

# Prospects of MeV telescopes in probing weak-scale Dark Matter

Arpan Kar



Based on: *SciPost Phys.* 19, 080 (2025) (*arXiv*: 2503.04907)

M. Cirelli, A. Kar

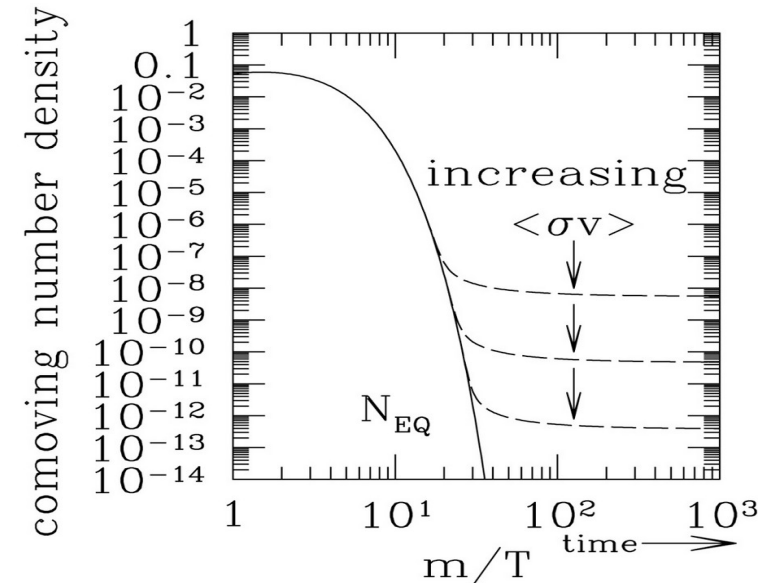


# Dark matter : WIMPs

- Dark Matter (DM) exists and provides  $\sim 25\%$  of the energy density of the Universe
- Weakly Interacting Massive Particles (WIMPs) : one of the most popular candidates for DM

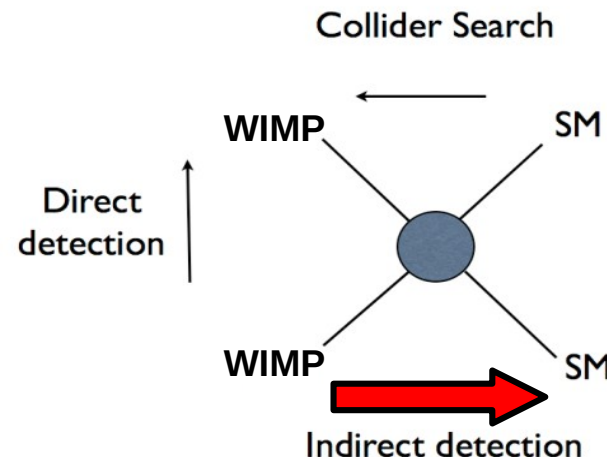
- Stable, no electric charge, no colors
- Mass at the weak scale (GeV – TeV)
- Weak interactions ( $\sigma v \sim 10^{-26} \text{ cm}^3 \text{ s}^{-1}$ ) keep WIMPs in thermal equilibrium in the early Universe and provide correct relic abundance through thermal decoupling

$$\Omega_{DM} h^2 = 0.12 \text{ (Obs.)} \quad \Omega_{DM} h^2 \propto 1/\langle \sigma v \rangle$$



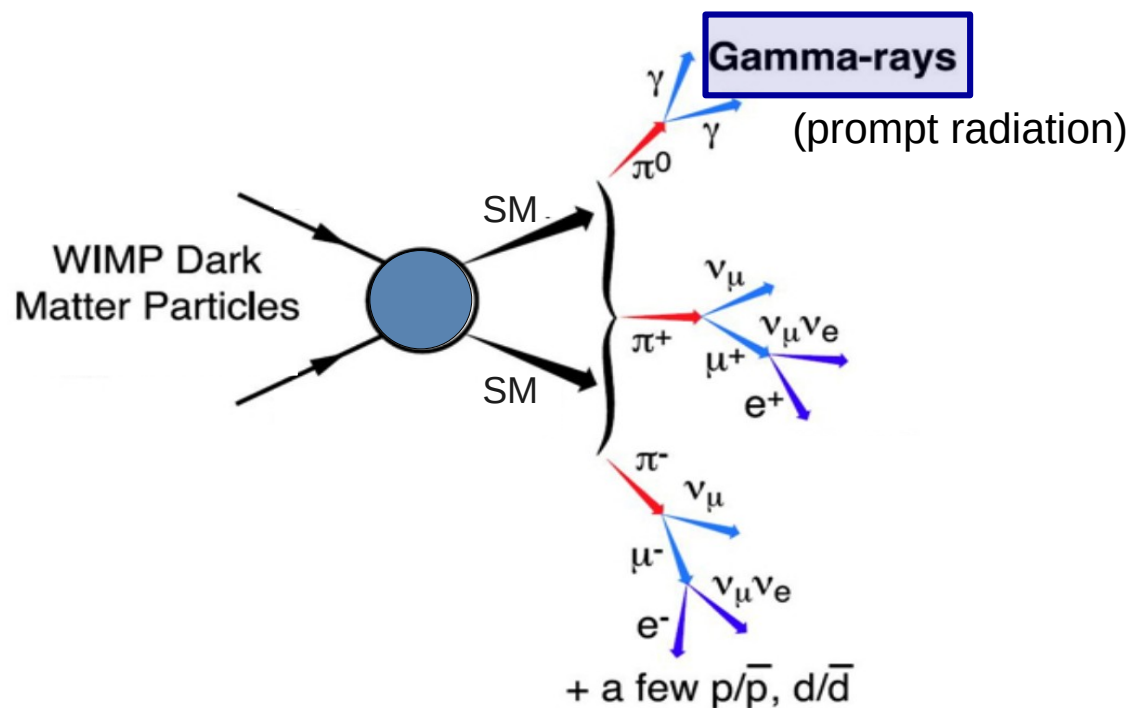
- WIMP searches :

- Direct detection
- Collider searches
- Indirect detection



## Searches for WIMP DM in photon observations from the inner Galaxy

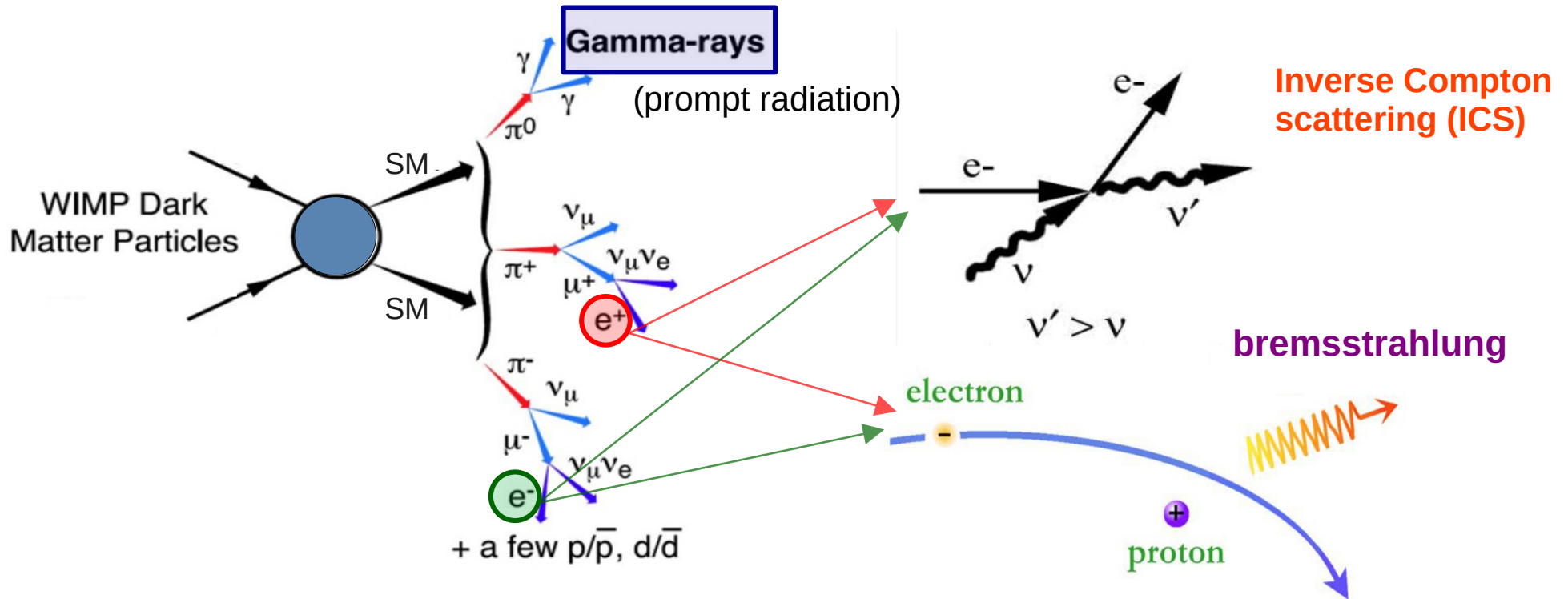
- **Photon signals from inner Galaxy:** the observables consist of mainly two types of photon signal



- **Prompt radiation:** High-energy  $\gamma$ - rays are produced directly in the WIMP annihilation process

## Searches for WIMP DM in photon observations from the inner Galaxy

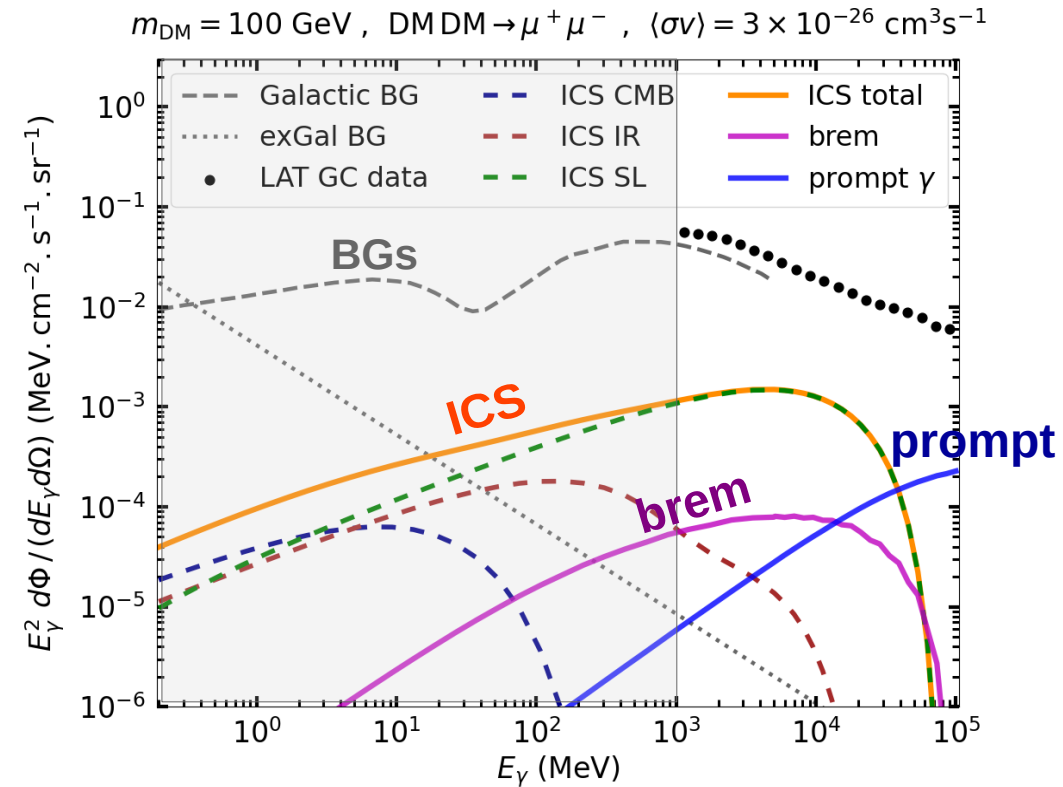
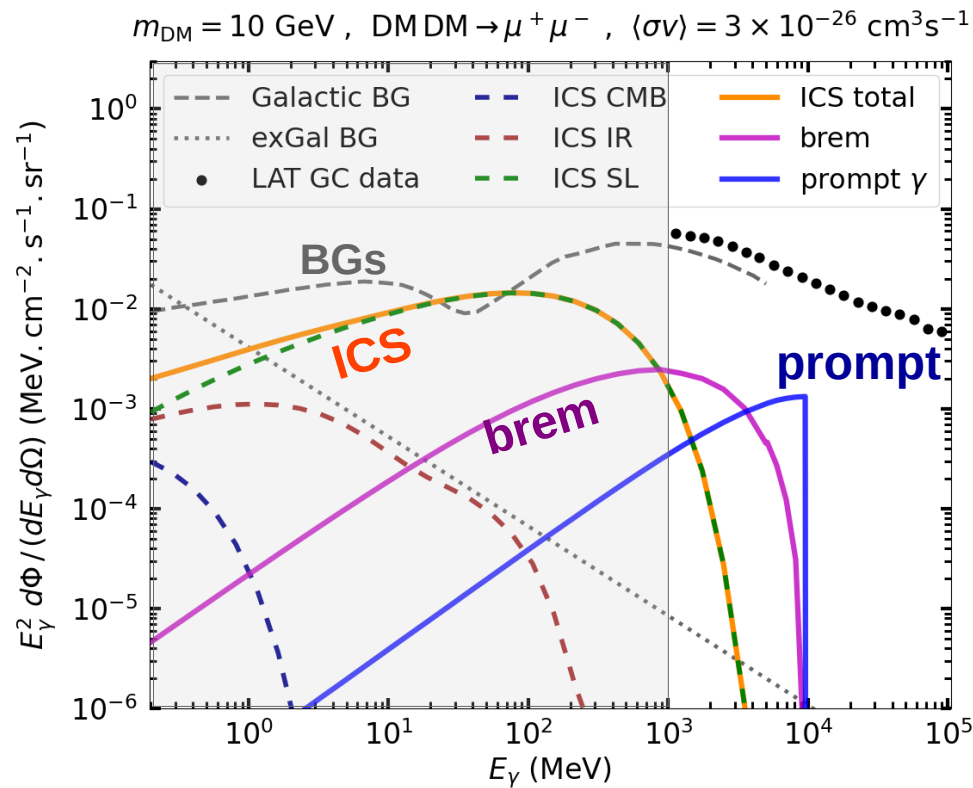
- **Photon signals from inner Galaxy:** the observables consist of mainly two types of photon signal



- **Prompt radiation:** High-energy  $\gamma$ - rays are produced directly in the WIMP annihilation process
- **Secondary radiation:** Galactic WIMP annihilations generate abundant energetic  $e^\pm$ , which subsequently emit through **Inverse Compton scattering (ICS)** , **bremsstrahlung**
  - **Comparatively lower-energetic gamma-rays photons**
  - **Enhanced for the lepton-rich annihilation channels of WIMPs**

e.g.,  $DM DM \rightarrow e^+ e^-$  ,  $DM DM \rightarrow \mu^+ \mu^-$

# Prompt and secondary photons from WIMP annihilations in the Galaxy

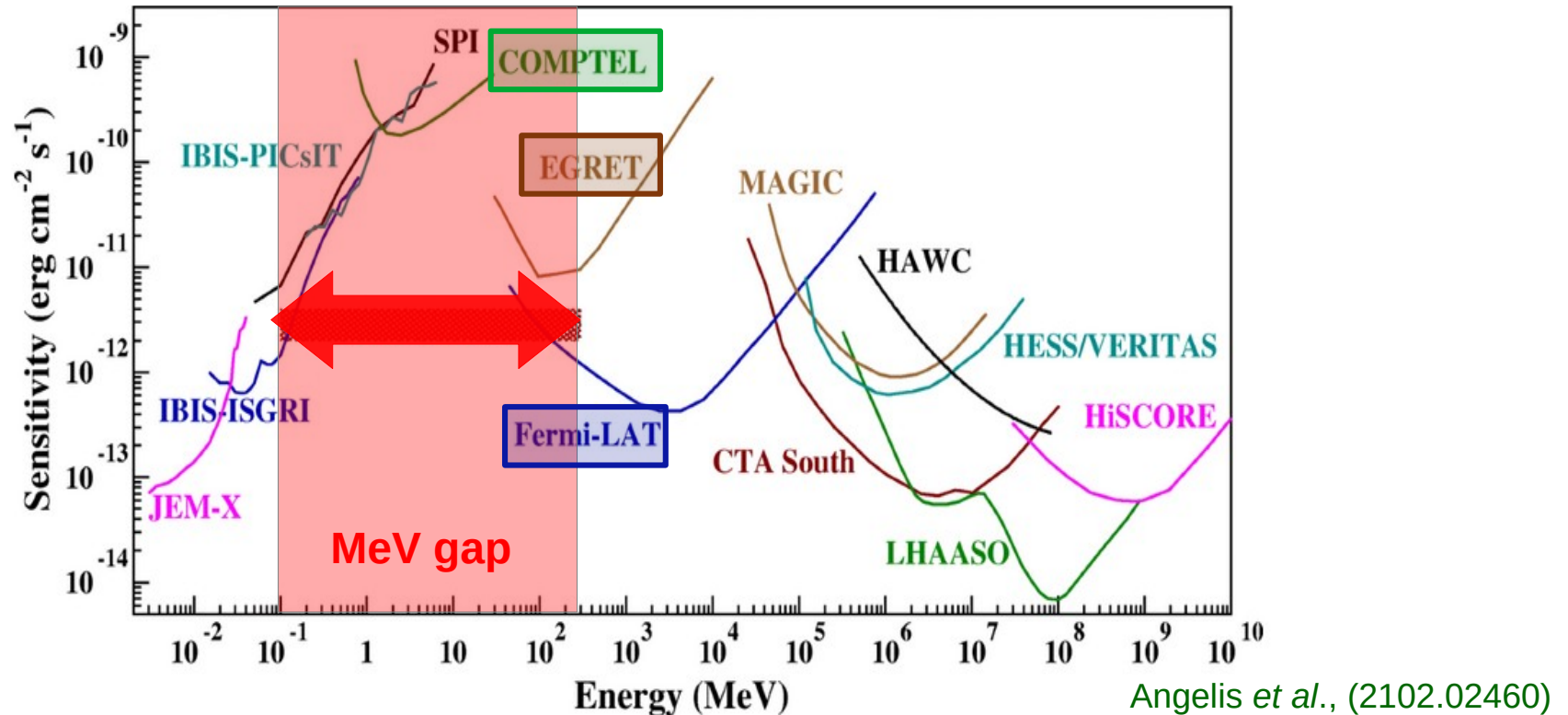


M. Cirelli, A.K.; (SciPost Phys. 19 (2025) 080)

- Secondary photons from WIMP annihilation in general populate the Sub-GeV energy range (MeV - GeV)

## MeV Gap and upcoming MeV $\gamma$ -ray telescopes

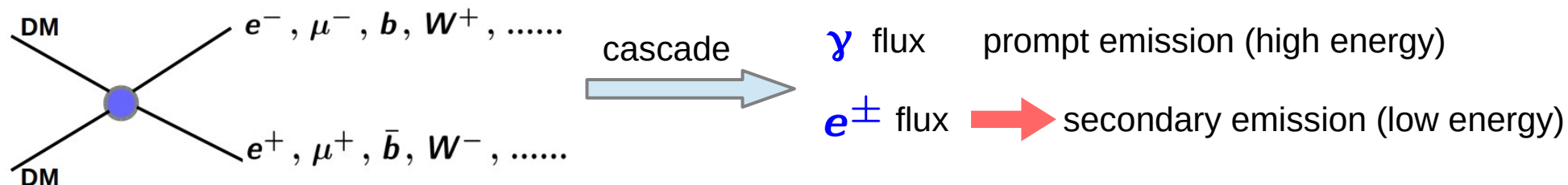
- One potential difficulty for detecting the secondary emissions: relatively poor sensitivity of existing telescopes (COMPTEL, EGRET, Fermi-LAT, etc.) in the sub-GeV range



- The upcoming space-based MeV telescopes will efficiently fill the MeV gap with *better sensitivity*
  - COSI, AMEGO, e-ASTROGAM, GECCO, AdEPT, PANGU, GRAMS, MAST, .....
- Potential of these MeV telescopes in probing WIMP DM, based on the secondary emission ?

## MeV-GeV photons from WIMP annihilations in the Galaxy

- Target region for observation: a disk of  $10^\circ$  radius around the galactic Center (GC)  
 → same order as the maximum angular width of the MeV telescopes



### ● Prompt $\gamma$ - ray emission flux:

$$\frac{d\Phi_{\text{prompt}}}{dE_\gamma d\Omega} = \frac{\langle \sigma v \rangle}{8\pi m_{\text{DM}}^2} \frac{dN_\gamma}{dE_\gamma} \frac{J_{\Delta\Omega}}{\Delta\Omega}$$

$\frac{dN_\gamma}{dE_\gamma}$  spectra produced per annihilation  
in a given annihilation channel

$$J_{\Delta\Omega} = \int_{\Delta\Omega} d\Omega \int_{\text{l.o.s.}} ds \rho_{\text{DM}}^2(r(s, \theta))$$

( $10^\circ$  around GC)

$s \rightarrow$  line-of-sight (l.o.s.)

$$d\Omega = 2\pi \sin\theta d\theta$$

NFW DM profile :

$$\rho_{\text{DM}}(r) = \frac{\rho_0}{\left(\frac{r}{r_s}\right) \left(1 + \frac{r}{r_s}\right)^2}$$

Salas *et al.*, (1906.06133)

Cirelli *et al.*, (2406.01705)

# MeV-GeV photons from WIMP annihilations in the Galaxy

## • Secondary $\gamma$ - ray emission flux:

$(s, b, l) \rightarrow$  Galactic coordinates  
 $\cos b \cos l = \cos \theta$

$$\frac{d\Phi_{2\text{ndary}}}{dE_\gamma d\Omega} = \frac{1}{\Delta\Omega} \int_{\Delta\Omega} d\Omega \left[ \frac{1}{E_\gamma} \int_{l.o.s.} ds \frac{j_{2\text{ndary}}(E_\gamma, \vec{x}(s, b, l))}{4\pi} \right]$$

$$j_{2\text{ndary}} = j_{\text{ICS}} + j_{\text{brem}}$$

$$j_{\text{ICS}}(E_\gamma, \vec{x}(s, b, l)) = 2 \int_{m_e}^{m_{\text{DM}}} dE_e \sum_{i \in \text{ISRF}} \mathcal{P}_{\text{ICS}}^i(E_\gamma, E_e, \vec{x}) \frac{dn_e}{dE_e}(E_e, \vec{x})$$

From DM annihilation

$$j_{\text{brem}}(E_\gamma, \vec{x}(s, b, l)) = 2 \int_{m_e}^{m_{\text{DM}}} dE_e \mathcal{P}_{\text{brem}}(E_\gamma, E_e, \vec{x}) \frac{dn_e}{dE_e}(E_e, \vec{x})$$



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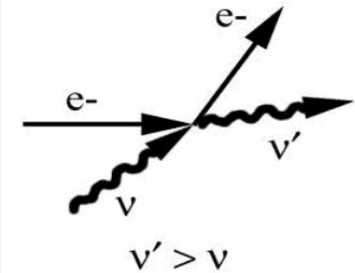
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$$\mathcal{P}_{\text{ICS}}^i(E_\gamma, E_e, \vec{x}) = c E_\gamma \int d\epsilon n_i^{\text{ISRF}}(\epsilon, \vec{x}) \sigma_{\text{IC}}(\epsilon, E_\gamma, E_e)$$

Inter-Stellar Radiation Field (ISRF) : CMB, infrared (IR), starlight (SL)

Buch, *et al.*, (PPPC 4 DM, [1505.01049]), (GALPROP)



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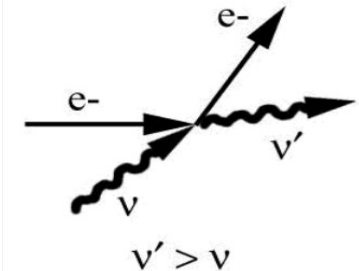
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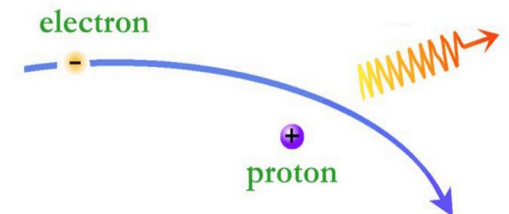
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Gas species : ionic, atomic and molecular



# MeV-GeV photons from WIMP annihilations in the Galaxy

## • Distribution of WIMP induced $e^\pm$ in the galaxy :

Source function from WIMP annihilation :

$$Q_e(E_e^S, r) = \frac{\langle \sigma v \rangle}{2 m_{\text{DM}}^2} \frac{dN_e}{dE_e^S} \rho_{\text{DM}}^2(r)$$

DM density

spectra produced per annihilation in a given annihilation channel

## ▪ Semi-Analytic :

$$\frac{dn_e}{dE_e}(E_e, \vec{x}) = \frac{1}{b_{\text{tot}}(E_e, \vec{x})} \int_{E_e}^{m_{\text{DM}}} dE_e^S Q_e(E_e^S, r)$$

$b_{\text{tot}}(E_e, \vec{x})$  : total energy loss rate of  $e^\pm$  ➡ Dominating process near the GC region

- |   |  |
|---|--|
| ➔ ICS on ambient photons                      | ➔ Coulomb interactions with interstellar gases |
| ➔ synchrotron emission in galactic $B$ -field | ➔ ionization of the same gases                 |
|   | ➔ bremsstrahlung on the same gases             |

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Buch, *et al.*, (PPPC 4 DM, [1505.01049])

## ▪ Full-propagation of $e^\pm$ :

( spatial difusion, advection/convection, re-acceleration, energy losses, various nuclear processes )

$$\nabla \cdot (\vec{J}_i - \vec{v}_w N_i) + \frac{\partial}{\partial p} \left[ p^2 D_{pp} \frac{\partial}{\partial p} \left( \frac{N_i}{p^2} \right) \right] - \frac{\partial}{\partial p} \left[ \dot{p} N_i - \frac{p}{3} (\vec{\nabla} \cdot \vec{v}_w) N_i \right] =$$

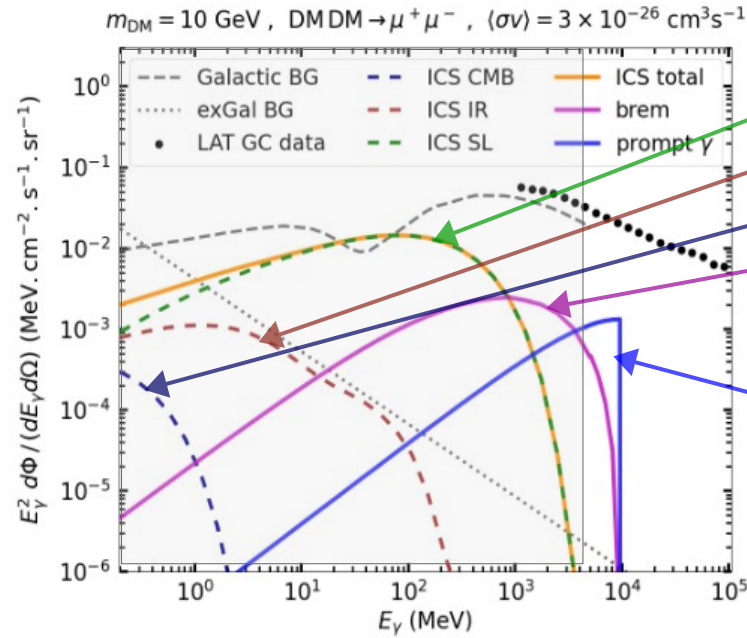
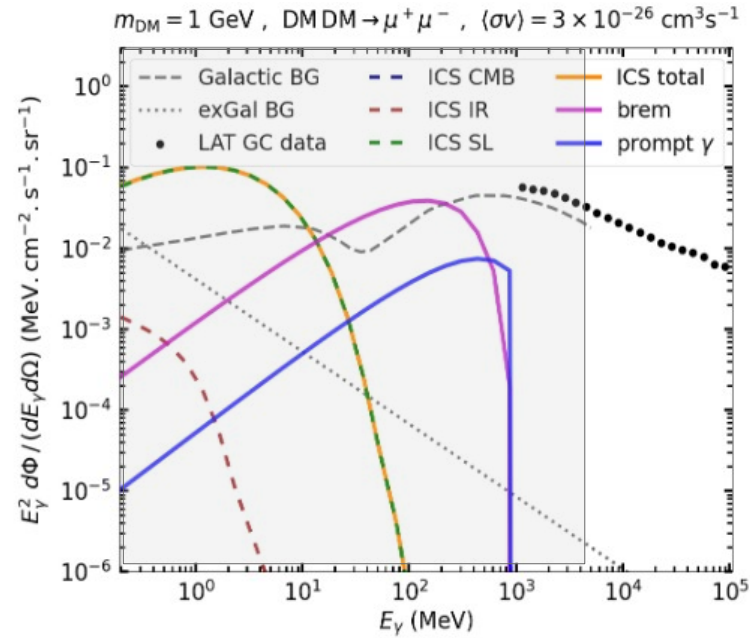
$$Q + \sum_{i < j} \left( c \beta n_{\text{gas}} \sigma_{j \rightarrow i} + \frac{1}{\gamma \tau_{j \rightarrow i}} \right) N_j - \left( c \beta n_{\text{gas}} \sigma_i + \frac{1}{\gamma \tau_i} \right) N_i$$

$$J_i = -D_{ij} \nabla_j N$$

Evoli *et al.*, (1607.07886)

$N_i(\vec{r}, p)$  : no. density of  $e^-/e^+$  per momentum  $p$

# MeV-GeV photons from WIMP annihilations in the Galaxy



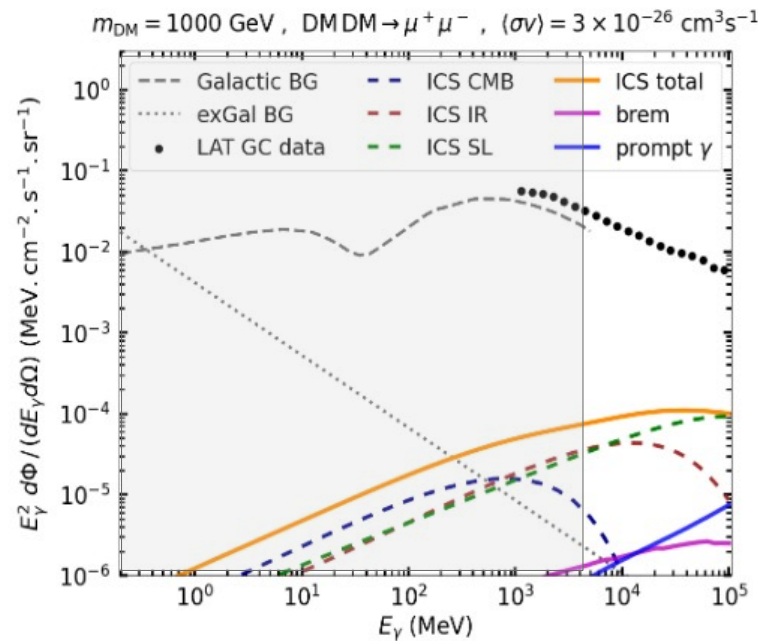
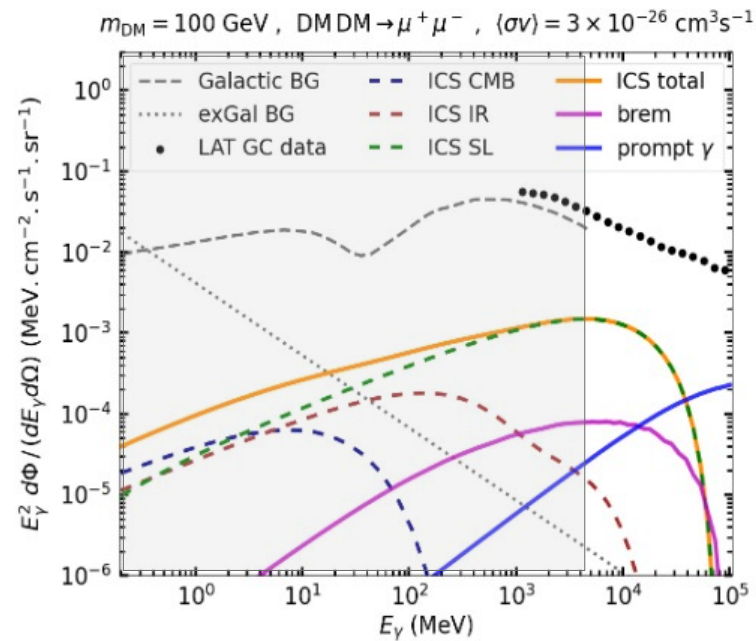
ICS on starlight

ICS on infrared

ICS on CMB

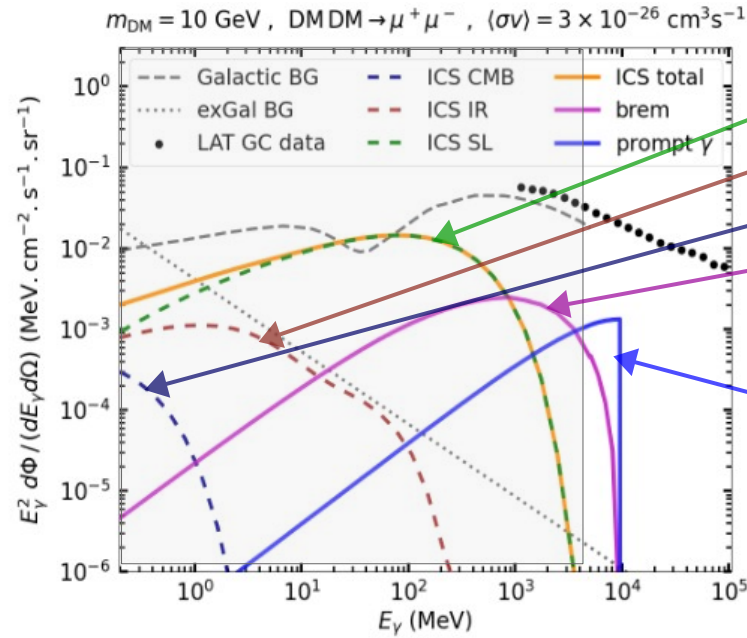
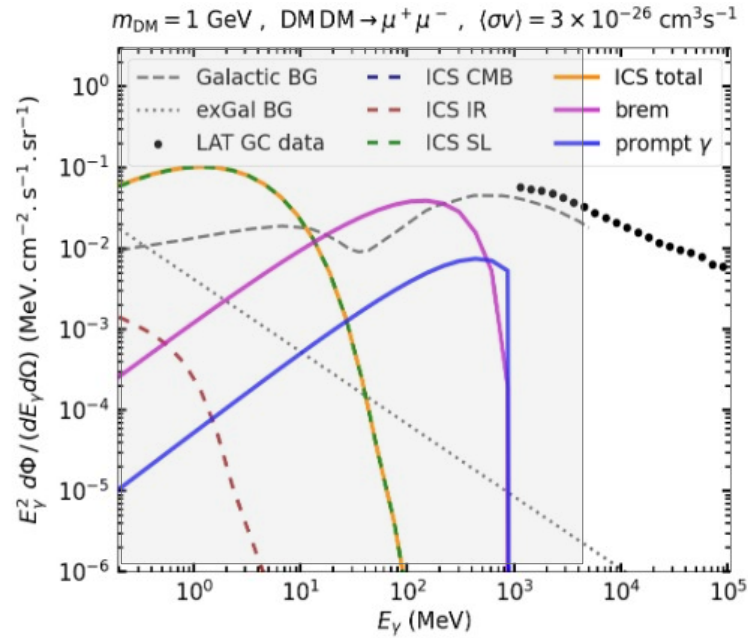
bremsstrahlung on gas

Prompt emission

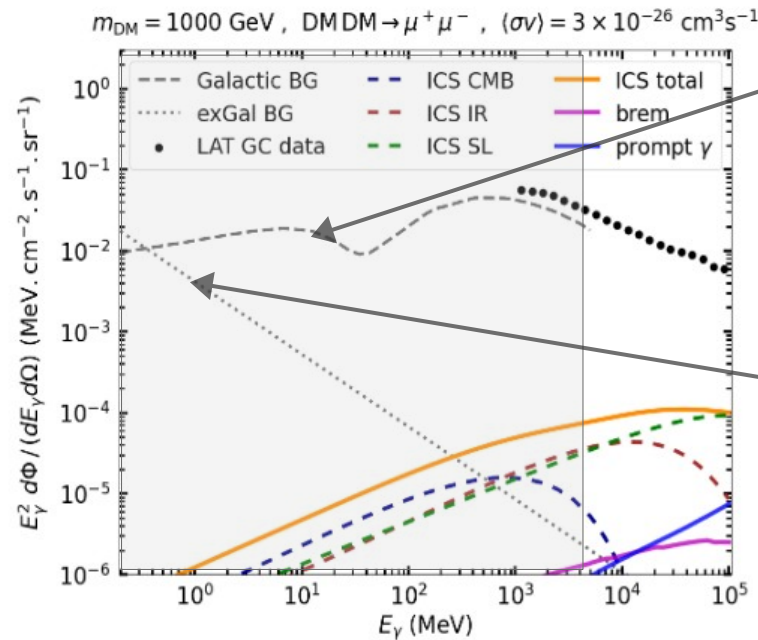
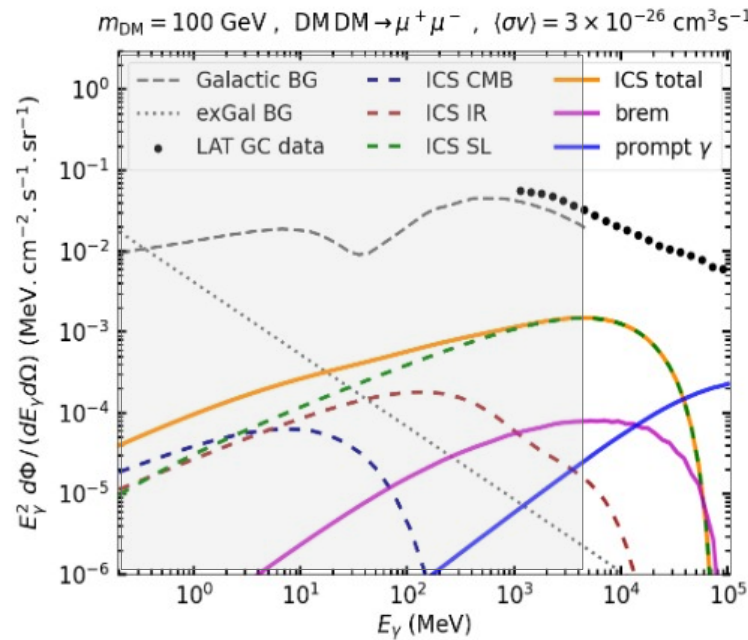




# MeV-GeV photons from WIMP annihilations in the Galaxy



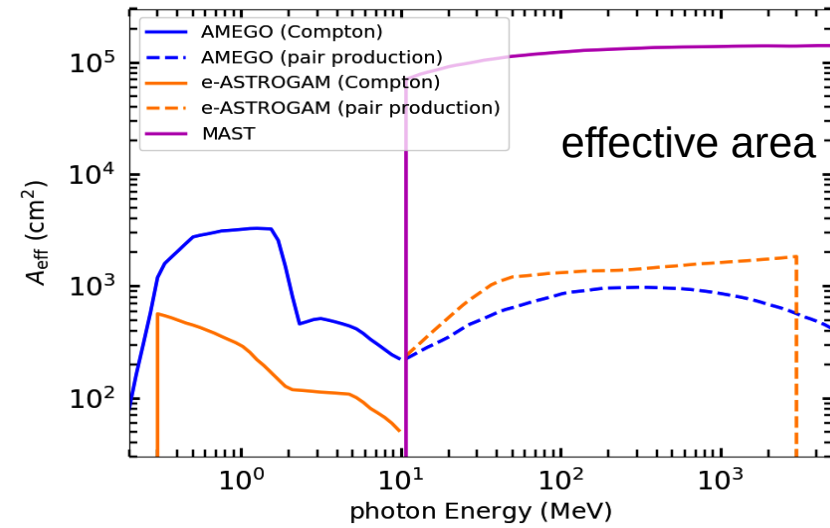
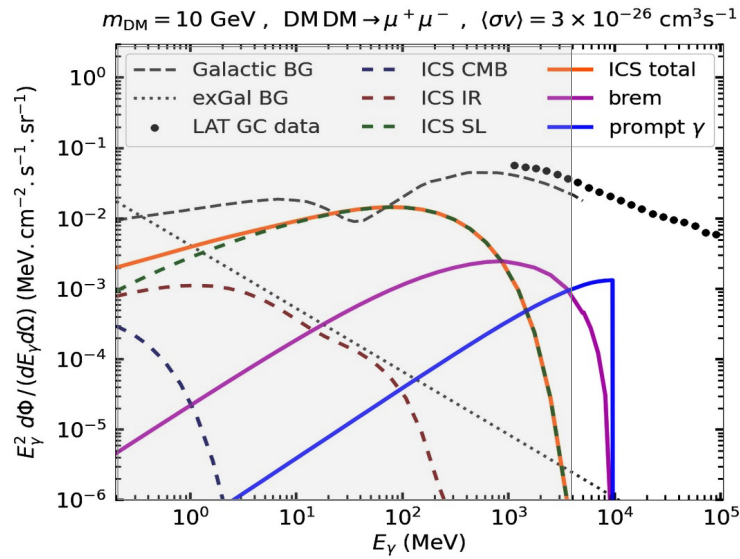
ICS on starlight  
ICS on infrared  
ICS on CMB  
bremsstrahlung on gas  
Prompt emission



**Galactic BGs**  
→ *bremsstrahlung*  
→  $\pi^0 \rightarrow \gamma \gamma$   
→ *ICS-high*  
→ *ICS-low*  
**extra-Galactic BG**

# WIMP annihilation signals at the MeV telescopes

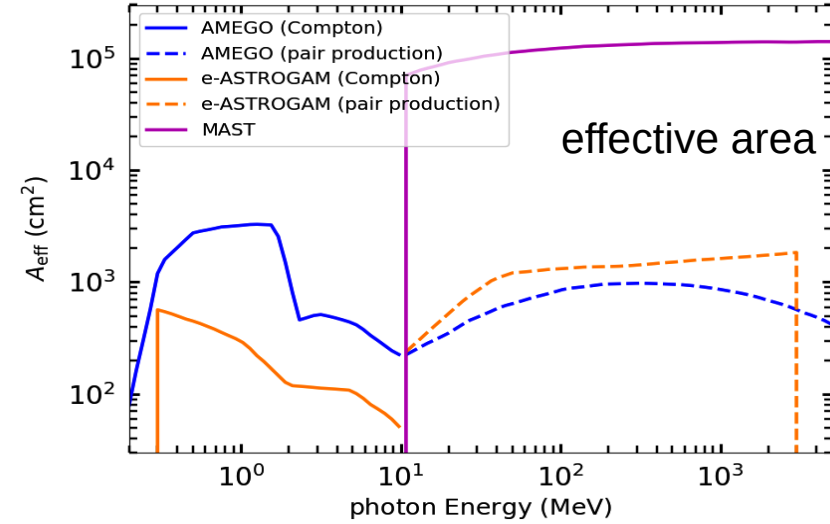
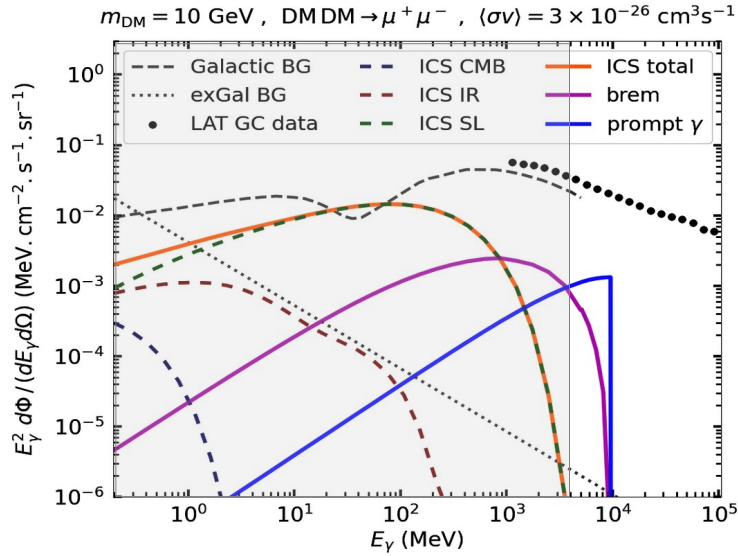
- MeV telescopes : **AMEGO**, **e-ASTROGAM** and **MAST**  $0.2 \text{ MeV} \lesssim E_\gamma \lesssim 5 \text{ GeV}$



M. Cirelli, A.K.; (SciPost Phys. 19 (2025) 080)

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Fisher-projections :

$$\mathcal{F}_{ij} = t_{\text{obs}} \int_{E_{\text{min}}}^{E_{\text{max}}} dE_\gamma A_{\text{eff}}(E_\gamma) \int_{\Delta\Omega} d\Omega \left( \frac{1}{\phi_{\text{tot}}} \frac{\partial \phi_{\text{tot}}}{\partial \theta_i} \frac{\partial \phi_{\text{tot}}}{\partial \theta_j} \right)_{\vec{\theta} = \vec{\theta}_{\text{fiducial}}}$$

$$\phi_{\text{tot}}(\vec{\theta}) = \frac{d\Phi^{\text{SIG}}}{dE_\gamma d\Omega}(\Gamma^{\text{SIG}}) + \sum_I \theta_I^{\text{BG}} \left\{ \frac{d\Phi_{\text{BG}}^I}{dE_\gamma d\Omega} \right\}_{\text{fiducial}} \quad \vec{\theta} = [\Gamma^{\text{SIG}}, \theta_{\text{brem}}^{\text{BG}}, \theta_{\pi^0}^{\text{BG}}, \theta_{\text{ICS}_{\text{hi}}}^{\text{BG}}, \theta_{\text{ICS}_{\text{lo}}}^{\text{BG}}, \theta_{\text{e.g.}}^{\text{BG}}]$$

$t_{\text{obs}} \rightarrow$  observation time

Energy-resolution

$$\frac{d\Phi}{dE_\gamma d\Omega} = \int dE'_\gamma R_\epsilon(E_\gamma, E'_\gamma) \frac{d\Phi}{dE'_\gamma d\Omega}$$

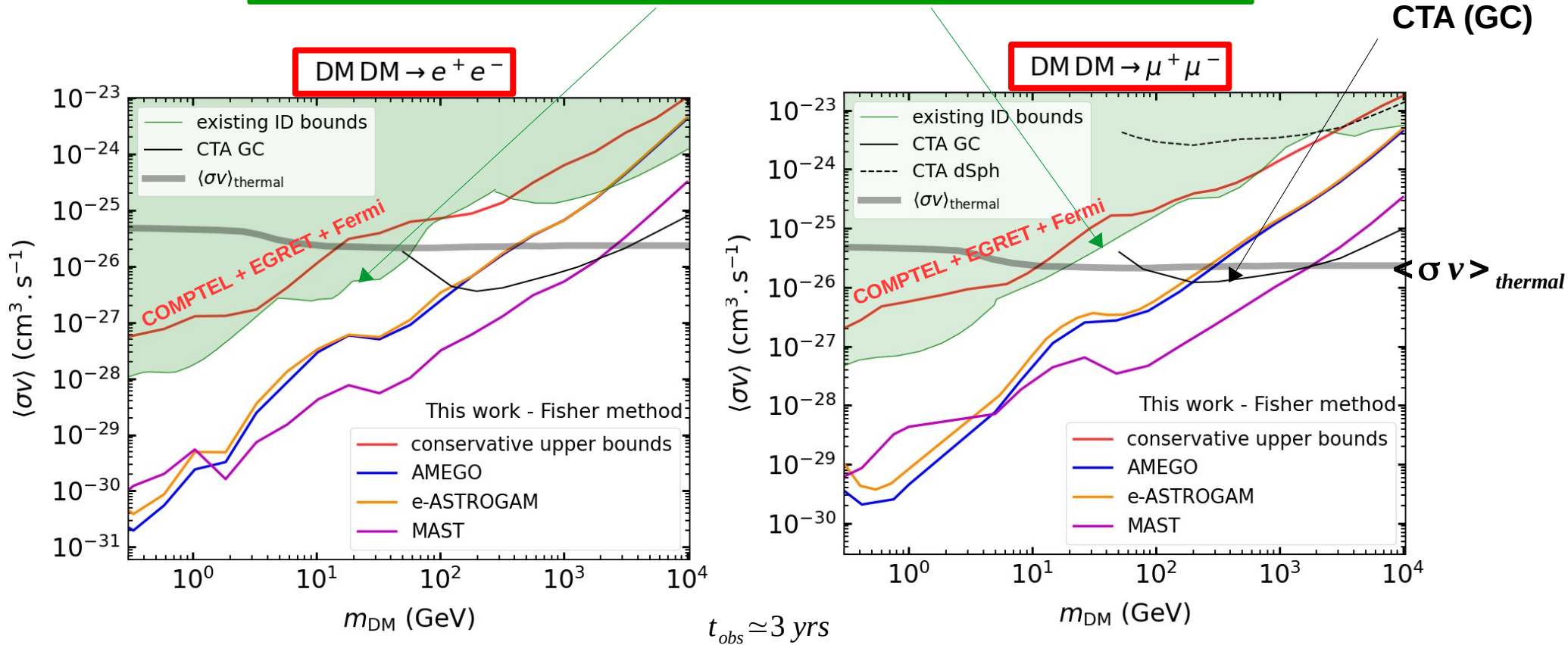
AMEGO (30%) , e-Astrogam (30%), MAST (30%)

( $2\sigma$  - projections)  $\Gamma_{\text{proj}}^{\text{SIG}} = 2 \sqrt{(\mathcal{F}^{-1})_{11}}$  (scales as  $\sqrt{t_{\text{obs}}}$ )



# Projected sensitivities for WIMP (leptonic annihilations)

Existing constraints: CMB + X-rays +  $\gamma$ -rays + AMS  $e^+$  + neutrinos

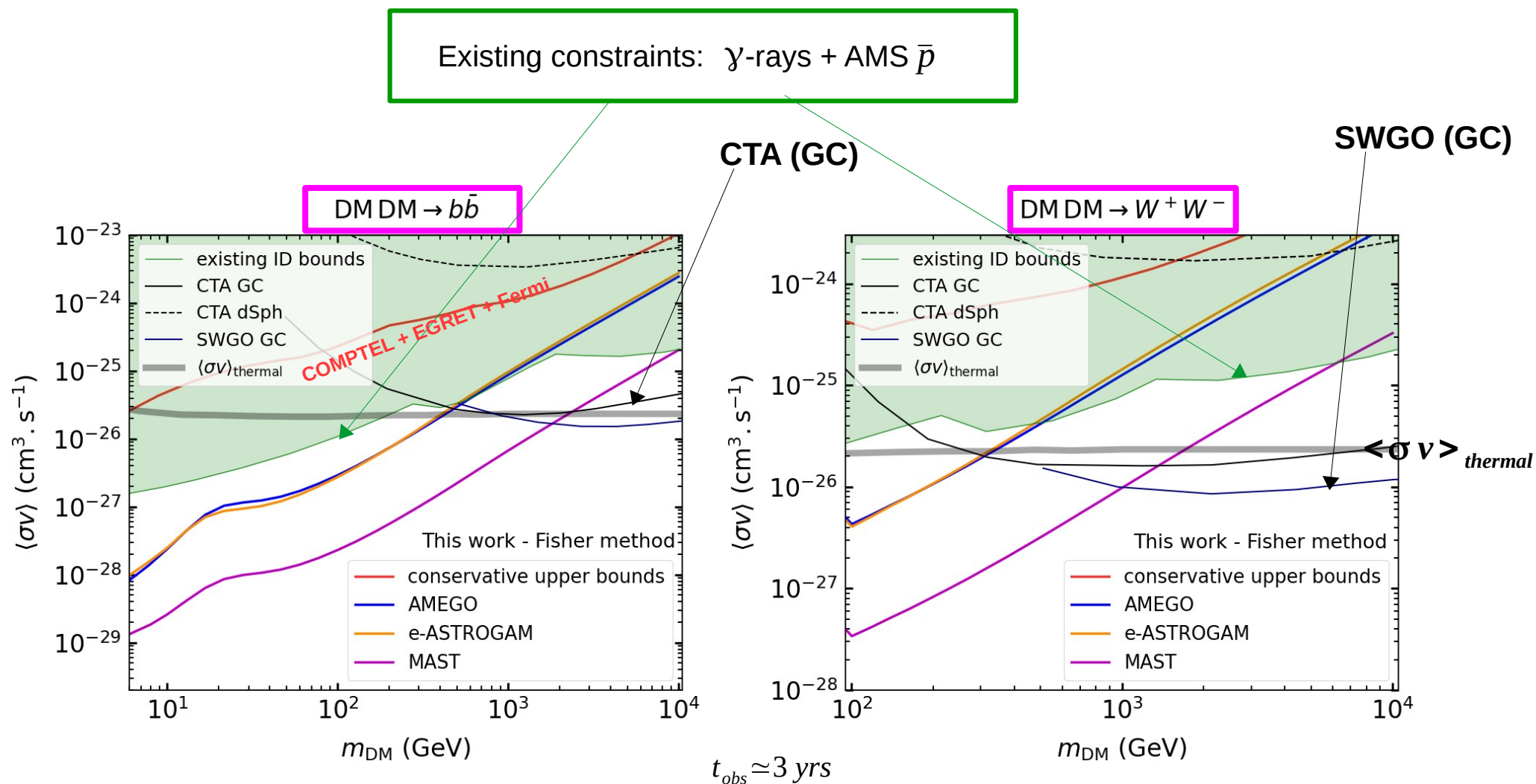


upper-bounds and projected sensitivities of the MeV telescopes AMEGO, e-ASTROGAM and MAST

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- Future space-based MeV gamma-ray telescopes will complement the ground-based high energy gamma-ray instruments in the indirect searches for weak-scale DM

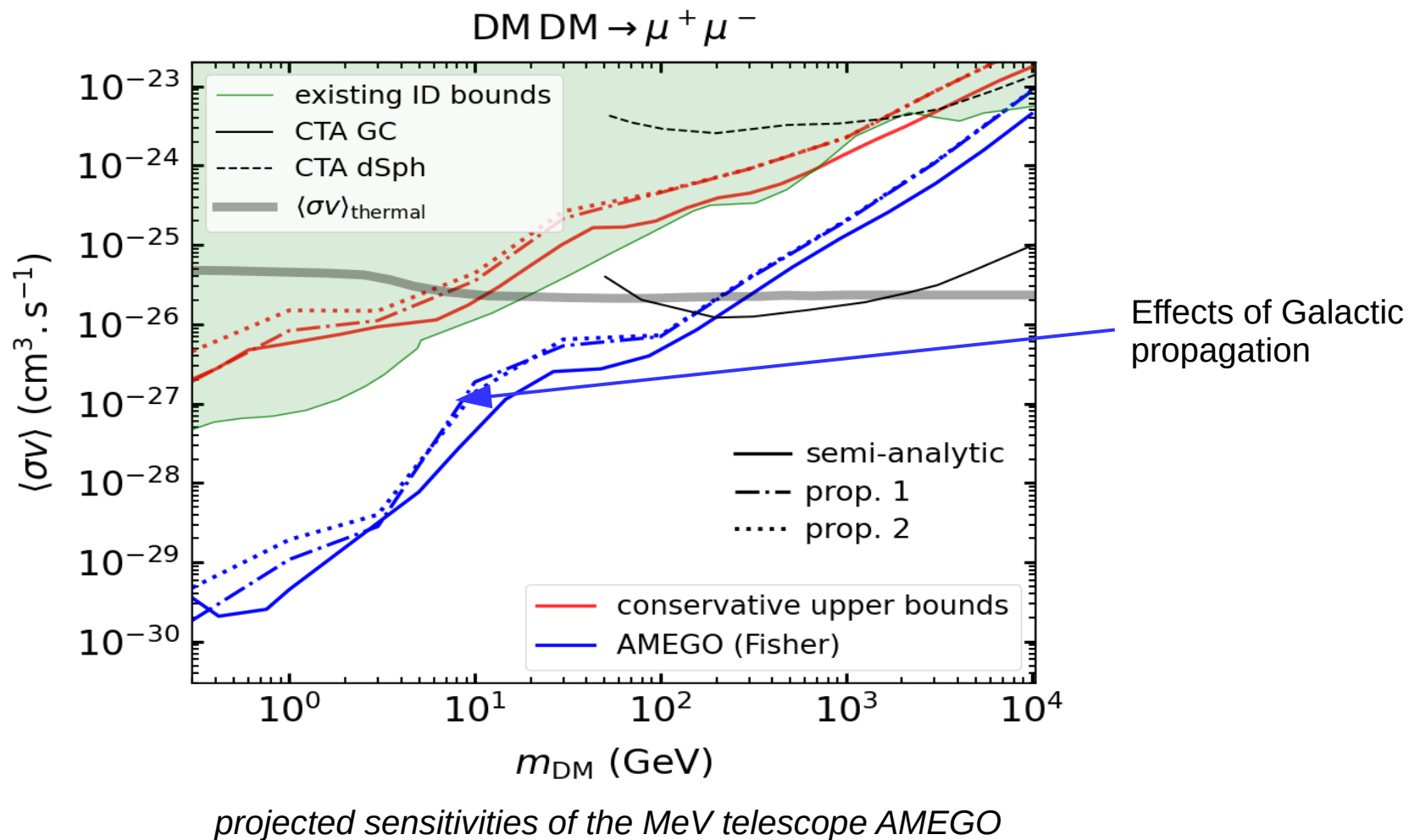
*Projected sensitivities for WIMP (hadronic annihilations)*



M. Cirelli, A.K.; (SciPost Phys. 19 (2025) 080)

- **Future space-based MeV gamma-ray telescopes will complement the ground-based high energy gamma-ray instruments in the indirect searches for weak-scale DM**

# Effects of propagation of $e^\pm$ in the Galaxy



Prop. 1 : propagation model from [Strong *et al.*, (1008.4330), (1101.1381)]

Prop. 2 : propagation model from [Calore *et al.*, (1409.0042)]

used also to estimate the  
**secondary photon BGs**  
towards the GC region

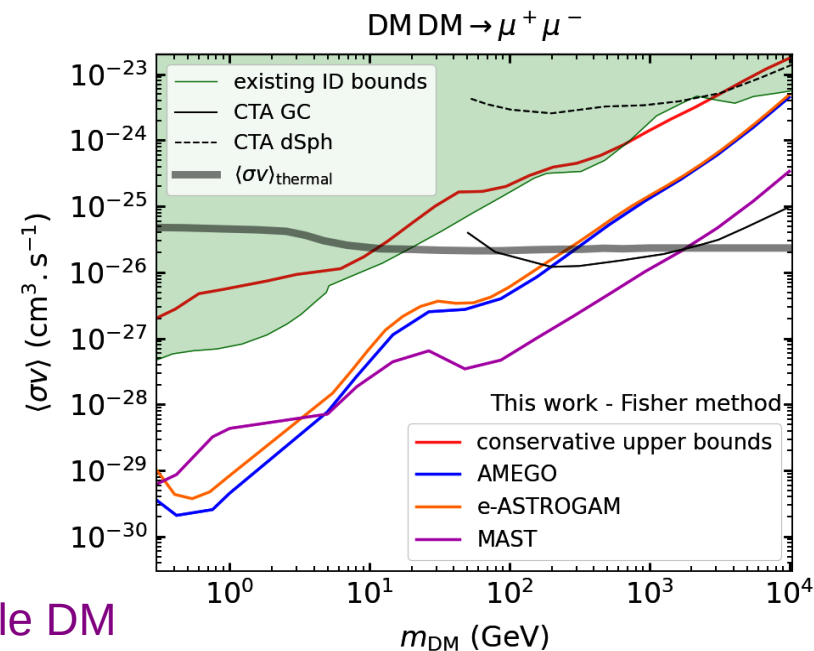
# Summary

- We explore the potential of the upcoming MeV telescopes in probing the photon signals from weak-scale DM annihilations in the Galaxy
- Low-energy secondary emissions (e.g., ICS and bremsstrahlung) produced by DM induced  $e^\pm$  significantly enhance the sub-GeV  $\gamma$ -ray signals of weak-scale DM annihilations

(for lepton-rich annihilation channels)

⇒ Significant enhancements in the sensitivity of MeV telescopes in probing weak-scale DM

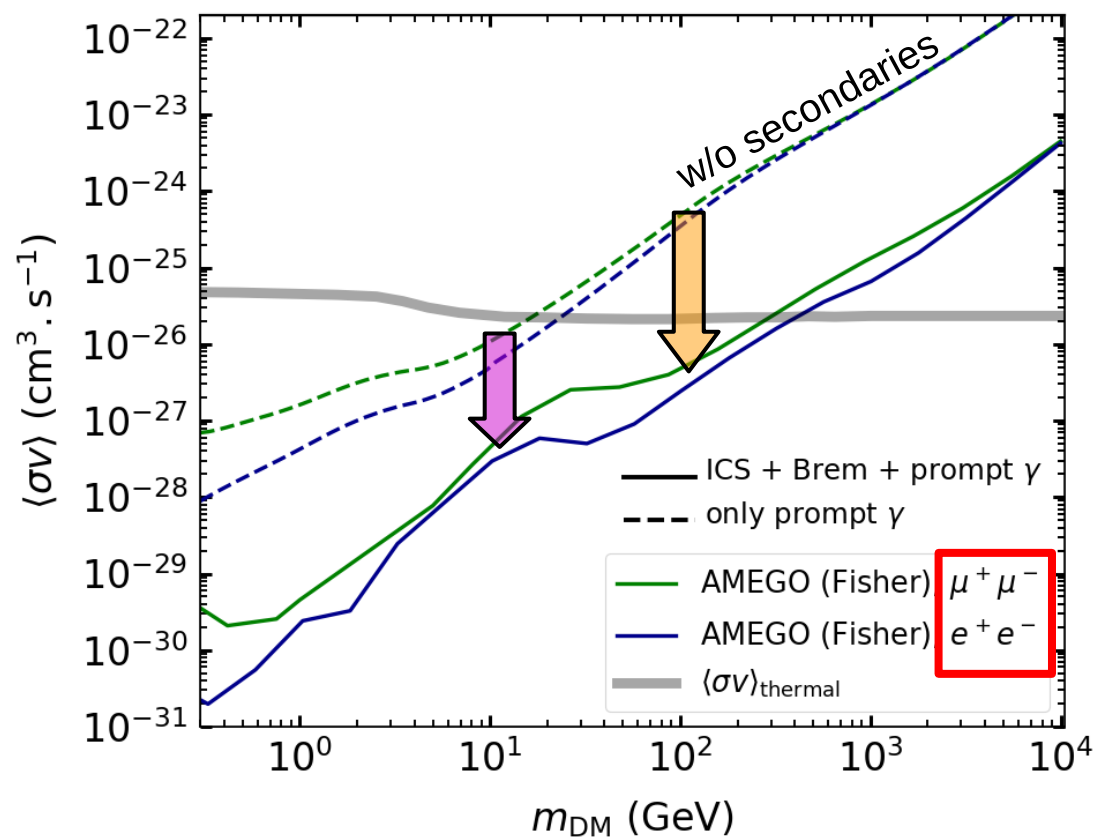
- Based on such signals, the MeV telescopes will be able to explore a wide area of the  $m_{\text{DM}}$ - $\langle\sigma v\rangle$  plane that is yet unconstrained
- MeV  $\gamma$ -ray telescopes can efficiently complement the ground-based high energy  $\gamma$ -ray instruments in the indirect searches for weak-scale DM
- Add an important tool in the indirect searches of weak-scale DM



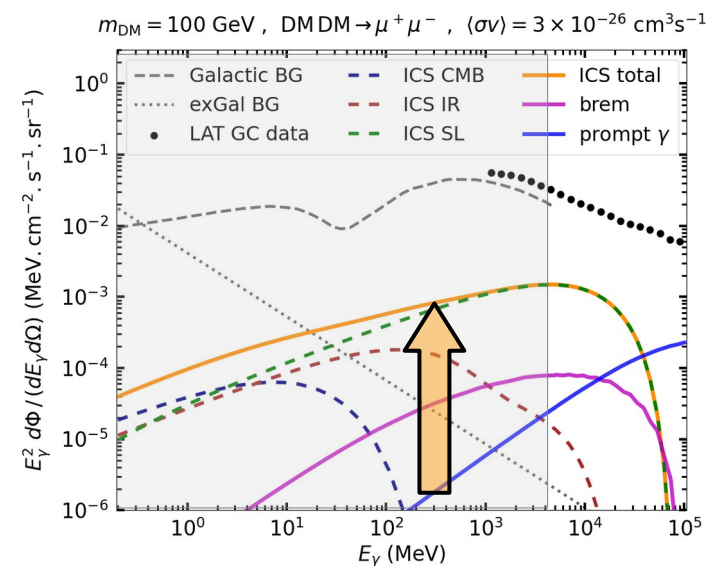
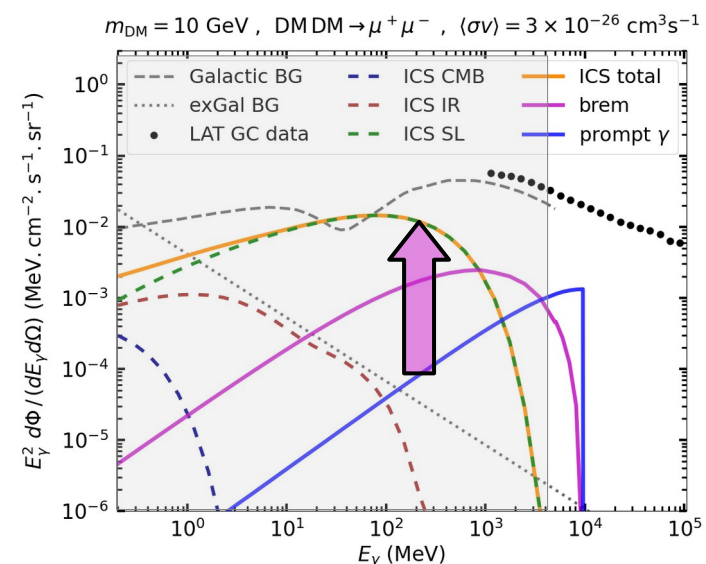
Thank  
you

**Backup slides**

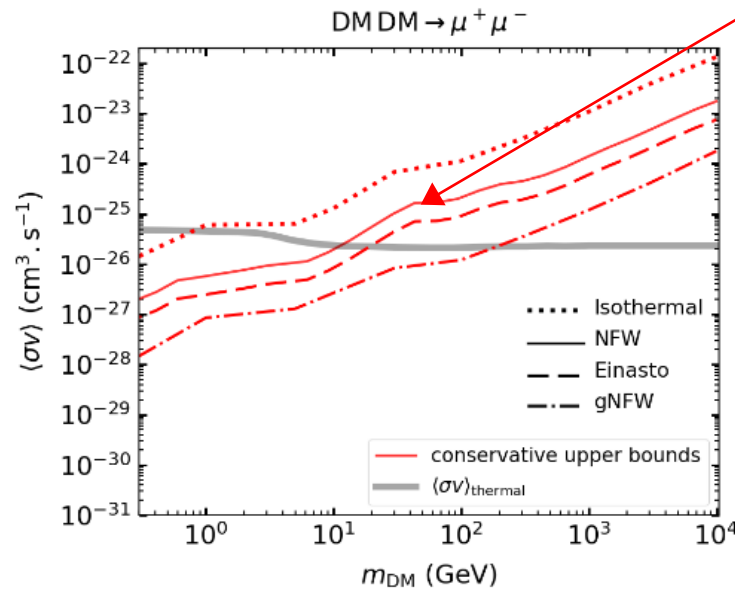
## Importance of considering the secondary signals for WIMPs



*projected sensitivities of the MeV telescope AMEGO*



# Impact of the choice of DM profile



NFW: 
$$\rho_{\text{DM}}(r) = \frac{\rho_0}{\left(\frac{r}{r_s}\right) \left(1 + \frac{r}{r_s}\right)^2}$$

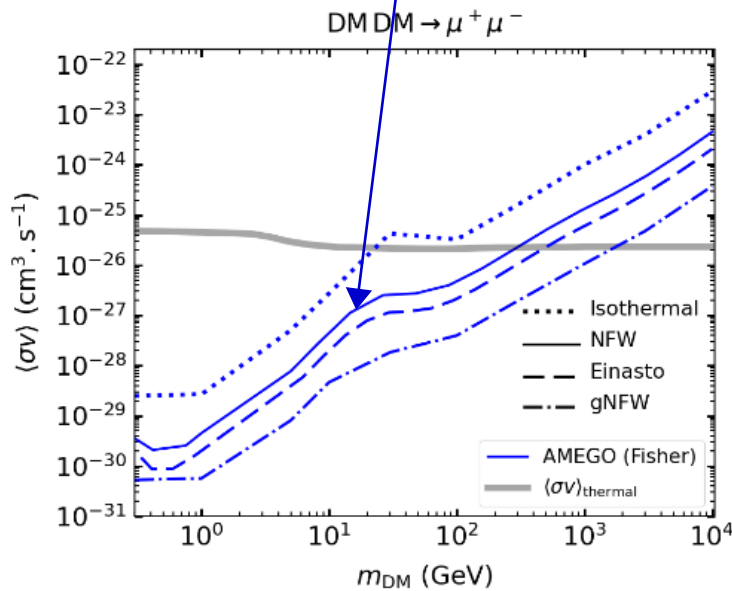
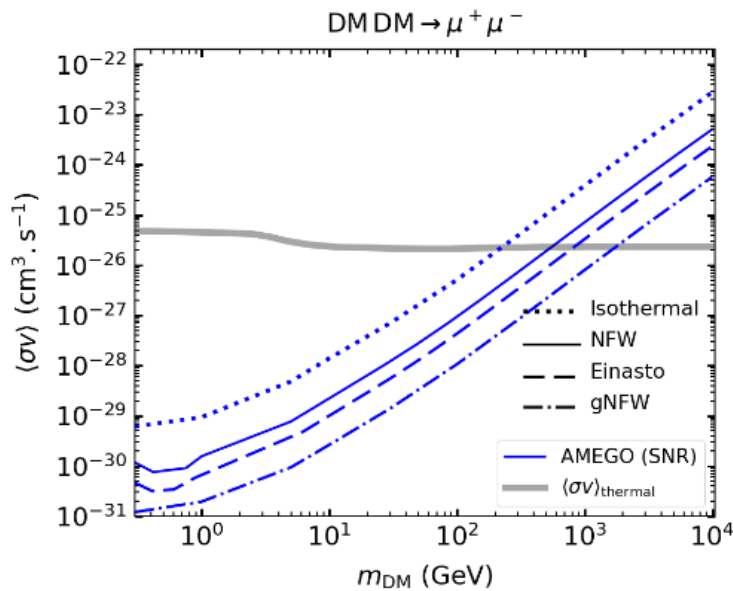
Einasto: 
$$\rho_{\text{DM}}^{\text{Ein}}(r) = \rho_0 \exp \left\{ -\frac{2}{\alpha} \left( \left( \frac{r}{r_s} \right)^\alpha - 1 \right) \right\}$$
  

$$\alpha = 0.18$$

gNFW: 
$$\rho_{\text{DM}}^{\text{gNFW}}(r) = \frac{\rho_0}{\left(\frac{r}{r_s}\right)^\gamma \left(1 + \frac{r}{r_s}\right)^{3-\gamma}}$$
  

$$\gamma = 1.3$$

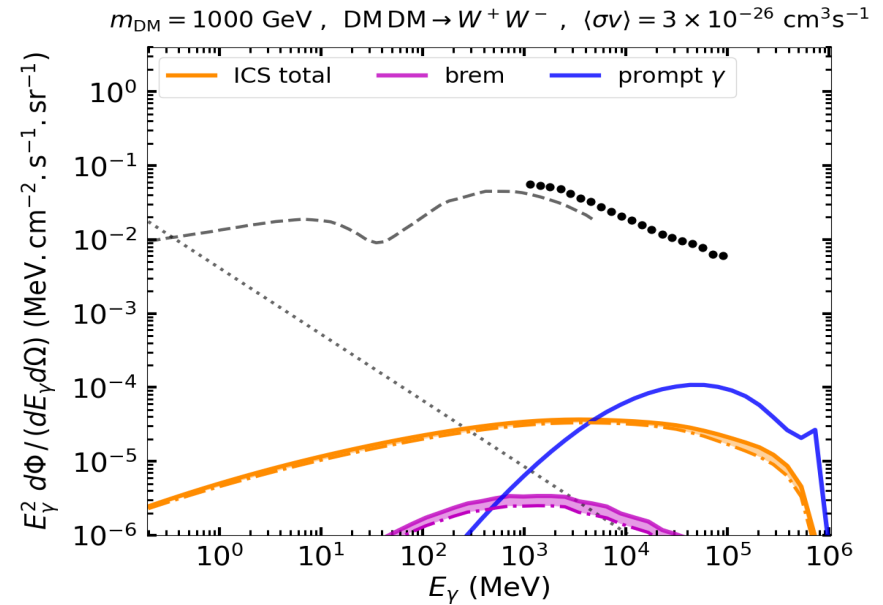
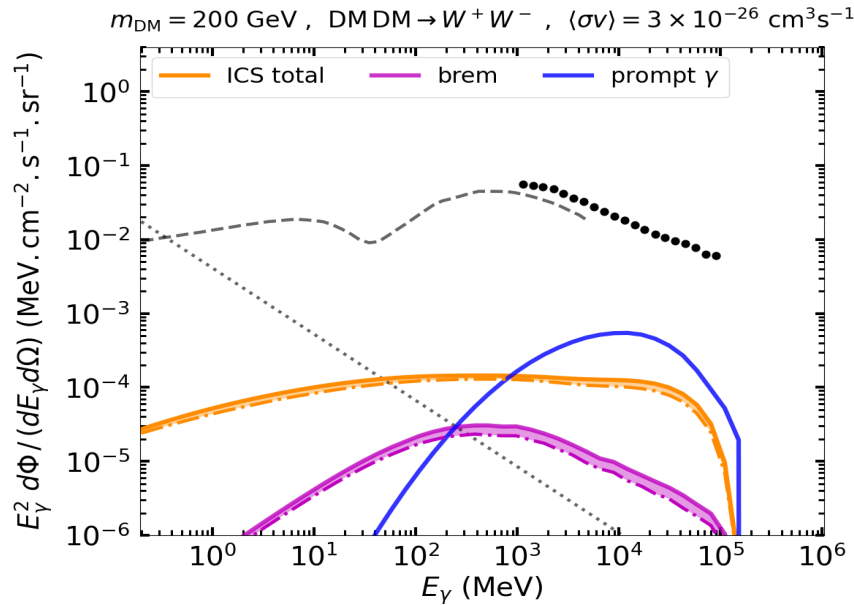
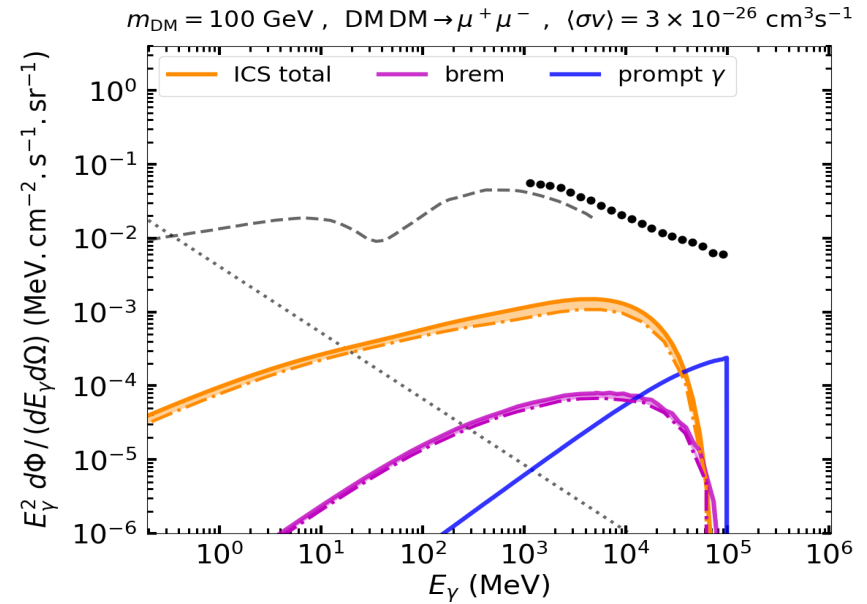
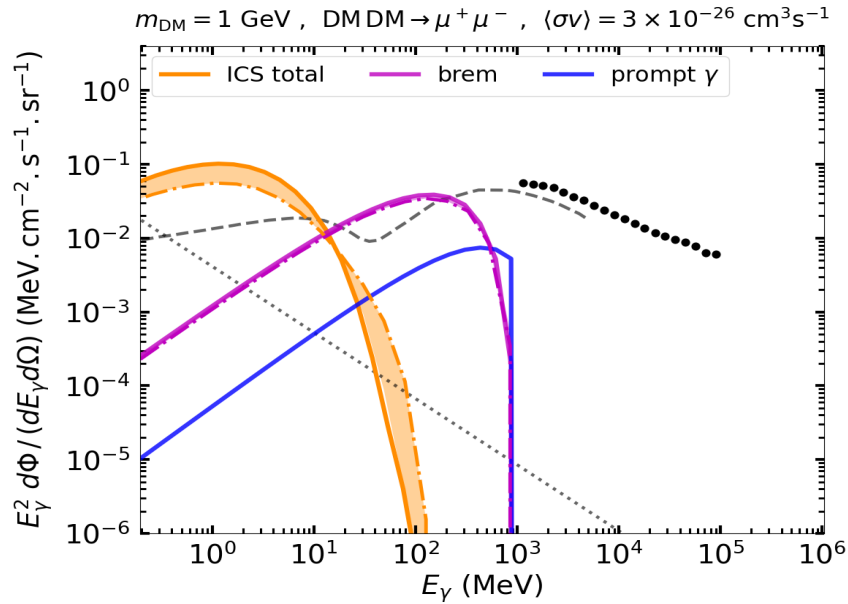
Isothermal (cored): 
$$\rho_{\text{DM}}^{\text{Iso}}(r) = \frac{\rho_0}{1 + \left(\frac{r}{r_s}\right)^2}$$



Salas *et al.*, (1906.06133)  
 Cirelli *et al.*, (2406.01705)



# ISRF models



M. Cirelli, A.K.; (SciPost Phys. 19 (2025) 080)

ISRF model 1 (solid lines):

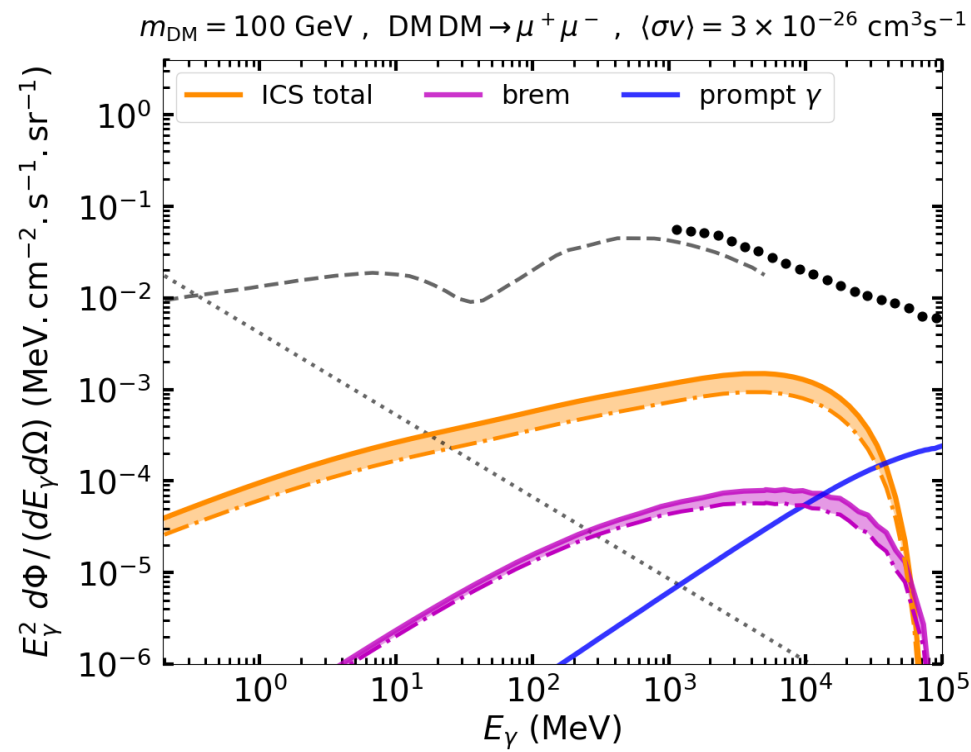
Buch, *et al.*, (PPPC 4 DM, [1505.01049]), (GALPROP)

ISRF model 2 (dashed-dotted lines):

Porter, *et al.*, (astro-ph/0507119) (used in DRAGON)



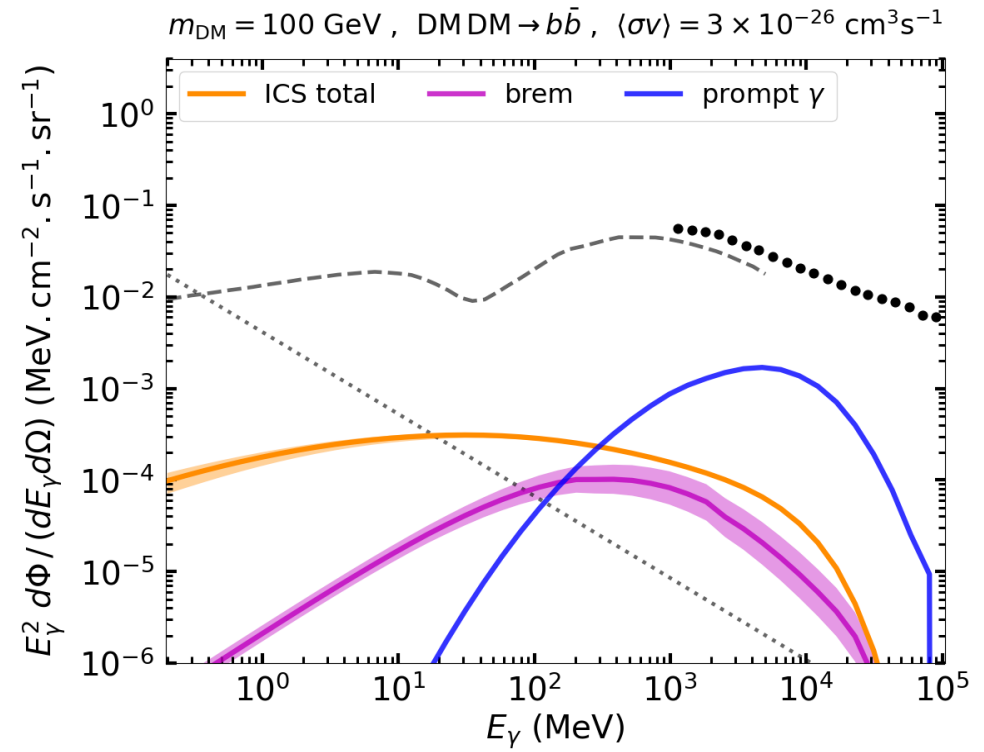
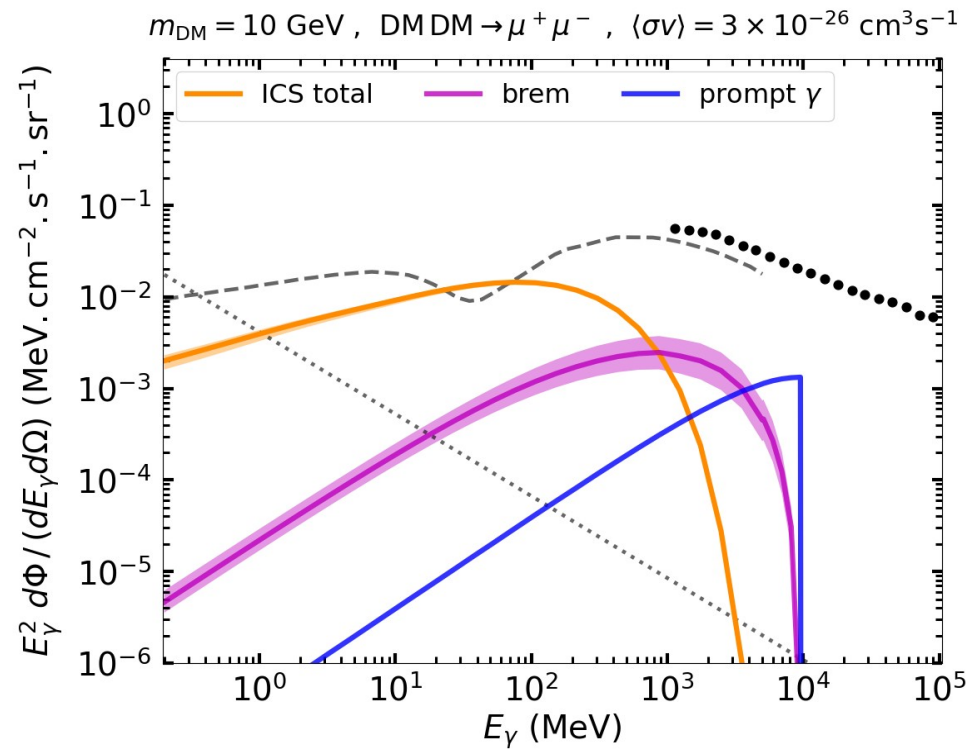
## Galactic B-field models



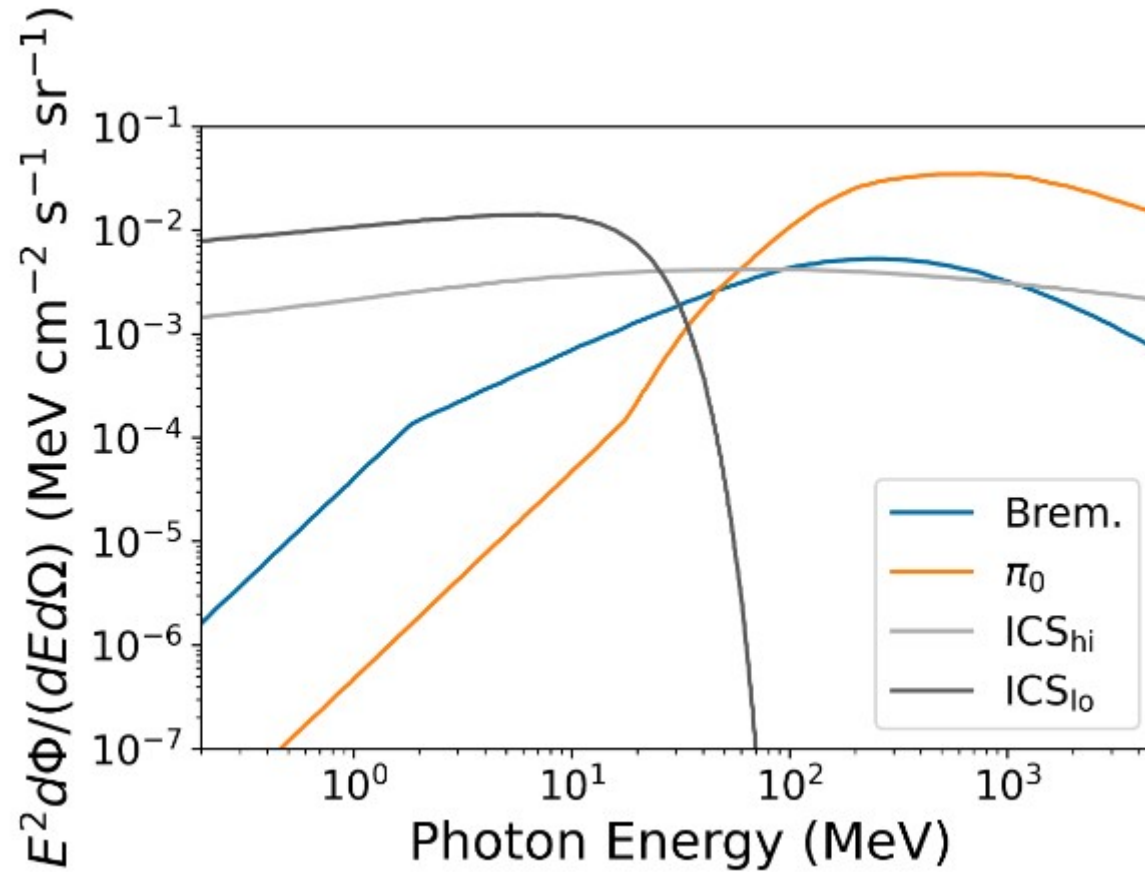
Different Galactic B-field models from:  
Buch, *et al.*, (PPPC 4 DM, [1505.01049])

M. Cirelli, A.K.; (SciPost Phys. 19 (2025) 080)

# Galactic gas map models



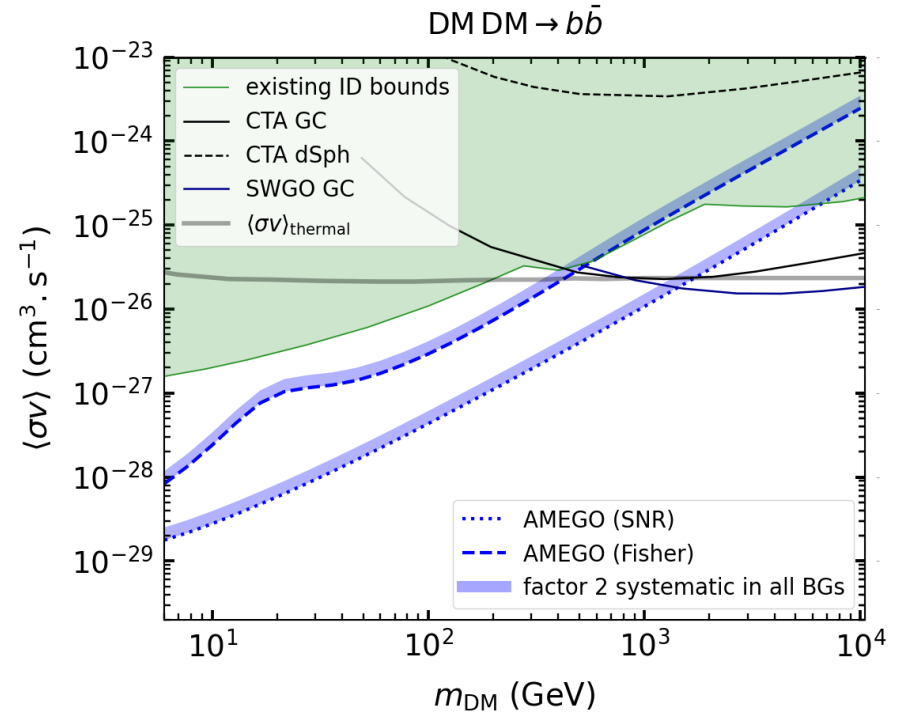
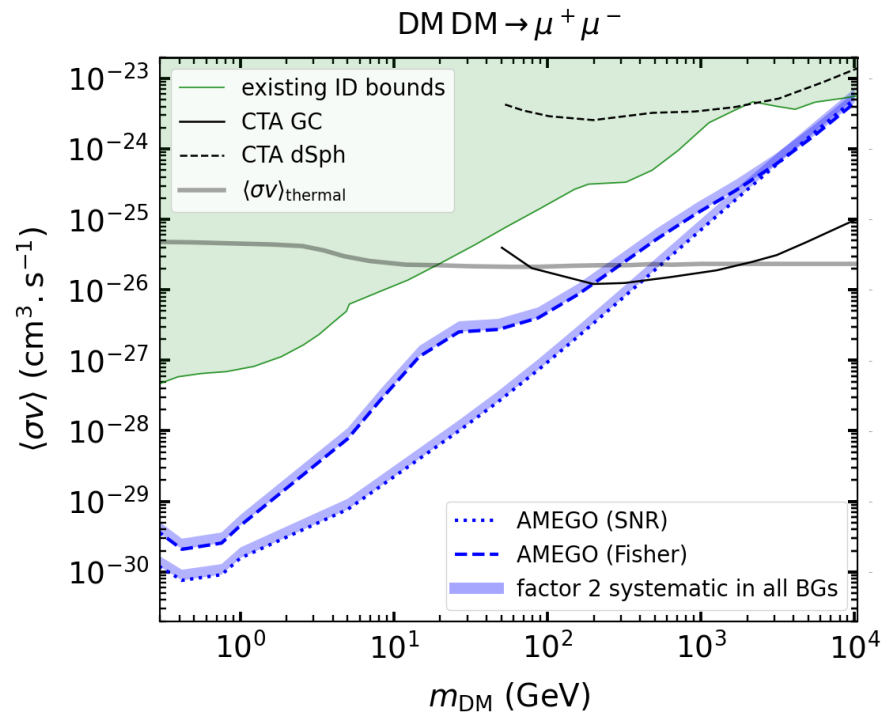
## Photon backgrounds from the inner Galaxy



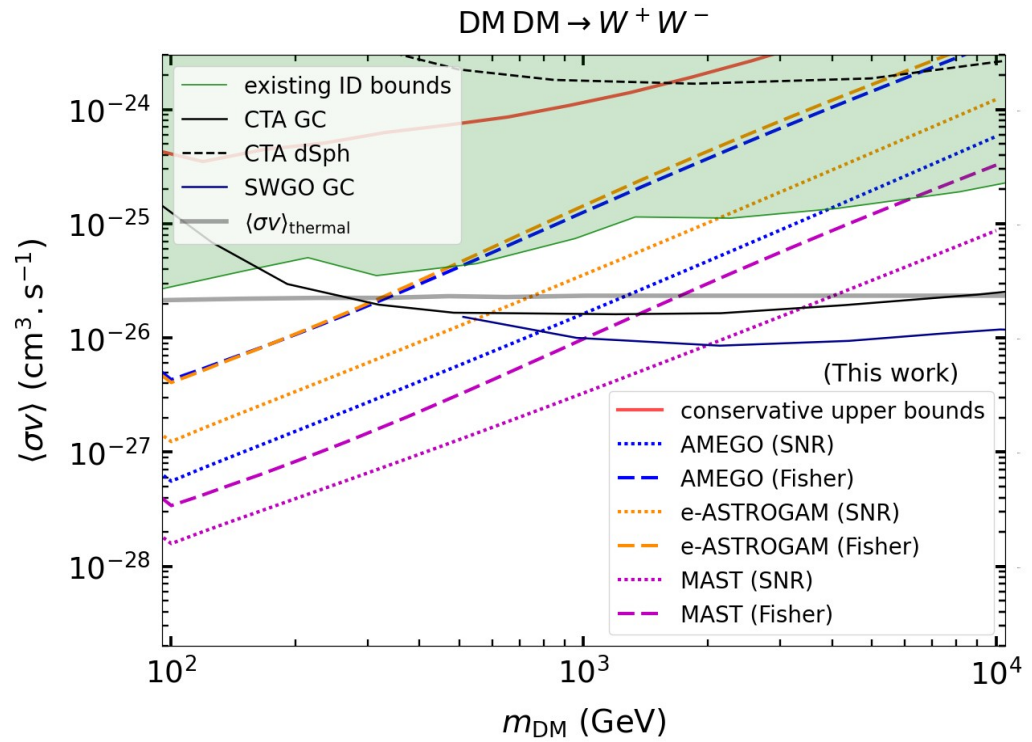
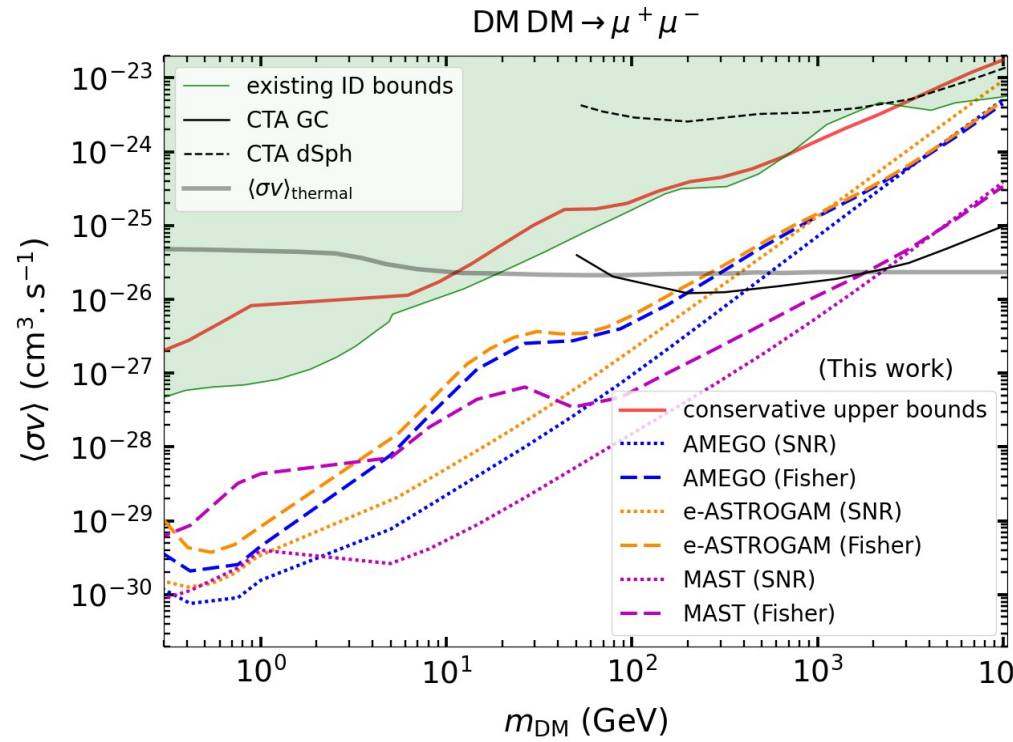
O'Donnell *et al.*, (2411.00087)

Galactic photon backgrounds from a region of  $10^\circ$  radius around the GC

# Systematic uncertainties in the backgrounds



# Signal-to-noise ratio and Fisher methods



Fisher-projections :

$$\mathcal{F}_{ij} = t_{\text{obs}} \int_{E_{\min}}^{E_{\max}} dE_{\gamma} A_{\text{eff}}(E_{\gamma}) \int_{\Delta\Omega} d\Omega \left( \frac{1}{\phi_{\text{tot}}} \frac{\partial \phi_{\text{tot}}}{\partial \theta_i} \frac{\partial \phi_{\text{tot}}}{\partial \theta_j} \right)_{\vec{\theta}=\vec{\theta}_{\text{fiducial}}}$$

$$\phi_{\text{tot}}(\vec{\theta}) = \frac{d\Phi^{\text{SIG}}}{dE_{\gamma} d\Omega}(\Gamma^{\text{SIG}}) + \sum_I \theta_I^{\text{BG}} \left\{ \frac{d\Phi_{\text{BG}}^I}{dE_{\gamma} d\Omega} \right\}_{\text{fiducial}}$$

$$\vec{\theta} = [\Gamma^{\text{SIG}}, \theta_{\text{brem}}^{\text{BG}}, \theta_{\pi^0}^{\text{BG}}, \theta_{\text{ICS}_{\text{hi}}}^{\text{BG}}, \theta_{\text{ICS}_{\text{lo}}}^{\text{BG}}, \theta_{\text{e.g.}}^{\text{BG}}]$$

$$\Gamma_{\text{proj}}^{\text{SIG}} = 2 \sqrt{(\mathcal{F}^{-1})_{11}}$$

signal-to-noise ratio (SNR) :

$$\frac{N_{\gamma}|_{\text{DM}}}{\sqrt{N_{\gamma}|_{\text{BG}}}} \geq 5$$

$$N_{\gamma} = t_{\text{obs}} \int_{E_{\min}}^{E_{\max}} dE_{\gamma} A_{\text{eff}}(E_{\gamma}) \int_{\Delta\Omega} d\Omega \frac{d\Phi}{dE_{\gamma} d\Omega}$$

## Atmospheric backgrounds

