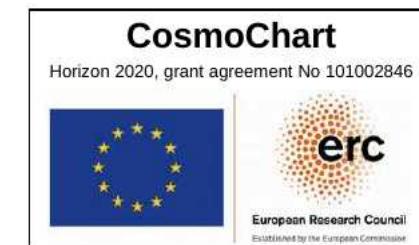


Self-interacting dark matter and inelasticity

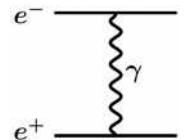
Kallia Petraki



Small-scale structure of the universe and self-interacting dark matter
Valencia, 11 June 2025

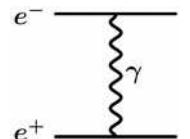
Self-interacting dark matter is great, *and complicated*

- Self-interactions refer to **elastic scattering** between DM particles
- **Light force mediators** are the canonical paradigm of SIDM
Ordinary matter analogue: $e^+ e^-$ scattering via virtual photon exchange.

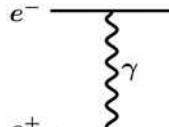


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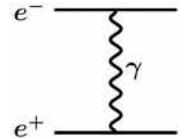
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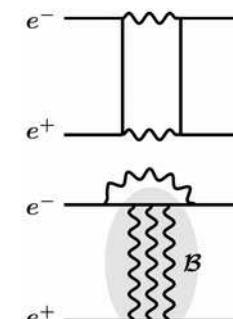
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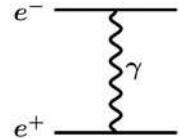
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or, more importantly,

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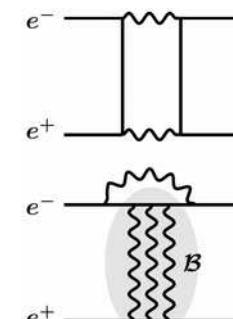
Important in many theories
with long-range interactions
due to BSF

(but not minimal (dark) QED)

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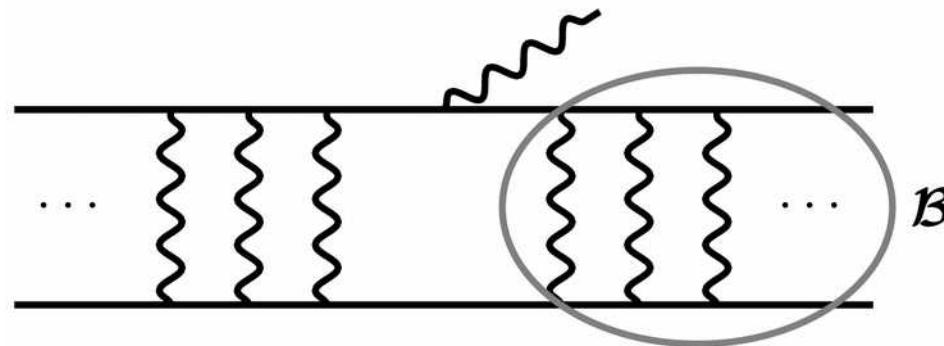
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Bound-state formation: standard calculation

Scattering state:
two free (DM) particles
interacting via a
long-range force

Transition to a lower energy state
requires dissipation of energy:
emission of a force mediator



The long-range force
captures particles into a
bound state

Cross-section depends on the
overlap of scattering and bound state wavefunctions, i.e. on
the potentials they experience

- If the potentials of the initial and final state are the same (e.g. minimal QED): cross-section well behaved.
- **If the potentials are different** → large wavefunction overlaps → large cross-sections that can **exceed unitarity limit, even for arbitrarily small couplings.**

Oncala, Petraki: 1911.02605, 2101.08666, 2101.08667; Binder, Garry, Heisig, Lederer, Urban: 2308.01336;
Beneke, Binder, de Ros, Garry, Lederer: 2411.08737; Petraki, Socha, Vasilaki: 2505.20443.



Bound-state formation: transitions between different potentials

- In all non-Abelian theories (e.g. QCD):
non-Abelian gauge bosons are charged under the gauge symmetry,
carry away conserved charge \rightarrow different in/out potentials.
- In (slightly non-minimal) Abelian theories:
e.g. darkly charged fermions coupled to a dark photon and a doubly charged light scalar mediator.

$$\mathcal{L} \supset \bar{X}X\gamma + XX\Phi^\dagger$$

Capture into bound states in two ways:

$$\begin{array}{ccc} \text{same potential} & & \text{well-behaved} \\ \overbrace{X + \bar{X}}^{\text{same potential}} \rightarrow \mathcal{B}(X\bar{X}) + \gamma & & \text{well-behaved} \\ \overbrace{X + X}^{\text{different potentials}} \rightarrow \mathcal{B}(X\bar{X}) + \Phi & & \text{unitarity violation!} \end{array}$$

But what is Unitarity, and what is it good for?

Cross-sections \sim (Probability for scattering) \times (phase space)

Cannot be
larger than 1 !

Unitarity implies
upper bounds on elastic and inelastic cross-sections

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Depend on
kinematic variables (momentum),
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If a cross-section
violates these bounds,
something has gone
terribly wrong!

Either the theory or
the calculation
need amendment.

Unitarity

$$S^\dagger S = 1 \quad \xrightarrow{S=1+iT} \quad -i(T - T^\dagger) = T^\dagger T$$

Project on a partial wave, insert complete set of states on RHS, and rescale amplitude, $M_\ell^{ab}(p^a, p^b) = \sqrt{4p^a p^b / s} \times \mathcal{M}_\ell^{ab}(p^a, p^b)$, to obtain

↓

Generalised optical theorem: on-/off-shell and partial-wave

$$\text{Im}[M_\ell^{aa}(p, p')] = \sum_{b: \text{ on-shell}} M_\ell^{ab*}(p, p^b) M_\ell^{ab}(p', p^b)$$

$$\text{Im } \begin{array}{c} \text{---} \\ \text{---} \\ M_\ell^{aa} \\ \text{---} \\ \text{---} \end{array} = \sum_b \left(\begin{array}{c} \text{---} \\ \text{---} \\ M_\ell^{ab*} \\ \text{---} \\ \text{---} \end{array} \right) \left(\begin{array}{c} \text{---} \\ \text{---} \\ M_\ell^{ab} \\ \text{---} \\ \text{---} \end{array} \right)$$

Unitarity

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⇓ on-shell

Unitarity circle:

$$|M_\ell^{aa} - i/2| \leq 1/2$$

Bounds on elastic and inelastic cross-sections:

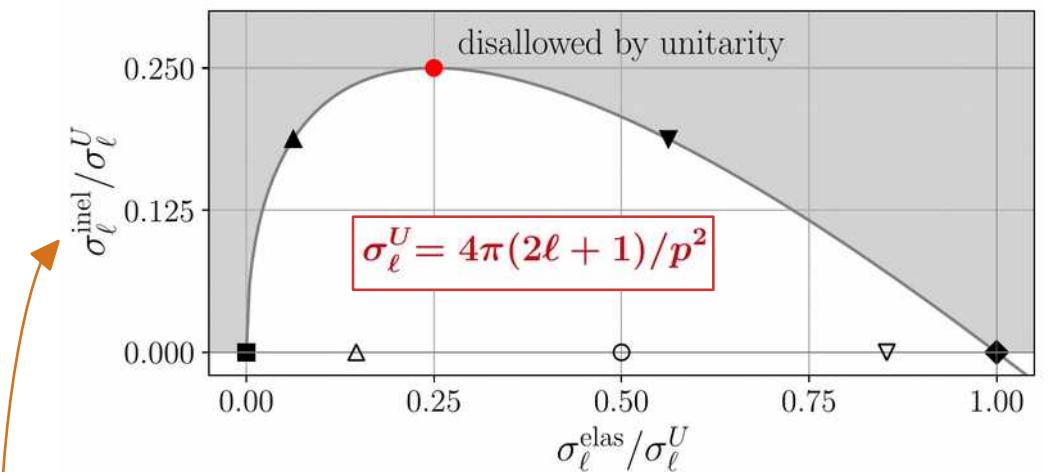
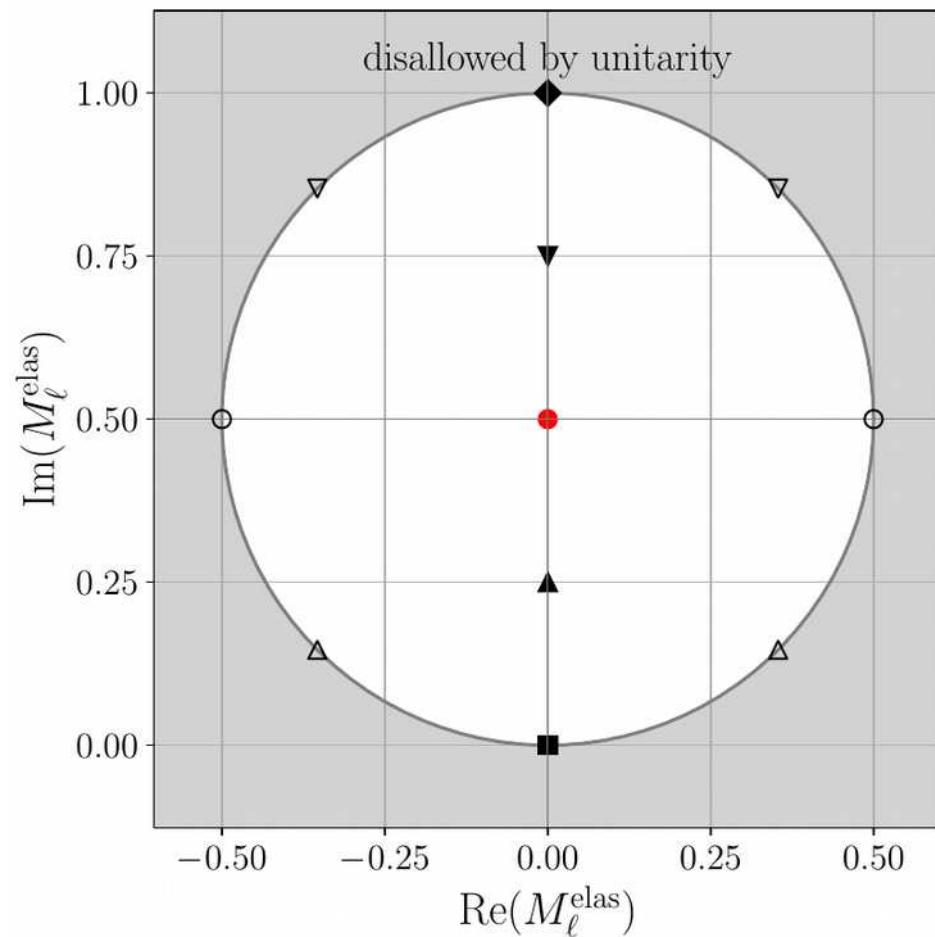
$$\sigma_\ell^{ab} = \sigma_\ell^U \times |M^{ab}|^2$$

$$\sigma_\ell^U = 4\pi(2\ell + 1)/p^2$$

$$(\sigma_\ell^{\text{elas}} + \sigma_\ell^{\text{inel}})/\sigma_\ell^U = \sum_b |M_\ell^{ab}|^2 = \text{Im}|M_\ell^{aa}| \leq |M_\ell^{aa}| \leq \sqrt{\sigma_\ell^{\text{elas}}/\sigma_\ell^U}$$

Unitarity

Bounds on partial-wave cross-sections



$$\sigma_\ell^{\text{inel}}/\sigma_\ell^U \leq \sqrt{\sigma_\ell^{\text{elas}}/\sigma_\ell^U} \left(1 - \sqrt{\sigma_\ell^{\text{elas}}/\sigma_\ell^U} \right)$$

$$\sigma_\ell^{\text{elas}}/\sigma_\ell^U \leq 1$$

$$\sigma_\ell^{\text{inel}}/\sigma_\ell^U \leq 1/4$$

Inelastic scattering generates elastic scattering!

Unitarity troubles

**Unitarity is a sacred principle in particle physics,
however**

$$\text{Im}[M_\ell^{aa}(p, p')] = \sum_{b: \text{ on-shell}} M_\ell^{ab*}(p, p^b) M_\ell^{ab}(p', p^b)$$

**any perturbative calculation at finite order of the
couplings violates it!**

Unitarity troubles

Resolution:

Calculate at infinite order!

How?

Resummation

What is resummation?

Something we learn in high-school
(but we did not know it was that)

Elastic unitarity and resummation

Consider, for example, the interaction of two charged particles

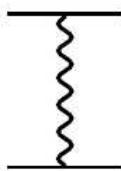
Tree-level



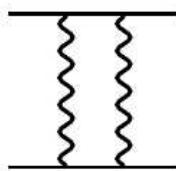
Rutherford scattering:
violates unitarity at low
enough velocities

$$\sigma_\ell^{\text{elas}} \propto \frac{\alpha^2}{m^2 v_{\text{rel}}^4}$$

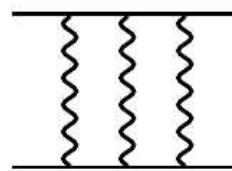
Resummation



+



+



+

...

$$\sigma_\ell^{\text{elas}} = \frac{4\pi(2\ell+1)}{m^2 v_{\text{rel}}^2} \times \sin^2 \left(\frac{1}{2i} \ln \left[\frac{\Gamma(1+\ell+i\alpha/v_{\text{rel}})}{\Gamma(1+\ell-i\alpha/v_{\text{rel}})} \right] \right)$$

Equivalent to solving
Schroedinger's eq.
for a Coulomb potential

Respects elastic unitarity at all velocities,
and spans the entire allowed range
(feature of long-range interactions)

What about inelastic unitarity?

Unitarity

Resummation of inelastic interactions

Flores, KP 2405.02222

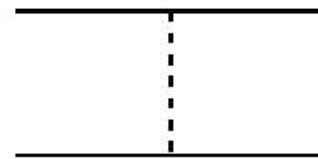
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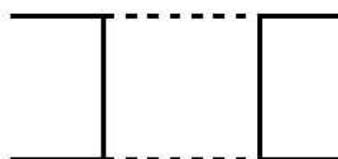
Purely elastic kernel

→ real potential

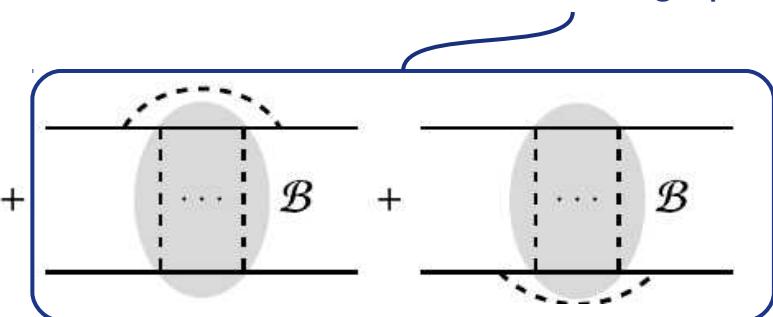


Squared inelastic kernel

→ imaginary potential



Contributions exist even for non-annihilating species



Solve Schrödinger's equation with

- (i) real potential only (→ unregulated cross-sections)
- (ii) complex potential (→ regulated cross-sections)

Unitarity

Resummation of inelastic interactions

The resummation of squared inelastic diagrams affects both elastic and inelastic cross-sections!

Flores, KP 2405.02222

$$\frac{\sigma_{\ell, \text{reg}}^{\text{inel}}}{\sigma_{\ell}^U} = \frac{\sigma_{\ell, \text{unreg}}^{\text{inel}} / \sigma_{\ell}^U}{\left(1 + \sigma_{\ell, \text{unreg}}^{\text{inel}} / \sigma_{\ell}^U\right)^2}$$

Corrected
cross-sections

$$\frac{\sigma_{\ell, \text{reg}}^{\text{elas}}}{\sigma_{\ell}^U} = \frac{\sigma_{\ell, \text{unreg}}^{\text{elas}} / \sigma_{\ell}^U}{\left(1 + \sigma_{\ell, \text{unreg}}^{\text{inel}} / \sigma_{\ell}^U\right)^2} + \left(1 - \sigma_{\ell, \text{unreg}}^{\text{elas}} / \sigma_{\ell}^U\right) \times \frac{\left(\sigma_{\ell, \text{unreg}}^{\text{inel}} / \sigma_{\ell}^U\right)^2}{\left(1 + \sigma_{\ell, \text{unreg}}^{\text{inel}} / \sigma_{\ell}^U\right)^2}$$

Unitarity

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regulating factor:
reduced flux due to inelastic scattering
(reflects continuity equation)

Important if $\sigma_{\text{unreg}} \gtrsim \sigma^U$.

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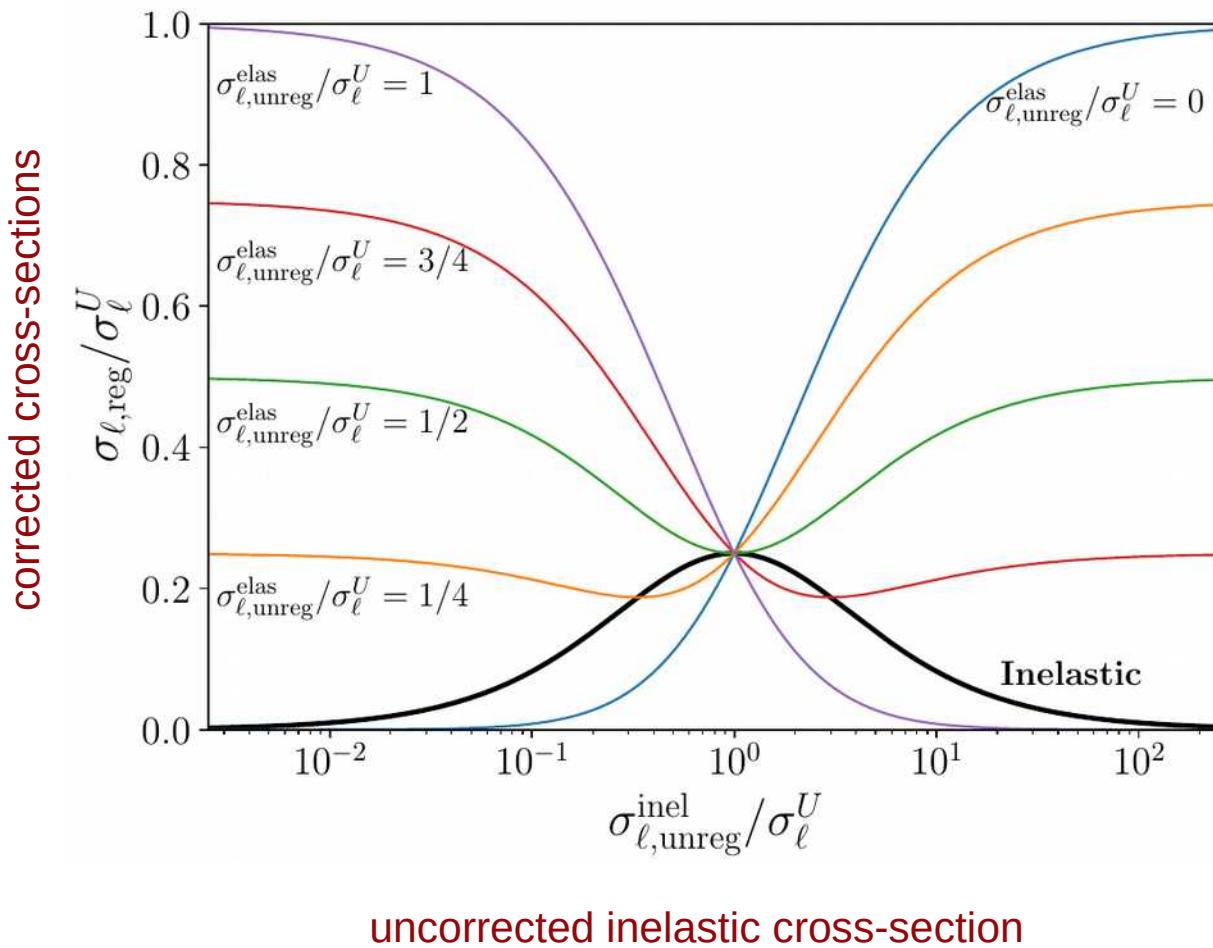
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original scattering probability,
suppressed by regulating factor

Additional contribution:
regenerated flux

Unitarity

Resummation of inelastic interactions



- Affects both strength and velocity dependence of the cross-sections.
- Can increase or decrease elastic scattering.
- Inelastic cross-sections: affect DM production in the early universe
 - correlation of parameters (mass-coupling relation)
 - DM self-interactions
- Elastic cross-section: affects DM scattering inside halos

[Flores, KP, Pignard: in progress]

Conclusions

Inelasticity is particularly important in the presence of light force mediators and affects elastic scattering in “real” and “virtual” ways:

- “Real” ways: formation of bound states, transitions between levels, dissipation.
- “Virtual” ways: inelastic scatterings feed into elastic scatterings.

- Unitarity has always been an important guide in finding new physics (e.g. discovery of electroweak interactions, Higgs etc).
Employ it for SIDM!
- Strong connection between unitarity limits and long-range interactions in the non-relativistic regime.