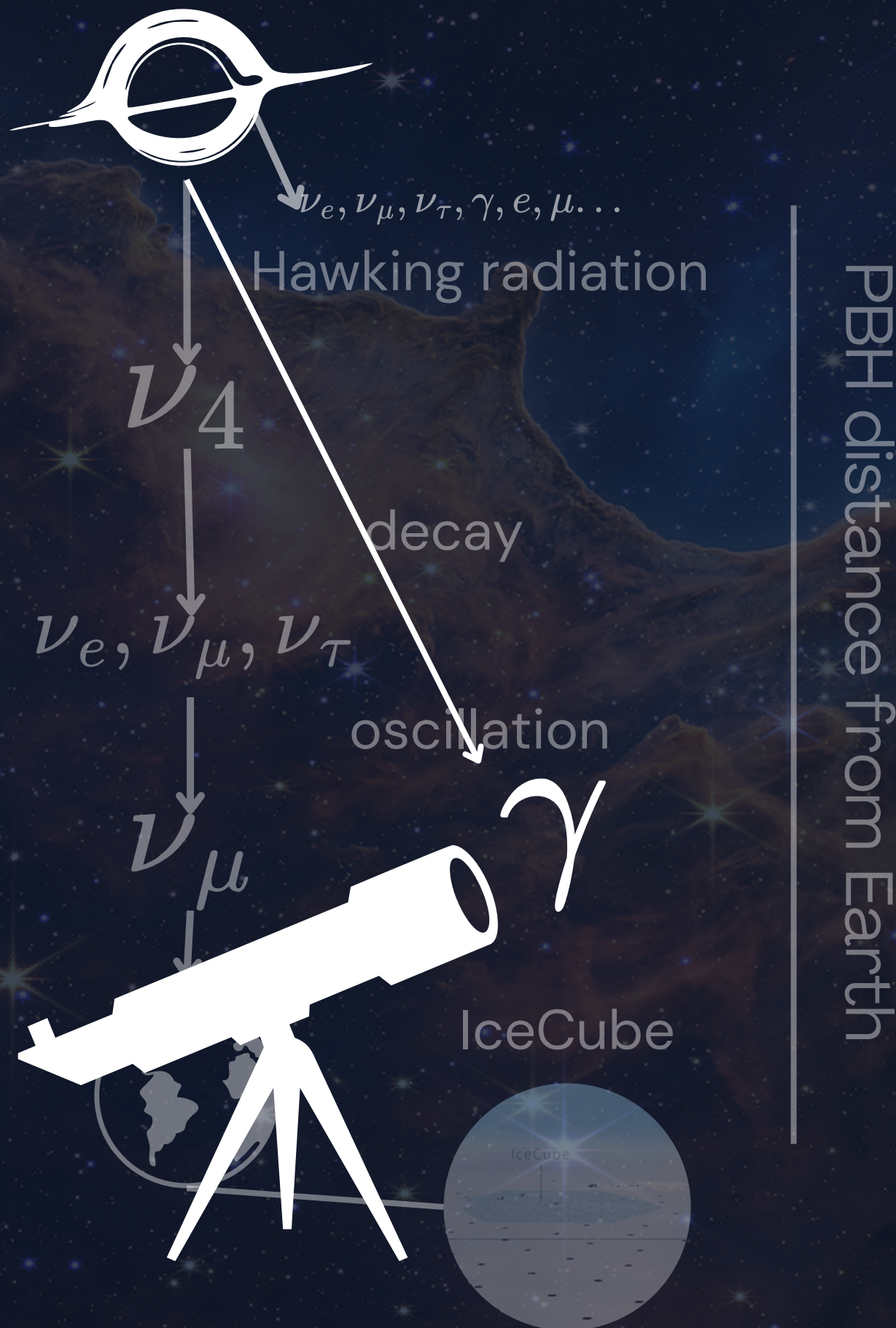


- ν_μ directly emitted by the PBH & through HNL decay are extremely boosted
- IceCube would be able to detect them, as it is sensitive to the right energy range :

100 GeV – 100 PeV

Signature: excess of muon neutrinos (with respect to the SM-only case) **from the PBH burst @IceCube due to HNL decay!**

... AND PHOTONS



- Photons are a smoking gun of PBH burst and can be produced in HNL decays as well
- Strongest constraints from H.E.S.S. & overdensities allow 1 exploding PBH at max 1 pc from us
- A combined $\nu + \gamma$ analysis is necessary to erase the **degeneracy** between the PBH position and the HNL contribution!
 - Photons @Hawk included in the statistical analysis

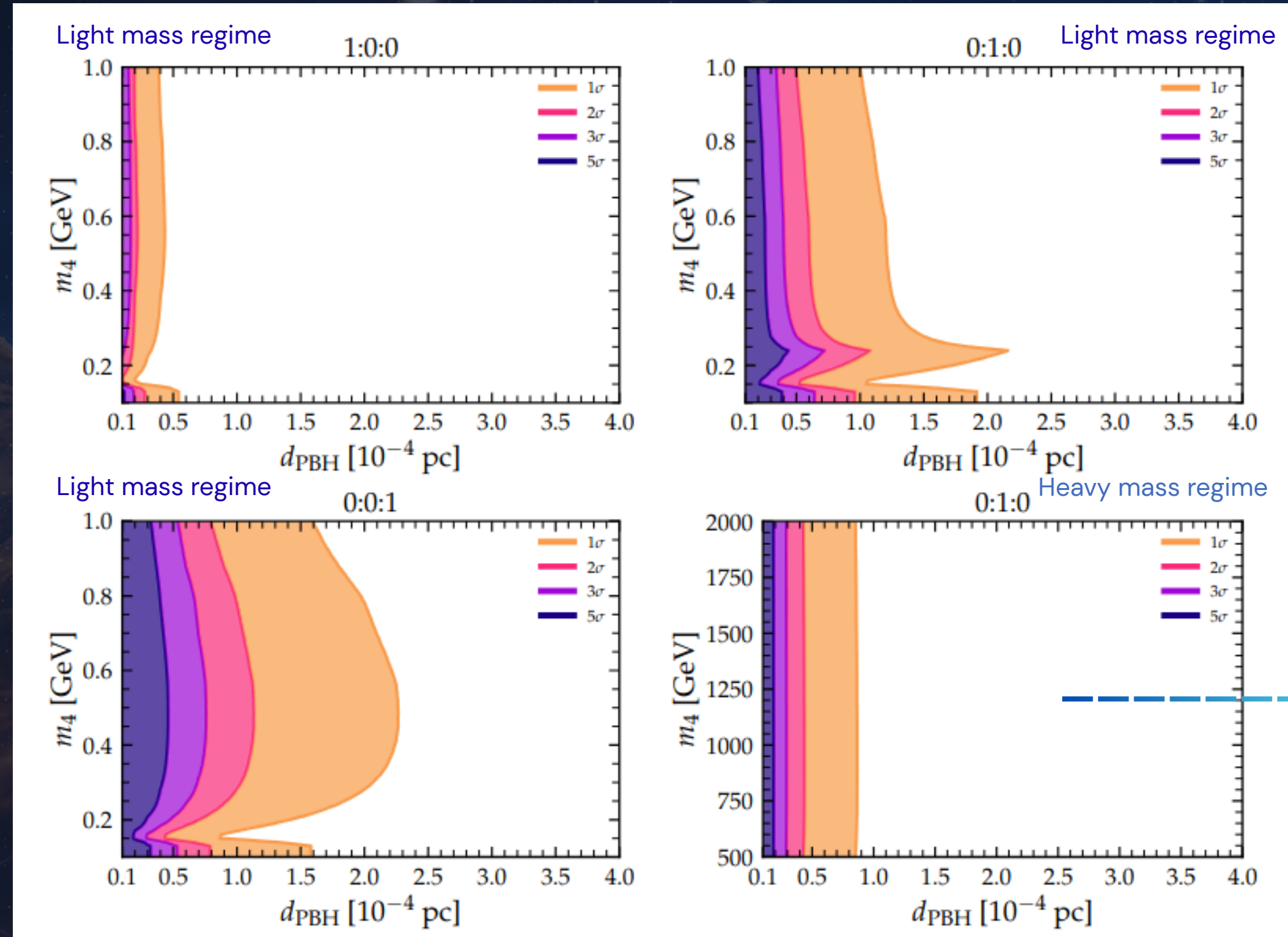
ANALYSIS SCHEME

Evaluated the expected number of muon ν at IceCube and γ at HAWK from northern hemisphere emitted in a **100s PBH burst** considering both **HNL + SM contributions**

Estimated the **IceCube & HAWK sensitivities** to HNL decays with a χ^2 analysis **including** the **uncertainty** over the **PBH distance** from Earth

SELECTED RESULTS

ICECUBE & HAWK SENSITIVITIES

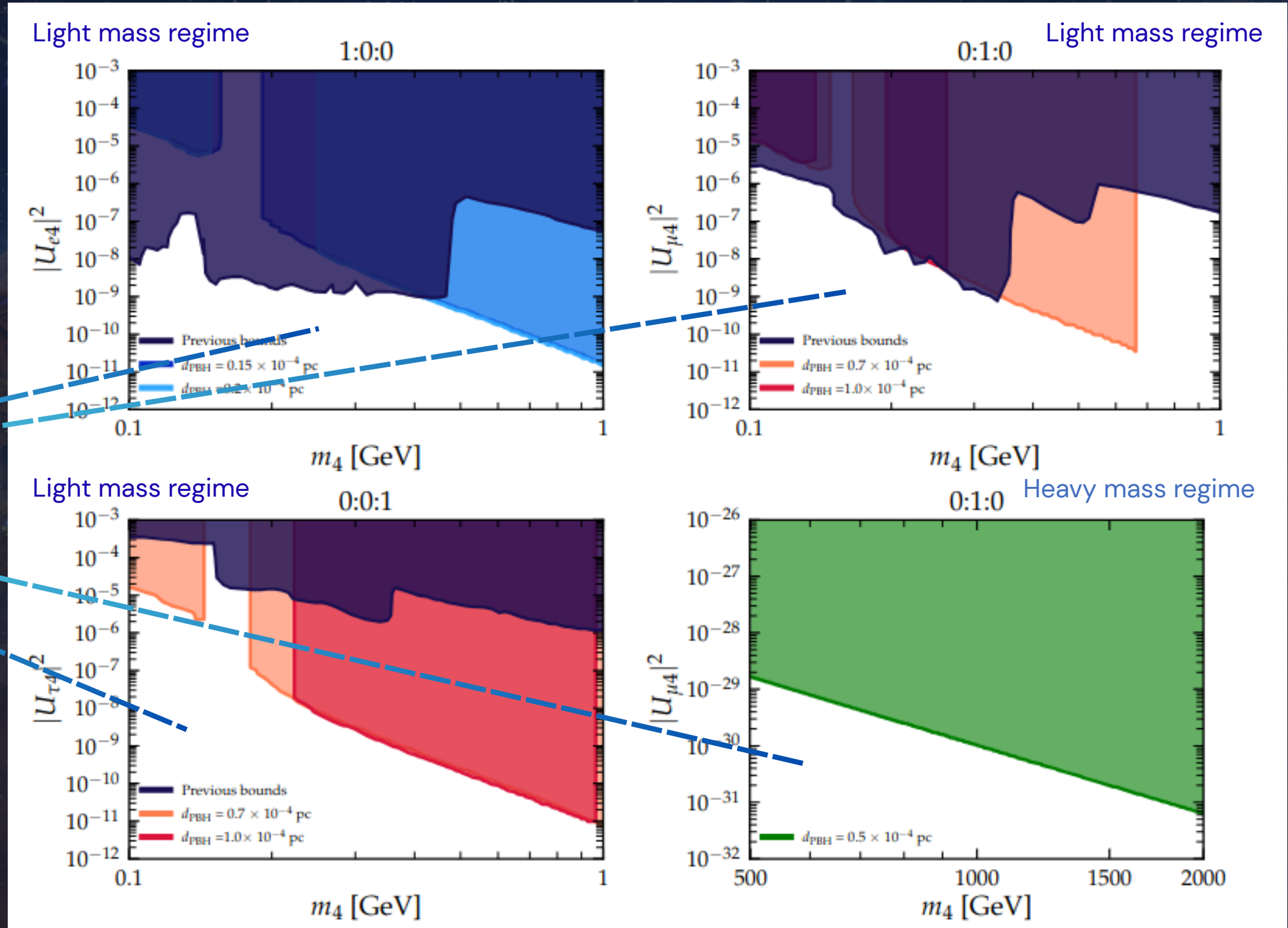


Combined IceCube & HAWK sensitivities to HNLs from a PBH burst lasting 100s

Expected IceCube & HAWK sensitivities at 90% CL for a 100s PBH burst

Same results can be casted in
the HNL mixing – mass plane

$$d_{\text{decay}} \leq d_{\text{PBH}}$$



CONCLUSIONS



1

We evaluated the muon ν signal at IceCube and γ at HAWK from 100s PBH burst

- included HNL decays
- considered for HNLs 2 mass ranges & 3 mixing scenarios

2

In the statistical analysis, we included both neutrinos @IceCube & photons @HAWK to break the degeneracy between the HNL contribution and the PBH position

3

We found that

- the HNL mass [0.1–1] GeV range could be probed at IceCube and Hawk if $d_{\text{PBH}} \sim 10^{-4} \text{ pc}$ for HNL mixings 0:1:0 and 0:0:1
- the [0.5–2] TeV range at $d_{\text{PBH}} \leq 10^{-4} \text{ pc}$ for 0:1:0



TAKE-HOME MESSAGE

Joint analyses between IceCube and gamma-ray observatories like HAWK would provide stringent limits on HNL masses and mixings!



TAKE-HOME MESSAGE

Joint analyses between IceCube and gamma-ray observatories like HAWK would provide stringent limits on HNL masses and mixings!

THANK YOU FOR YOUR ATTENTION!
QUESTIONS?

BACKUP SLIDES

PBH EVAPORATION

A PBH that is exploding now had at the time of its formation (during radiation domination) a mass equal to 10^{15} g



At the start of the 100s of the burst, i.e. the last 100s of its life, the PBH has a mass of 6.2×10^9 g

So we compute the time-integrated spectra with BlackHawk for a PBH that has an initial mass of 6.2×10^9 g, as we are interested only in the last 100s of its life:

$$\frac{dN_i}{dE} = \int_0^\tau dt \frac{d^2 N_i}{dE dt}(M(t)) .$$

NEUTRINO SPECTRA FROM HNL DECAYS

[0.1 – 1] GeV:

$$\begin{aligned}\nu_4 &\rightarrow \nu_\alpha \pi^0, \\ \nu_4 &\rightarrow \nu_\alpha \nu_\ell \bar{\nu}_\ell \quad (\ell = e, \mu, \tau),\end{aligned}$$

[0.5 – 2] TeV

$$\begin{aligned}\nu_4 &\rightarrow W^\pm \mu^\mp, \\ \nu_4 &\rightarrow Z^0 \nu_\mu, \\ \nu_4 &\rightarrow H^0 \nu_\mu.\end{aligned}$$

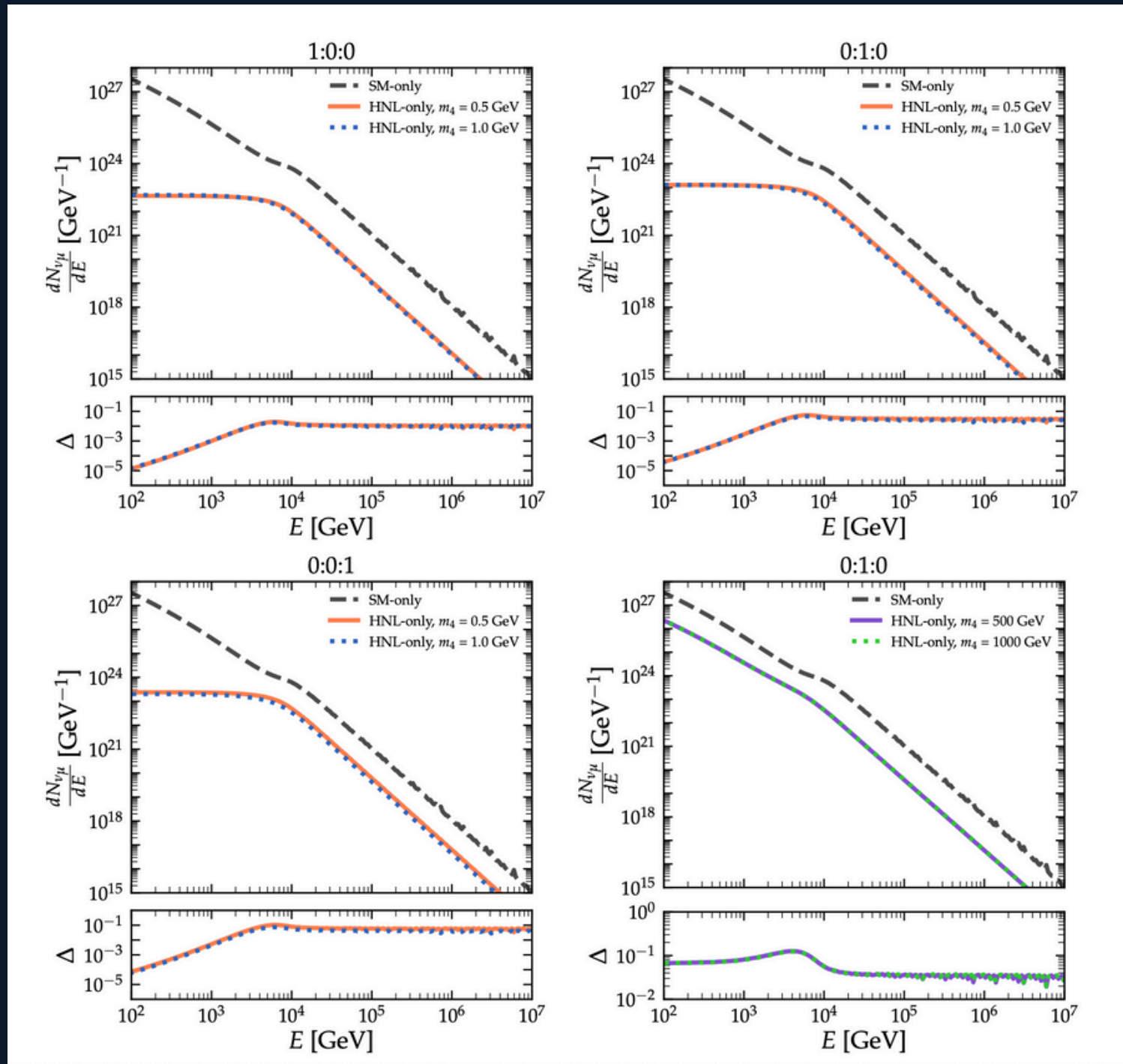
+ charge-conjugated channels

Time-integrated spectra of active neutrinos from an HNL decay:

$$\frac{dN_\alpha}{dE} = \mathcal{B}_\alpha \int d\cos\theta \int_{E_{s,\min}}^{E_{s,\max}} dE_s \frac{1}{\gamma_s (1 + \beta_s \cos\theta)} \frac{dN_s}{dE_s} \mathcal{F}_\alpha \left[\frac{E}{\gamma_s (1 + \beta_s \cos\theta)}, \cos\theta \right], \quad (20)$$

If only one active neutrino mixes with the HNL, the branching ratio does not depend on the mixing

NEUTRINO SPECTRA FROM HNL DECAYS



Light mass regime:

Angular distributions depend only on the neutrino energy in the HNL rest frame \rightarrow contribution peaks near the primary bump.

Heavy mass regime:

Angular distributions arise from multiple secondary processes \rightarrow spectrum resembles the SM-like shape.

Expected muon neutrino spectrum from HNL-only (color) and SM-only (grey) contributions in a 100s PBH burst

PHOTON SPECTRA

[0.1 – 1] GeV

$$\nu_4 \rightarrow \pi^0 \nu_\alpha \rightarrow \gamma \gamma \nu_\alpha.$$



$$\left. \frac{dN_\gamma}{dE} \right|_{\text{HNL}} = 2 \int_{E_{\min}}^{\infty} dE_{\pi^0} \frac{dN_{\pi^0}}{dE_{\pi^0}} \frac{1}{p_{\pi^0}},$$

where the pion spectra is

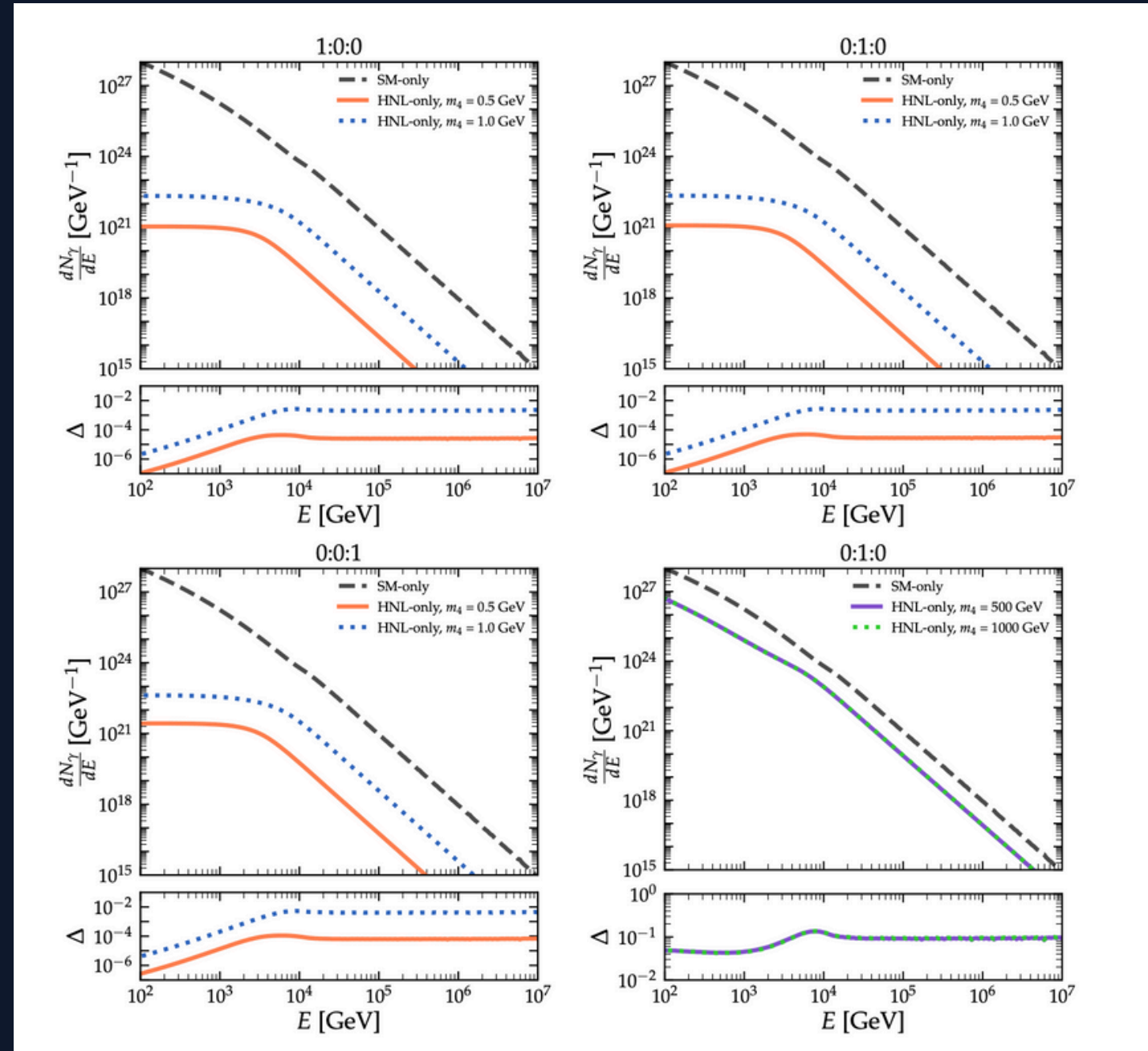
$$\frac{dN_{\pi^0}}{dE} = \frac{1}{2} \mathcal{B}(\nu_4 \rightarrow \pi^0 \nu_\alpha) \int dE_s \left(\frac{E\bar{E} - \gamma_s m_{\pi^0}^2}{\gamma_s \beta_s \bar{p}^3} \right) \frac{\bar{p}}{\gamma_s(\bar{p} - \bar{E}) + E} \frac{dN_s}{dE_s},$$

[0.5 – 2] TeV



as in the neutrino case

PHOTON SPECTRA



Light mass regime:

As in the neutrino case, the angular distribution of the photons that come from the pion decay only depend on the pion energy distribution, which depends only on its energy in the HNL rest frame.

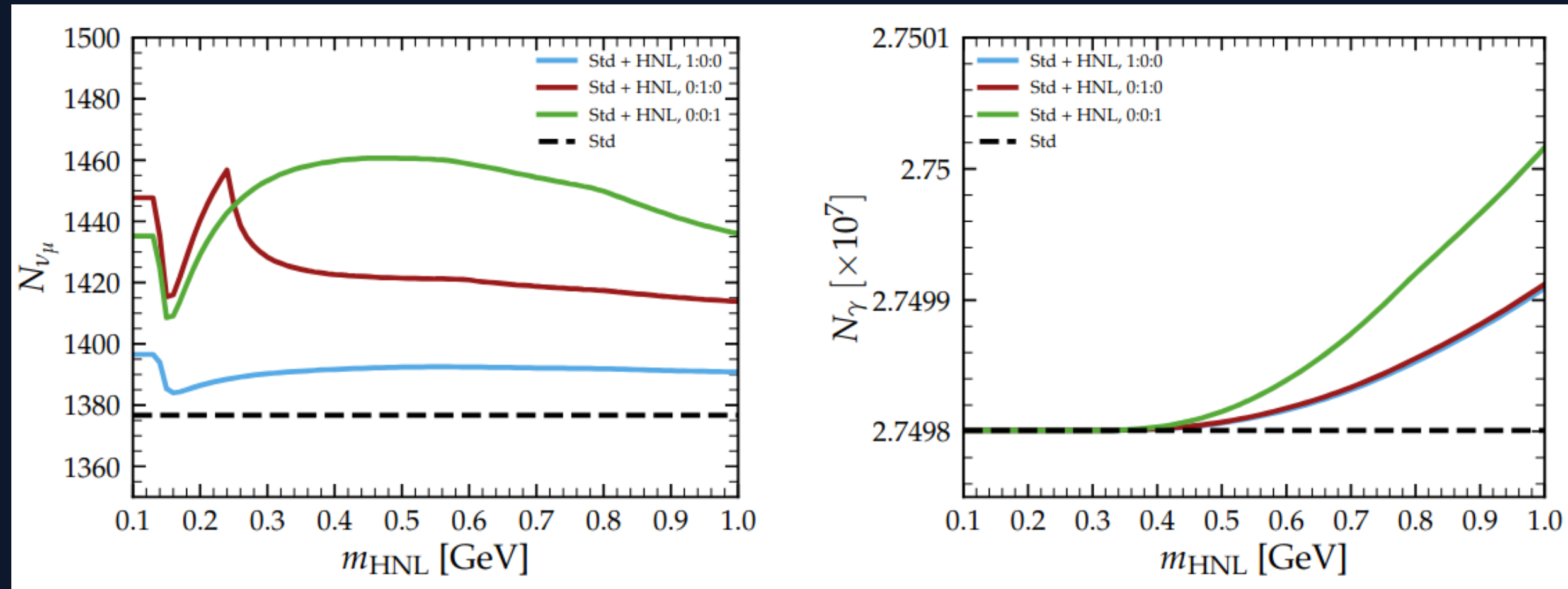
Heavy mass regime:

As in the neutrino case, the angular distributions arise from multiple secondary processes → SM-like shape.

Expected photon neutrino spectrum from HNL-only (color) and SM-only (grey) contributions in a 100s PBH burst

NUMBER OF EVENTS

[0.1 – 1] GeV regime



Expected muon-neutrino events at IceCube (left panel) and photon events at HAWC (right panel) as a function of the HNL mass. We assume a 100 s burst and a PBH at 0.0001 pc from Earth located in the north hemisphere (declination angle $[30^\circ < \delta < 90^\circ]$ for IceCube, zenith angles $[0^\circ < \zeta < 26^\circ]$ for HAWK).

NUMBER OF EVENTS

[0.5 – 2] TeV regime

We predict around 1470 neutrino events from both SM and HNL decays,
independently of the HNL mass

In the same way, we predict 2.87×10^7 events at HAWK regardless of the
HNL mass

ANALYSIS

To estimate the combined sensitivity at IceCube and HAWC we rely on the χ^2 test statistics

$$\chi^2(d_{\text{PBH}}, m_4) = \min_{\alpha} \left\{ \frac{[N_{\gamma}^{\text{SM+HNL}}(m_4, d_{\text{PBH}}(1 + \alpha)) - N_{\gamma}^{\text{SM}}(d_{\text{PBH}})]^2}{\sigma_{\gamma}^2} + \frac{[N_{\nu_{\mu}}^{\text{SM+HNL}}(m_4, d_{\text{PBH}}(1 + \alpha)) - N_{\nu_{\mu}}^{\text{SM}}(d_{\text{PBH}})]^2}{\sigma_{\nu_{\mu}}^2} \right\}, \quad (49)$$

that we minimize over the nuisance parameter α

We find $\alpha \sim 10^{-7} - 10^{-5}$ in the
light-mass regime, $\alpha \sim 10^{-2}$
in the high-mass regime