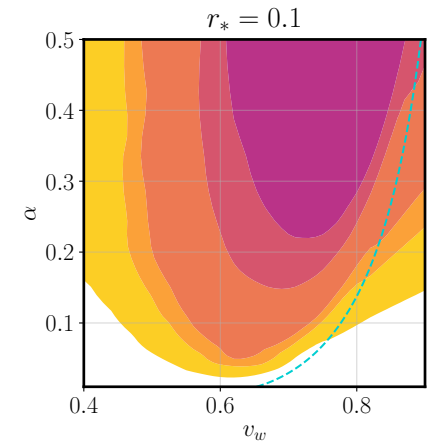
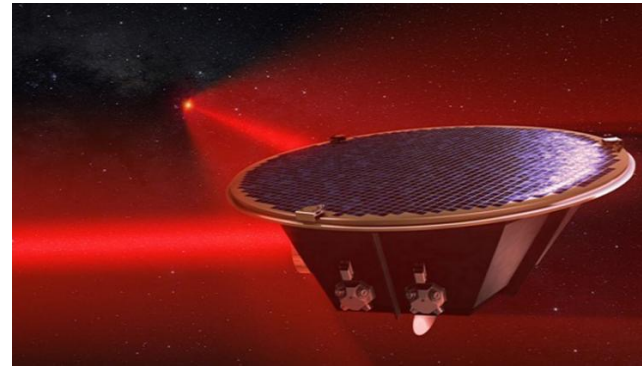
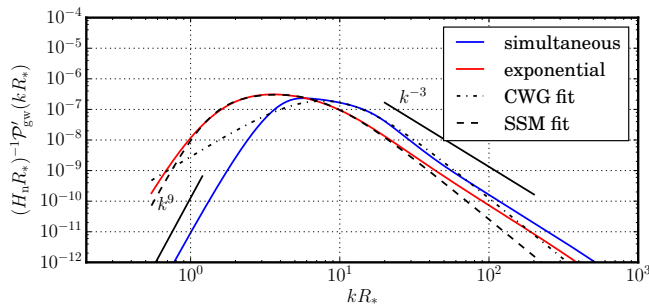
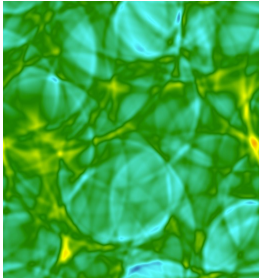




Gravitational waves from first order phase transitions

Mark Hindmarsh

Helsinki Institute of Physics & Dept of Physics,
University of Helsinki

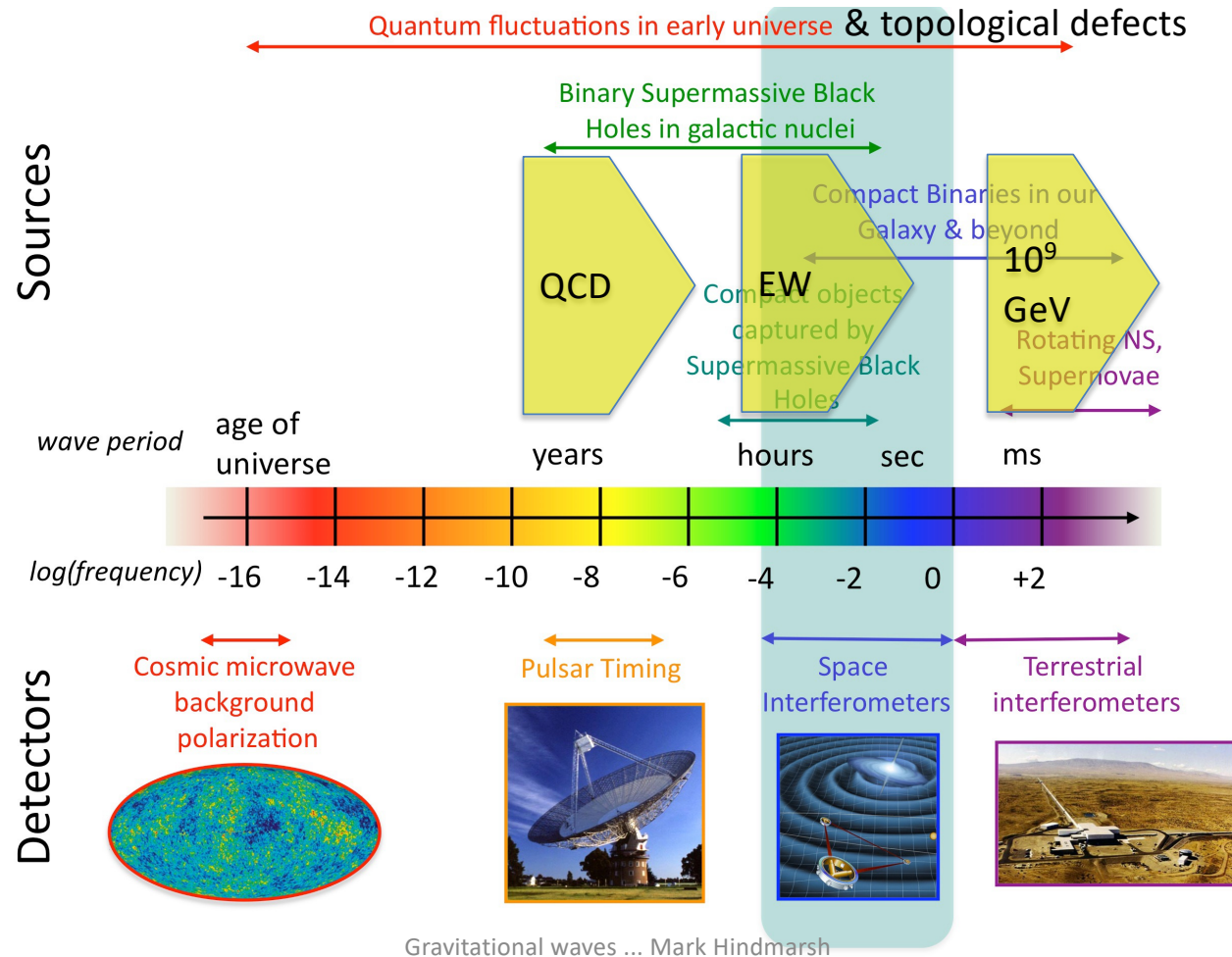
and

Department of Physics & Astronomy,
University of Sussex

Towards ... Valencia
11. joulukuuta 2024

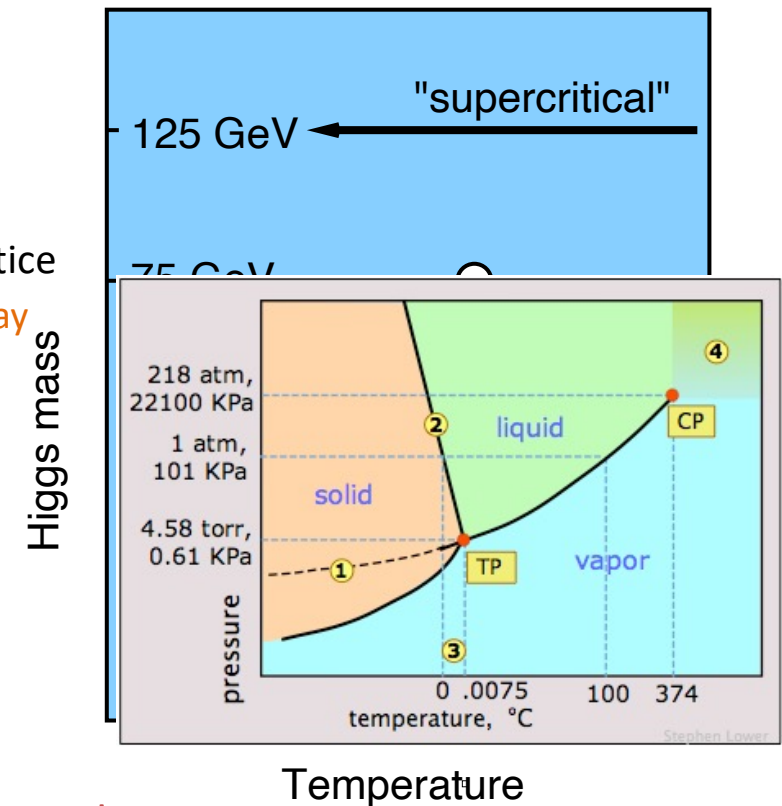
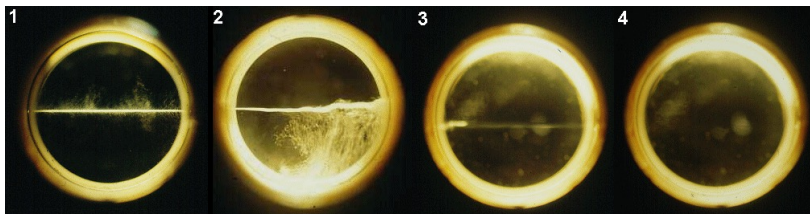


Gravitational wave spectrum



Electroweak transition: 100 GeV, 10 ps

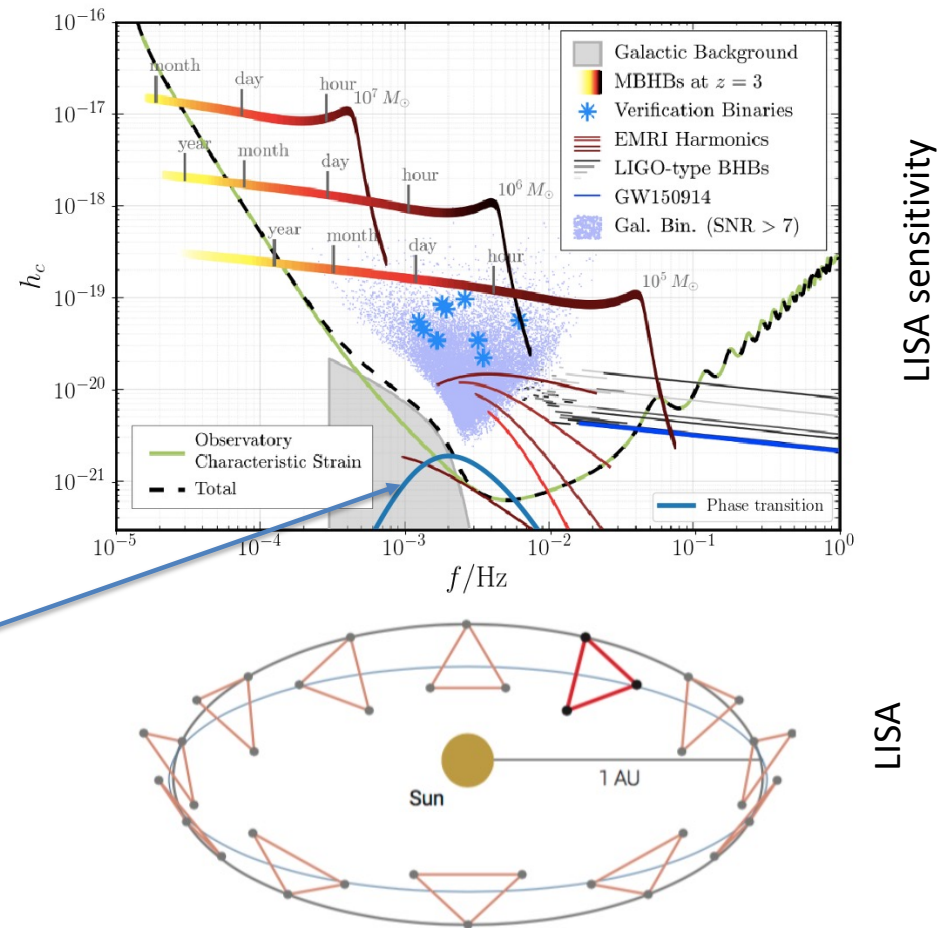
- Perturbative: weakly first order transition
Kirzhnits, Linde (1972,4)
- But: SM is not weakly coupled at high T
Linde (1980)
- Non-perturbative techniques:
 - Dimensional reduction to 3D effective field theory + 3D lattice
Kajantie, Laine, Rummukainen, Shaposhnikov (1995,6) P Schicho talk, Friday
 - SU(2)-Higgs on 4D lattice
Czikor, Fodor, Heitger (1998)
- SM transition at $m_h \approx 125$ GeV is a cross-over
- a **supercritical fluid**



- Search for 1st order transition is a search for physics beyond SM

Laser Interferometer Space Antenna

- Launch mid 2030s
- 4-year mission (up to 10 years)
- 2.5M km arms
- Science objectives:
 - White dwarves
 - Black holes
 - Galaxy mergers
 - Extreme gravity
 - TeV-scale early Universe
- Other missions: Taiji, TianQin
- Proposals: DECIGO, BBO



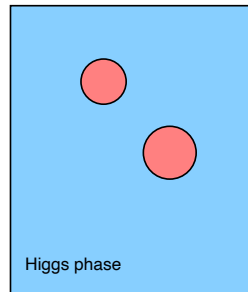
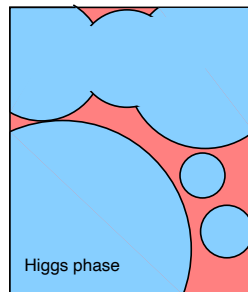
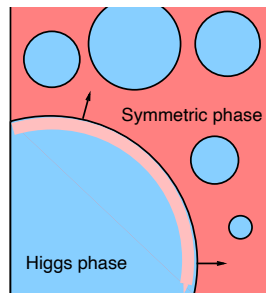
Little bangs in the Big Bang

- 1st order transition by nucleation of bubbles of low- T phase
Langer 1969, Coleman 1974, Linde 1983
- Nucleation rate/volume $p(t)$ rapidly increases below T_c
- Expanding bubbles generate pressure waves in hot fluid
- Gravitational wave (GW) production
- GW spectrum has information about phase transition
- Departure from equilibrium: needed for baryogenesis

Sakharov (1967)

Kuzmin, Rubakov, Shaposhnikov (1985)

M. Postma, Friday



Steinhardt (1982); Hogan (1983,86);
Gyulassy et al (1984); Witten (1984)

Gravitational waves ... Mark Hindmarsh

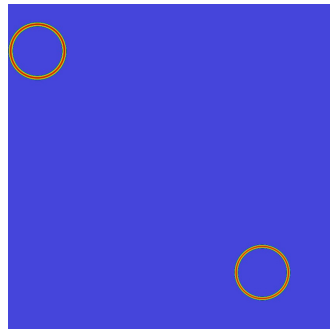
Fluid kinetic energy



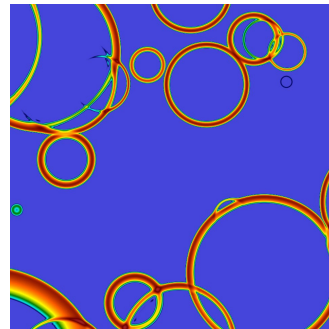
MH, Huber, Rummukainen, Weir (2013,5,7)
Cutting, MH, Weir (2018,9)

David Weir

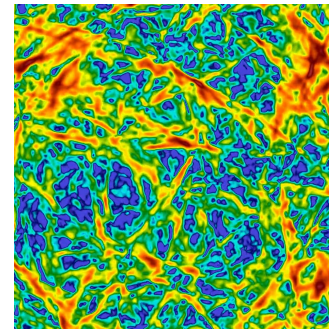
Phases of a phase transition



1



2



3



4

1. Bubble nucleation and expansion
2. Collision
3. Acoustic waves (vorticity)
4. Non-linear (shocks, turbulence)

$$\tau_{nl} \sim L_f / \bar{U}_f$$

L_f – fluid flow length scale

U_f – RMS fluid velocity

‘exponential’ nucleation rate/volume $p(t)$

$$p(t) = p_n e^{\beta(t-t_n)}$$

$$\tau_{co} = \beta^{-1} \text{ Ekstedt, Gould, Hirvonen 2023}$$

β – transition rate parameter

$\beta > H$ for successful transition

Guth, Weinberg 1983; Enqvist et al 1992;

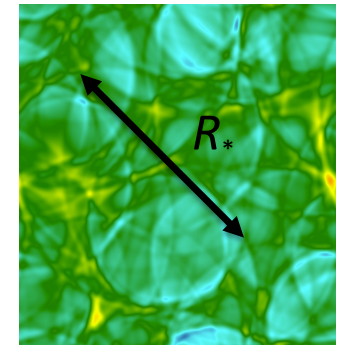
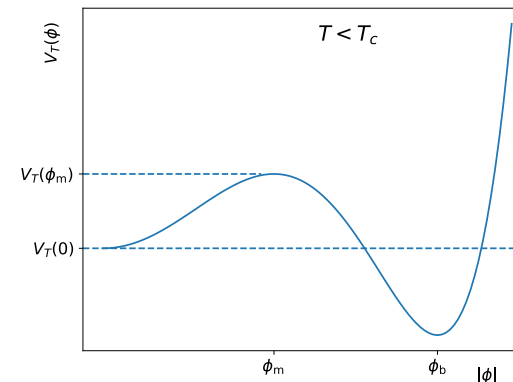
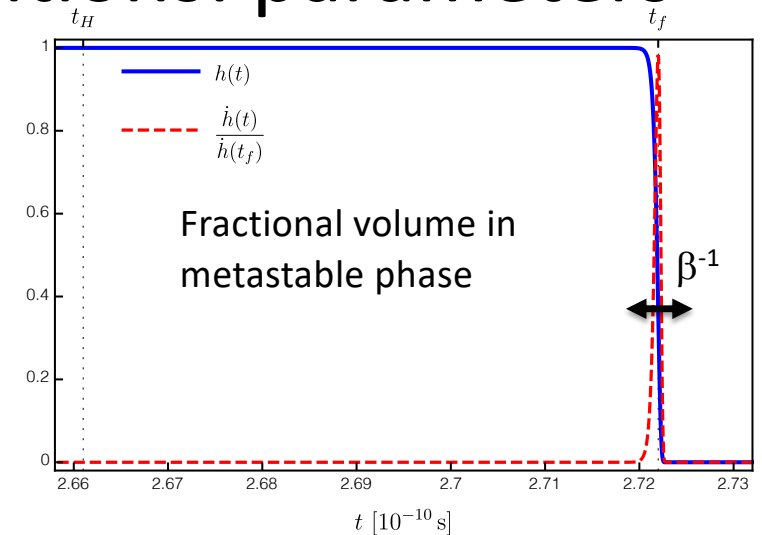
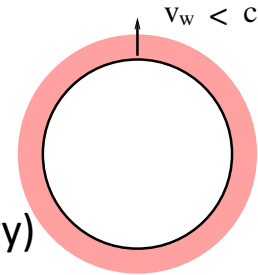
Turner, Weinberg, Widrow 1992

Review: MH, Lüben, Lumma, Pauly 2021

GWs from first order phase transitions: parameters

- Parameters of transition:
 - T_n = Temperature at nucleation
 - β = transition rate ($= -d \log p / dt$)
 - v_w = Bubble wall speed
 - α = (potential energy change^{*})/(heat energy)
 - c_{s1}, c_{s2} = sound speeds Giese et al 2020
- Useful derived parameters:
 - r_* = (bubble centre spacing R_*)/Hubble length
 - K = fluid kinetic energy fraction Steinhardt '84
Espinosa et al 2010
- Fluid kinetic energy makes GWs
 - Energy release via self-similar bubble solutions

$$* \frac{1}{4} \Delta(e - 3p)$$



GWs from an early universe phase transition

Assume rapid transition, $\beta \gg H$, neglect expansion of universe

- Effective theory:

Ignatius et al (1994), Kurki-Suonio, Laine (1996)

- **Higgs field**

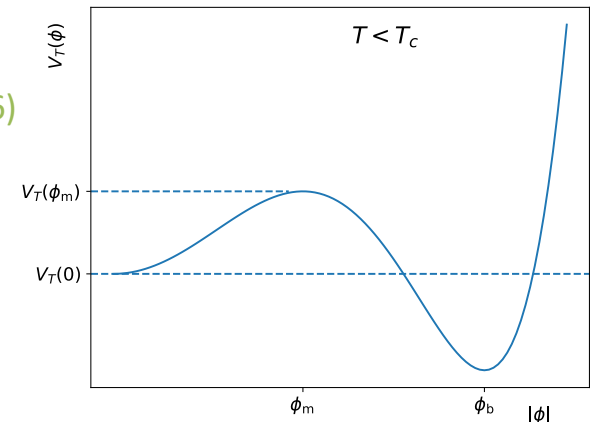
$$\square\phi - V'_T(\phi) = \eta_T(\phi)U \cdot \partial\phi$$

- $V_T(\phi)$ equation of state
- $\eta_T(\phi)$ field-fluid coupling (models friction)

- **Relativistic fluid (ideal limit)**

$$T_f^{\mu\nu} = (e + P)U^\mu U^\nu + P g^{\mu\nu}$$

$$\partial_\mu T_f^{\mu\nu} + \partial^\nu \phi V'_T(\phi) = \eta_T(\phi)(U \cdot \partial\phi)\partial^\nu \phi$$



- **Metric perturbation (GW strain)**

$$\ddot{u}_{ij} - \nabla^2 u_{ij} = 16\pi G T_{ij} \quad \longrightarrow \quad \tilde{h}_{ij}(\mathbf{k}) = \Lambda_{ij,kl}^{TT} u_{kl}(\mathbf{k}) \quad \text{Garcia-Bellido, Figueroa, Sastre (2008)}$$

GWs from an early universe phase transition

Assume rapid transition, $\beta \gg H$, neglect expansion of universe

- Effective theory: Ignatius et al (1994), Kurki-Suonio, Laine (1996)

– **Higgs field** $-\ddot{\phi} + \nabla^2 \phi - V'_T(\phi) = \eta_T(\phi)W(\dot{\phi} + V^i \partial_i \phi)$

- $V_T(\phi)$ equation of state
- $\eta_T(\phi)$ field-fluid coupling (“friction”, high T: $\eta_T(\phi) \propto \phi^2/T$)

– **Relativistic fluid**

$$\dot{E} + \partial_i(EV^i) + P[\dot{W} + \partial_i(WV^i)] - \frac{\partial V}{\partial \phi} W(\dot{\phi} + V^i \partial_i \phi) = \eta W^2(\dot{\phi} + V^i \partial_i \phi)^2.$$

$$\dot{Z}_i + \partial_j(Z_i V^j) + \partial_i P + \frac{\partial V}{\partial \phi} \partial_i \phi = -\eta W(\dot{\phi} + V^j \partial_j \phi) \partial_i \phi.$$

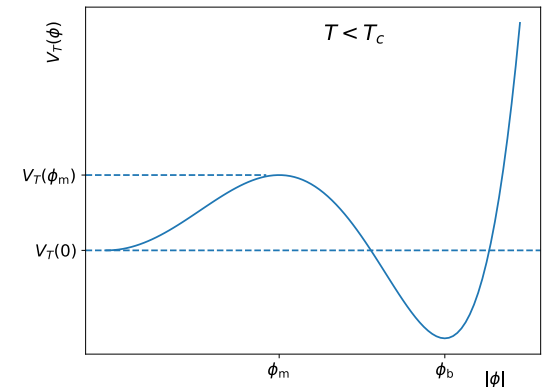
- $E = W^*(\text{energy density})$, $Z_i = \text{momentum density}$, $V_i = 3\text{-velocity}$, $W = \text{Lorentz factor}$

– Discretisation * Wilson & Matthews (2003)

Also: Brandenburg, Enqvist, Olesen (1996); Giblin, Mertens (2013); Jinno, Konstandin, Rubira (2020)

– **Metric perturbation** (GW strain)

$$\ddot{u}_{ij} - \nabla^2 u_{ij} = 16\pi G T_{ij} \quad \longrightarrow \quad \tilde{h}_{ij}(\mathbf{k}) = \Lambda_{ij,kl}^{TT} u_{kl}(\mathbf{k}) \quad \text{Garcia-Bellido, Figueroa, Sastre (2008)}$$



Connection to fundamental theory

• Scalar hydrodynamics $-\ddot{\phi} + \nabla^2 \phi - V'_T(\phi) = \eta_T(\phi)W(\dot{\phi} + V^i \partial_i \phi)$

• Scalar effective potential $V_T(\phi)$ $\rightarrow$

equilibrium, quasi-eqm. $(T_n, \alpha, \beta, c_s, g_{\text{eff}})$

• Scalar-fluid coupling $\eta_T(\phi)$ $\rightarrow$

non-equilibrium (v_w)

Ekstedt, Schicho, Tenkanen 2024

Ekstedt, Gould, Hirvonen 2023

Ekstedt et al 2024

$$\Omega_{\text{gw}}(f) = \frac{2\pi^2}{3H_0^2} f^2 h_c^2(f)$$

Phase transition parameters :

T_n = nucleation temperature

g_{eff} = effective d.o.f. in plasma

$\alpha \sim (\text{potential energy})/(\text{thermal energy})$

c_{s1}, c_{s2} = sound speed(s)

β = transition rate

v_w = bubble wall speed

Simulations, Modelling

$H_n(T_n, g_{\text{eff}})$ (Hubble rate)

$K(v_w, \alpha, c_s)$ (kinetic energy fraction)

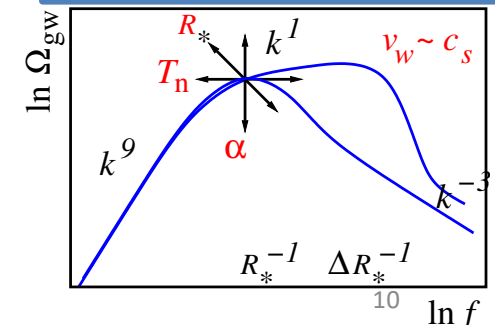
$R_*(\beta, v_w)$ (mean bubble separation)

GW spectrum

Ω_p = peak amplitude

f_p = peak frequency

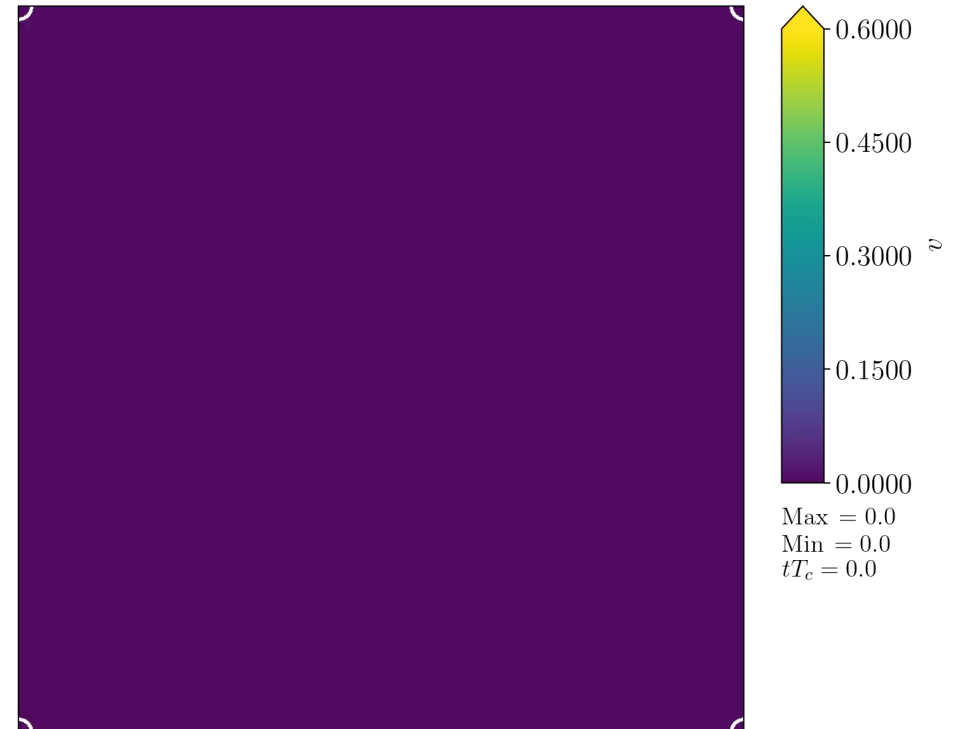
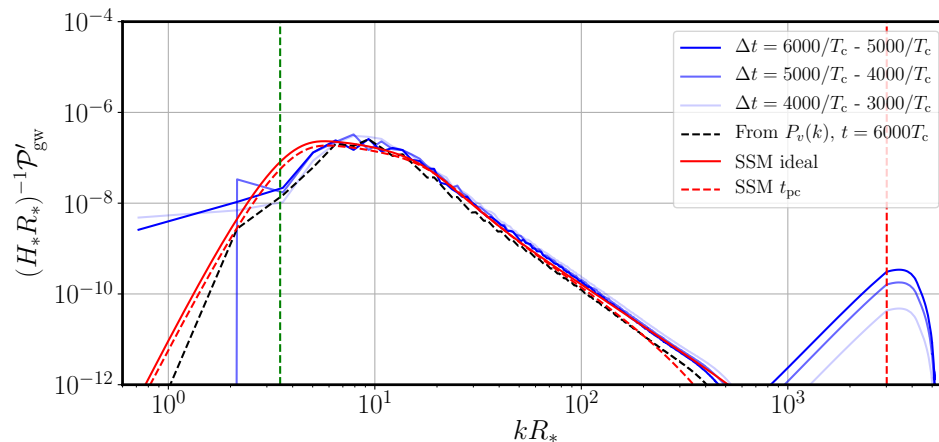
σ_i = shape parameters



3D hydrodynamic simulations of phase transitions

Hindmarsh et al 2013, 2015, 2017, 2019; Jinno et al 2022; Caprini et al 2024

- Relativistic fluid + scalar order parameter (“Higgs”)
- Linearised gravitational wave production
- Key output: GW power spectrum
(fractional GW energy density per log wavenumber)



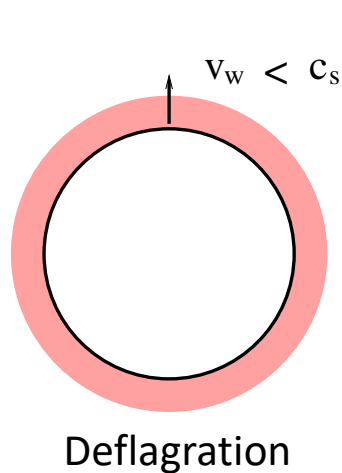
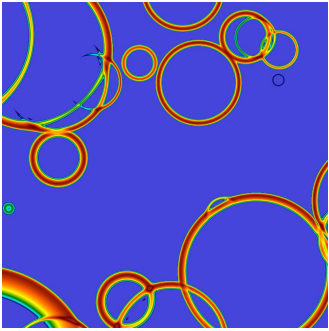
- While sound waves persist, GW power spectrum grows
- Plot: GW power spectrum **growth rate** (scaled by H_n)



Gravitational waves ... Mark Hindmarsh

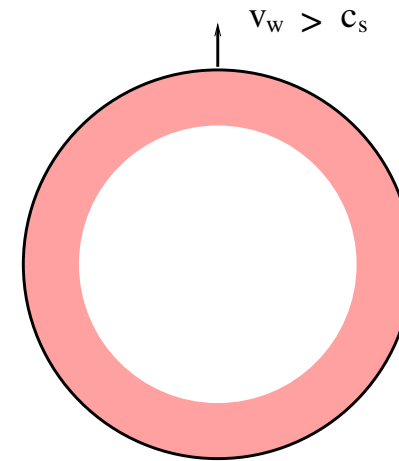
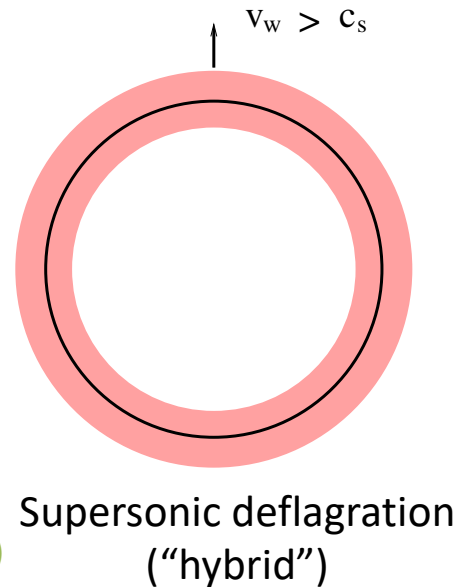


Towards a model: relativistic combustion

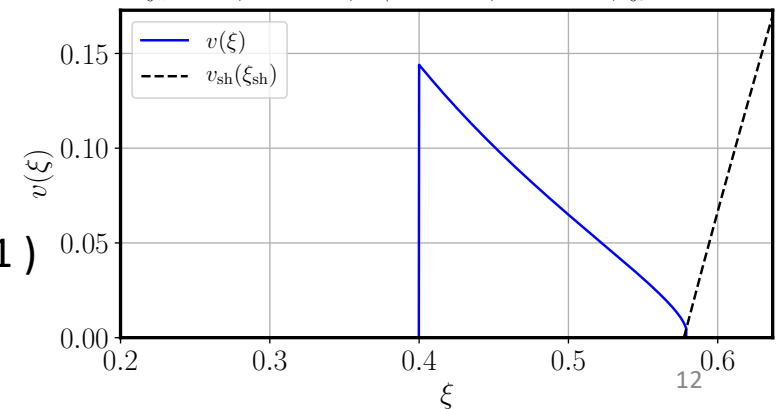


Landau & Lifshitz; Steinhardt (1984)

Kurki-Suonio, Laine (1991), Espinosa et al (2010)

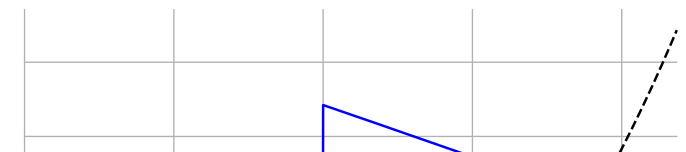


$\xi_w = 0.4, \alpha_n = 0.1, \alpha_+ = 0.078, r = 1.624, \xi_{sh} = 0.579$



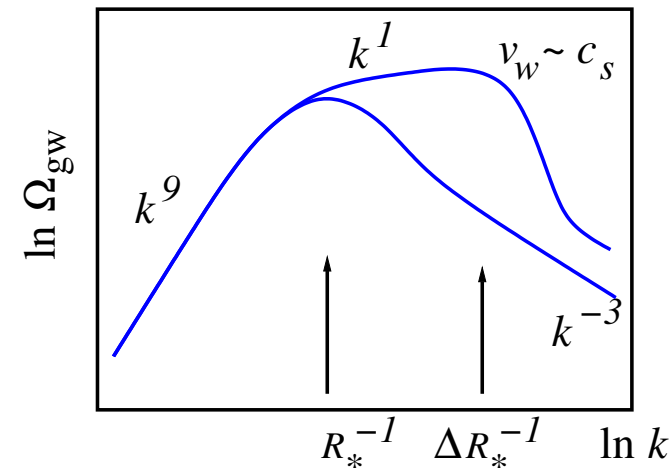
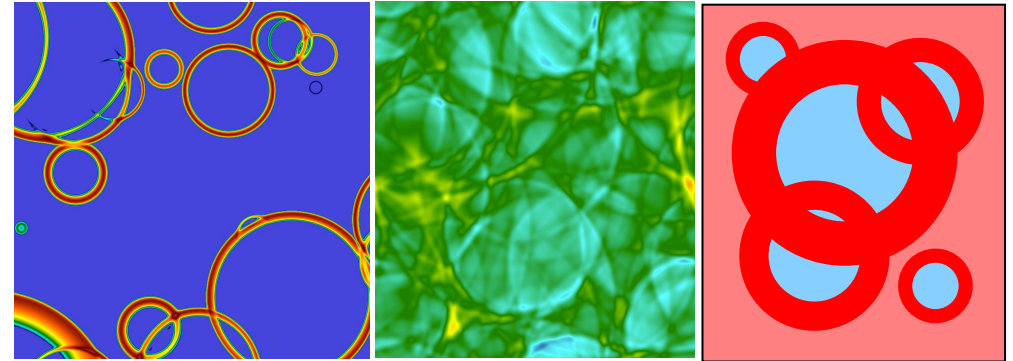
- Large scales: ideal relativistic hydrodynamics
- Microphysics: $e(T)$, $p(T)$, v_w
- Radial fluid velocity $v(r,t)$ and enthalpy distribution $w(r,t)$
 - **Similarity solution $v(x)$, $w(x)$: $\xi = r/t$**
- Low-friction or ultra-strong transition: can be "runaway" ($v_w \rightarrow 1$)
 - (not considered here) Bodeker Moore 2010, 2017

Gravitational waves ... Mark Hindmarsh



GWs from phase transitions: Sound shell model

- GWs from Gaussian velocity field
Caprini, Durrer, Servant (2007,2009)
- Velocity field: weighted addition of self-similar sound “shells” $\mathbf{v}_q(t_i)$ from bubbles
MH 2016, MH, Hijazi (2019)
- Two length scales:
 - Bubble spacing R_*
 - Shell width $R_*|v_w - c_s|/c_s$
- Double broken power law
 - $P_{gw} \sim k^9, k^1, k^{-3}$
- Amplitude proportional to:
 - Bubble spacing
 - Shear stress lifetime
 - (Kinetic energy)²
- Similar: bulk flow model (k^1 at low k)
Jinno, Takimoto 2017; Konstandin 2018



Estimating GW power

$$\square h \sim T \longrightarrow P_h(t, k) \sim \int^t dt_1 \int^t dt_2 \cos[k(t - t_1)] \cos[k(t - t_2)] \langle T_k(t_1) T_k^*(t_2) \rangle$$

- GW energy fraction: $\longrightarrow \Omega_{\text{gw}} \sim (H_n \tau_v)(H_n \tau_c) K^2$
 - H_n Hubble rate (max nucleation)
 - τ_v duration of stresses
 - τ_c coherence time
- Coherence time:
 - $\tau_c \sim R_*$ (bubble spacing)
- Shear stress lifetime (shocks):
 - $\tau_v \sim H_n^{-1}/(1 + U_f/R_* H_n)$
 - U_f ($\sim$ RMS velocity) $\sim \sqrt{K}$
 - or $K = (4/3)U_f^2$

$R_*(\beta, v_w, \text{also } \alpha_n)$ (mean bubble spacing)
from bubble growth dynamics

$K(v_w, \alpha_n)$ (kinetic energy fraction)
from self-similar hydro solution

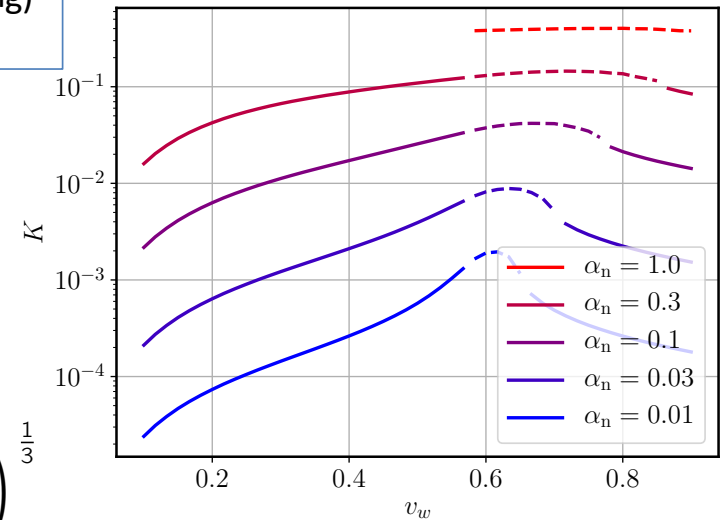
$$\Omega_{\text{gw},0} \simeq F_{\text{gw},0} \frac{r_*}{1 + \sqrt{K}/r_*} K^2 \tilde{\Omega}_{\text{gw}}$$

Numerical simulations:

$$\tilde{\Omega}_{\text{gw}} = \mathcal{O}(10^{-2})$$

Standard cosmology:

$$F_{\text{gw},0} = 3.6 \times 10^{-5} \left(\frac{100}{g_{\text{eff}}} \right)^{1/3}$$

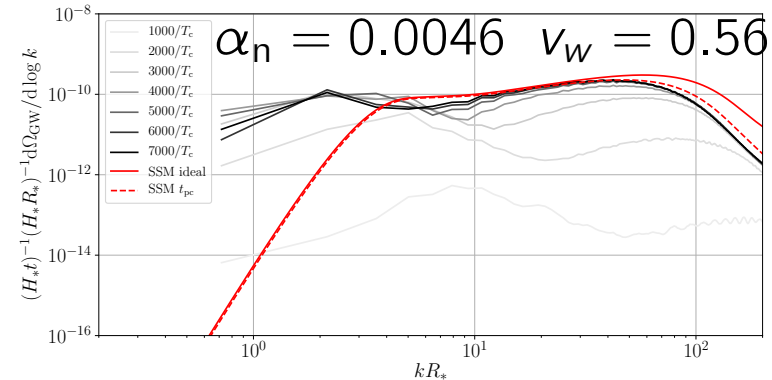
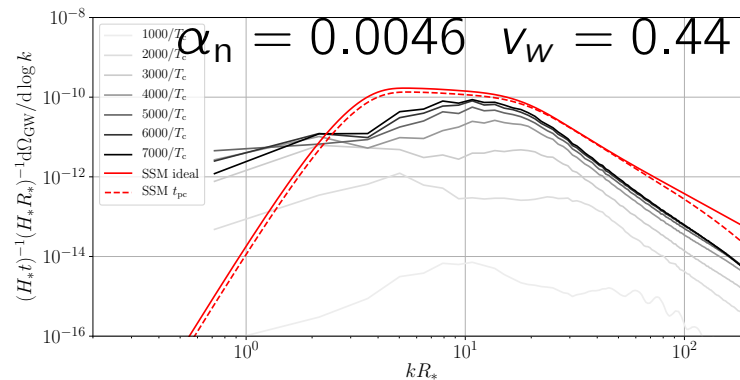
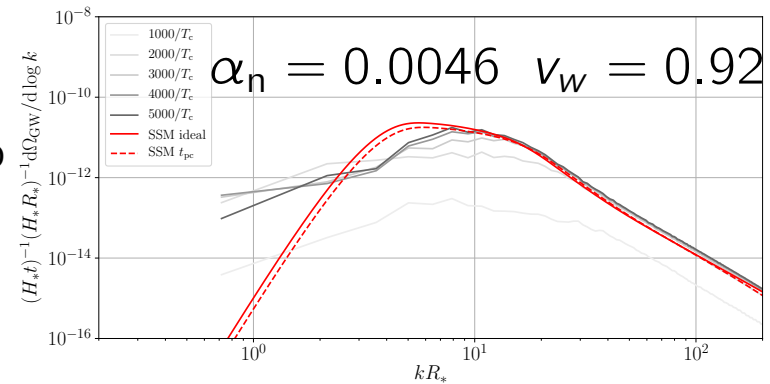


Sound shell model vs. simulations P_{gw}

GW power growth rate

- **Solid:** ideal self-similar sound shell
- **Dash:** evolving sound shell at peak collision time in 1+1D scalar hydro
- **Grey:** simulations:
 - simultaneous nucleation of bubbles

Hindmarsh, Huber, Rummukainen, Weir 2017

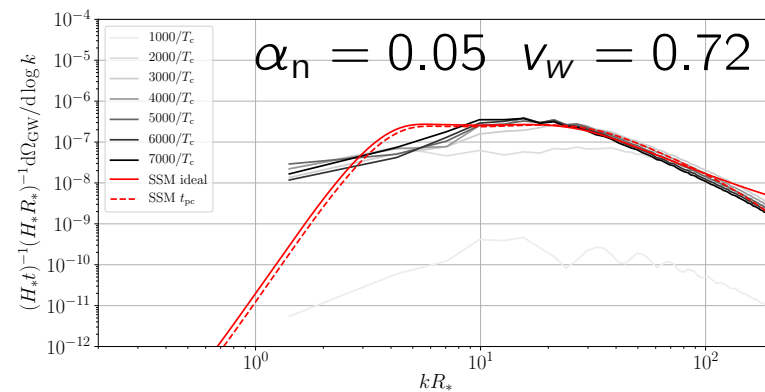
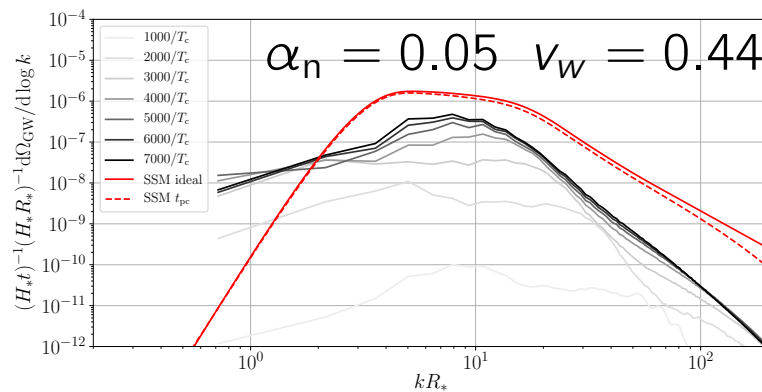
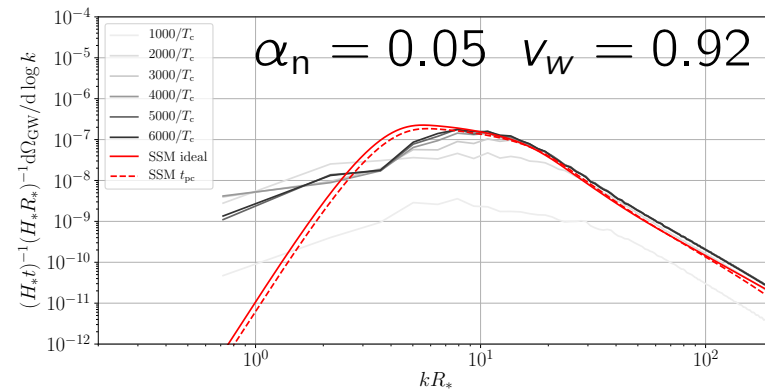


Sound shell model vs. simulations P_{gw}

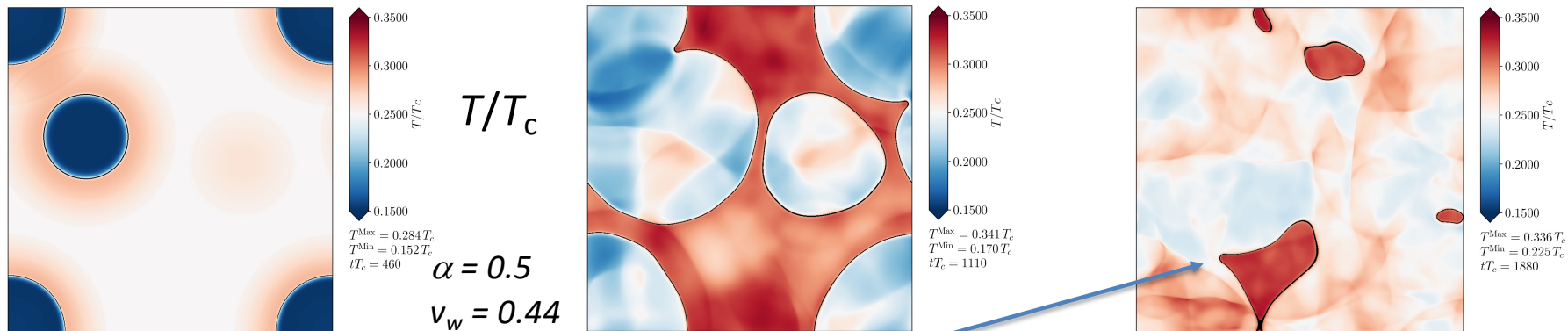
GW power growth rate

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Hindmarsh, Huber, Rummukainen, Weir 2017

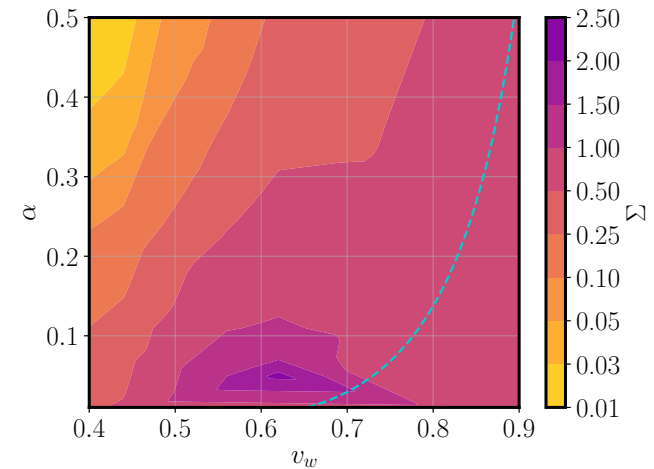


Nonlinearities 1: Kinetic energy & GW suppression



Cutting, MH, Weir, 2020

- **Deflagrations:** heat up fluid in front
- Pressure in front of wall increases, walls slow down
- Formation of hot droplets Cutting, Vilhonen, Weir (2022)
- Less transfer into kinetic energy, more into heat.
- Include **GW suppression factor** as a numerical parameter (**right**)
- Also: nucleation suppression Al-Ajmi, MH (2023)
 - bigger bubbles, boosts signal

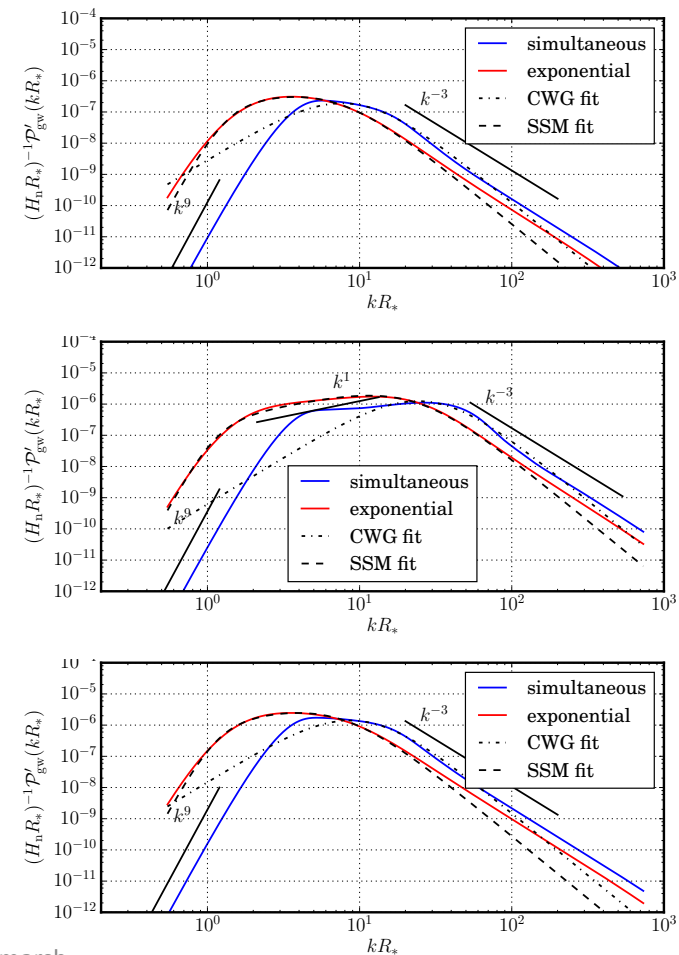


Gowling, MH (2021)

SSM toolbox: PTtools

Cutting, Gowling, MH, Mäki

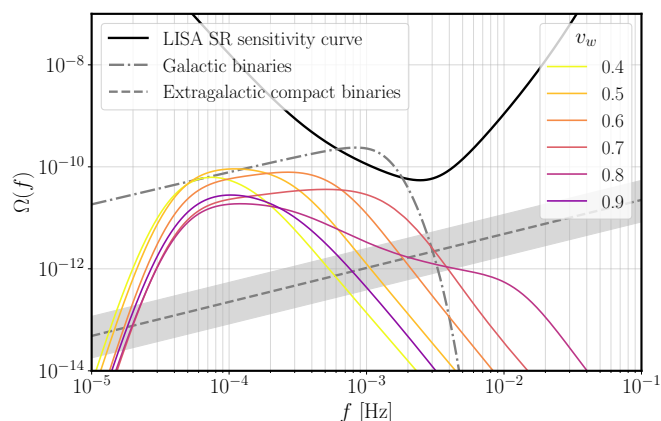
- Python package for GW spectra from 1st order phase transitions
- Currently: linearised approx with phenomenological corrections
 - intermediate strength ($\alpha \leq 0.3$)
 - Wall speeds $0.4 < v_w < 0.9$
 - Sub-Hubble bubble separations ($r_* \ll 1$)
- Speed-up with numba
 - 50 ms for 100 frequencies
- Support for general eqn of state
- Public release early 2025



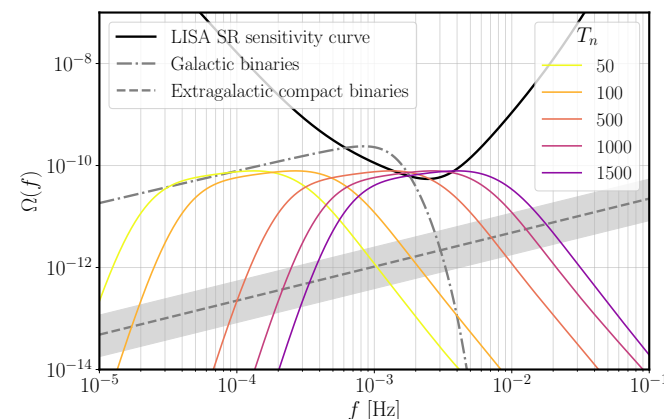
GW power spectra in the SSM

- **Sound shell model predictions: accuracy ~20% at peak for near-linear flows**
 - strengths $\alpha \leq 0.3$; wall speeds $0.4 \leq v_w \leq 0.9$; sub-Hubble bubble separations $r_* \ll 1$

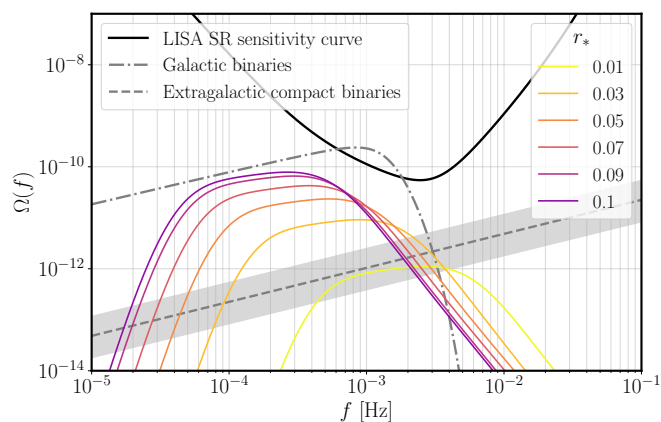
Gowling, MH (2021)



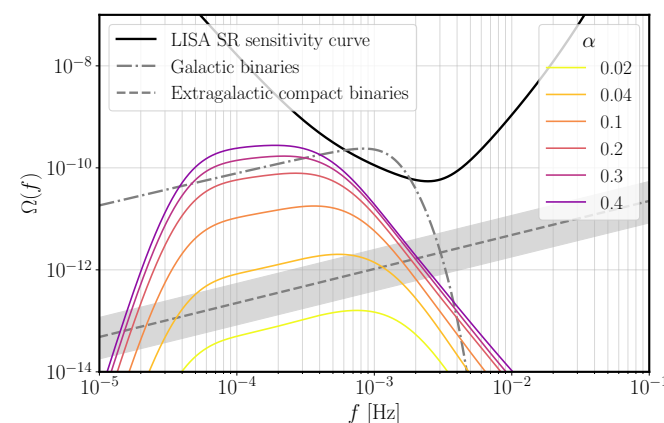
Vary wall speed



Vary transition temp

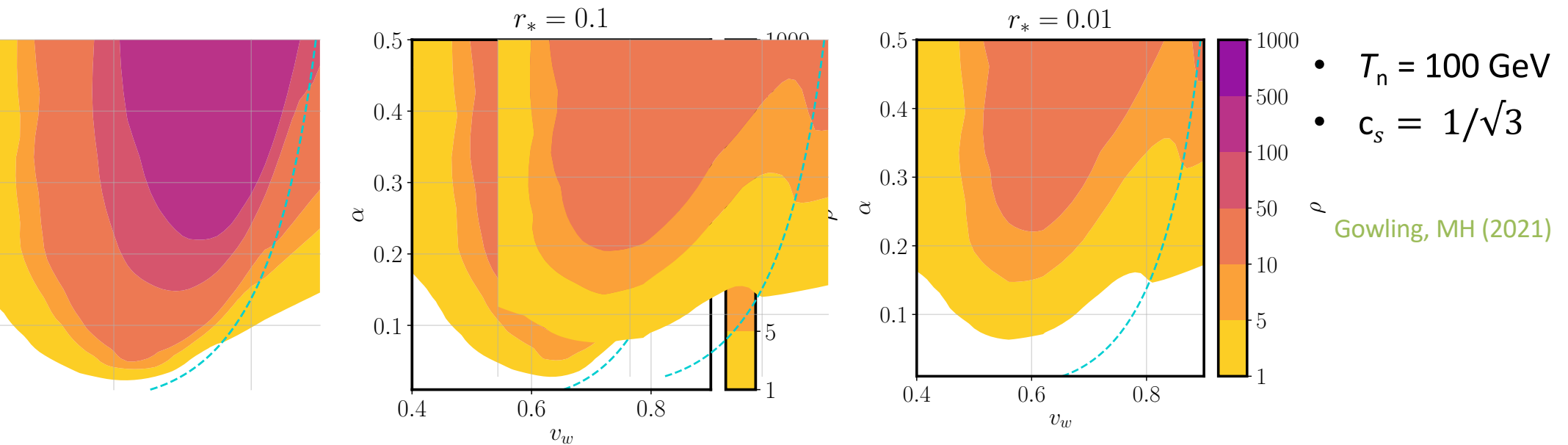


Vary bubble separation



Vary transition strength

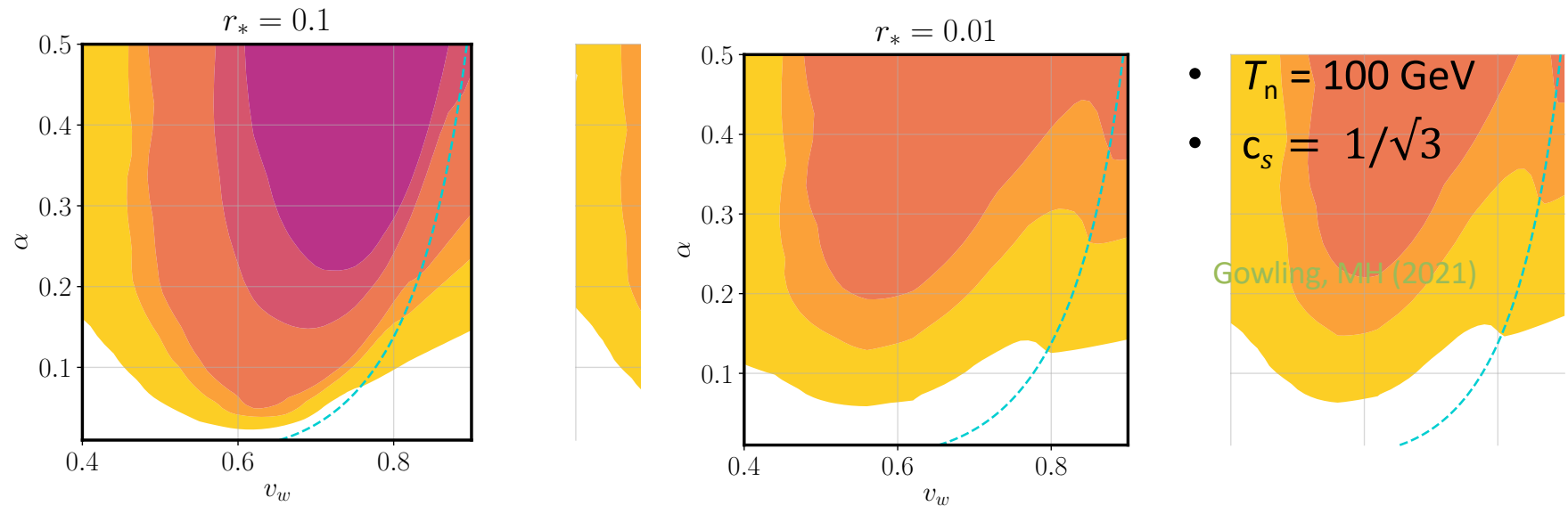
Signal-to-noise ratios (LISA)



- Signal-to-noise ratio ρ ($t_{\text{obs}} = 4 \text{ years}$)
- “Worst case” galactic binary foreground
 - (NB annual variation aids removal) [MH, Hooper, Minkkinen, Weir \(2024\)](#)
- “LISA science requirements” instrument noise

$$\rho^2(\vec{\theta}) = t_{\text{obs}} \int df \left(\frac{\Omega_{\text{gw}}(f; \vec{\theta})}{\Omega_{\text{noise}}(f)} \right)^2$$

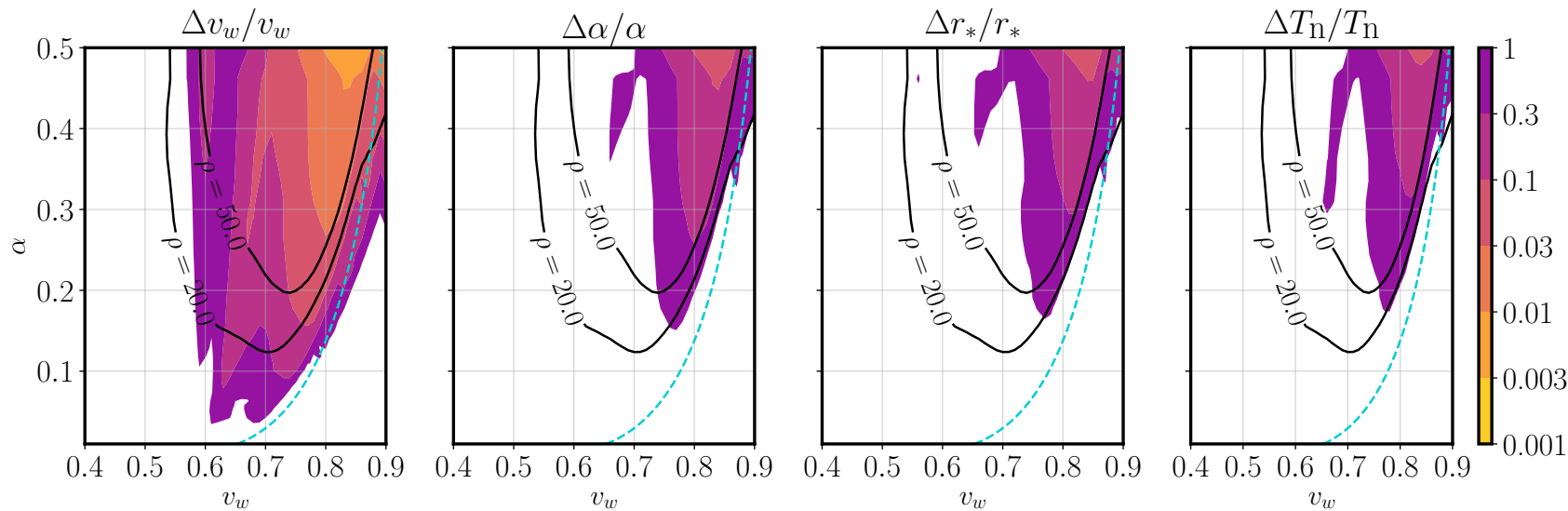
Signal-to-noise ratios (LISA)



- Signal-to-noise ratio r ($t_{\text{obs}} = 4$ years)
- Perfect removal of GB foreground
- “LISA science requirements” instrument noise

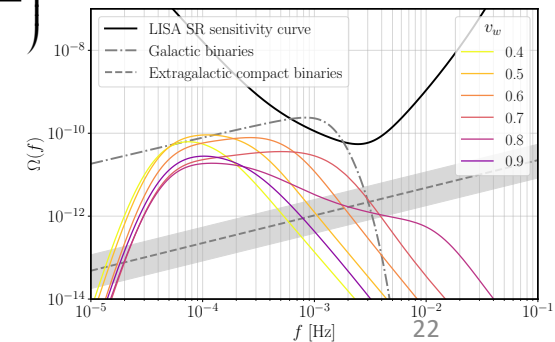
$$\rho^2(\vec{\theta}) = t_{\text{obs}} \int df \left(\frac{\Omega_{\text{gw}}(f; \vec{\theta})}{\Omega_{\text{noise}}(f)} \right)^2$$

Observability of PT parameters: Fisher analysis



- $T_c = 100$ GeV
 - $c_s = 1/\sqrt{3}$
 - $r_* = 0.1$
- Gowling, MH (2021)

- Covariance matrix $\mathbf{C}$ ($t_{\text{obs}} = 4$ years) $C_{ij}^{-1}(\vec{\theta}) = 2t_{\text{obs}} \int df \frac{1}{\Omega_{\text{noise}}^2(f)} \left(\frac{\partial \Omega_{\text{gw}}(f; \vec{\theta})}{\partial \theta_i} \frac{\partial \Omega_{\text{gw}}(f; \vec{\theta})}{\partial \theta_j} \right)$
- Uncertainty from eigenvalues of $\mathbf{C}$
 - “noise” = GB foreground + “LISA science requirements” instrument noise
- **Wall speed** is best determined parameter
 - shape dependence of power spectrum

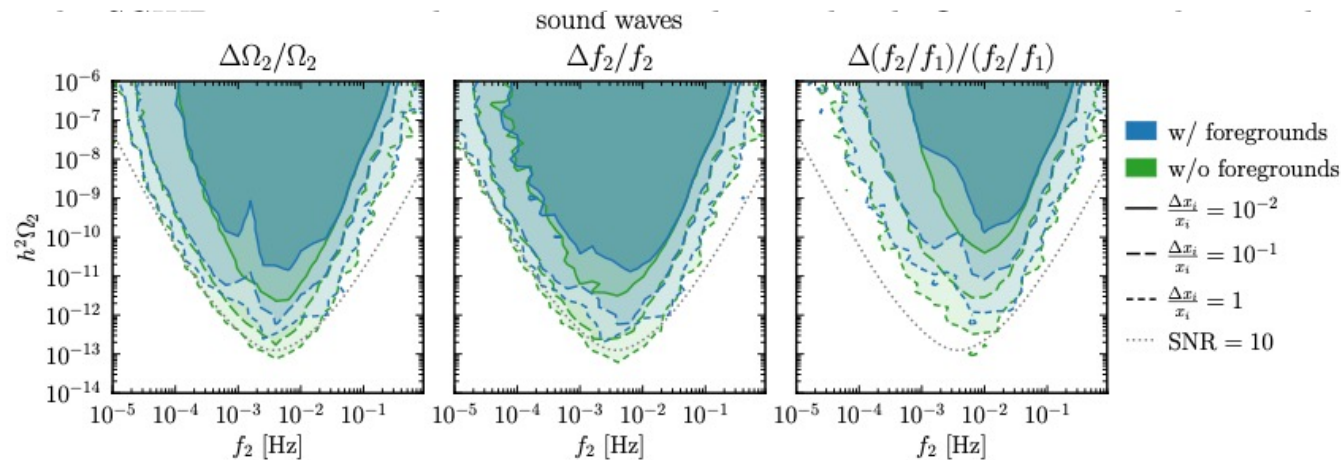


Observability of parametrized spectra

Caprini et al (LISA CWG, 2024)

$$\Omega_{\text{GW}}^{\text{DBPL}}(f, \vec{\theta}_{\text{Cosmo}}) = \Omega_{\text{int}} \times S(f) = \Omega_2 \times S_2(f), \quad (2.8)$$

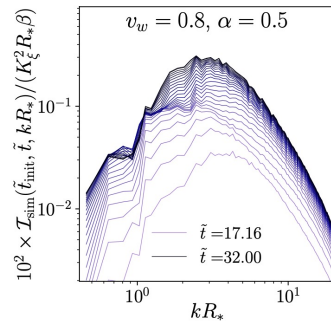
$$S(f) = N \left(\frac{f}{f_1} \right)^{n_1} \left[1 + \left(\frac{f}{f_1} \right)^{a_1} \right]^{\frac{-n_1+n_2}{a_1}} \left[1 + \left(\frac{f}{f_2} \right)^{a_2} \right]^{\frac{-n_2+n_3}{a_2}},$$



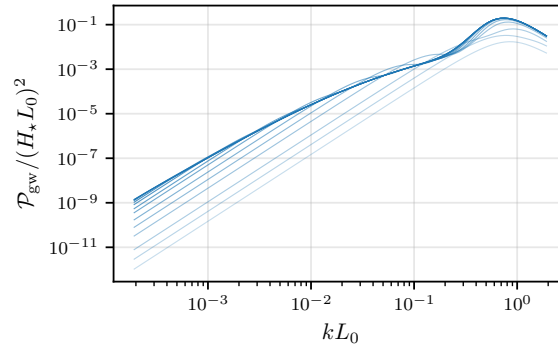
- Marginalised posteriors from Polychord
 - “noise” = GB foreground + “LISA science requirements” instrument noise

Recent advances in hydrodynamics

- Shape of GW spectrum at low k : k^3 and k^1
Sharma, Dahl, Brandenburg, MH 2023; Roper-Pol, Procacci, Caprini 2023
- Effect of sound speed on GW spectrum
Giese et al 2020, Giombi, Dahl, MH 2024 Racco, Poletti 2022
- GWs from strong transitions ($\alpha \sim 1$)
 - $v_{\text{rms}} \sim 0.1 - 0.3$, decaying flow (shock dissipation)
 - peak shape change from shock evolution Dahl et al 2024
 - Decrease of GW growth rate, convergence Caprini et al 2024

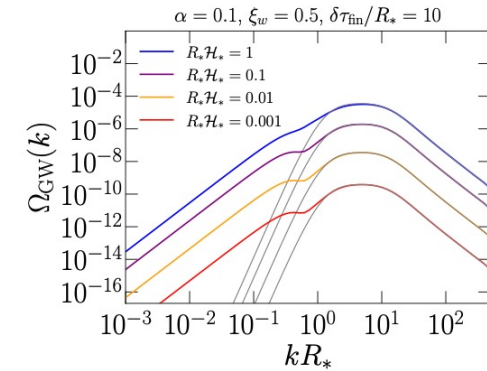


Caprini et al 2024

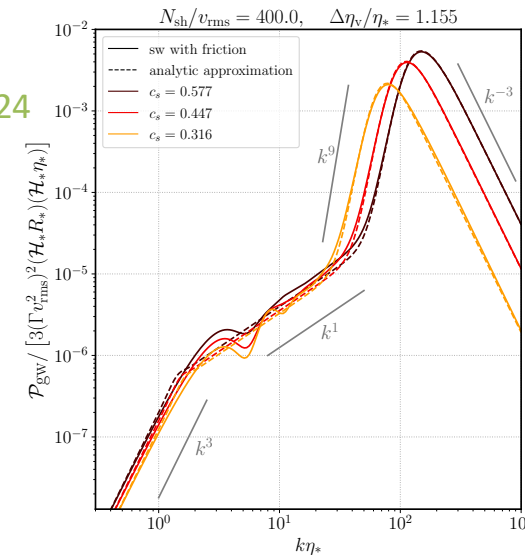


Dahl et al 2024

Gravitational waves ... Mark Hindmarsh



Roper Pol, Procacci, Caprini 2023



Giombi, Dahl, MH 2024

Conclusions

- LISA and other missions will probe physics of electroweak-scale transitions from mid-2030s
 - Measure/constrain phase transition parameters
 - Wall speed likely to be best determined (if signal seen)
 - Parameters from underlying particle physics models
 - Recent progress with equilibrium parameters
 - Wall speed the hardest (non-equilibrium)
- Towards accurate calculations of GW power spectrum from parameters
 - Now: $< 20\%$ accuracy for strength $\alpha \leq 0.3$; speed $0.4 < v_w < 0.9$; bubble separation $r_* \ll 1$; sound speed $c_s = 1/\sqrt{3}$
 - Non-linear evolution ($\alpha \sim 1$) beginning to be investigated
 - Strongly supercooled transitions ($\alpha \gg 1$)?
- Ambition: make GWs as good a probe of the electroweak era as CMB is for the decoupling era

Jinno et al 2019, Lewicki, Vaskonen 2022

