

# Probing Gauge-Higgs Unification models at the ILC with $q\bar{q} A_{FB}$ at c.m.e. above the Z mass

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Gen=T



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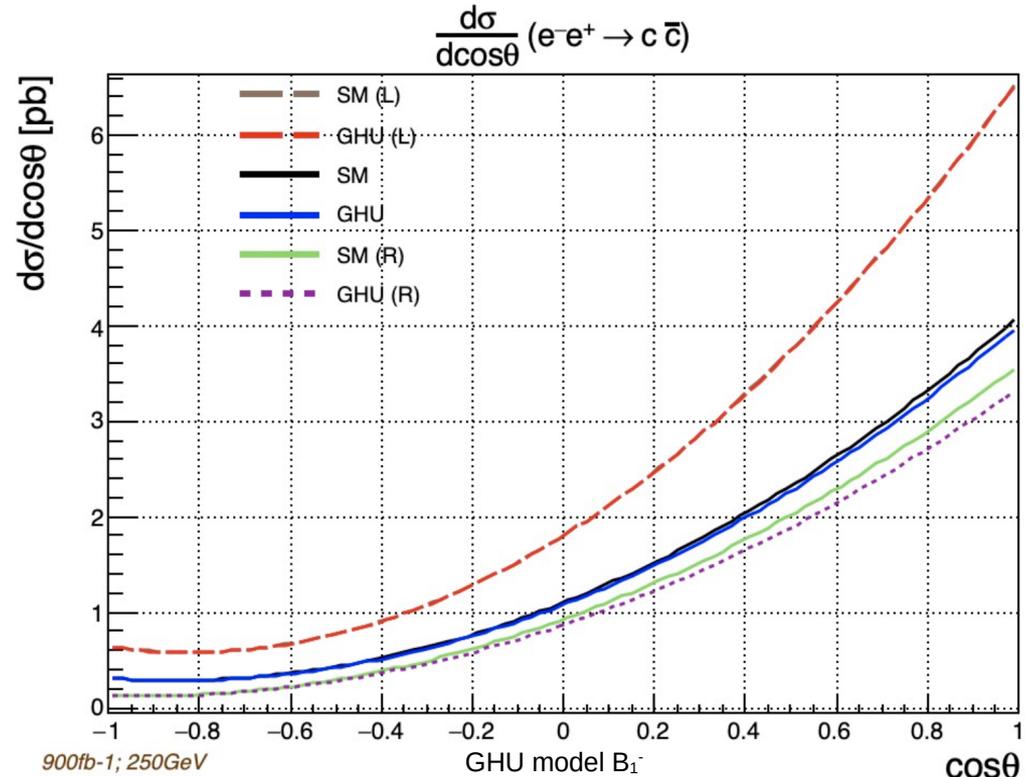
# **Introduction & Motivation**



# Gauge-Higgs Unification Models (GHU)

- ▶ Randall-Sundrum metric (5D).
- ▶ The symmetry breaking pattern is different than in the SM and features the *Hosotani mechanism*:
  - Masses are generated dynamically from the extra-dimension properties.
- ▶ Only one parameter, *Hosotani's angle*  $\theta_H$ , determines the projection of the 5D fields, fixing all physical effects:
  - KK resonances of the Z/ $\gamma$  with  $m_{\text{KK}} \sim 10\text{-}25$  TeV.
  - **Modifications and new EW couplings/helicity amplitudes.**
  - Already visible effects at 250GeV.

As **Benchmark**, we will use the [Funatsu, Hatanaka, Hosotani, Orikasa, Yamatsu] models.



**Probing such scenarios require at least per mil level experimental precision**  
tt/bb/cc/ss/... **Achievable at future colliders?**

▶ A models: ([arxiv:1705.05282](https://arxiv.org/abs/1705.05282))

$$A_1 : \theta_H = 0.0917, m_{KK} = 8.81 \text{ TeV} \rightarrow m_{Z^1} = 7.19 \text{ TeV};$$

$$A_2 : \theta_H = 0.0737, m_{KK} = 10.3 \text{ TeV} \rightarrow m_{Z^1} = 8.52 \text{ TeV},$$

▶ B models: ([arxiv:2309.01132](https://arxiv.org/abs/2309.01132)) ([arxiv:2301.07833](https://arxiv.org/abs/2301.07833))

$$B_1^\pm : \theta_H = 0.10, m_{KK} = 13 \text{ TeV} \rightarrow m_{Z^1} = 10.2 \text{ TeV};$$

$$B_2^\pm : \theta_H = 0.07, m_{KK} = 19 \text{ TeV} \rightarrow m_{Z^1} = 14.9 \text{ TeV};$$

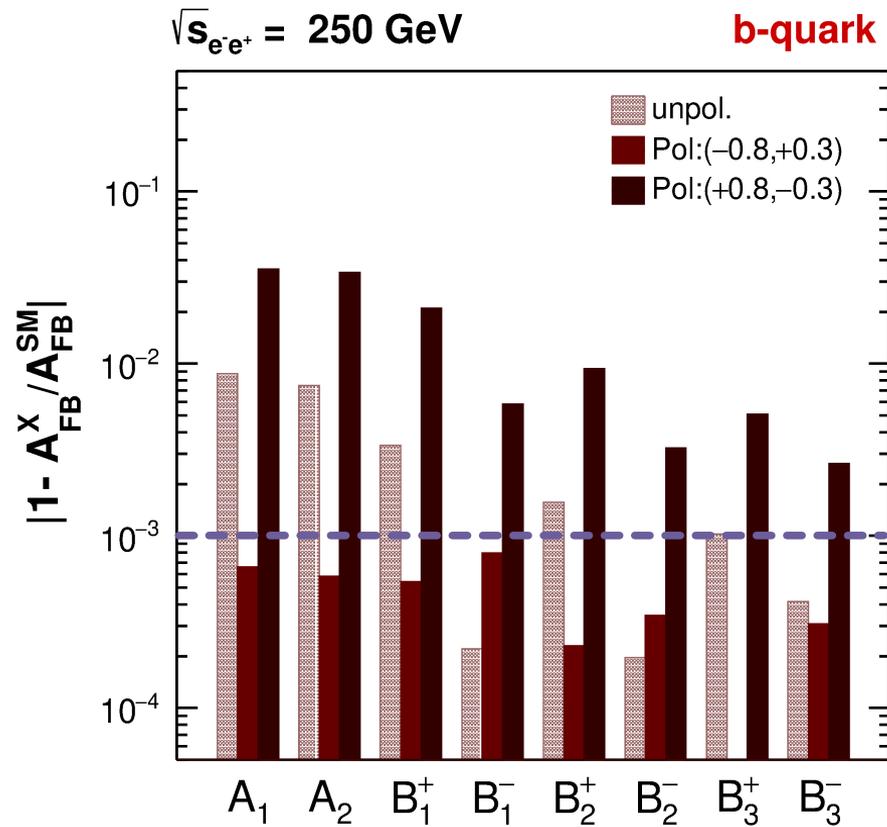
$$B_3^\pm : \theta_H = 0.05, m_{KK} = 25 \text{ TeV} \rightarrow m_{Z^1} = 19.6 \text{ TeV};$$

Resonances of O(10) TeV: Only indirect measurements are possible!

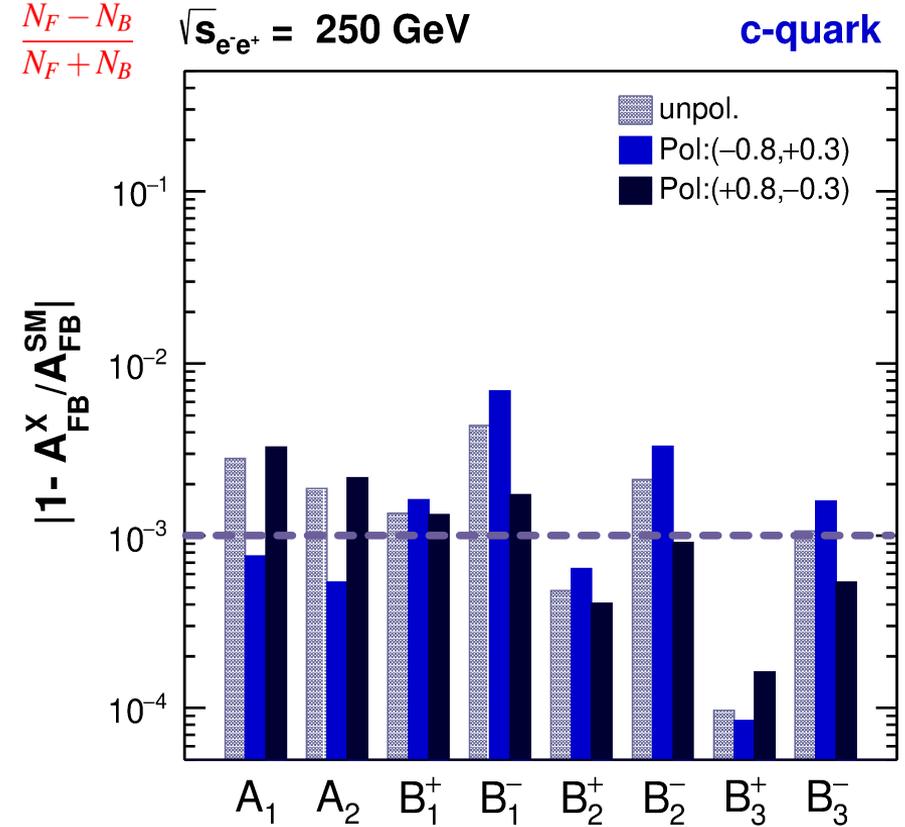
**This talk: Phenomenology at ILC H20-staged program.**

- Runs at 250, 500, 1000 GeV.
- Polarized  $e^-$  and  $e^+$  beams.

# GHU vs SM (250 GeV)



$$A_{FB} = \frac{N_F - N_B}{N_F + N_B}$$



Deviations at the **per mil** level!

# **Experimental study with full simulation**

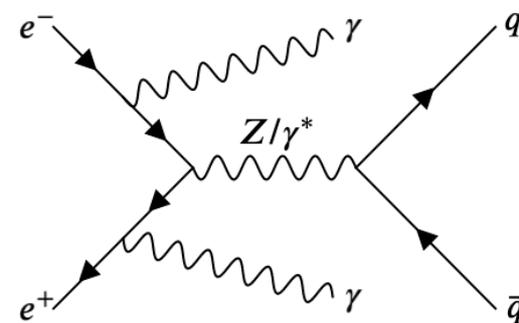


# Heavy flavor production in $e^-e^+$ collisions

► We work with  $A_{FB}$  for b and c quarks.

$$A_{FB} = \frac{N_F - N_B}{N_F + N_B}$$

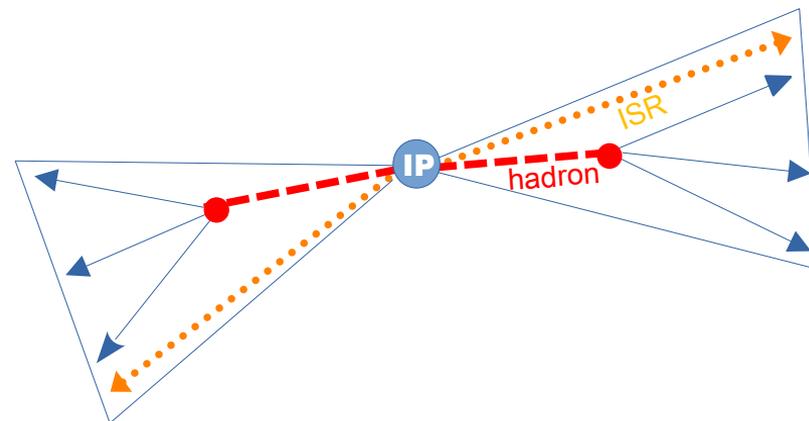
- MC simulations at 250 and 500 GeV.
  - International Linear Collider (ILC) run plan.
- **Full simulation** of the International Large Detector (ILD).



► Topology: Two back-to-back jets.

► Procedure (plots in back-up):

- 1 Background suppression → Selection of  $q\bar{q}$  events.
- 2 Flavor tagging → Selection of  $b\bar{b}$  &  $c\bar{c}$  events.
  - Double tagging.
- 3 Charge measurement → Quark-Antiquark identification.
  - Double charge.



**High-purity & independent samples for each quark flavour.**

# Results for ILC250 & ILC500

►  $A_{FB}$  definition:

$$A_{FB} = \frac{N_F - N_B}{N_F + N_B}$$

► At least 4 observables for  $A_{FB}$  at ILC per energy point

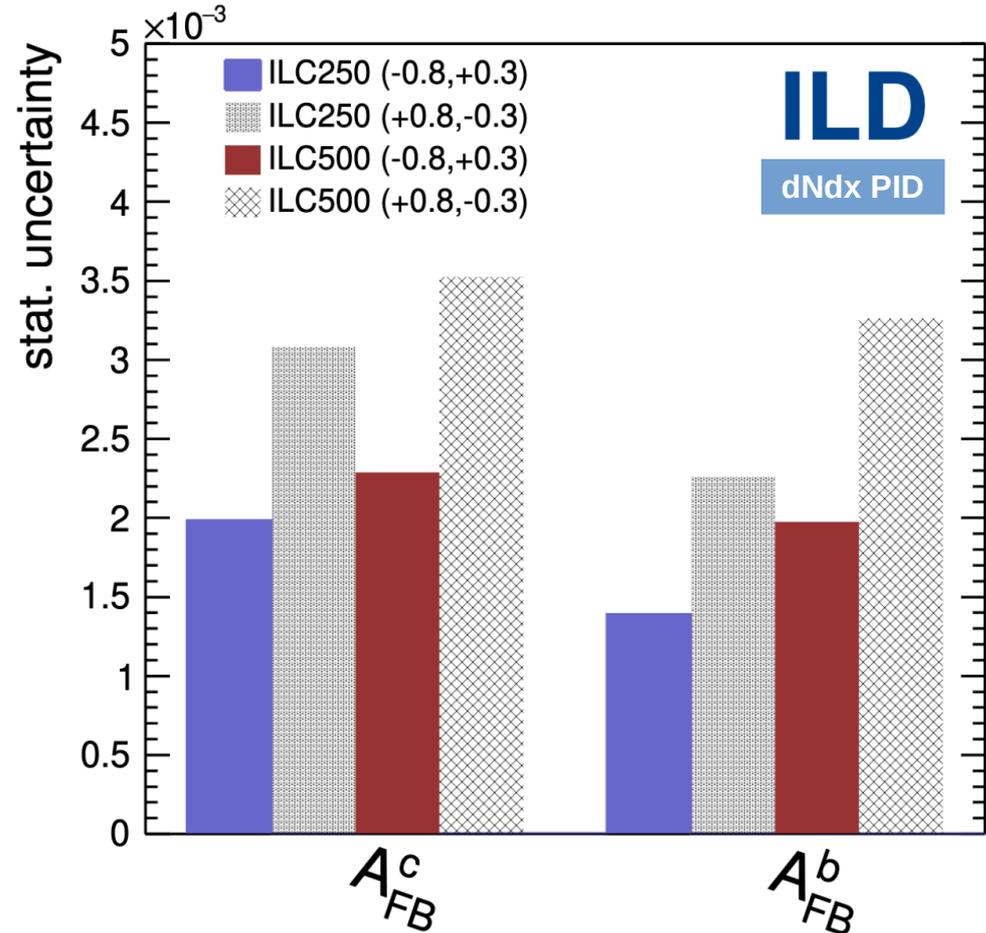
- 2 quarks (b and c).
- 2 polarizations ( $e_{LpR}$ ,  $e_{RpL}$ ).

► **Per mil level statistical uncertainties** reachable for the nominal ILC program

- **Smaller exp. syst. Uncertainties**

► Running at ILC500

- Similar uncertainties but bigger deviations.
- Possibility of combining with the ILC250 results.

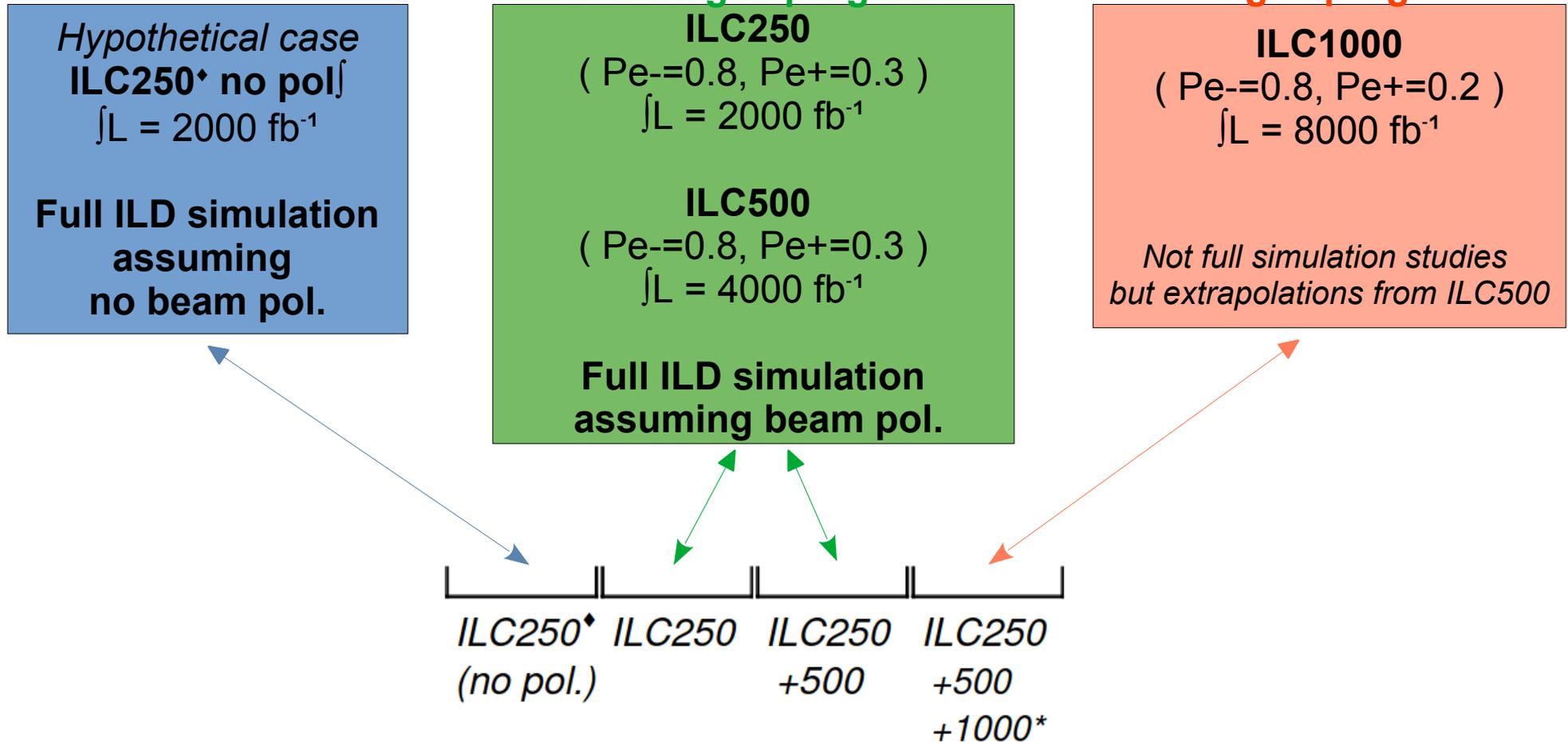


**Discrimination  
power between  
GHU & SM**

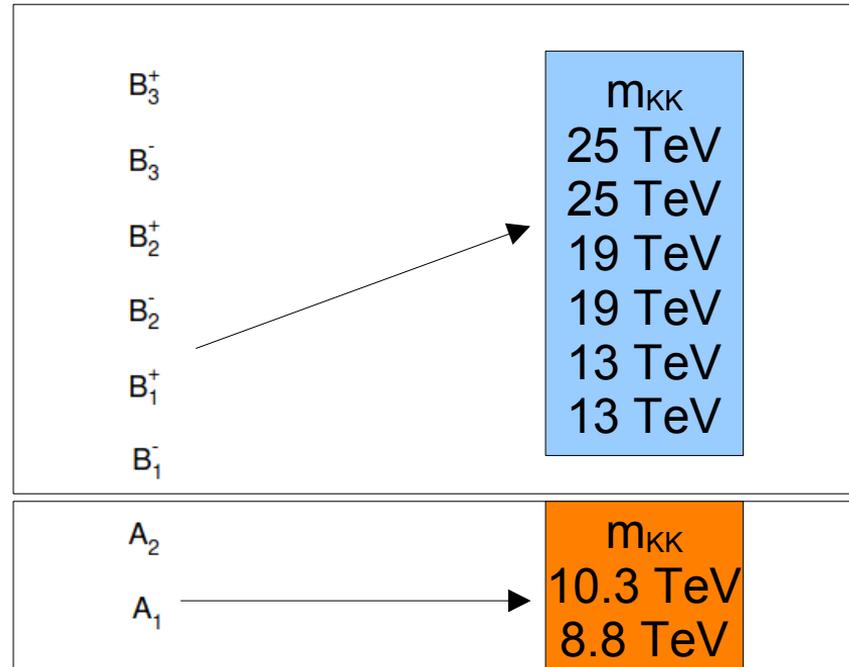


- ▶ Procedure: Testing the statistical significance of model  $AFB_{\text{test}}$  vs a reference model  $AFB_{\text{ref}}$  assuming that one of them is measured.
- ▶ The uncertainties are considered normally distributed:
  - Significance in  $\sigma$ .
  - P-value: Gaussian at  $d_\sigma$ .
$$d_\sigma = \frac{\|AFB_{\text{test}} - AFB_{\text{ref}}\|}{\Delta_{AFB_{\text{ref}}}}$$
- ▶ Combination of multiple measurements is done with a *multivariate gaussian*.
  - Assuming no correlations for AFB.
- ▶ We also assumed different precisions for the SM Z boson couplings:
  - Current precision, ILC250 and Giga-Z (ILC run at the Z-Pole).

# GHU vs SM: Beam scenarios



# GHU vs SM: GHU energy scale



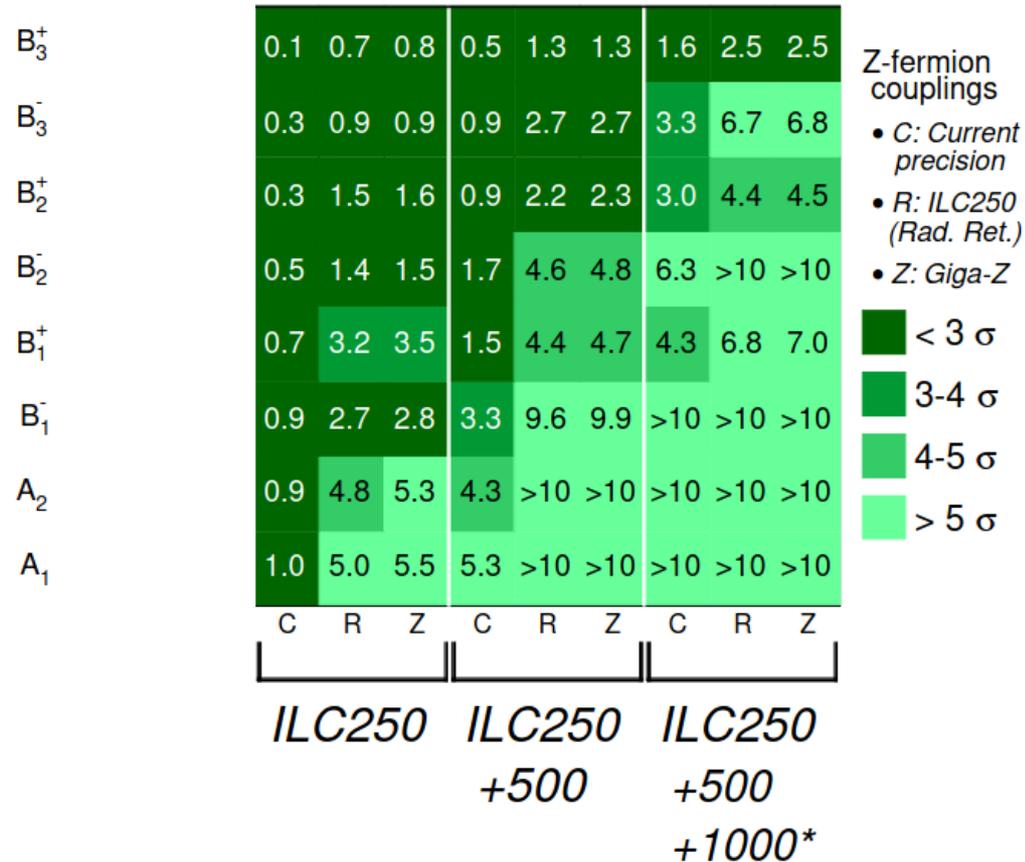
## Similar structure for all plots:

- More massive resonances (harder to detect models) as we move up.
- Higher energy accessed by the ILC runs as we move to the right.

# GHU vs SM: Precision on Z-couplings

GHU vs SM discrimination power ( $\sigma$ -level)

ILD



# GHU vs SM: Beam(s) polarization

GHU vs SM discrimination power ( $\sigma$ -level)

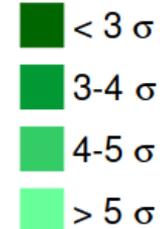
$B_3^+$	0.1	0.4	0.5	0.1	0.7	0.8
$B_3^-$	0.1	0.4	0.5	0.3	0.9	0.9
$B_2^+$	0.2	0.7	0.8	0.3	1.5	1.6
$B_2^-$	0.2	0.7	0.7	0.5	1.4	1.5
$B_1^+$	0.3	1.6	1.7	0.7	3.2	3.5
$B_1^-$	0.5	1.4	1.4	0.9	2.7	2.8
$A_2$	0.6	3.3	3.6	0.9	4.8	5.3
$A_1$	0.8	3.9	4.2	1.0	5.0	5.5
	C	R	Z	C	R	Z

ILC250<sup>♦</sup> ILC250  
(no pol.)

## ILD

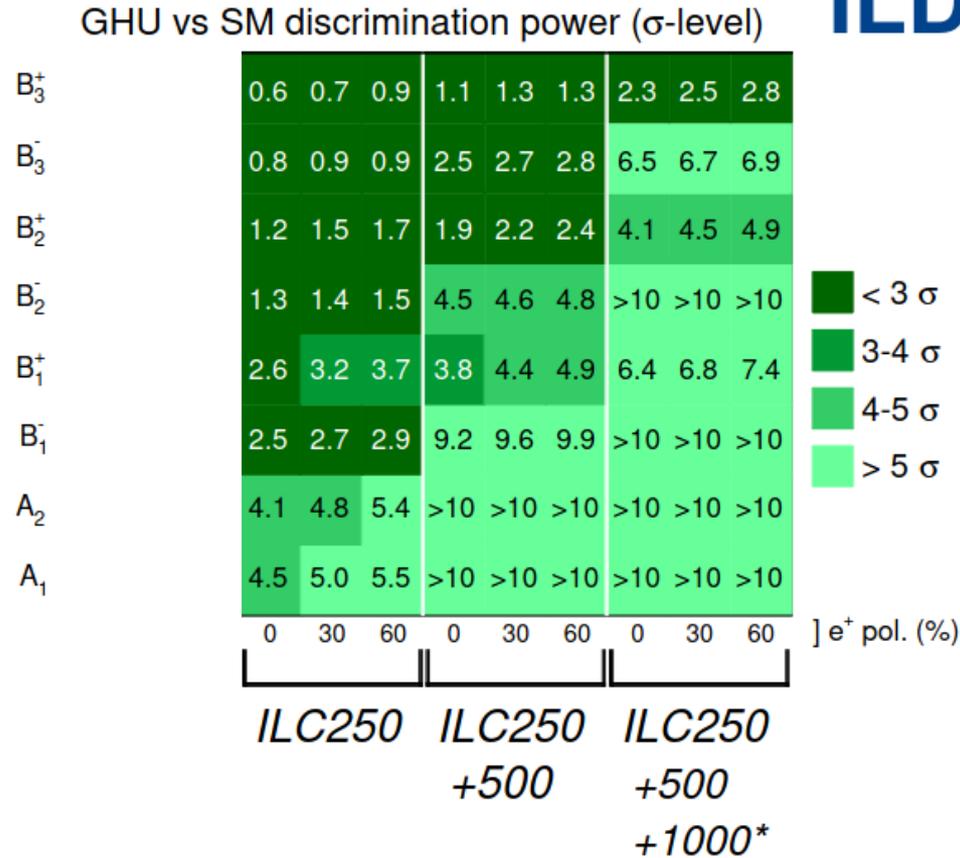
Z-fermion couplings

- C: Current precision
- R: ILC250 (Rad. Ret.)
- Z: Giga-Z



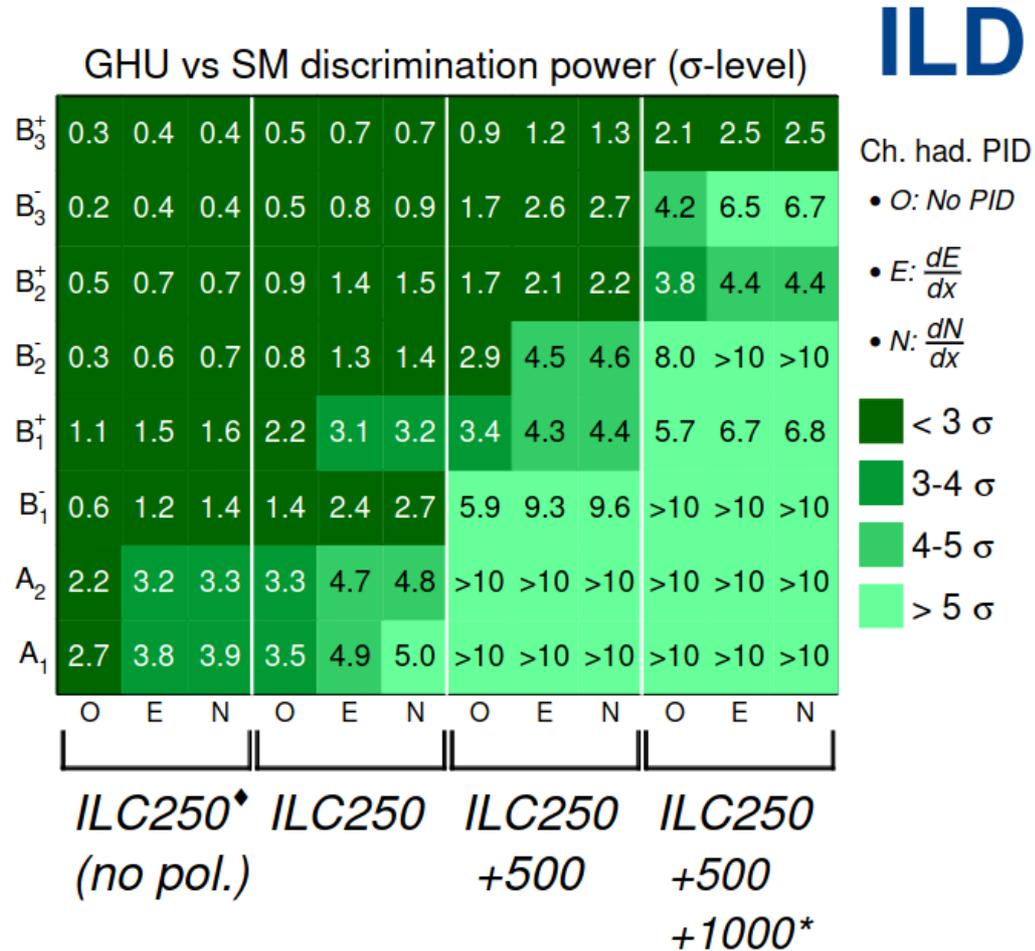
# GHU vs SM: Positron beam polarization

ILD

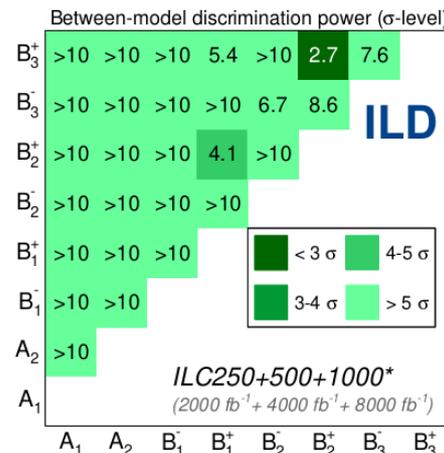
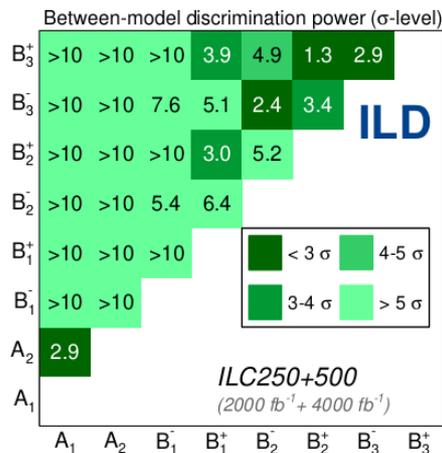
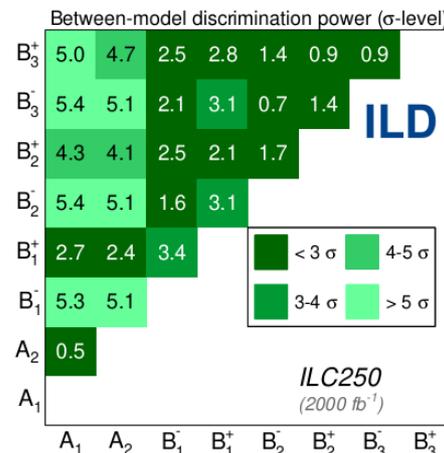
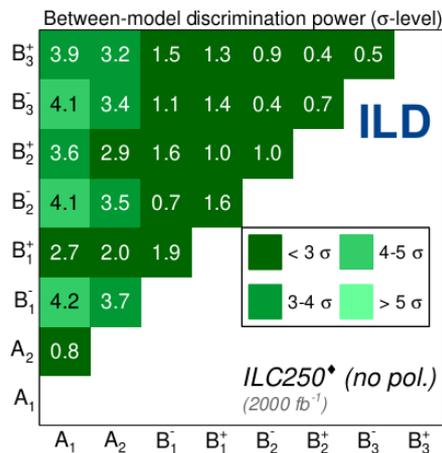


e- polarization at 80% in all of these scenarios

# GHU vs SM: Particle ID dependence



# GHU between model discrimination



# **Conclusion/ summary**



- ▶ Comprehensive study done at ILC250/ILC500 with ILD simulations:
  - Backgrounds, beam features, polarization, realistic reconstruction tools.
  - Uncertainties dominated by statistics, above the Z-pole.
  - Room for improvement (modern algorithms for flavour tagging, event selection, etc.)
- ▶ ILC offers unique capabilities to explore these signatures and discriminate GHU vs SM:
  - **High energy reach.**
  - Electron and positron **beam polarization** → enhancing the sensitivity but also allowing combination of measurements with different BSM sensitivity (for control of systematics).
  - Optimal use of PID via  **$dN/dx$** .
- ▶ **Full discrimination of almost all of the proposed models (and within models) is possible with the H20-staged (baseline) run plan for ILC!**

## ▶ Paper recently published in EPJ-C!

- DOI: [10.1140/epjc/s10052-024-12918-z](https://doi.org/10.1140/epjc/s10052-024-12918-z)

## ▶ Auxiliary and support studies:

- ILD note (2022) [2306.11413](#)
  - ▶ ILC250, b and c studies. (A. Irles, F. Richard, R. Pöschl).
- Proceeding LCWS (2023) [2307.14888](#)
  - ▶ Optimization of flavor tagging, use of dNdx PID and extension to 500 GeV. (J.P. Marquez)
- Proceeding EPS-HEP (2023) [2310.17617](#)
  - ▶ First phenomenology prospects. (J.P. Marquez)

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<https://doi.org/10.1140/epjc/s10052-024-12918-z>

THE EUROPEAN  
PHYSICAL JOURNAL C



Regular Article - Experimental Physics

## Probing gauge-Higgs unification models at the ILC with quark-antiquark forward-backward asymmetry at center-of-mass energies above the Z mass

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**Abstract** The International Linear Collider (ILC) will allow the precise study of  $e^-e^+ \rightarrow q\bar{q}$  interactions at different center-of-mass energies from the Z-pole to 1 TeV. In this paper, we discuss the experimental prospects for measuring differential observables in  $e^-e^+ \rightarrow b\bar{b}$  and  $e^-e^+ \rightarrow c\bar{c}$  at the ILC baseline energies, 250 and 500 GeV. The study is based on full simulation and reconstruction of the International Large Detector (ILD) concept. Two gauge-Higgs unification models predicting new high-mass resonances beyond the Standard Model are discussed. These models predict sizeable deviations of the forward-backward observables at the ILC running above the Z mass and with longitudinally polarized electron and positron beams. The ability of the ILC to probe these models via high-precision measurements of the forward-backward asymmetry is discussed. Alternative scenarios at other energies and beam polarization schemes are also discussed, extrapolating the estimated uncertainties from the two baseline scenarios.

### 1 Introduction

The Standard Model (SM) is a successful theory, well-established experimentally and theoretically. With the discovery of the Higgs boson [1, 2], the structure of the SM seems to be confirmed. However, the SM cannot explain many of its seemingly arbitrary features. An example is the striking mass hierarchy in the fermion sector. Moreover, while the dynamics of the SM gauge bosons, the photon,

W and Z bosons, and gluons are governed by the gauge principle, the dynamics of the Higgs boson are different and unique in the SM. The SM does not predict the strength of the Higgs couplings of quarks and leptons, nor the Higgs self-couplings. Large quantum corrections must be canceled by fine-tuning the parameters to match the measured Higgs boson mass. One possible solution to this issue, achieving stabilization of the Higgs mass against quantum corrections, appears when the Higgs boson is associated with the zeroth mode of a dimension-five component of extensions of the SM gauge group. These models are referred to as gauge-Higgs unification (GHU) models.

The two most precise determinations of  $\sin^2 \theta_{\text{eff}}$  by the LEP and SLC differ by 3.7 standard deviations, and neither agrees with the SM prediction [3, 4]. In particular, the LEP value was extracted from the forward-backward asymmetry measurement for b-quarks in LEP1 data, and is nearly three standard deviations away from the value predicted by the SM. Clarifying this anomaly and exploring the possibility of BSM physics motivates the study of quark pair production in high energy  $e^-e^+$  collisions at future colliders both at the Z boson mass and higher energies. In the SM, these interactions are mediated by the photon, Z boson, and their interference. Some BSM theories predict deviations of these bosons' couplings or even sizable new contributions to these processes from new mediators (such as heavy Z' resonances). These deviations would be accessible experimentally by performing high precision measurements of  $e^-e^+ \rightarrow q\bar{q}$  observables at different center-of-mass energies ( $\sqrt{s}$ ). The work presented here is based on the study of such processes at the ILC.

In parallel to the exploitation of data from the Large Hadron Collider (LHC), the high-energy accelerator-based

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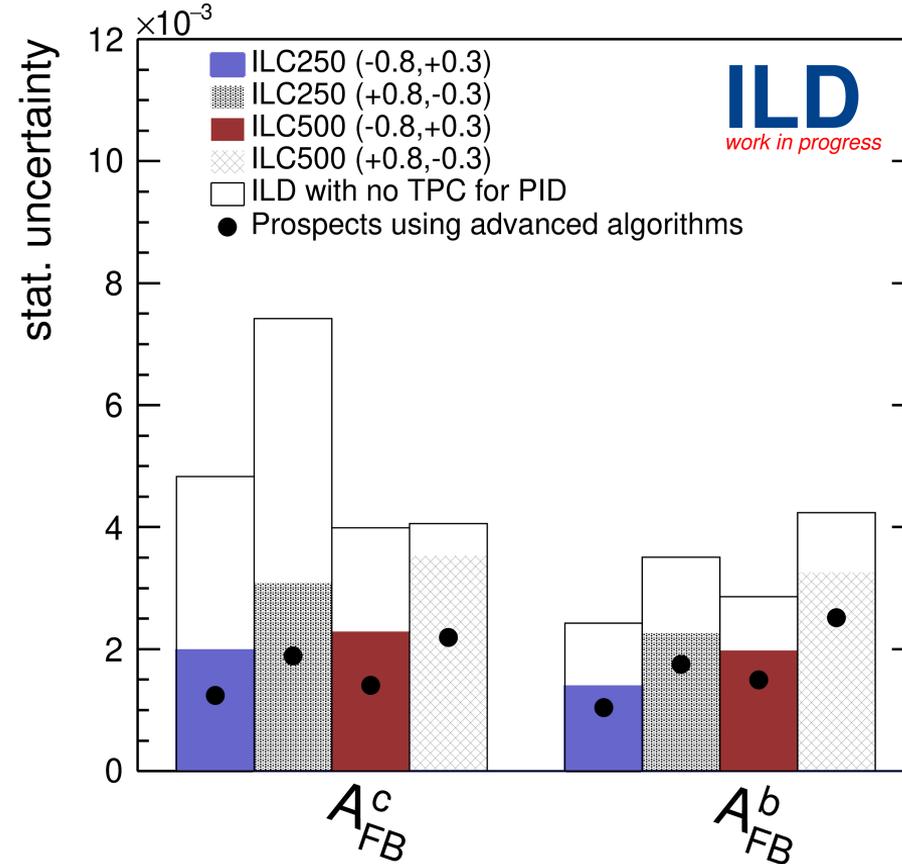
**Thanks for your attention!**

**back-up**



# Prospects for new flavor tagging algorithms

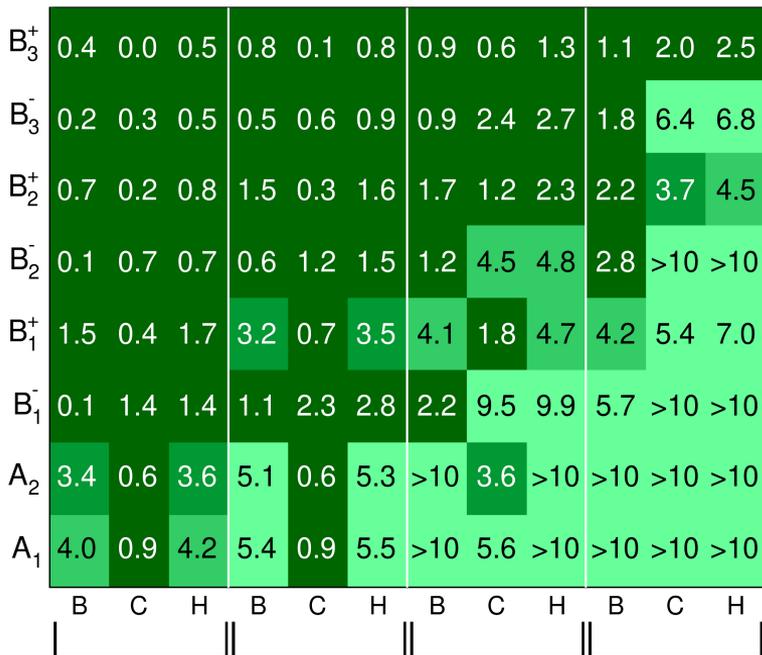
► Expected precisions when using a ParticleTransformer (GNN) algorithm for flavor tagging.



# Prospects for new flavor tagging algorithms

[Current ILD but exploiting dNdx for PID]

GHU vs SM discrimination power ( $\sigma$ -level)



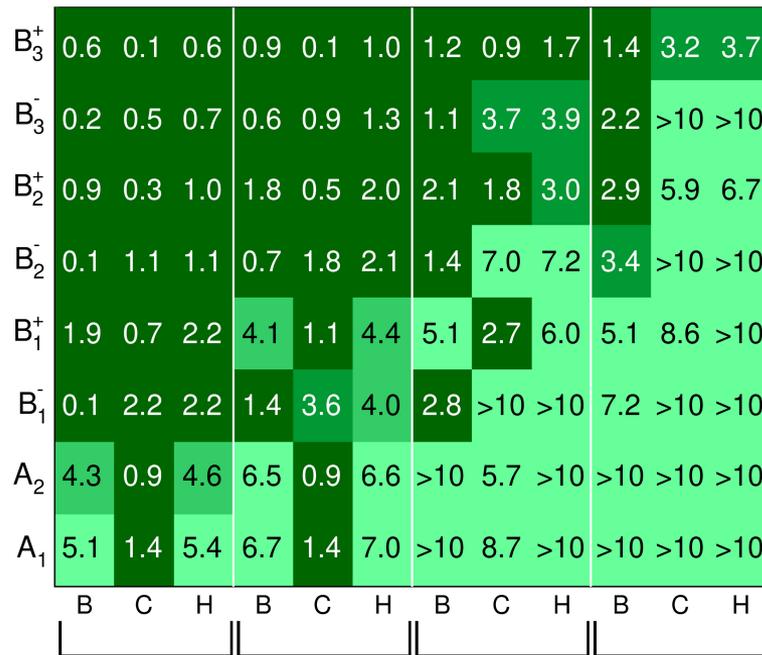
ILC250<sup>♦</sup> (no pol.)  
ILC250  
ILC250 +500  
ILC250 +500 +1000\*

ILD

Quark flavor  
• B: b-quark  
• C: c-quark  
• H: b & c  
■ < 3  $\sigma$   
■ 3-4  $\sigma$   
■ 4-5  $\sigma$   
■ > 5  $\sigma$

[Current ILD but exploiting ParticleTransformer algorithms]

GHU vs SM discrimination power ( $\sigma$ -level)



ILC250<sup>♦</sup> (no pol.)  
ILC250  
ILC250 +500  
ILC250 +500 +1000\*

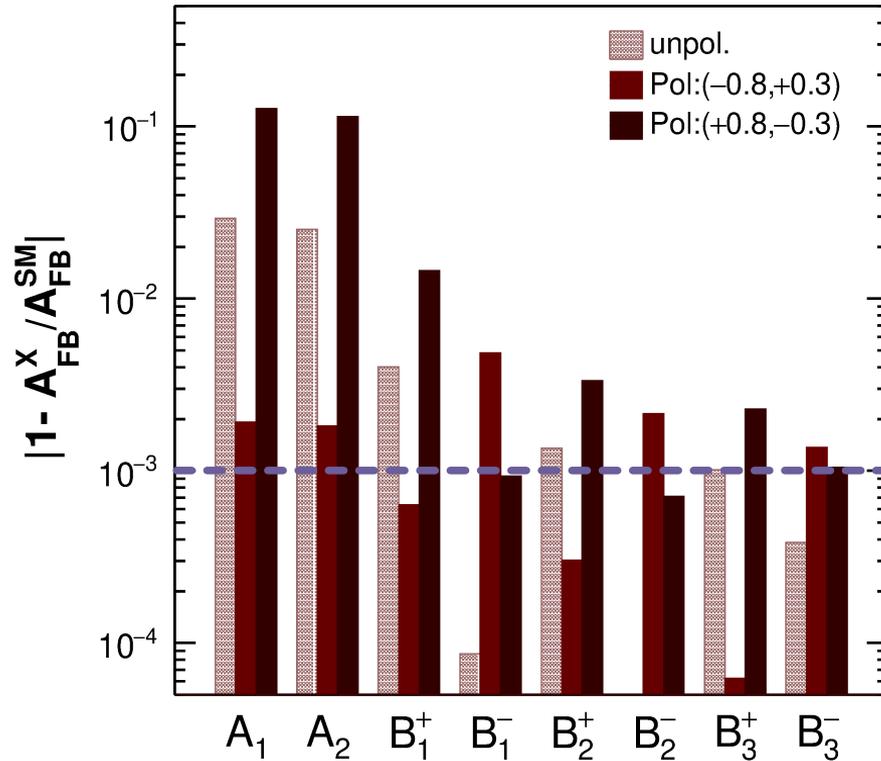
ILD

Work in progress  
Quark flavor  
• B: b-quark  
• C: c-quark  
• H: b & c  
■ < 3  $\sigma$   
■ 3-4  $\sigma$   
■ 4-5  $\sigma$   
■ > 5  $\sigma$

# GHU vs SM (500 GeV)

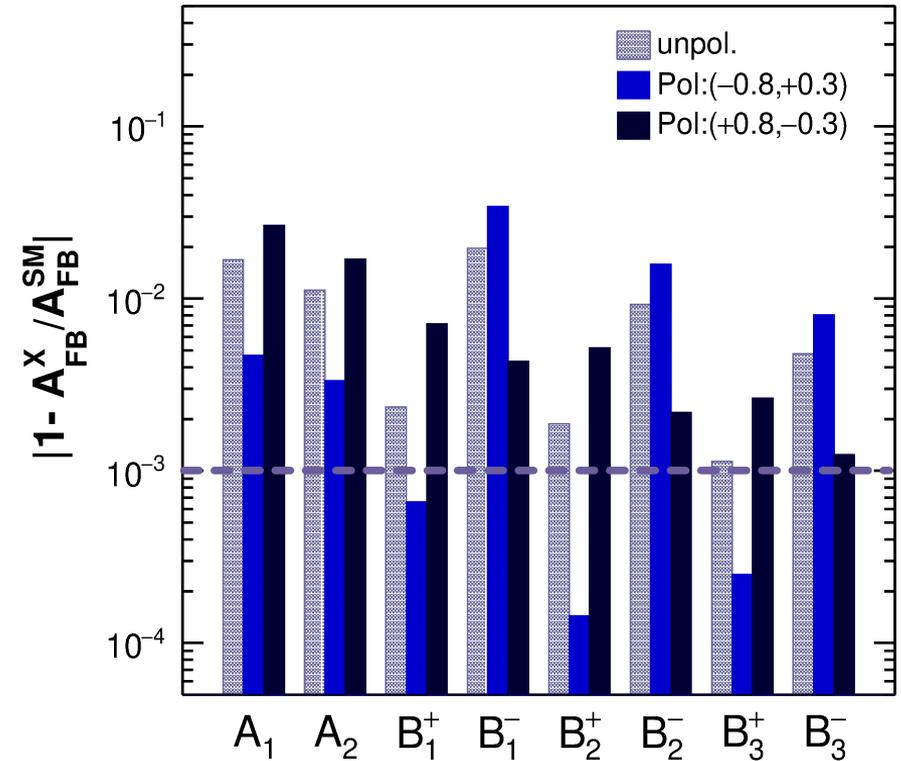
$\sqrt{s_{e^+e^-}} = 500 \text{ GeV}$

**b-quark**



$\sqrt{s_{e^+e^-}} = 500 \text{ GeV}$

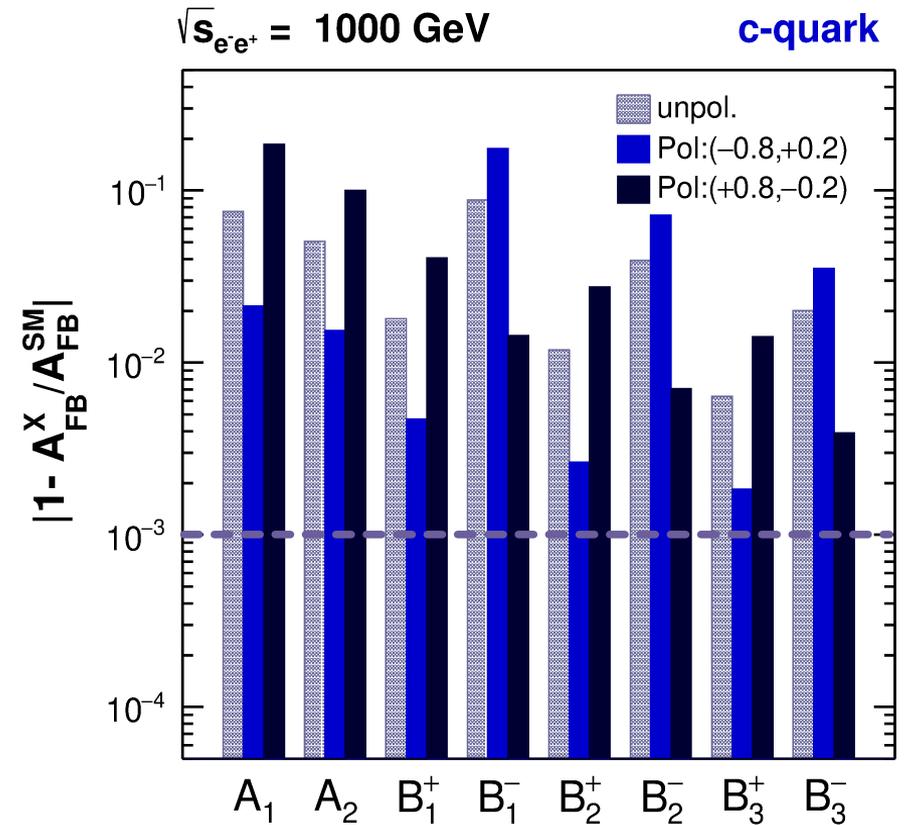
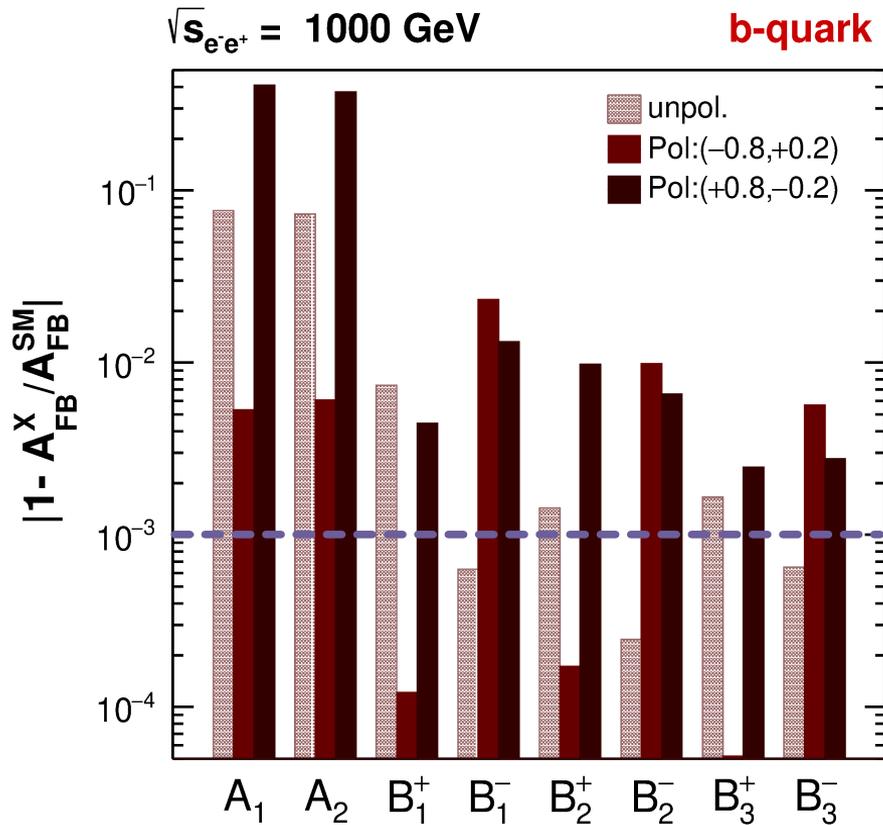
**c-quark**



$$A_{FB} = \frac{N_F - N_B}{N_F + N_B}$$

Deviations at the **per mil** level

# GHU vs SM (1 TeV)



Deviations at the per mil level

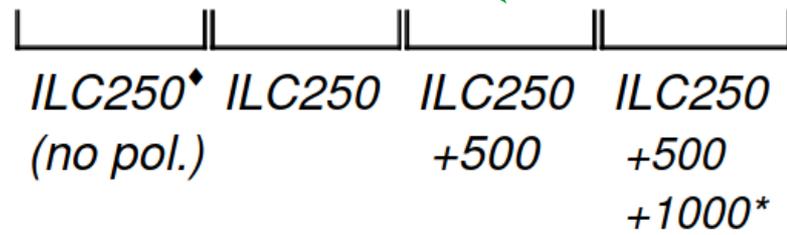
$$A_{FB} = \frac{N_F - N_B}{N_F + N_B}$$

# GHU vs SM: Beam scenarios

*Hypothetical case*  
**ILC250\* no pol**  
 $\int L = 2000\text{fb}^{-1}$   
OSP|SSP [%] = 45 | 5  
**Full ILD simulation**  
**assuming**  
**no beam pol**

**H20 nominal program**  
**ILC250**  
( $P_{e-}=0.8, P_{e+}=0.3$ )  
 $\int L = 2000\text{fb}^{-1}$   
OSP|SSP [%] = 45 | 5  
**ILC500**  
( $P_{e-}=0.8, P_{e+}=0.3$ )  
 $\int L = 4000\text{fb}^{-1}$   
OSP|SSP [%] = 40 | 10  
**Full ILD simulation**  
**assuming beam pol.**

**H20 nominal program**  
**ILC1000**  
( $P_{e-}=0.8, P_{e+}=0.2$ )  
 $\int L = 8000\text{fb}^{-1}$   
OSP|SSP [%] = 40 | 10  
  
*Not full simulation studies*  
*but extrapolations from ILC500*



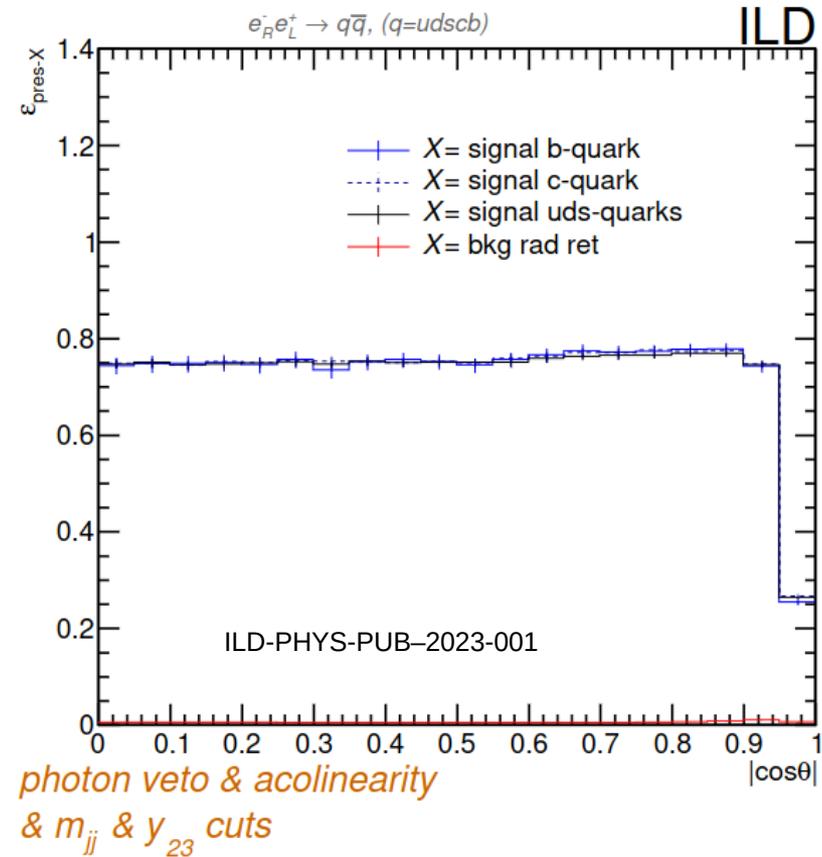
► <https://arxiv.org/pdf/2203.07622.pdf>

Quantity	Value	current $\delta[10^{-4}]$	Z pole		ILC250	
			$\delta_{stat}[10^{-4}]$	$\delta_{sys}[10^{-4}]$	$\delta_{stat}[10^{-4}]$	$\delta_{sys}[10^{-4}]$
boson properties						
$m_W$	80.379	1.5	-	-	0.08	0.3
$m_Z$	91.1876	0.23		0.022		-
$\Gamma_Z$	2.4952	9.4	0.5	-	6	-
$\Gamma_Z(had)$	1.7444	11.5		4.	-	-
Z-e couplings						
$1/R_e$	0.0482	24.	2.	5	5.5	10
$A_e$	0.1513	139.	1.5	1.2	12.	9.
$g_L^e$	-0.632	16.	1.0	3.2	2.8	7.6
$g_R^e$	0.551	18.	1.0	3.2	2.9	7.6
Z- $\ell$ couplings						
$1/R_\mu$	0.0482	16.	2.	2.	5.5	10
$1/R_\tau$	0.0482	22.	2.	2.	5.7	10
$A_\mu$	0.1515	991.	2.	5	54.	3.
$A_\tau$	0.1515	271.	2.	5.	57.	3
$g_L^\mu$	-0.632	66.	1.0	2.3	4.5	7.6
$g_R^\mu$	0.551	89.	1.0	2.3	5.5	7.6
$g_L^\tau$	-0.632	22.	1.0	2.8	4.7	7.6
$g_R^\tau$	0.551	27.	1.0	3.2	5.8	7.6
Z-b couplings						
$R_b$	0.2163	31.	0.4	7.	3.5	10
$A_b$	0.935	214.	1.	5.	5.7	3
$g_L^b$	-0.999	54.	0.32	4.2	2.2	7.6
$g_R^b$	0.184	1540	7.2	36.	41.	23.
Z-c couplings						
$R_c$	0.1721	174.	2.	30	5.8	50
$A_c$	0.668	404.	3.	5	21.	3
$g_L^c$	0.816	119.	1.2	15.	5.1	26.
$g_R^c$	-0.367	416.	3.1	17.	21.	26.

**back-up**  
**( $q\bar{q}$  analysis)**



- ▶ **Topology: 2 back-to-back jets (pencil-like topology)**
- ▶ **Preselection aiming for high background rejection and high efficiency.**
- ▶ **Main bkg  $ee \rightarrow Z\gamma$  (radiative return through ISR)**
  - $\sim x10$  larger than signal
  - **$\sim 90\%$  of such ISR photons are lost in the beam pipe**  $\rightarrow$  events filtered by energy & angular mom. conservation arguments
  - The **remaining  $\sim 10\%$  are filtered by identifying photons** in the detector (efficiency of  $>90\%$ )
  - PFA detector!!
- ▶ Other backgrounds from diboson production decaying hadronically are removed with extra topological cuts.



- Compare samples with 1 tag vs 2 tags (after preselection)

$$f_{1b} = \overline{\varepsilon_c} \overline{R_b} + \overline{\tilde{\varepsilon}_c} \overline{R_c} + \overline{\tilde{\varepsilon}_{uds}} (1 - \overline{R_b} - \overline{R_c})$$
$$f_{2b} = \overline{\varepsilon_b^2} (1 + \overline{\rho}) \overline{R_b} + \overline{\varepsilon_c^2} \overline{R_c} + \overline{\varepsilon_{uds}^2} (1 - \overline{R_b} - \overline{R_c})$$

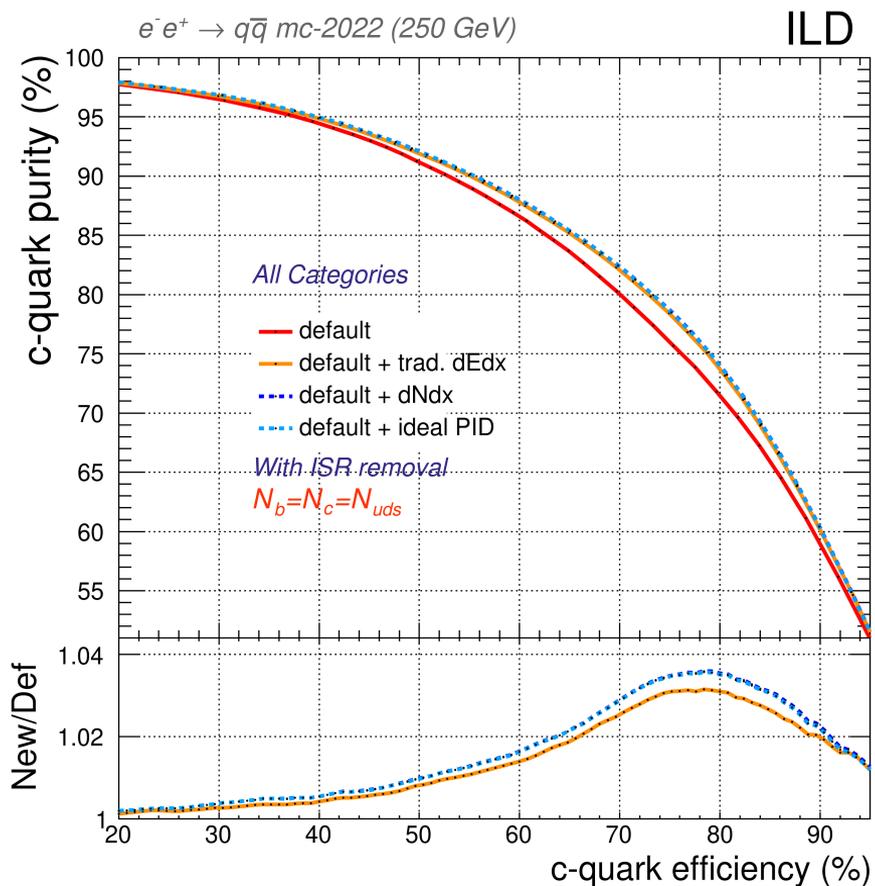
Measured observables

PHYSICS!  
Indirect observables

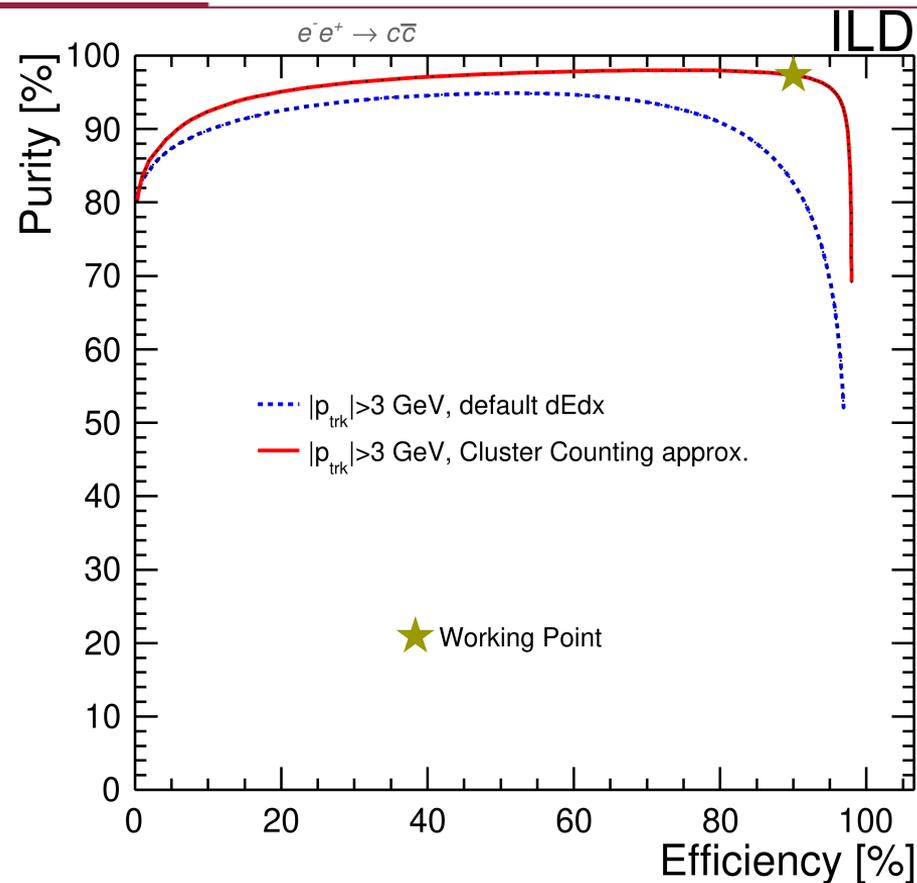
Inputs (MC or independent measurements)

**Similar set of equations  
for the c-quark  
solved simultaneously**

# PID: From dEdx to dNdx

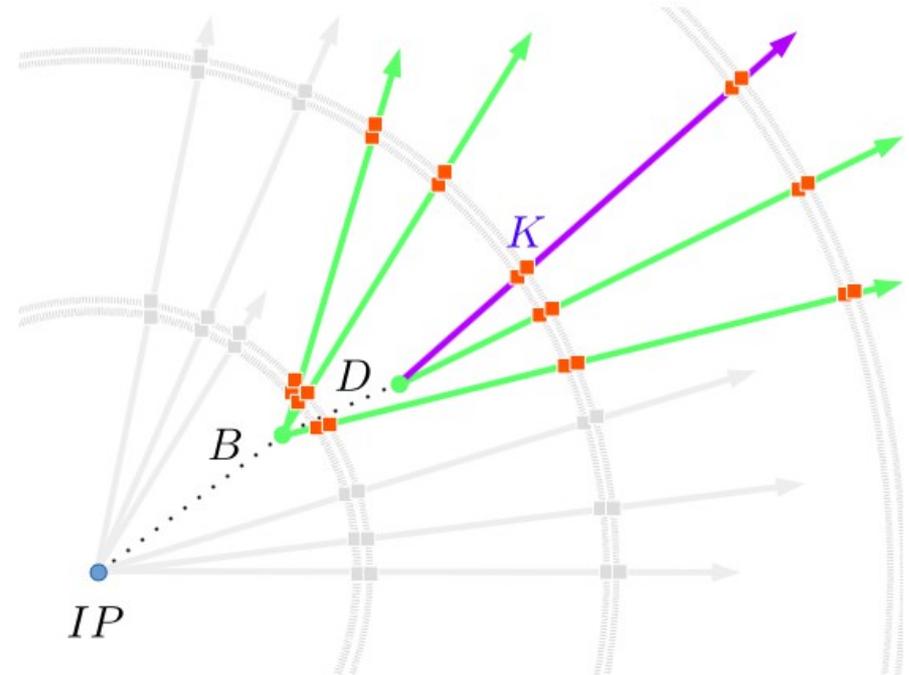


Effects in **Flavour Tagging**



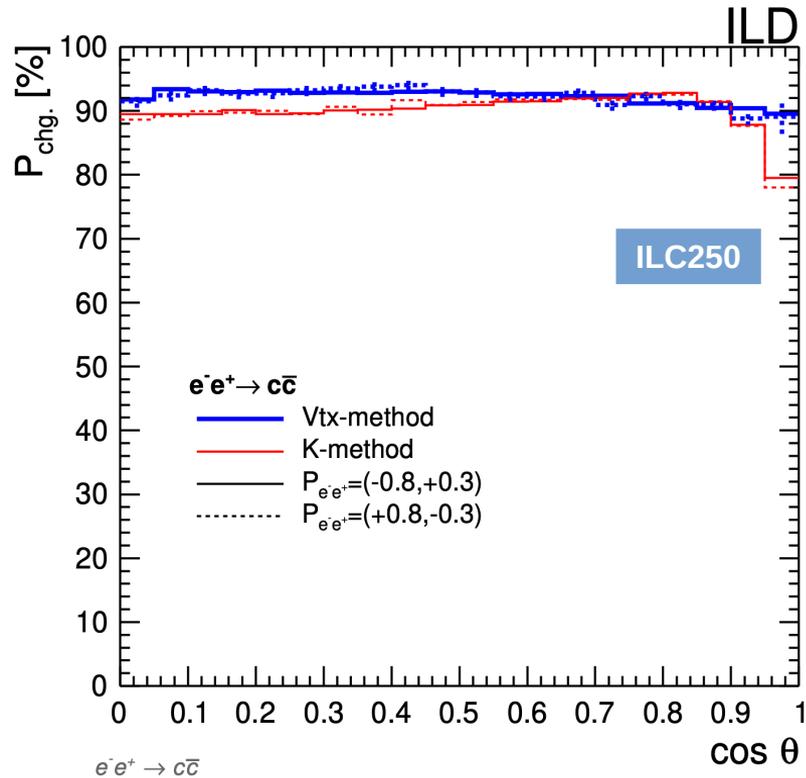
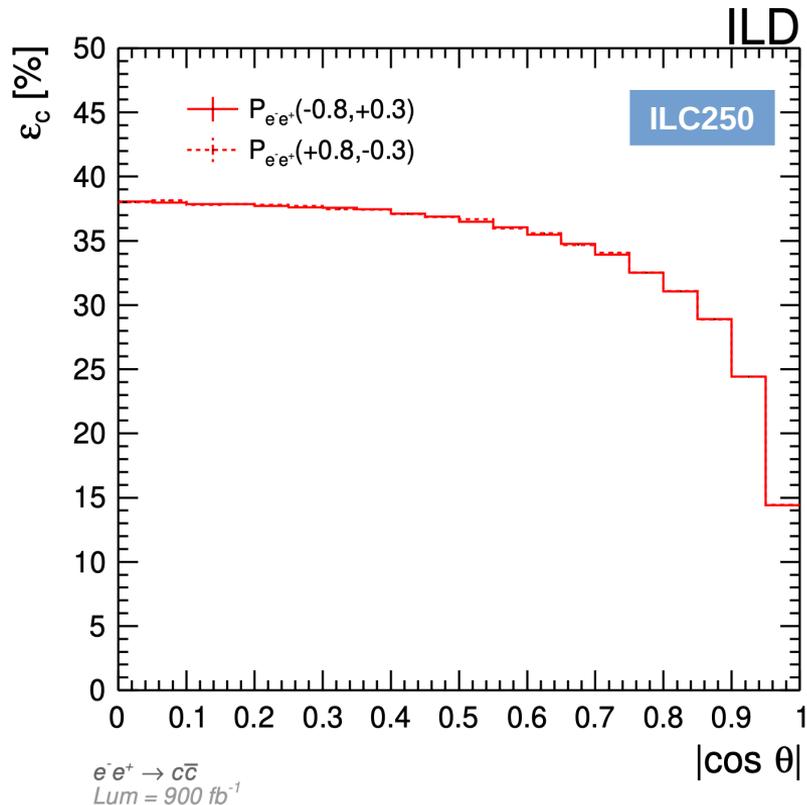
Effects in **Kaon ID** for charge reco.

- ▶ We start from a very pure & background-free **double tagged** sample
- ▶ We are required to **measure the jet charge**
  - Using K-ID and/or full Vtx charge measurement
  - K-ID is better suited for the C-quark (Vtx is better suited for b-quark)
- ▶ We use the **double charge** measurements
  - To control / reduce the systematic uncertainties



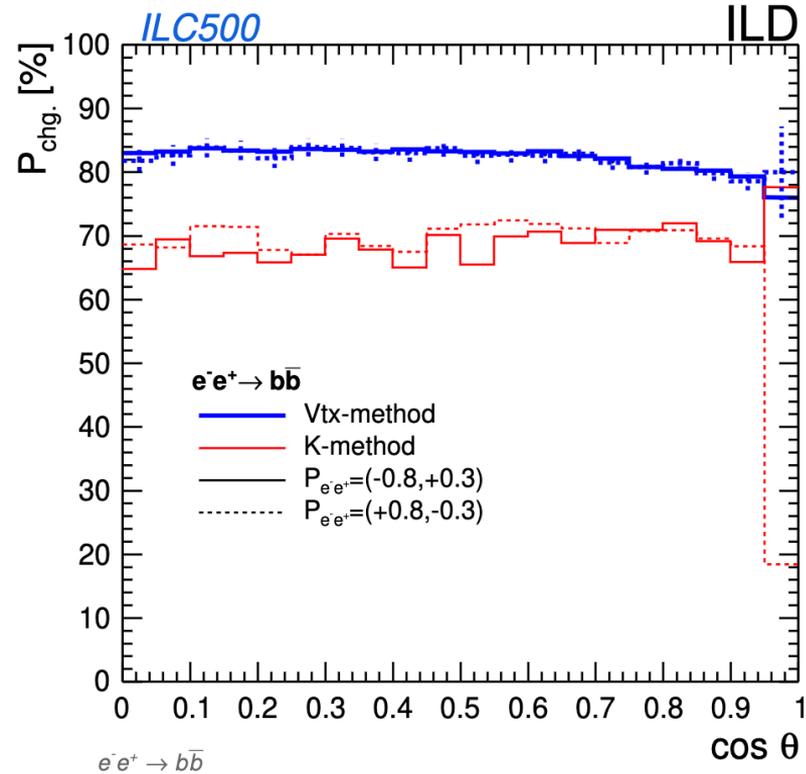
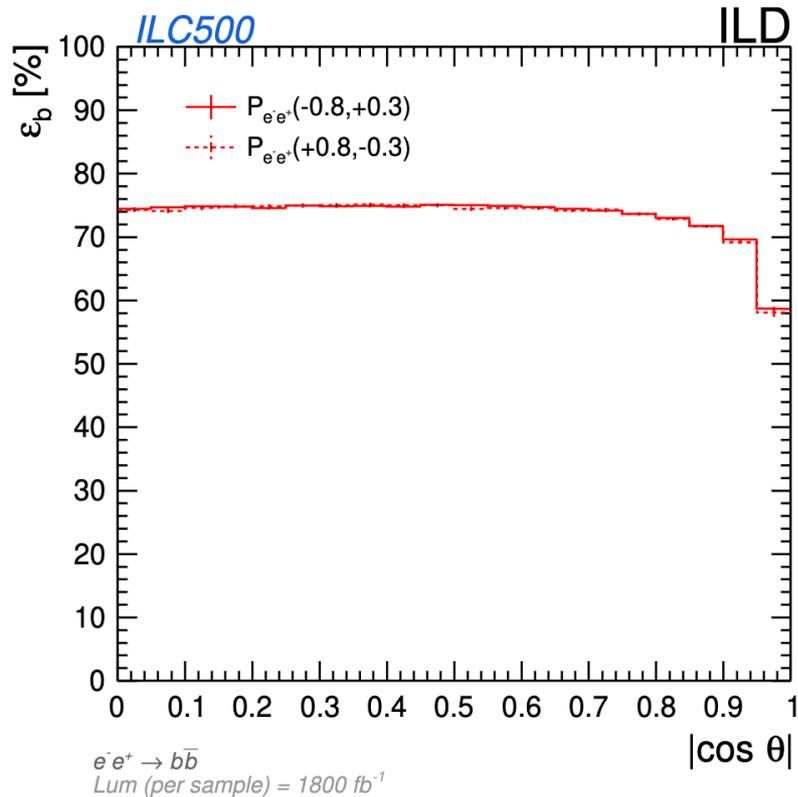
# Jet flavour tagging & charge measurement

- ▶ Double tagging & double charge measurement methods. ( described in previous ILD Note [2306.11413 \(2022\)](#) )
  - To maximally reduce the usage of MC tools (control of fragmentation, QCD correlations... uncertainties)



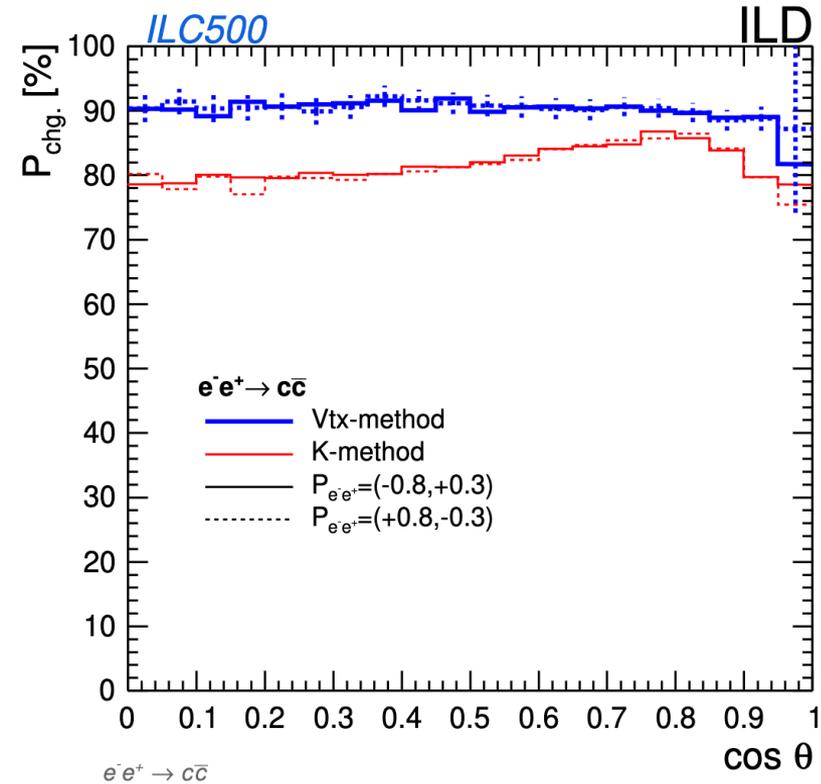
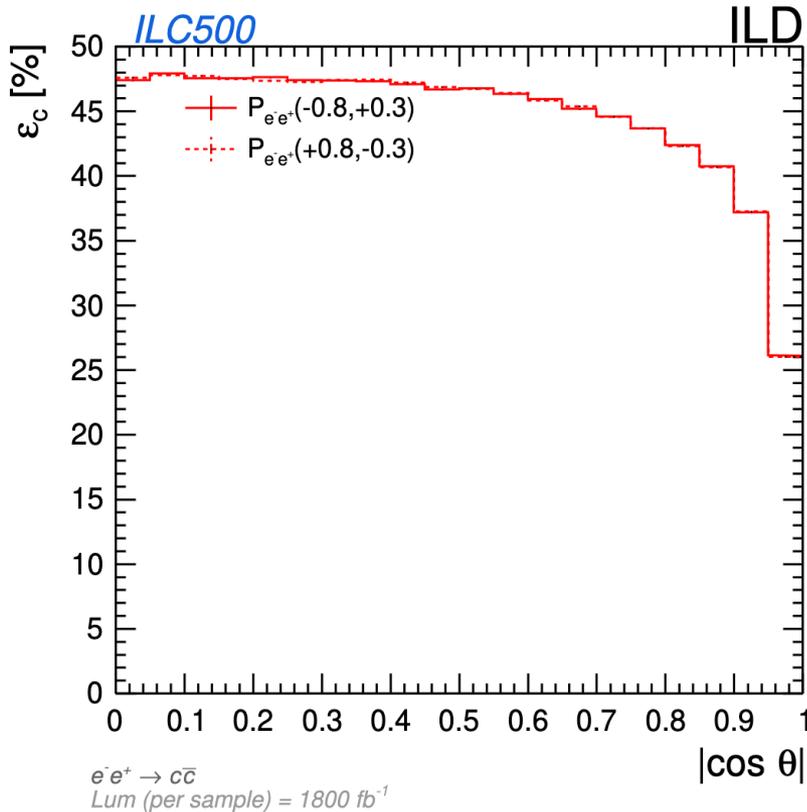
# Jet flavour tagging & charge measurement

- ▶ Double tagging & double charge measurement methods ( Plots from LCWS proceeding: 2307.14888 ).



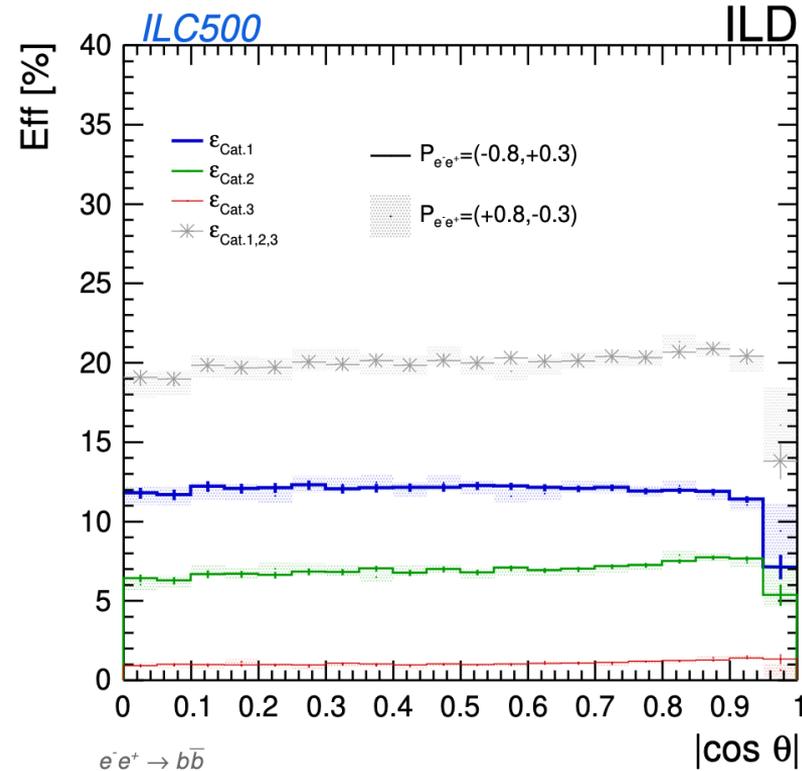
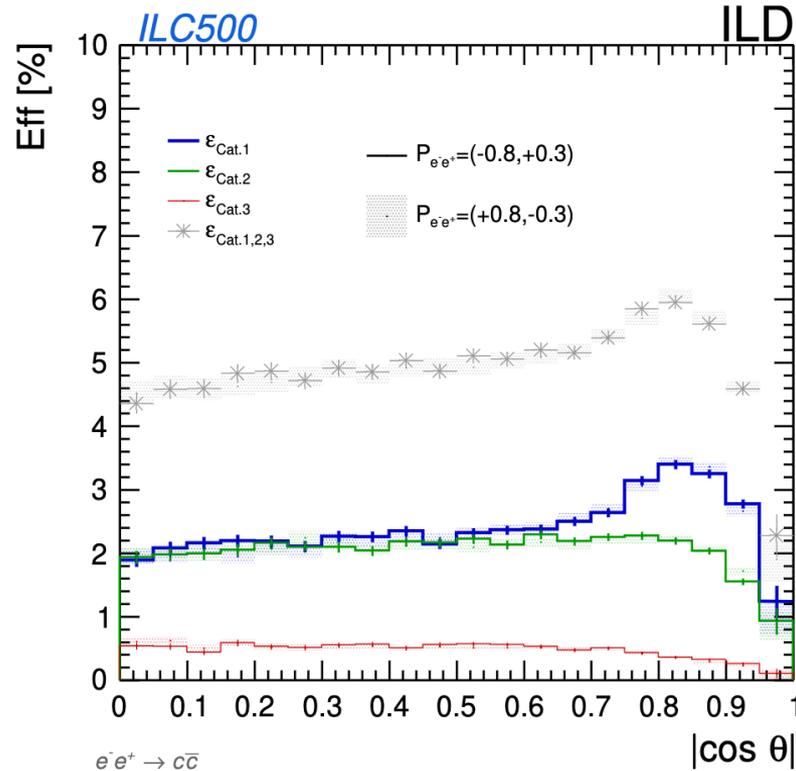
# Jet flavour tagging & charge measurement

► Double tagging & double charge measurement methods ( Plots from LCWS proceeding: [2307.14888](#) ).



# Jet flavour tagging & charge measurement

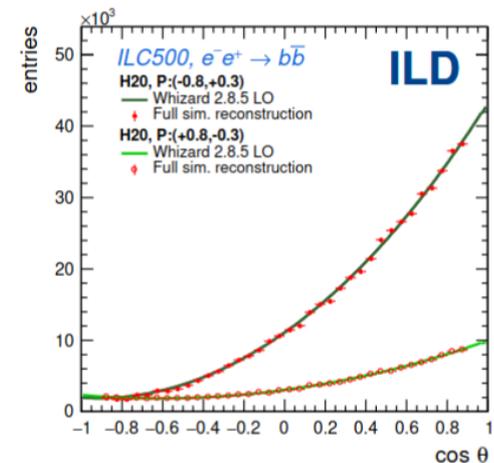
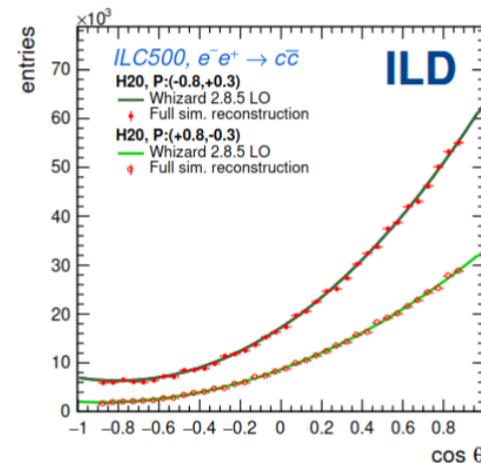
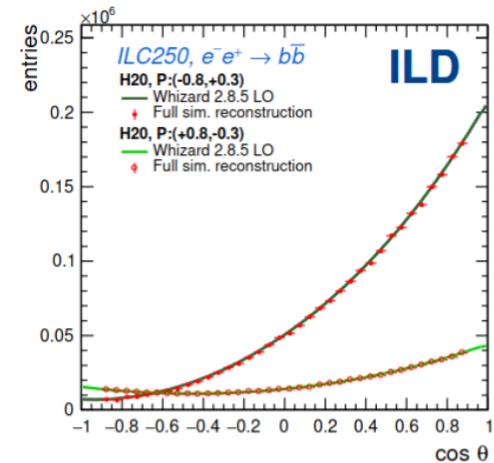
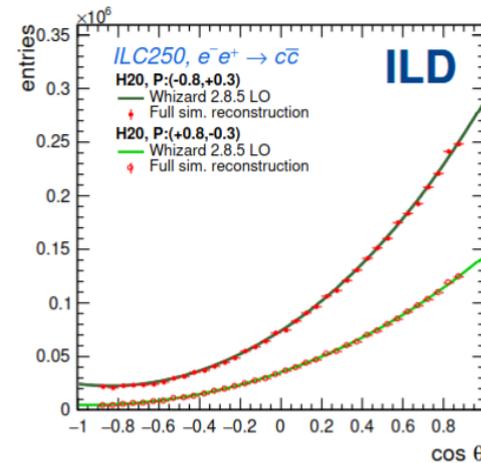
► Double tagging & double charge measurement methods ( Plots from LCWS proceeding: [2307.14888](#) ).



**Selection for the  $A_{\text{FB}}$  measurement at 500 GeV**

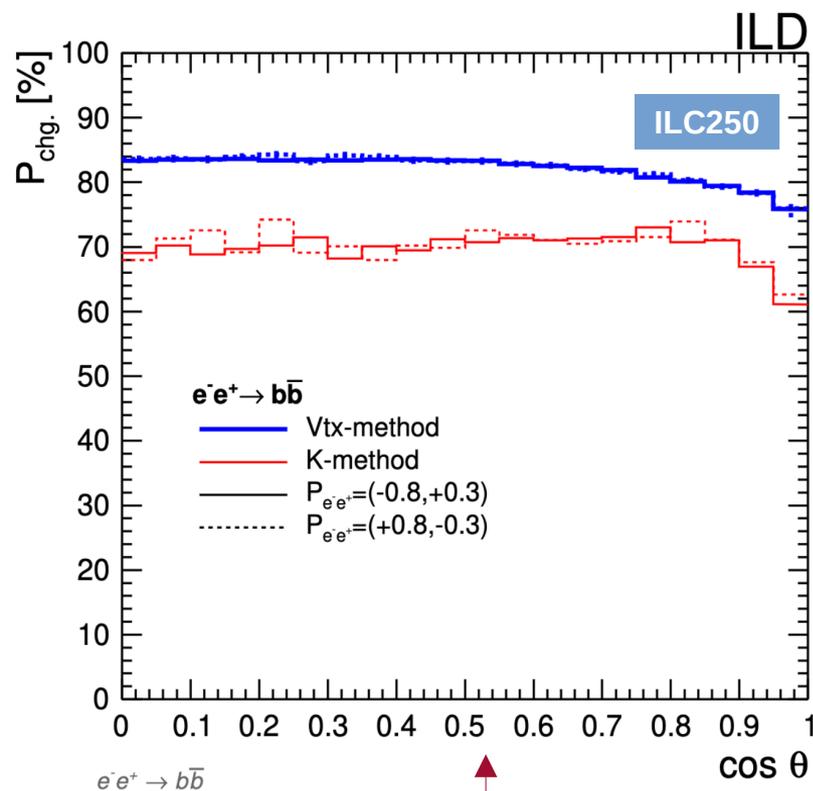
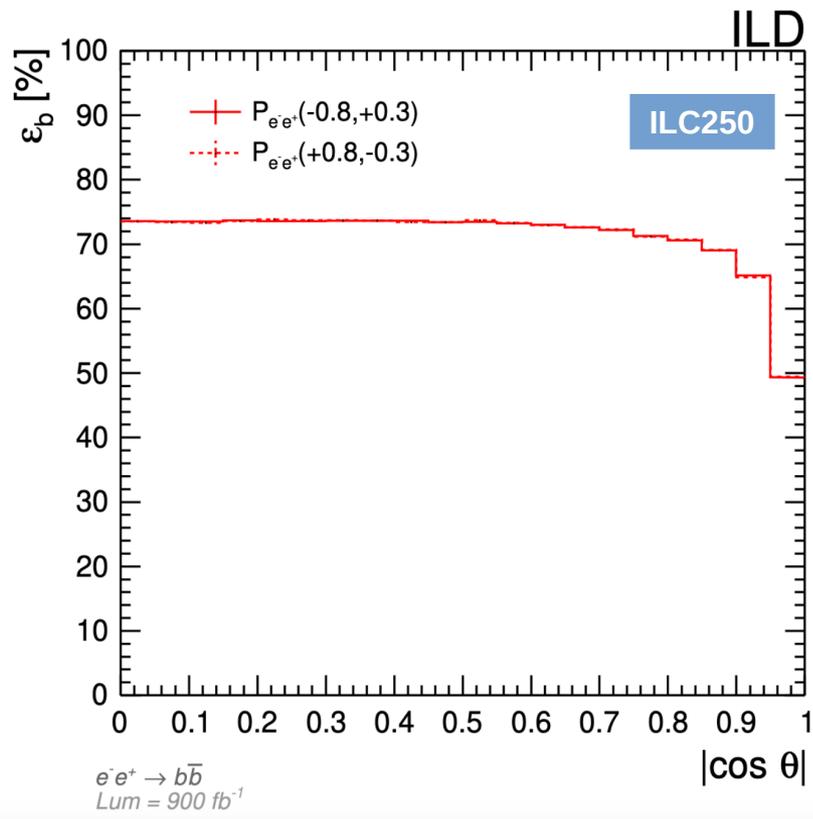
# Fit and results

- ▶ **Results for ILC250 and ILC500:**
  - Presented in LCWS (2023) [2307.14888](https://arxiv.org/abs/2307.14888)
- ▶ **At least 4 observables for AFB at ILC per energy point**
  - 2 quarks (b and c).
  - 2 polarizations ( $e_{LP_R}$ ,  $e_{RP_L}$ ).
- ▶ **Per mil level statistical uncertainties** reachable for the nominal **ILC program**
  - **Smaller exp syst. Uncertainties.**
    - Fragmentation, angular correlations, preselection efficiency, mistag, etc.



# Jet flavour tagging & charge measurement

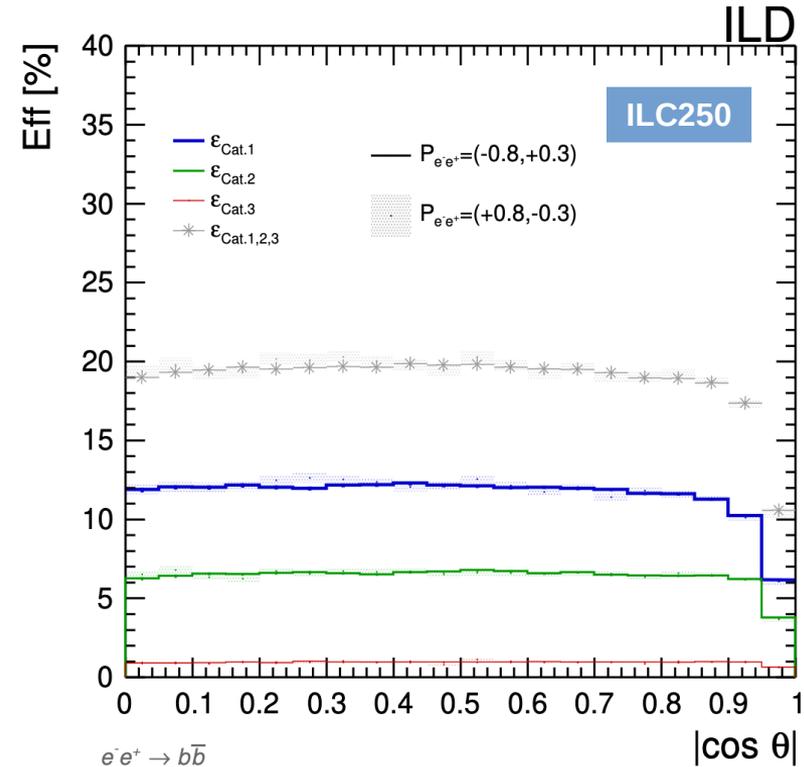
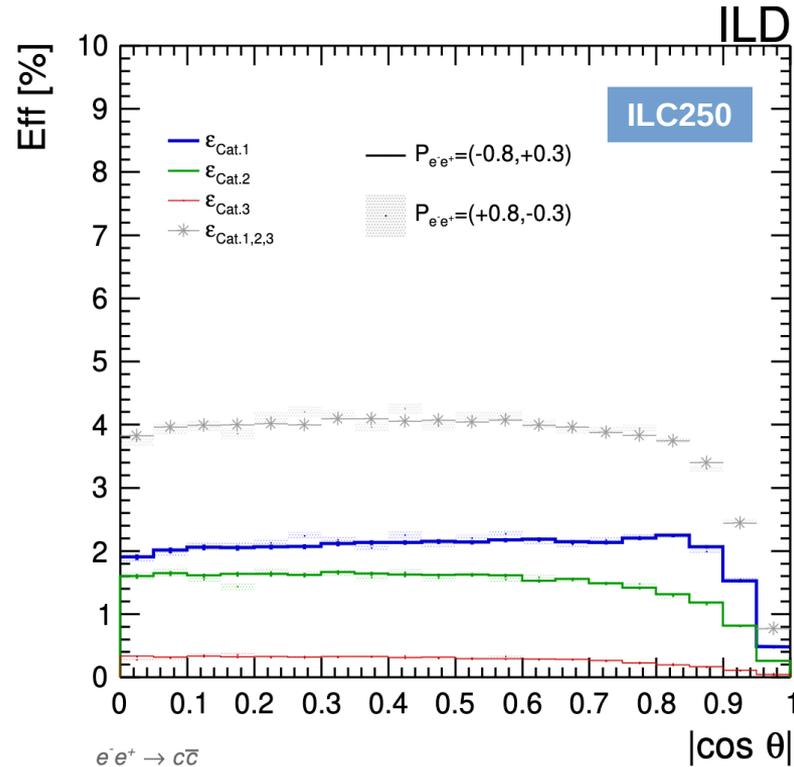
- ▶ Double tagging & double charge measurement methods. ( described in previous ILD Note [2306.11413 \(2022\)](#) )
  - To maximally reduce the usage of MC tools (control of fragmentation, QCD correlations... uncertainties)



PID is very important  
for this measurement!

# Jet flavour tagging & charge measurement

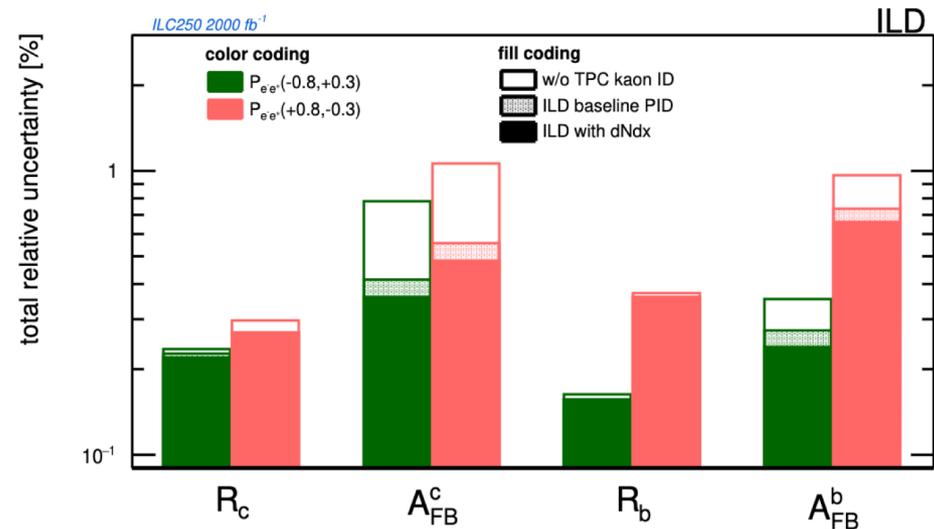
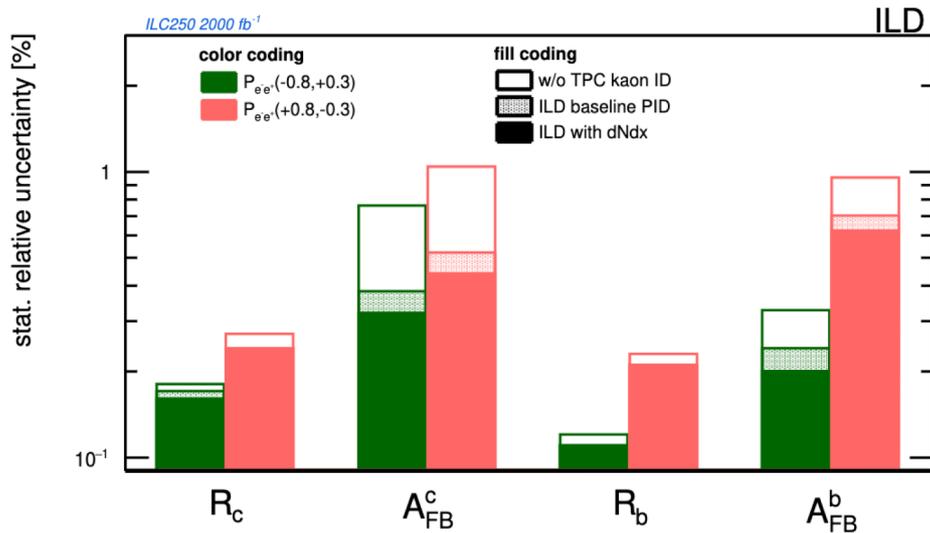
- Double tagging & double charge measurement methods. ( described in previous ILD Note [2306.11413 \(2022\)](#) )



Selection for the  $A_{\text{FB}}$  measurement at 250 GeV

# Uncertainties ILC250

► Presented in LCWS (2023) [2307.14888](#)

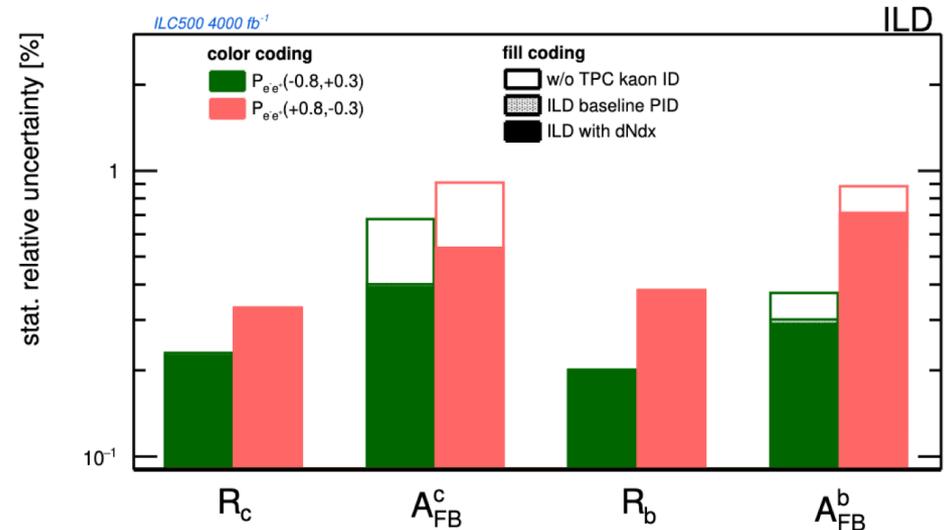
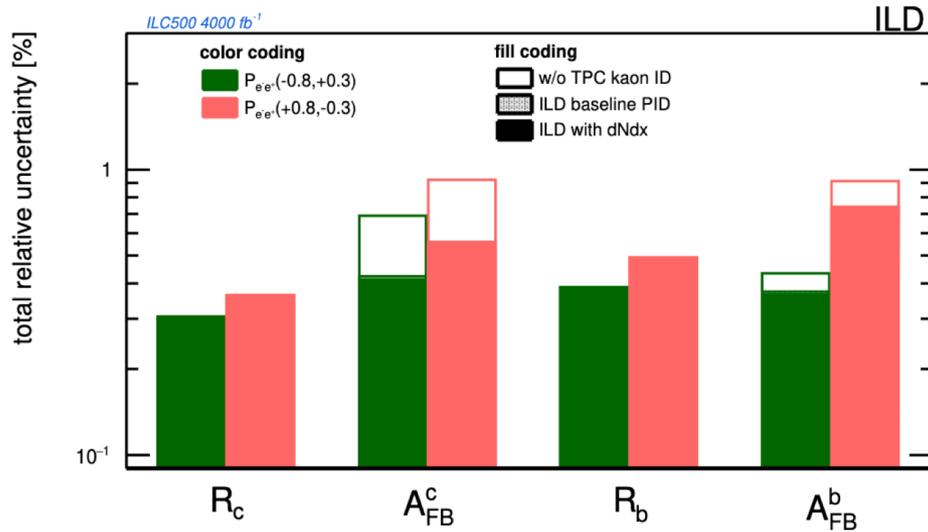


Statistical uncertainties dominate over systematic uncertainties

# Uncertainties ILC500

► Presented in LCWS (2023) [2307.14888](#)

- Less benefit from the use of PID, but the  $A_{FB}$  uncertainties are in the same level.



Statistical uncertainties dominate over systematic uncertainties