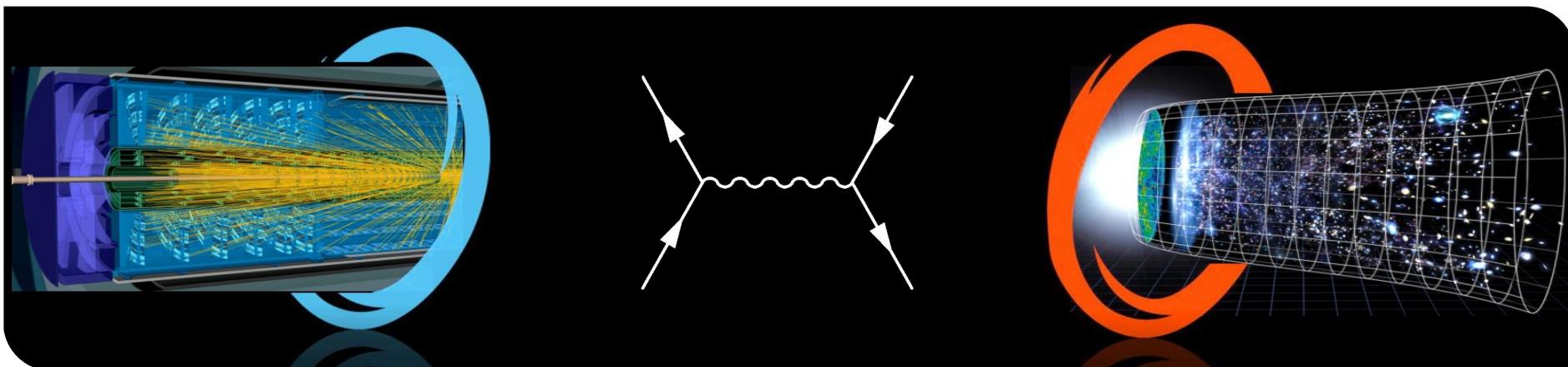


What is the mass scale of self-interacting dark matter?

Felix Kahlhoefer

Workshop on the Small-Scale Structure of the Universe
and Self-Interacting Dark Matter, Valencia, 13 June 2025



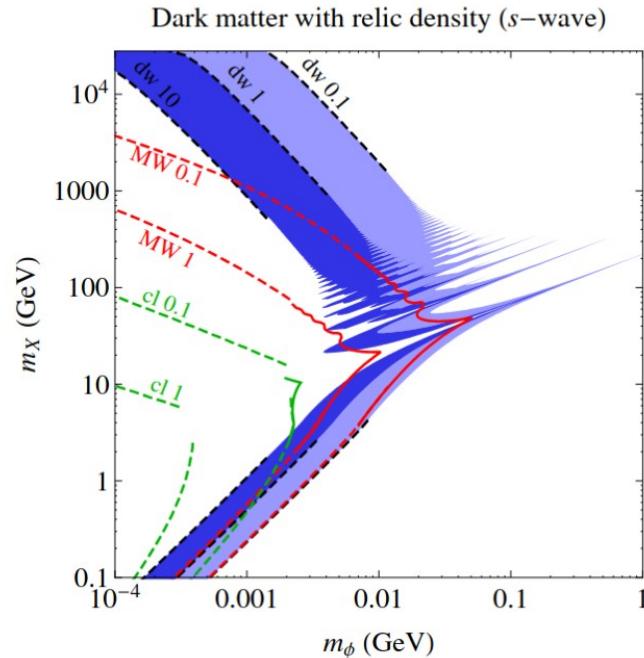
Back-of-the envelope calculation

- Typical cross section of interest:
 $\sigma / m_{\text{DM}} > 0.5 \text{ cm}^2 / \text{g} \sim 1 \text{ barn} / \text{GeV}$
- Assume there is only a single mass/energy scale v in the dark sector
- Dimensional analysis: $m_{\text{DM}} \sim v$ and $\sigma \sim v^{-2}$
 - Required scale: $v \sim 100 \text{ MeV}$



SIDM with several mass scales

- Most popular SIDM model (long-range interaction via light mediator exchange) requires two separate scales: **DM mass and mediator mass**
- Imposing relic density constraint (DM production via freeze-out mechanism) fixes the interaction strength, but neither of the mass scales
 - DM mass can vary over **many orders of magnitude**
 - Potentially **large hierarchy** between DM and mediator mass
 - Large self-interactions **accommodated but not predicted**

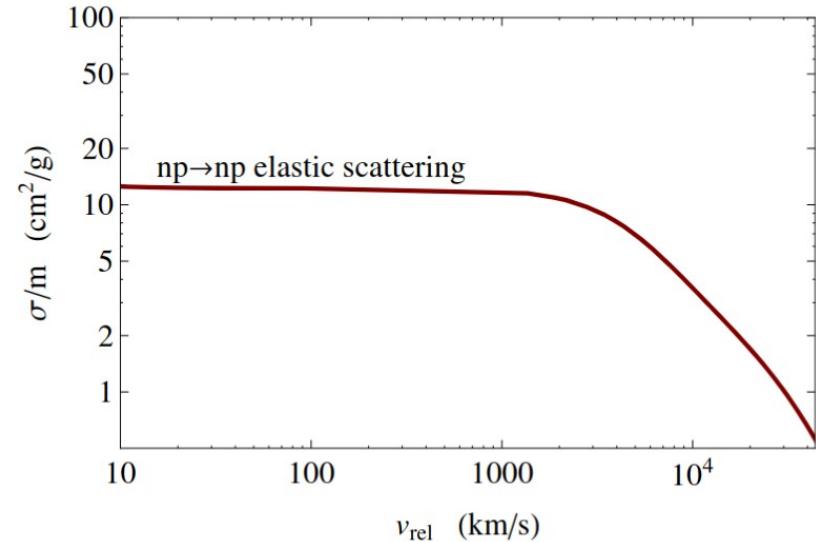


Tulin et al., arXiv:1302.3898

SIDM without a large hierarchy of scales

- **Curious observation:** $v \sim 100$ MeV is close to QCD confinement scale $\Lambda_{\text{QCD}} \sim 200$ MeV and to the pion decay constant $f_\pi \sim 92$ MeV
- Indeed, 1 barn/GeV is roughly the nucleon-nucleon scattering cross section
 - Compelling reason to think about DM particles with interactions similar to QCD

Can we get SIDM from new strong dynamics?



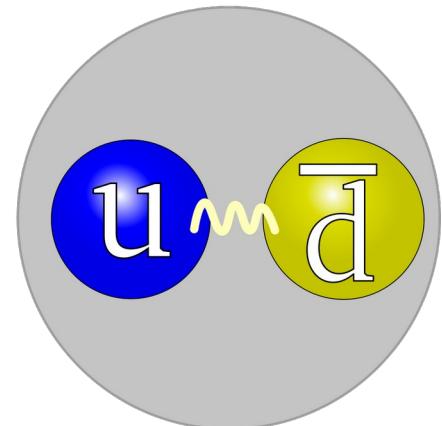
Tulin & Yu, arXiv:1705.02358

Strongly-interacting dark sectors

- Consider a dark sector that contains dark gluons and dark quarks:

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu}^a F^{\mu\nu a} + \bar{q}_d i \not{D} q_d - \bar{q}_d M_q q_d$$

- For energies below some scale Λ_d the dark sector confines, giving rise to dark mesons and dark baryons
- In contrast to the SM, the lightest dark mesons (i.e. the dark pions) may be stable and possible DM candidates



Chiral symmetry breaking

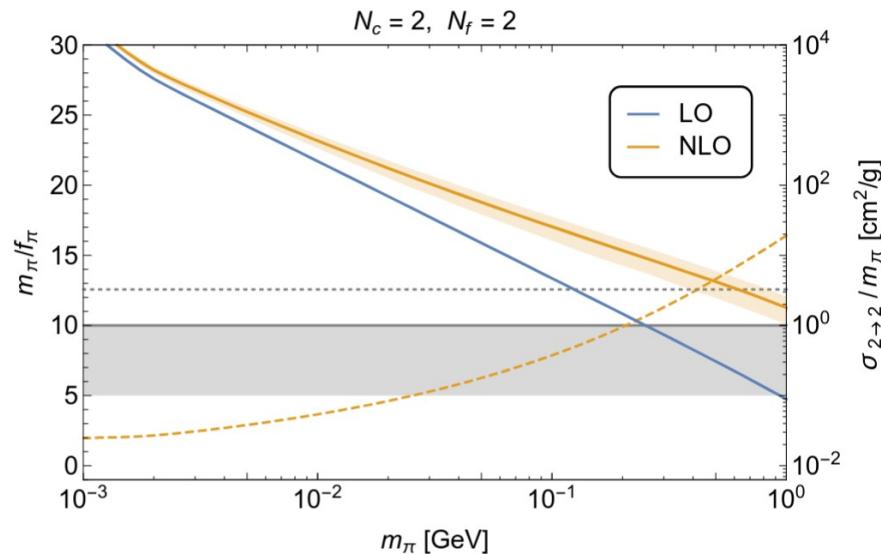
- In the absence of a mass term for the dark quarks, the dark pions are **Goldstone bosons of chiral symmetry breaking** and therefore massless
- A strongly-interacting dark sector therefore requires two independent mass scales
 - M_q and Λ_d (at high energies)
 - m_π and f_π (at low energies)
- SIDM cross section: $\sigma / m_{DM} \sim m_\pi / f_\pi^4$
- In the Standard Model, m_π and f_π are quite close, but with no fundamental reason
- Can we narrow down the range of m_π / f_π ?

Relic density requirement

- Fix combination of m_π and f_π from the relic density requirement
- **Simplest case:** Consider $3 \rightarrow 2$ annihilations via Wess-Zumino-Witten anomaly

See Xiaoyong's talk

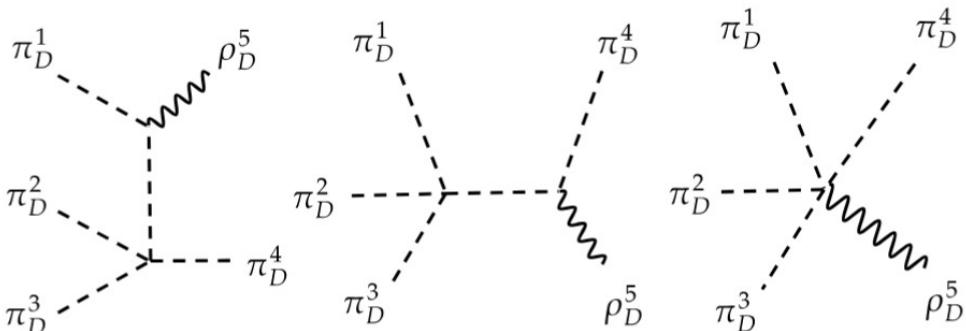
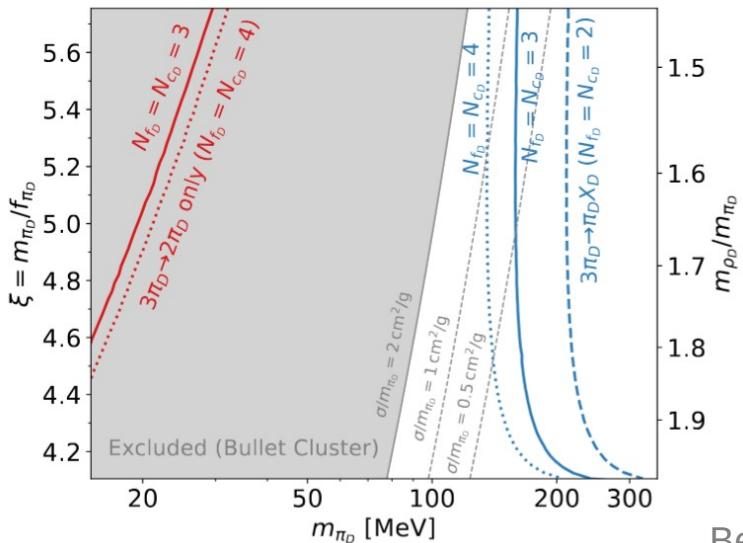
- Perturbativity ($m_\pi/f_\pi < 4\pi$) implies $m_\pi < 1$ GeV (dashed line)
- Predicted self-interaction cross section too large (solid line)



Hansen et al., arXiv:1507.01590

Strongly interacting DM with light vector mesons

- **Lattice simulations:** For $m_\pi/f_\pi > 4$, expect light vector mesons with $m_\rho < 2 m_\pi$
→ Relevant for relic density calculation



- Relic density requirement points to $m_\pi \sim 200$ MeV
- Predicted self-interaction cross section $\sim 1 \text{ cm}^2 / \text{g}$
→ Exactly in the interesting range!
- **But:** no significant velocity dependence expected

Bernreuther, FK et al., arXiv:2311.17157

Dark sectors with only a single mass scale

Key idea: Start with a dark sector without any mass scale

- Consider dark sector with a scalar field Φ and **conformal symmetry**
 - Invariance under rescaling $x^\mu \rightarrow \alpha x^\mu$, $\Phi(x) \rightarrow \alpha^{-1} \Phi(\alpha x)$
 - All mass terms in the Lagrangian forbidden, all particles massless at tree-level
- **But:** Radiative corrections introduce scale dependence
- **Coleman Weinberg mechanism:** effective potential has minimum at non-zero field value

$$V_{\text{eff}}(\phi_b, T = 0) = \sum_a \frac{g_a \eta_a}{64\pi^2} \frac{m_a^4(v)}{v^4} \phi_b^4 \left[\log \left(\frac{\phi_b^2}{v^2} \right) - \frac{1}{2} \right]$$

- Vacuum expectation value generates mass terms and breaks conformal symmetry

Connection to dark matter

- Assume dark scalar Φ and fermion fields χ charged under $U(1)'$ gauge symmetry

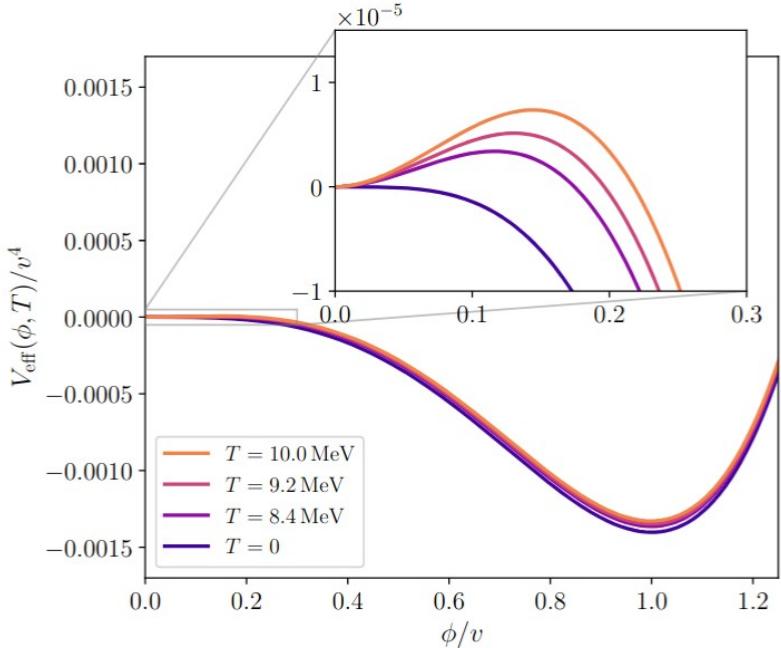
$$\mathcal{L} = |D_\mu \Phi|^2 - \frac{1}{4} F'_{\mu\nu} F'^{\mu\nu} + \bar{\chi}_1 i \not{D} \chi_1 + \bar{\chi}_2 i \not{D} \chi_2 - \left(\frac{y_1}{2} \Phi \bar{\chi}_1^c \chi_1 + \frac{y_2}{2} \Phi^* \bar{\chi}_2^c \chi_2 + \text{h.c.} \right) - V(\Phi)$$

- Conformal symmetry forbids all mass terms
- All masses must be proportional to vacuum expectation value of dark scalar
 - Gauge interactions: Dark photon mass
 - Yukawa interactions: Dark matter mass
- Interaction terms:

$$\mathcal{L}_{\text{int}} \supset -\lambda v \phi^3 - \frac{\lambda}{4} \phi^4 + g^2 v A'_\mu A'^\mu \phi + \frac{g^2}{2} \phi^2 A'_\mu A'^\mu + \sum_i \frac{g}{2} Q_i \bar{\chi}'_i \gamma^\mu \gamma^5 \chi'_i A'_\mu - \frac{y}{2\sqrt{2}} \phi \bar{\chi}'_i \chi'_i$$

Conformal dark sector at finite temperature

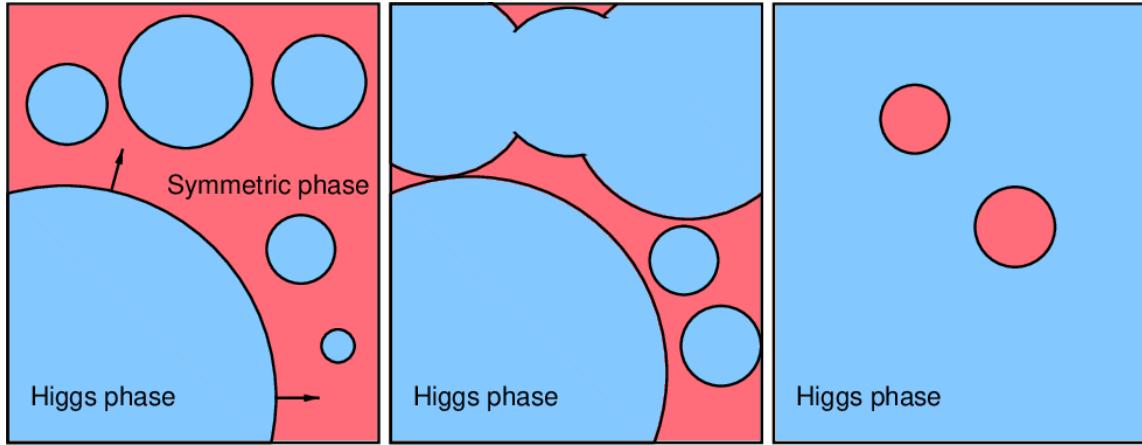
- At high temperatures and densities (early universe), conformal symmetry is restored
- As the universe cools down, symmetric minimum becomes metastable
- Potential barrier prevents transition to energetically favourable ground state
 - Conformal symmetry breaking happens through a (strongly) supercooled first-order phase transition



Balan, FK et al., arXiv:2502.19478

First-order phase transition

- Symmetry breaking occurs through bubble nucleation (local transition to ground state)
- Bubbles expand into symmetric phase and collide
- Expect gravitational waves from sound waves, turbulence and bubble collisions
- Typical peak frequency:



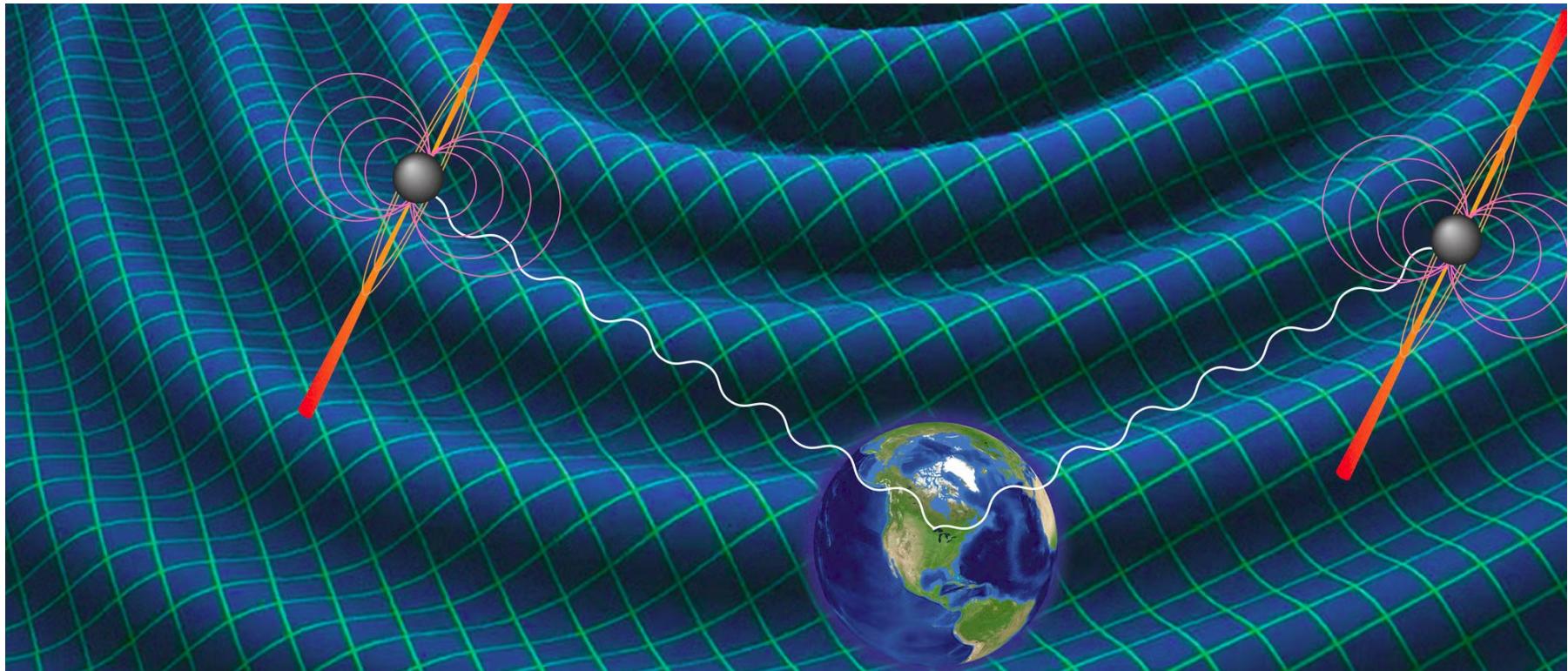
$$f_{\text{peak}} \simeq 10 \text{ mHz} \left(\frac{\beta/H}{100} \right) \left(\frac{T_p}{1 \text{ TeV}} \right)$$

Bringmann, FK et al.,
arXiv:2311.06346

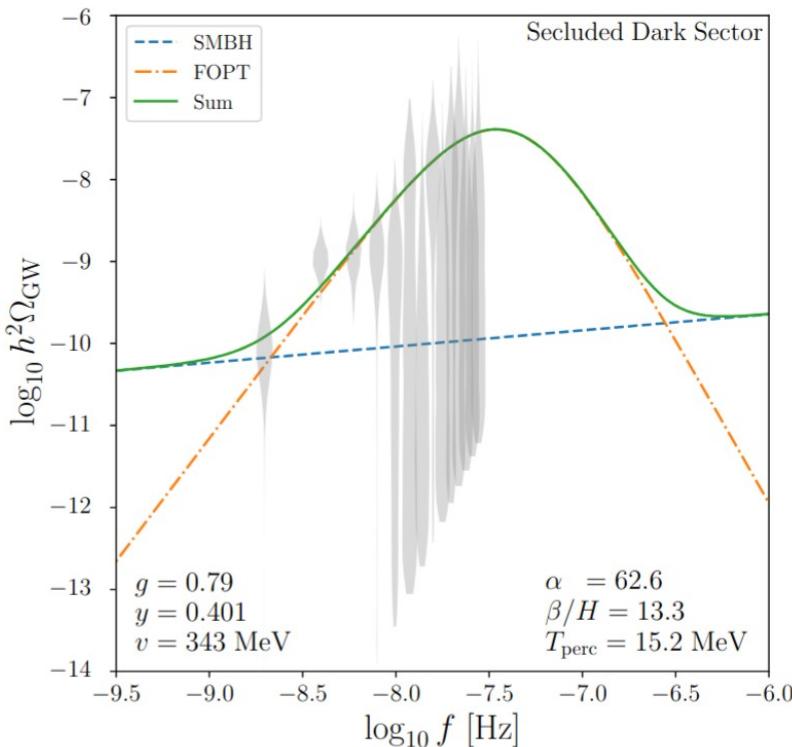
Phase transition temperature

Phase transition speed

Pulsar Timing Arrays (PTAs)



Stochastic gravitational wave background



- Pulsar timing array data strongly suggests presence of a stochastic GW background
- Peak frequency around 10 nHz
- Typical values:
 - $\beta/H \sim 10$
 - $T_p \sim 10 \text{ MeV}$
 - Dark Higgs vev $v \sim 100 \text{ MeV}$

PTAs prefer same mass scale as SIDM!

Balan, FK et al., arXiv:2502.19478

Challenge 1: Portal interactions

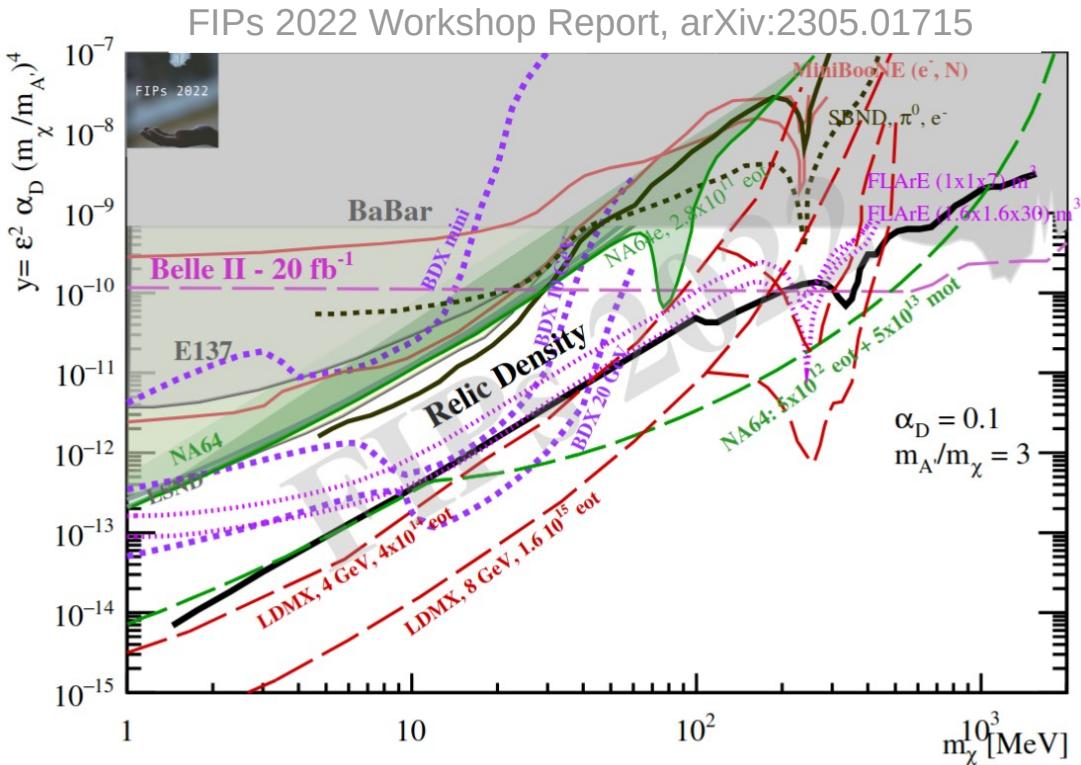
- Phase transition requires dark sector in equilibrium with Standard Model
 - Need “portal” interactions for interactions with SM particles

Option 1: Higgs mixing

- Spoils conformal symmetry after electroweak symmetry breaking

Option 2: Kinetic mixing

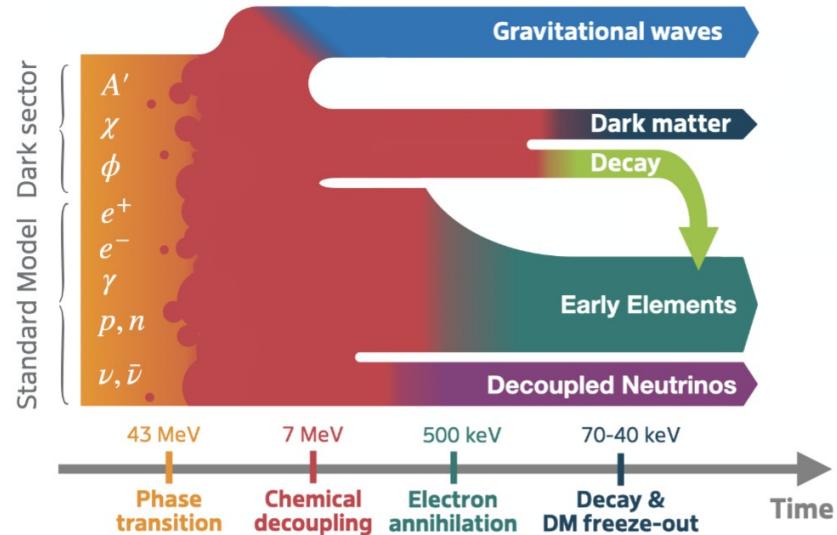
- Subject to strong experimental constraints



Challenge 1: Portal interactions

- Need to transfer energy density stored in the dark sector back into the Standard Model
- Non-negligible abundance of decoupled dark Higgs bosons remains after BBN
- Decay into e^+e^- via dark-photon loops:

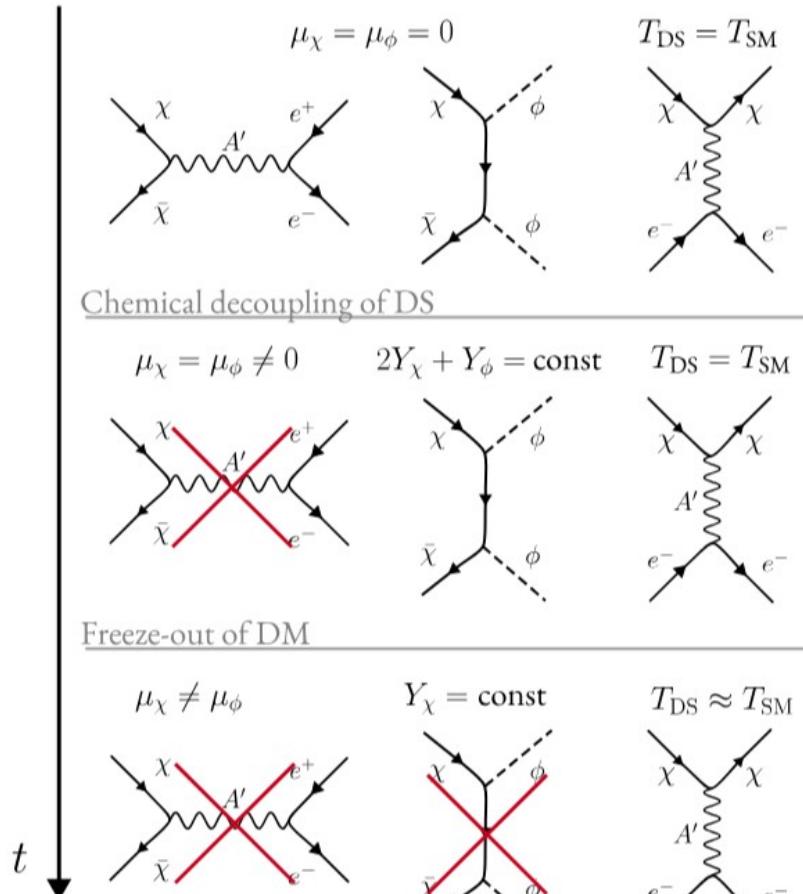
$$\tau_\phi \approx 2500 \text{ s} \left(\frac{\kappa}{10^{-4}} \right)^4 \left(\frac{g}{0.75} \right)^2 \frac{m_\phi}{30 \text{ MeV}} \left(\frac{m_{A'}}{100 \text{ MeV}} \right)^{-2}$$



Balan, FK et al., arXiv:2502.19478

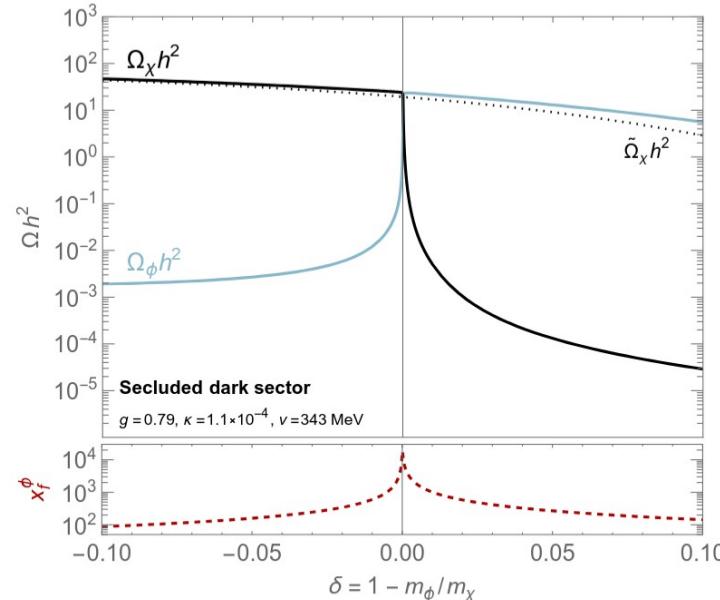
Challenge 2: Dark matter relic density

- Naive estimate of dark matter relic density: $\Omega_{\text{DM}} \simeq 0.1 \frac{10^{-8} \text{ GeV}^{-2}}{\langle \sigma_{\text{ann}} v \rangle}$
- For direct annihilation of DM into dark Higgs bosons: $\langle \sigma_{\text{ann}} v \rangle \sim \frac{y^4}{m_{\text{DM}}^2} \sim \frac{y^2}{v_\phi^2}$
- First-order phase transition requires dark sector couplings of order unity
 - Observed DM relic density implies $v_\phi \sim \text{TeV}$
 - Perfect for LISA, but not for pulsar timing arrays
 - Need to find a way to suppress annihilation cross section



Detailed evolution

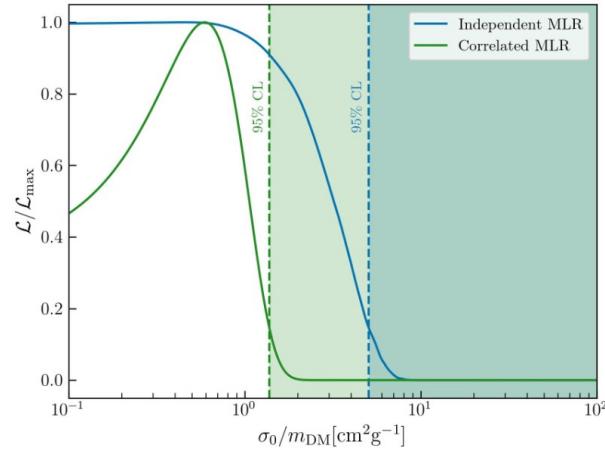
- Reproducing DM relic density requires m_χ and m_ϕ to be very close



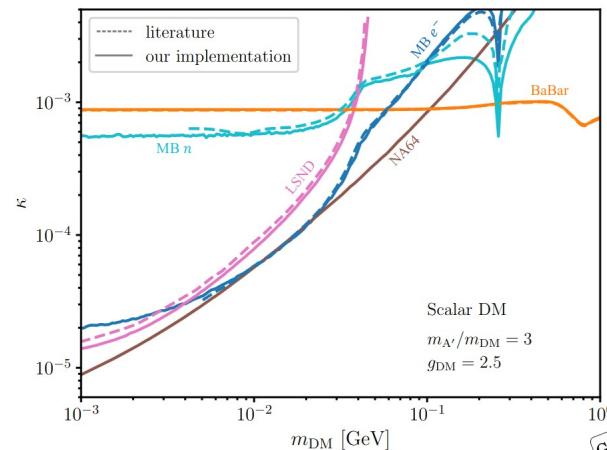
Global fit of a classically conformal dark sector

- Need to take into account many different constraints for global analysis

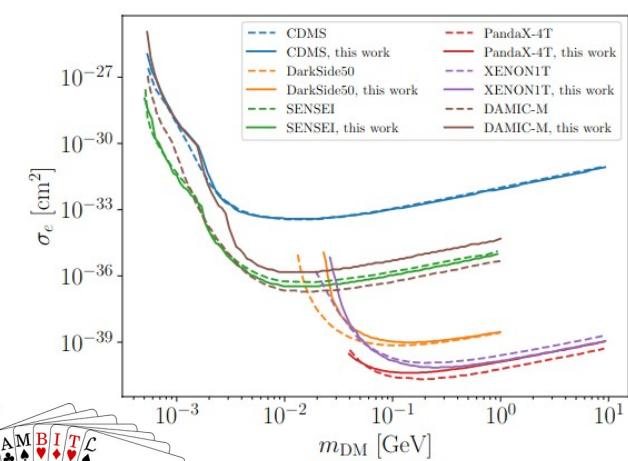
Dark matter self-interactions



Accelerator searches



Scattering off electrons

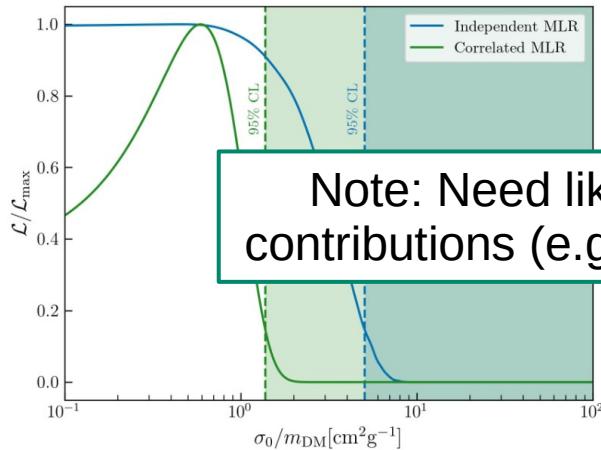


Balan, FK et al., arXiv:2405.17548

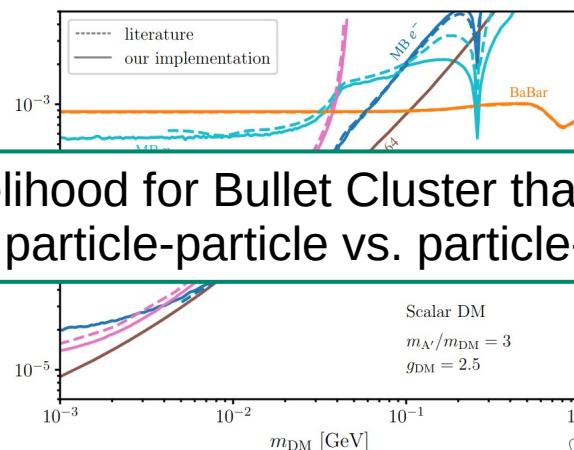
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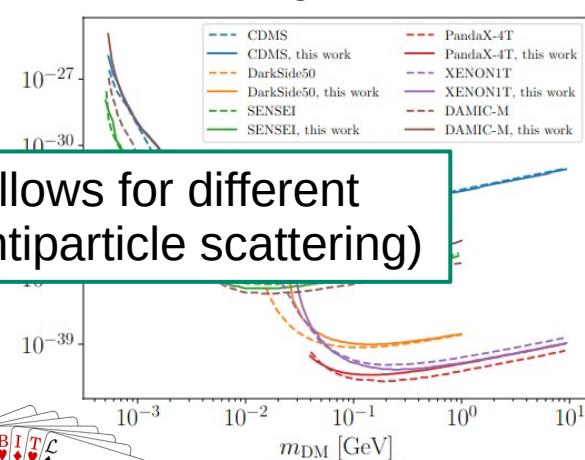
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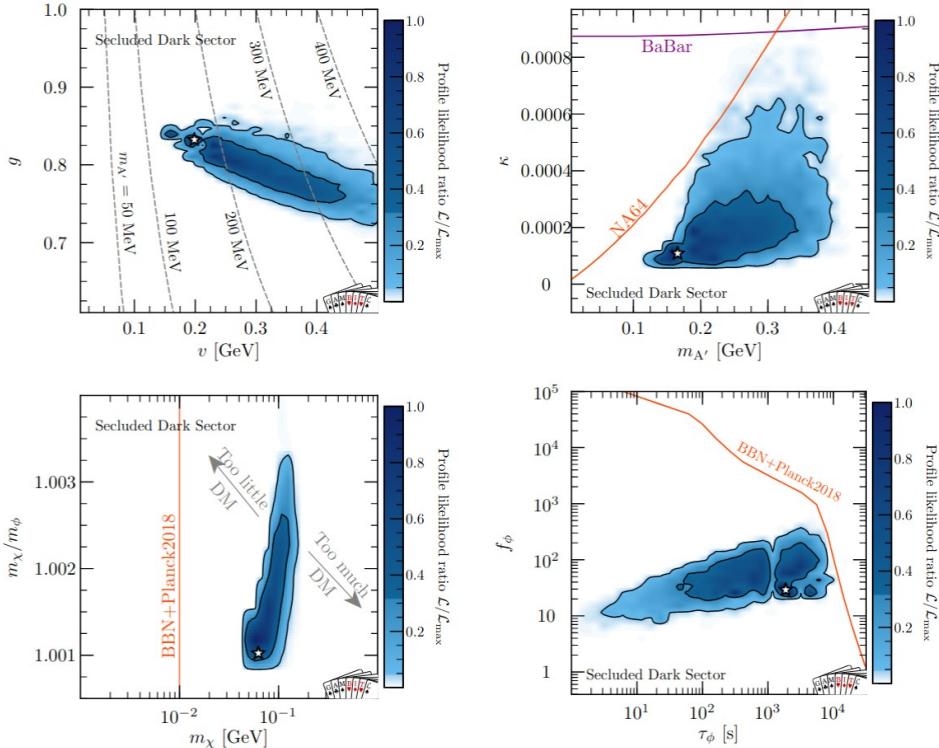


Note: Need likelihood for Bullet Cluster that allows for different contributions (e.g. particle-particle vs. particle-antiparticle scattering)



Balan, FK et al., arXiv:2405.17548

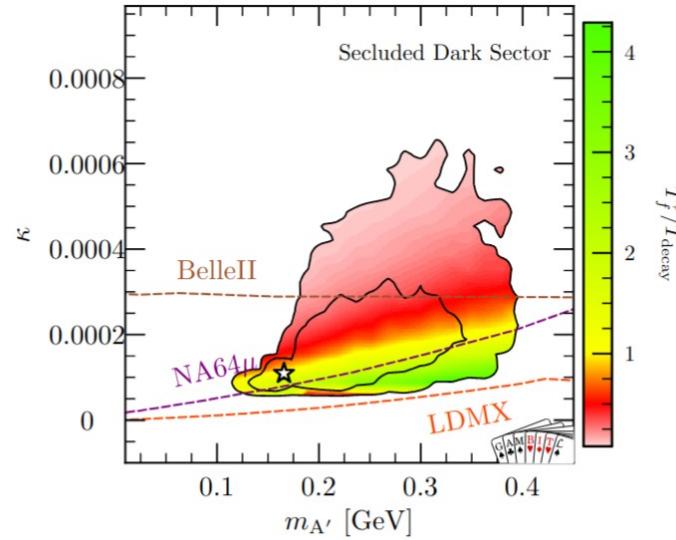
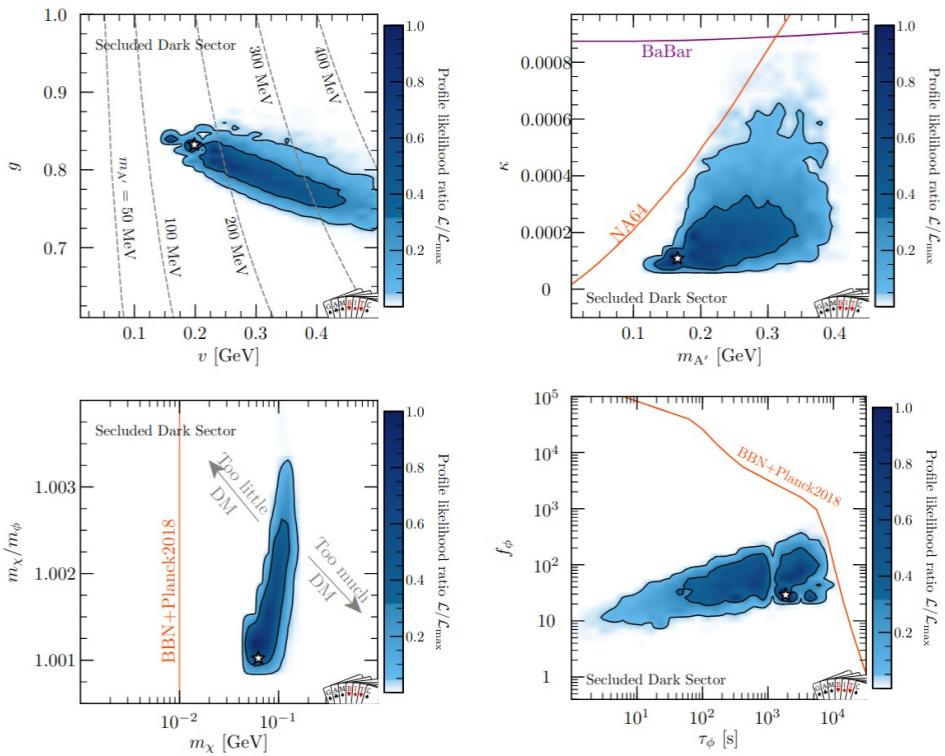
Results



- Possible to fit PTA signal and reproduce observed DM relic abundance while satisfying bounds from dark photon searches
- Abundance and lifetime of dark Higgs bosons small enough to comply with BBN constraints
- Well-defined allowed parameter regions
→ Predictions for future searches

Balan, FK et al., arXiv:2502.19478

Results



Model can be tested with future searches for invisibly decaying dark photons

Balan, FK et al., arXiv:2502.19478

Implications for self-interacting dark matter

- Dark matter self-scattering dominated by dark Higgs boson exchange
- Dark photon too heavy to give relevant contribution
- Typical self-interaction cross sections:
 $\sigma / m_{\text{DM}} \sim 0.1\text{--}1 \text{ cm}^2 / \text{g}$
- No velocity dependence...
- Maybe interesting for future improvements on galaxy cluster scales?

$$\sigma_{\chi\chi \rightarrow \chi\chi} = \frac{y^2}{256\pi v^2} \left(3 - 4 \frac{m_\chi^2}{m_\phi^2} \right)^2$$
$$\sigma_{\chi\bar{\chi} \rightarrow \chi\bar{\chi}} = \frac{y^2}{128\pi v^2} \left(7 - 4 \frac{m_\chi^2}{m_\phi^2} + 16 \frac{m_\chi^4}{m_\phi^4} \right)$$

Conclusions

- Not all models of self-interacting DM actually **predict** observable effects
- Many models are underconstrained and allow for a wide range of cross sections
- Two notable exceptions:

1. Self-interacting dark matter from strong dynamics

- Two independent mass scales, but additional input from relic density requirement
- Preferred parameter region: $m_{\text{DM}} \sim 200 \text{ MeV}$, $\sigma / m_{\text{DM}} \sim 1 \text{ cm}^2 / \text{g}$

2. Self-interacting dark matter from spontaneous breaking of conformal symmetry

- Mass scale generated dynamically through first-order phase transition
- PTA signal prefers $m_{\text{DM}} \sim 100 \text{ MeV}$, $\sigma / m_{\text{DM}} \sim 0.1 \text{ cm}^2 / \text{g}$
- Challenges: Energy transfer and dark matter relic density