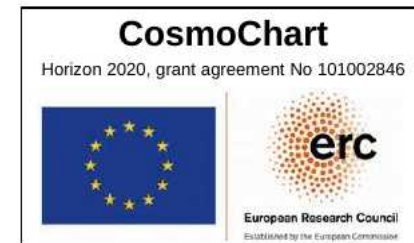


Self-interacting dark matter and inelasticity

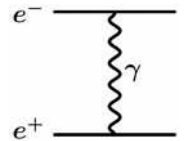
Kallia Petraki



Small-scale structure of the universe and self-interacting dark matter
Valencia, 11 June 2025

Self-interacting dark matter is great, *and complicated*

- Self-interactions refer to **elastic scattering** between DM particles
- **Light force mediators** are the canonical paradigm of SIDM
Ordinary matter analogue: $e^+ e^-$ scattering via virtual photon exchange.

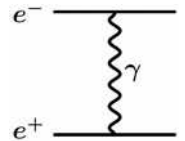


Self-interacting dark matter is great, *and complicated*

- Self-interactions refer to **elastic scattering** between DM particles

- **Light force mediators** are the canonical paradigm of SIDM

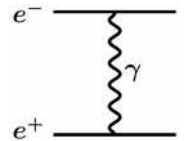
Ordinary matter analogue: $e^+ e^-$ scattering via virtual photon exchange.



- **Light force mediators imply inelastic processes:**
 - @ astrophysicists: bremsstrahlung \rightarrow dissipation
 - @ particle physicists: annihilation, bremsstrahlung, **bound-state formation**

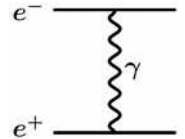
Self-interacting dark matter is great, *and complicated*

- Self-interactions refer to **elastic scattering** between DM particles
- **Light force mediators** are the canonical paradigm of SIDM
Ordinary matter analogue: $e^+ e^-$ scattering via virtual photon exchange.
- **Light force mediators imply inelastic processes:**
 - @ astrophysicists: bremsstrahlung \rightarrow dissipation
 - @ particle physicists: annihilation, bremsstrahlung, **bound-state formation**
- **Inelastic interactions change the state of DM** and affect its self interactions.
Bound states: ion-ion scattering very different from atom-atom scattering!



Self-interacting dark matter is great, *and complicated*

- Self-interactions refer to **elastic scattering** between DM particles
- Light force mediators** are the canonical paradigm of SIDM
Ordinary matter analogue: $e^+ e^-$ scattering via virtual photon exchange.
- Light force mediators imply inelastic processes:**
 - @ astrophysicists: bremsstrahlung \rightarrow dissipation
 - @ particle physicists: annihilation, bremsstrahlung, **bound-state formation**
- Inelastic interactions change the state of DM** and affect its self interactions.
Bound states: ion-ion scattering very different from atom-atom scattering!

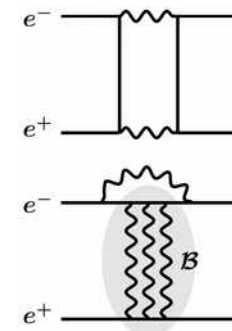


- Inelastic interactions generate elastic scattering**

$$e^+ e^- \rightarrow \gamma \gamma \rightarrow e^+ e^-$$

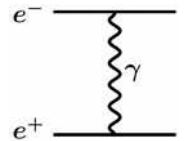
or, more importantly,

$$e^+ e^- \rightarrow \text{Positronium} + \gamma \rightarrow e^+ e^-$$



Self-interacting dark matter is great, *and complicated*

- Self-interactions refer to **elastic scattering** between DM particles
- Light force mediators** are the canonical paradigm of SIDM
Ordinary matter analogue: $e^+ e^-$ scattering via virtual photon exchange.
- Light force mediators imply inelastic processes:**
 - @ astrophysicists: bremsstrahlung \rightarrow dissipation
 - @ particle physicists: annihilation, bremsstrahlung, **bound-state formation**
- Inelastic interactions change the state of DM** and affect its self interactions.
Bound states: ion-ion scattering very different from atom-atom scattering!

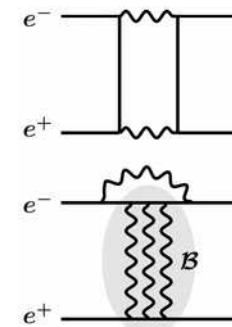


- Inelastic interactions generate elastic scattering**

$$e^+ e^- \rightarrow \gamma \gamma \rightarrow e^+ e^-$$

or, more importantly,

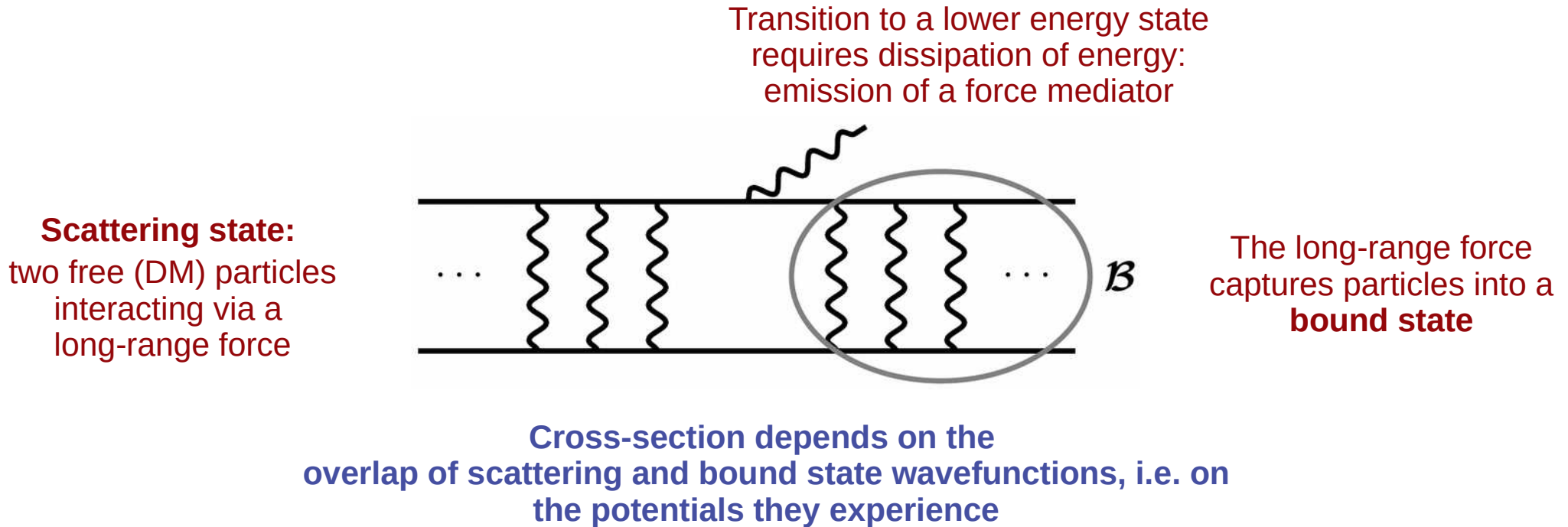
$$e^+ e^- \rightarrow \text{Positronium} + \gamma \rightarrow e^+ e^-$$



Important in many theories
with long-range interactions
due to BSF

(but not minimal (dark) QED)

Bound-state formation: standard calculation



- If the potentials of the initial and final state are the same (e.g. minimal QED): cross-section well behaved.
- **If the potentials are different** → large wavefunction overlaps → large cross-sections that can **exceed unitarity limit, even for arbitrarily small couplings.**



Oncala, Petraki: 1911.02605, 2101.08666, 2101.08667; Binder, Garny, Heisig, Lederer, Urban: 2308.01336;
Beneke, Binder, de Ros, Garny, Lederer: 2411.08737; Petraki, Socha, Vasilaki: 2505.20443.

Bound-state formation: transitions between different potentials

- In all non-Abelian theories (e.g. QCD):
non-Abelian gauge bosons are charged under the gauge symmetry, carry away conserved charge → different in/out potentials.
- In (slightly non-minimal) Abelian theories:
e.g. darkly charged fermions coupled to a dark photon and a doubly charged light scalar mediator.

$$\mathcal{L} \supset \bar{X} X \gamma + X X \Phi^\dagger$$

Capture into bound states in two ways:

$$\begin{array}{lcl}
 \underbrace{X + \bar{X}}_{\text{same potential}} \rightarrow \underbrace{\mathcal{B}(X \bar{X})}_{\text{same potential}} + \gamma & \text{well-behaved} \\
 \underbrace{X + X}_{\text{different potentials}} \rightarrow \underbrace{\mathcal{B}(X \bar{X})}_{\text{different potentials}} + \Phi & \text{unitarity violation!}
 \end{array}$$

But what is Unitarity, and what is it good for?

Cross-sections \sim (Probability for scattering) \times (phase space)

**Cannot be
larger than 1 !**

**Unitarity implies
upper bounds on elastic and inelastic cross-sections**

But what is Unitarity, and what is it good for?

Cross-sections \sim (Probability for scattering) \times (phase space)

Cannot be
larger than 1 !

Unitarity implies
upper bounds on elastic and inelastic cross-sections

Physical meaning:
Saturation of probability for scattering

But what is Unitarity, and what is it good for?

Cross-sections \sim (Probability for scattering) \times (phase space)

Cannot be
larger than 1 !

Unitarity implies
upper bounds on elastic and inelastic cross-sections

Physical meaning:
Saturation of probability for scattering

Depend on
kinematic variables (momentum),
but not on any coupling

But what is Unitarity, and what is it good for?

Cross-sections \sim (Probability for scattering) \times (phase space)

Cannot be
larger than 1 !

Unitarity implies
upper bounds on elastic and inelastic cross-sections

Physical meaning:
Saturation of probability for scattering

Depend on
kinematic variables (momentum),
but not on any coupling

If a cross-section
violates these bounds,
something has gone
terribly wrong!

Either the theory or
the calculation
need amendment.

Unitarity

$$S^\dagger S = 1 \xrightarrow{S=1+iT} -i(T - T^\dagger) = T^\dagger T$$

Project on a partial wave, insert complete set of states on RHS, and rescale amplitude, $M_\ell^{ab}(p^a, p^b) = \sqrt{4p^a p^b / s} \times \mathcal{M}_\ell^{ab}(p^a, p^b)$, to obtain

\Downarrow

Generalised optical theorem: on-/off-shell and partial-wave

$$\text{Im}[M_\ell^{aa}(p, p')] = \sum_{b: \text{ on-shell}} M_\ell^{ab*}(p, p^b) M_\ell^{ab}(p', p^b)$$

$$\text{Im} \left(M_\ell^{aa} \right) = \sum_b \left(M_\ell^{ab*} \right) \left(M_\ell^{ab} \right)$$

Unitarity

Generalised optical theorem: on-/off-shell and partial-wave

$$\text{Im}[M_\ell^{aa}(p, p')] = \sum_{b: \text{ on-shell}} M_\ell^{ab*}(p, p^b) M_\ell^{ab}(p', p^b)$$

⇓ **on-shell**

Unitarity circle:

$$|M_\ell^{aa} - \text{i}/2| \leq 1/2$$

Bounds on elastic and inelastic cross-sections:

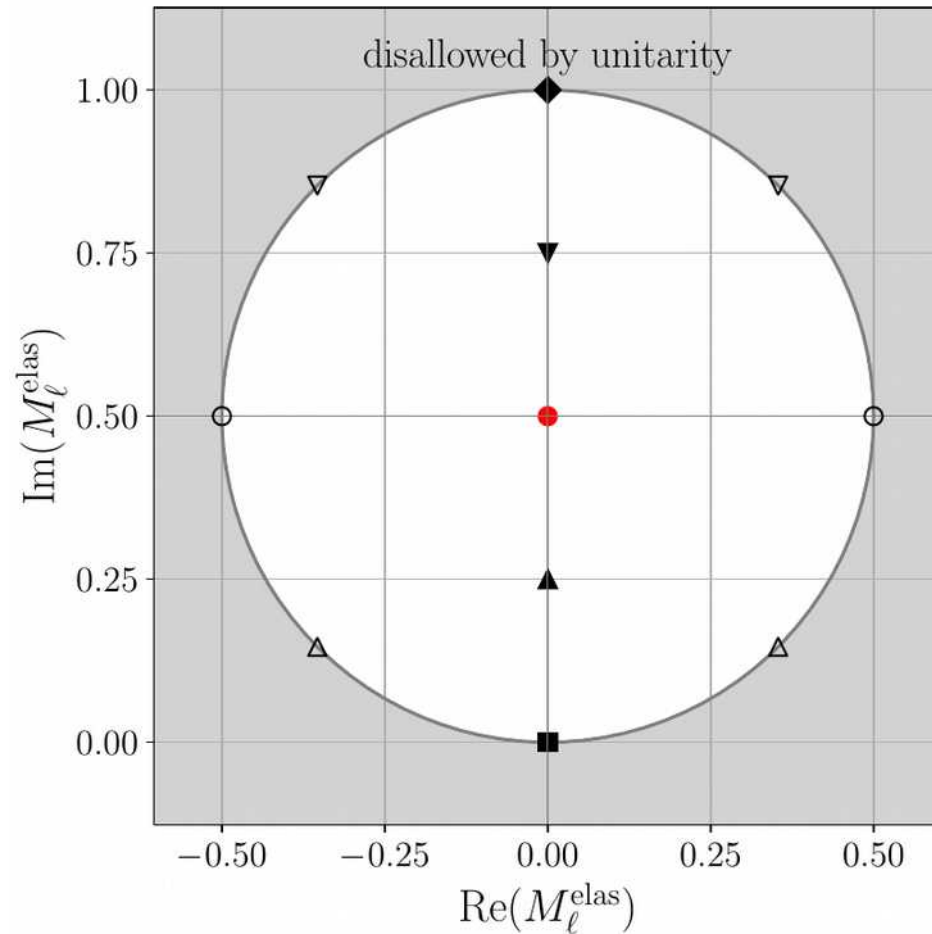
$$\sigma_\ell^{ab} = \sigma_\ell^U \times |M^{ab}|^2$$

$$\sigma_\ell^U = 4\pi(2\ell + 1)/p^2$$

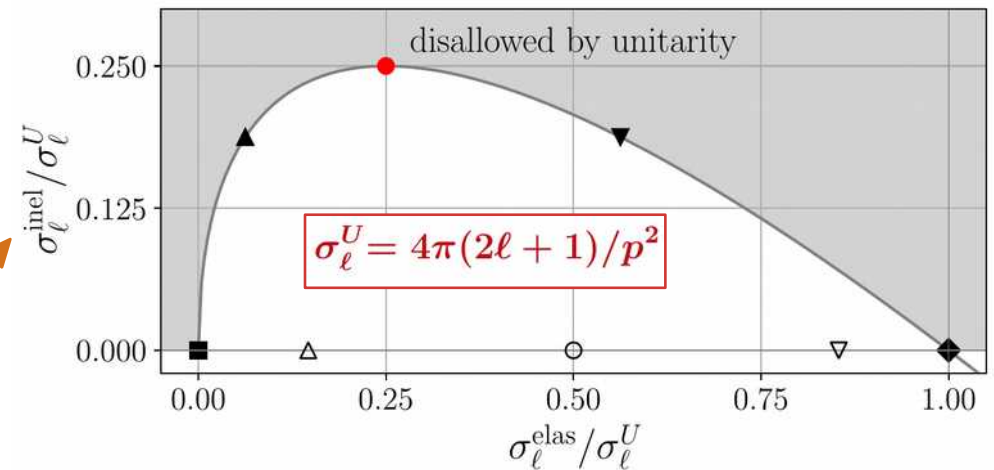
$$(\sigma_\ell^{\text{elas}} + \sigma_\ell^{\text{inel}})/\sigma_\ell^U = \sum_b |M_\ell^{ab}|^2 = \text{Im}|M_\ell^{aa}| \leq |M_\ell^{aa}| \leq \sqrt{\sigma_\ell^{\text{elas}}/\sigma_\ell^U}$$

Unitarity

Bounds on partial-wave cross-sections



Flores, KP: 2405.02222



$$\sigma_{\ell}^{\text{inel}} / \sigma_{\ell}^U \leq \sqrt{\sigma_{\ell}^{\text{elas}} / \sigma_{\ell}^U} \left(1 - \sqrt{\sigma_{\ell}^{\text{elas}} / \sigma_{\ell}^U} \right)$$

$$\sigma_{\ell}^{\text{elas}} / \sigma_{\ell}^U \leq 1$$

$$\sigma_{\ell}^{\text{inel}} / \sigma_{\ell}^U \leq 1/4$$

Inelastic scattering
generates elastic scattering!

Unitarity troubles

Unitarity is a sacred principle in particle physics,
however

$$\text{Im}[M_\ell^{aa}(p, p')] = \sum_{b: \text{ on-shell}} M_\ell^{ab*}(p, p^b) M_\ell^{ab}(p', p^b)$$

any perturbative calculation at finite order of the
couplings violates it!

Unitarity troubles

Resolution:

Calculate at infinite order!

How?

Resummation

What is resummation?

Something we learn in high-school

(but we did not know it was that)

Elastic unitarity and resummation

Consider, for example, the interaction of two charged particles

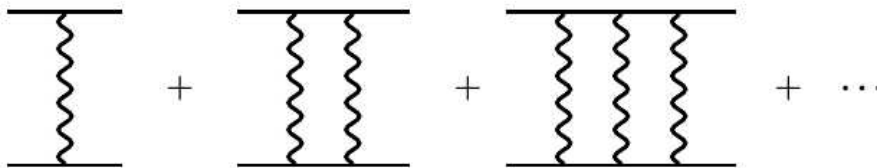
Tree-level



$$\sigma_{\ell}^{\text{elas}} \propto \frac{\alpha^2}{m^2 v_{\text{rel}}^4}$$

Rutherford scattering:
violates unitarity at low
enough velocities

Resummation



$$\sigma_{\ell}^{\text{elas}} = \frac{4\pi(2\ell + 1)}{m^2 v_{\text{rel}}^2} \times \sin^2 \left(\frac{1}{2i} \ln \left[\frac{\Gamma(1 + \ell + i\alpha/v_{\text{rel}})}{\Gamma(1 + \ell - i\alpha/v_{\text{rel}})} \right] \right)$$

Equivalent to solving
Schroedinger's eq.
for a Coulomb potential

Respects elastic unitarity at all velocities,
and spans the entire allowed range
(feature of long-range interactions)

What about inelastic unitarity?

Unitarity

Resummation of inelastic interactions

Flores, KP 2405.02222

$$\text{Im} \left(M_{\ell}^{aa} \right) = \sum_b \left(M_{\ell}^{ab*} \right) \left(M_{\ell}^{ab} \right)$$

Generalised optical theorem: on-/off-shell and partial-wave

$$\text{Im}[M_{\ell}^{aa}(p, p')] = \sum_{b: \text{on-shell}} M_{\ell}^{ab*}(p, p^b) M_{\ell}^{ab}(p', p^b)$$

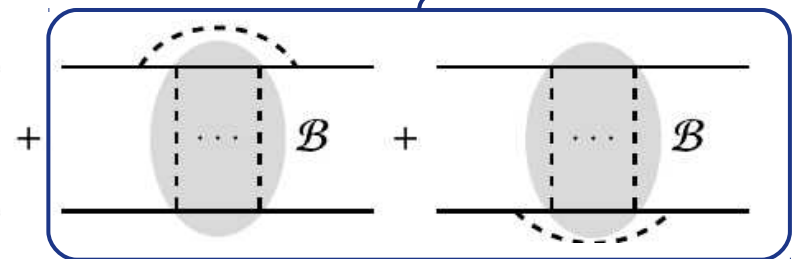
Purely
elastic
kernel

→ real potential



Squared
inelastic
kernel

→ imaginary potential



Solve Schroedinger's equation with

- (i) real potential only (→ unregulated cross-sections)
- (ii) complex potential (→ regulated cross-sections)


Unitarity


Resummation of inelastic interactions

The resummation of squared inelastic diagrams affects both elastic and inelastic cross-sections!

Flores, KP 2405.02222

Corrected
cross-sections


$$\frac{\sigma_{\ell,\text{reg}}^{\text{inel}}}{\sigma_{\ell}^U} = \frac{\sigma_{\ell,\text{unreg}}^{\text{inel}}/\sigma_{\ell}^U}{\left(1 + \sigma_{\ell,\text{unreg}}^{\text{inel}}/\sigma_{\ell}^U\right)^2}$$


$$\frac{\sigma_{\ell,\text{reg}}^{\text{elas}}}{\sigma_{\ell}^U} = \frac{\sigma_{\ell,\text{unreg}}^{\text{elas}}/\sigma_{\ell}^U}{\left(1 + \sigma_{\ell,\text{unreg}}^{\text{inel}}/\sigma_{\ell}^U\right)^2} + \left(1 - \sigma_{\ell,\text{unreg}}^{\text{elas}}/\sigma_{\ell}^U\right) \times \frac{\left(\sigma_{\ell,\text{unreg}}^{\text{inel}}/\sigma_{\ell}^U\right)^2}{\left(1 + \sigma_{\ell,\text{unreg}}^{\text{inel}}/\sigma_{\ell}^U\right)^2}$$

Unitarity

Resummation of inelastic interactions

The resummation of squared inelastic diagrams affects both elastic and inelastic cross-sections!

Flores, KP 2405.02222

Corrected
cross-sections

$$\frac{\sigma_{\ell,\text{reg}}^{\text{inel}}}{\sigma_{\ell}^U} = \frac{\sigma_{\ell,\text{unreg}}^{\text{inel}}/\sigma_{\ell}^U}{\left(1 + \sigma_{\ell,\text{unreg}}^{\text{inel}}/\sigma_{\ell}^U\right)^2}$$

regulating factor:
reduced flux due to inelastic scattering
(reflects continuity equation)

Important if $\sigma_{\text{unreg}} \gtrsim \sigma^U$.

$$\frac{\sigma_{\ell,\text{reg}}^{\text{elas}}}{\sigma_{\ell}^U} = \frac{\sigma_{\ell,\text{unreg}}^{\text{elas}}/\sigma_{\ell}^U}{\left(1 + \sigma_{\ell,\text{unreg}}^{\text{inel}}/\sigma_{\ell}^U\right)^2} + \left(1 - \sigma_{\ell,\text{unreg}}^{\text{elas}}/\sigma_{\ell}^U\right) \times \frac{\left(\sigma_{\ell,\text{unreg}}^{\text{inel}}/\sigma_{\ell}^U\right)^2}{\left(1 + \sigma_{\ell,\text{unreg}}^{\text{inel}}/\sigma_{\ell}^U\right)^2}$$

Unitarity

Resummation of inelastic interactions

The resummation of squared inelastic diagrams affects both elastic and inelastic cross-sections!

Flores, KP 2405.02222

Corrected cross-sections

$$\frac{\sigma_{\ell,\text{reg}}^{\text{inel}}}{\sigma_{\ell}^U} = \frac{\sigma_{\ell,\text{unreg}}^{\text{inel}}/\sigma_{\ell}^U}{\left(1 + \sigma_{\ell,\text{unreg}}^{\text{inel}}/\sigma_{\ell}^U\right)^2}$$

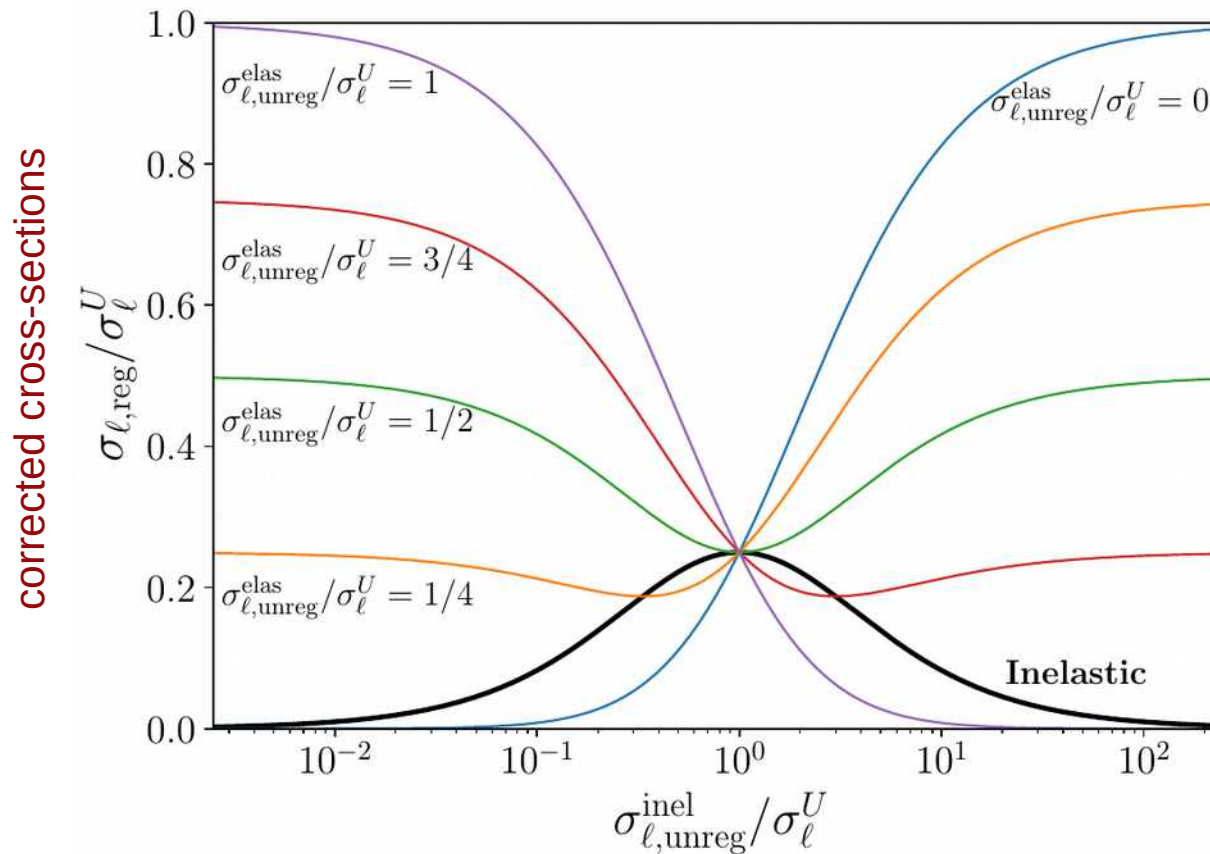
$$\frac{\sigma_{\ell,\text{reg}}^{\text{elas}}}{\sigma_{\ell}^U} = \frac{\sigma_{\ell,\text{unreg}}^{\text{elas}}/\sigma_{\ell}^U}{\left(1 + \sigma_{\ell,\text{unreg}}^{\text{inel}}/\sigma_{\ell}^U\right)^2} + \left(1 - \sigma_{\ell,\text{unreg}}^{\text{elas}}/\sigma_{\ell}^U\right) \times \frac{\left(\sigma_{\ell,\text{unreg}}^{\text{inel}}/\sigma_{\ell}^U\right)^2}{\left(1 + \sigma_{\ell,\text{unreg}}^{\text{inel}}/\sigma_{\ell}^U\right)^2}$$

original scattering probability, suppressed by regulating factor

Additional contribution: regenerated flux

Unitarity

Resummation of inelastic interactions



uncorrected inelastic cross-section

- Affects both strength and velocity dependence of the cross-sections.
- Can increase or decrease elastic scattering.
- Inelastic cross-sections: affect DM production in the early universe
 - correlation of parameters (mass-coupling relation)
 - DM self-interactions
- Elastic cross-section: affects DM scattering inside halos

[KP, Socha, Vasilaki 2505.20443]

[Flores, KP, Pignard: in progress]

Conclusions

Inelasticity is particularly important in the presence of light force mediators and affects elastic scattering in “real” and “virtual” ways:

- “Real” ways: formation of bound states, transitions between levels, dissipation.
- “Virtual” ways: inelastic scatterings feed into elastic scatterings.

- Unitarity has always been an important guide in finding new physics (e.g. discovery of electroweak interactions, Higgs etc).

Employ it for SIDM!

- Strong connection between unitarity limits and long-range interactions in the non-relativistic regime.