

The importance of being warm (during inflation)

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- Inflation: basics
- Warm inflation
- Contact with observations
- $V = \lambda\phi^4$ is alive!

V CPAN Days
Santiago de Compostela,
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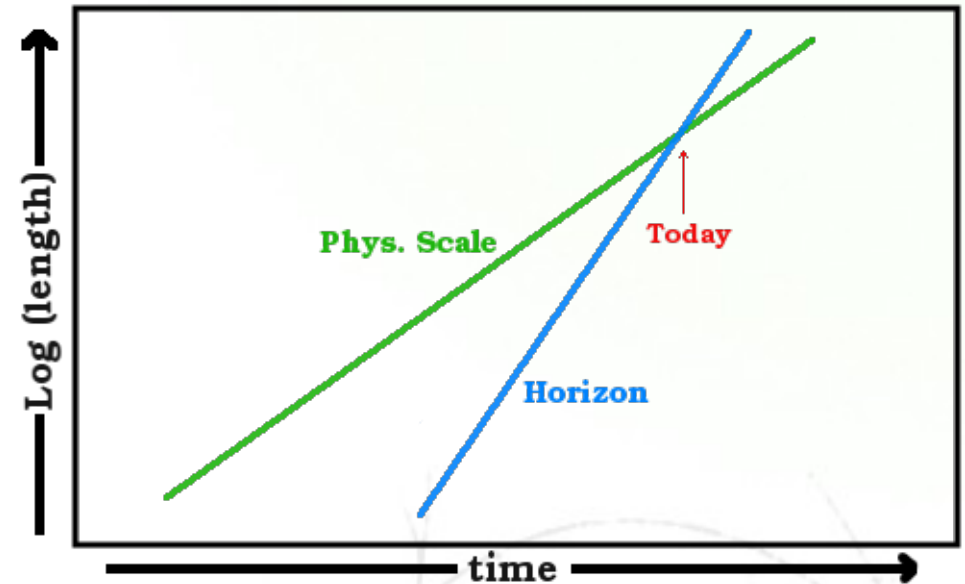
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Cosmological Puzzles

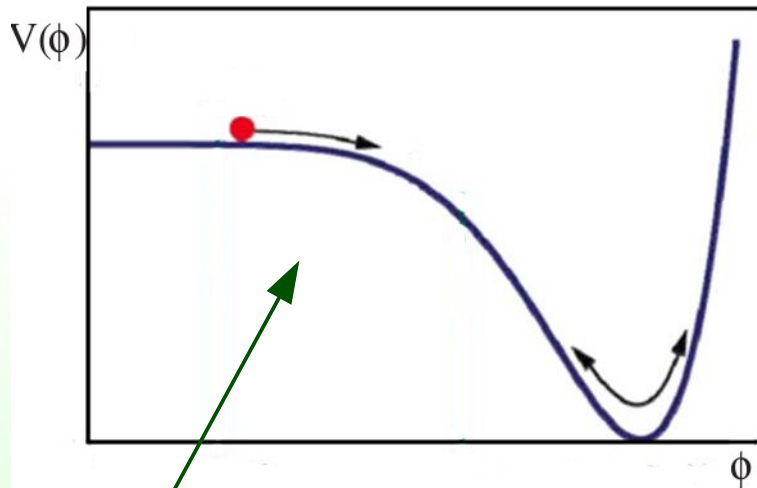
- Horizon problem
Why the Universe is homogeneous?
- Flatness problem
Why it is flat?
 $|\Omega(1s) - 1| < O(10^{-16})$
- Unwanted relics
How to avoid them?
- Primordial density perturbation
What's the origin of structure?



Inflation

$$\ddot{a} \sim -(\rho + 3p)$$

- Constant ✗
- Scalar field ✓



Inflation interactions?

Reheating interactions required!

$$V(\phi) \gg \dot{\phi}^2$$

$$\ddot{\phi} + \underbrace{3H\dot{\phi}}_{\text{friction}} + V_{\phi} = 0$$

Flat potential

Inflaton coupled to other fields $\mathcal{L} \sim g^2 \phi^2 \chi^2$



Equation of motion for the quantum effective action

$$\frac{\delta\Gamma}{\delta\phi(x)} = \ddot{\phi} + 3H\dot{\phi} + \frac{\partial V}{\partial\phi} + \int dx'^4 \Sigma_R(x-x')\delta\phi(x) = 0$$



Adiabatic evolution

$$\begin{aligned}\ddot{\phi} + 3H\dot{\phi} + V_\phi &= -\Upsilon(t)\dot{\phi} \\ \dot{\rho}_r + 4H\rho_r &= +\Upsilon(t)\dot{\phi}^2\end{aligned}$$

Dissipation!

- Extra friction
- Smooth graceful exit
- Observational effects

$$T > H$$

Which interactions?

Thermal corrections to the inflaton mass

$$\mathcal{L} \sim g^2 \phi^2 \chi^2 \longrightarrow m_\phi^2 \sim \begin{cases} g^2 T^2 & m_\chi \ll T \quad \times \\ g^2 e^{-m_\chi/T} & m_\chi \gg T \end{cases}$$

[Berera, Gleiser & Ramos '98; Yokoyama & Linde '98]

Pattern of interactions

[Berera & Ramos '02]

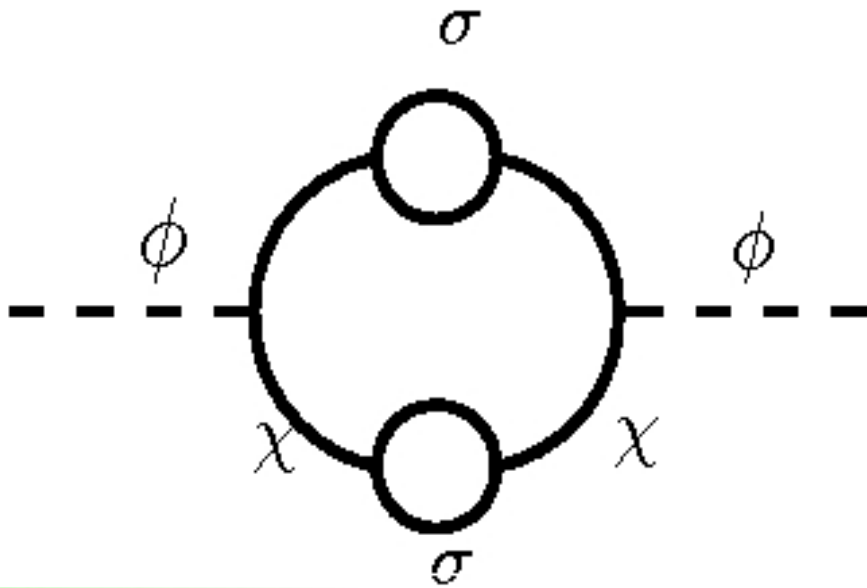
$$\mathcal{L} \sim g^2 \phi^2 \chi^2 + hM\chi\sigma^2 \begin{cases} m_\chi \gg T & \text{heavy catalyst field} \\ m_\sigma \ll T & \text{radiation} \end{cases}$$

Supersymmetric version

$$W = f(\phi) + \frac{g}{2} \Phi X^2 + \frac{h}{2} XY^2 \quad \text{small radiative corrections}$$
$$g^2 N_X, h^2 N_Y \lesssim 1$$

One loop correction to the inflaton self-energy

$$\Upsilon = \frac{4}{T} g^4 \phi^2 \int \frac{d^4 p}{(2\pi)^4} \rho_\chi^2 n_B (1 + n_B)$$



$$\rho_\chi = \frac{4w_p \Gamma_\chi}{(p_0^2 - w_p^2)^2 + 4w_p^2 \Gamma_\chi^2}$$

$$w_p^2 = p^2 + m_\chi^2$$

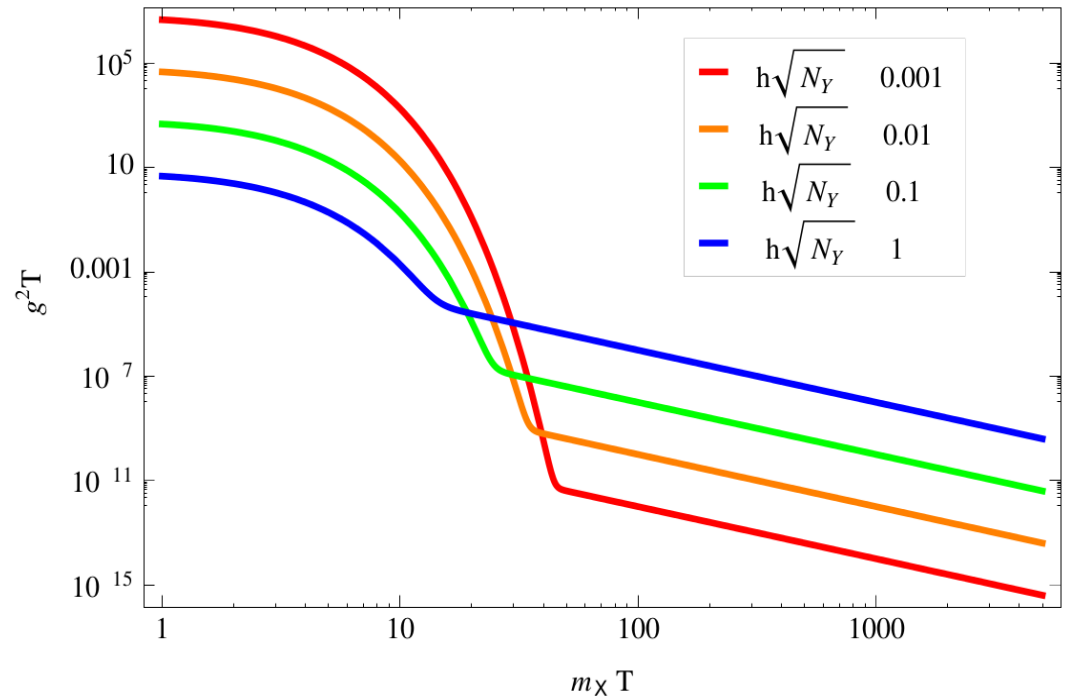
[Moss&Xiong'08; BasteroGil,Berera&Ramos'11]

[BasteroGil,Berera,Ramos&Rosa'12]

- Virtual modes (low-momentum)

$$\Upsilon = C_\phi \frac{T^3}{\phi^2}$$

$$C_\phi = 0.02 h^2 N_X N_Y$$



- On-shell modes (pole)

$$\Upsilon = \frac{32}{\sqrt{2\pi}} \frac{g^2 N_X}{h^2 N_Y} \sqrt{m_\chi T} e^{-m_\chi/T}$$

[BasteroGil, Berera, Ramos & Rosa '12]

Interactions \longrightarrow Energy transfer

$$\begin{aligned}\ddot{\phi} + 3H\dot{\phi} + V_{\phi} &= -\Upsilon\dot{\phi} \\ \dot{\rho}_r + 4H\rho_r &= +\Upsilon\dot{\phi}^2\end{aligned}$$

Conditions

- $V(\phi) \gg \dot{\phi}^2$
- $\rho_{\phi} > \rho_r$
- $T > H$

$$\rho_r = \frac{\pi^2}{30} g_* T^4$$

Strong dissipation

- Smaller scales
- No η -problem
- Steeper potentials

\longrightarrow inflection point inflation

$$V(\phi) \simeq V_0 \left[1 + 3\beta^2 \left(\frac{\phi - \phi_0}{\phi_0} \right) + 4 \left(\frac{\phi - \phi_0}{\phi_0} \right)^3 \right]$$

Warm inflection \longrightarrow no fine tuning

For low-momentum dissipation

strong dissipation $\longrightarrow N_X \gtrsim 10^5 - 10^6$

- Larger GUTs, Branes, extra dimensions?

[BasteroGil, Berera & Rosa '12]

[Matsuda '13; Bhattacharjee & Deshamukhya '13]

- Bulk viscosity effects

[DelCampo, Herrera, Pavón, Villanueva '10]

[BasteroGil, RCB, Ramos & Vicente '12]

- Pole dominated dissipation

$$h^2 N_Y < g^2 N_X \lesssim 1$$

[BasteroGil, Berera, Ramos & Rosa '12]

[Berera&Fang'95; Berera'95;Berera'00;Moss&Berera'04]

$$\delta\ddot{\phi}_k + 3H(1 + Q)\delta\dot{\phi}_k + \frac{k^2}{a^2}\delta\phi_k \simeq \sqrt{2\Upsilon T}a^{-3/2}\xi_k$$

$$Q = \frac{\Upsilon}{3H}$$

Weak dissipation

$$\Delta_{\mathcal{R}}^2 \simeq \left(\frac{H_*}{\phi_*}\right)^2 \left(\frac{H}{2\pi}\right)^2 [1 + 2n_* + \kappa_*]$$

$$r \simeq \frac{16\epsilon_\phi}{1 + 2n_* + \kappa_*}$$

dissipation

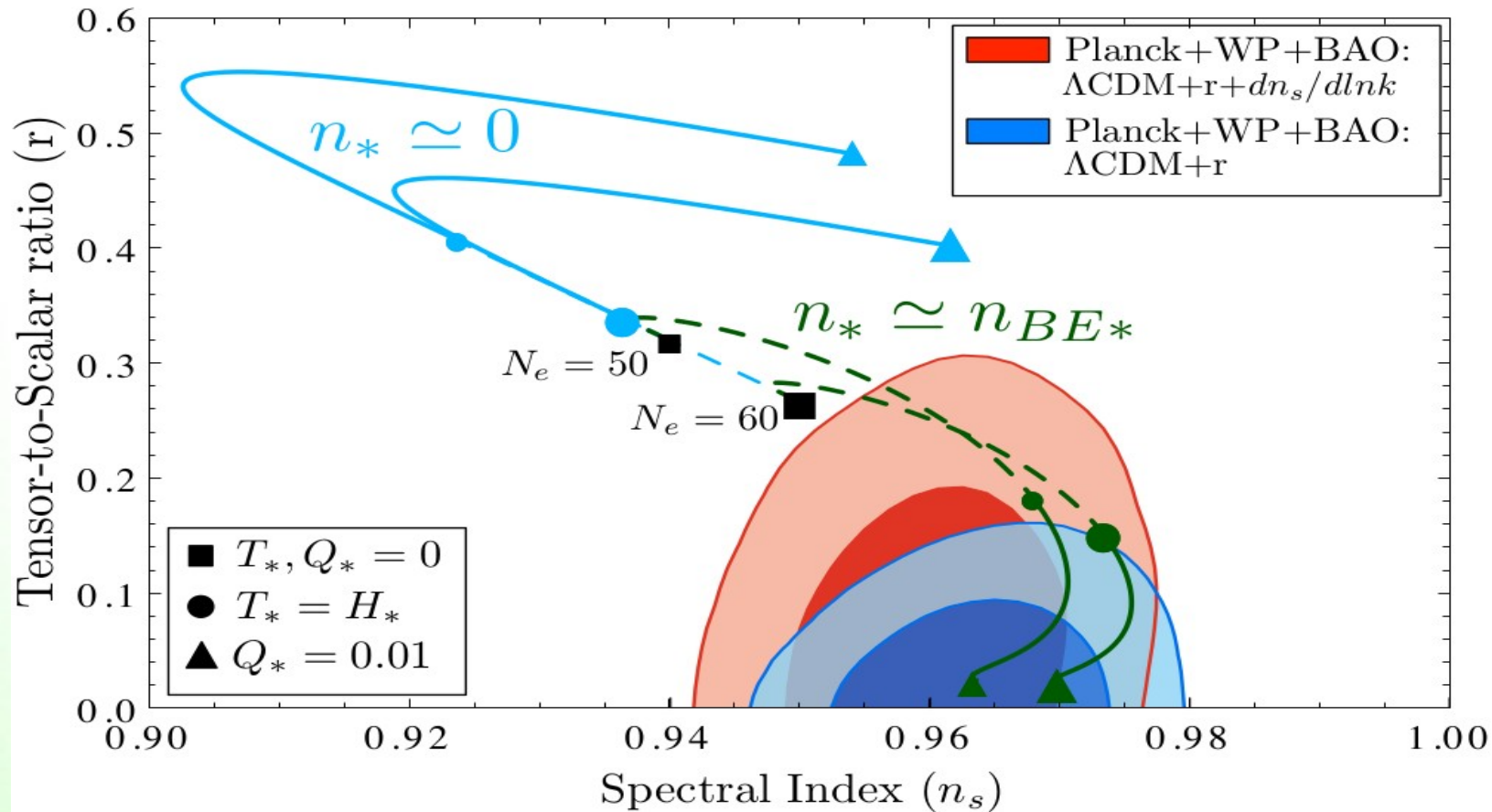
$$\kappa = 2\pi Q \frac{T}{H}$$

inflaton state

$$n_* \begin{cases} \ll 1 & \text{vacuum} \\ (e^{k/aT} - 1)^{-1} & \text{thermal} \end{cases}$$

[Bartrum,BasteroGil,Berera,RC,Ramos&Rosa'13]

Example: quartic potential



$$V = \frac{\lambda}{4} \phi^4$$

$$g_* = 228, 75$$

- $T \ll H$ ill-defined thermal equilibrium
- $n_*, \kappa_* \ll 1$ larger field values than CI
- $n_* \simeq n_{BE}$ strong suppression

$$\chi \rightarrow \sigma\sigma\phi$$

Thermalization

- decays
- inverse decays
- thermal scatterings

[Anisimov, Buchmuller, Drewes, Mendizabal'09]

$$\frac{\Gamma_{\phi^*}}{H_*} \simeq 9(\alpha_h \alpha_g)^{3/2} \left(\frac{1 - n_s}{0.04} \right) \left(\frac{0.01}{r} \right)^{3/2} \left(\frac{0.005}{Q_*} \right)^{1/2}$$

$$\alpha_h = \frac{h^2 N_Y}{4\pi}$$

$$\alpha_g = \frac{g^2 N_X}{4\pi}$$

not too small

Out-of-equilibrium dissipation + B,C,CP violation



baryon asymmetry

[BasteroGil,Berera,Ramos&Rosa'12]

- Baryon-to-entropy ratio depends on inflaton

(no GUT relics or gravitinos) [Bartrum,Berera,&Rosa'12]

- Baryon isocurvature and adiabatic modes anticorrelated

$$B_B \simeq 3(n_s - 1) \simeq -0.12 \quad n_{iso} \simeq (3 - n_s)/2 \simeq 1.02$$

- Relative matter isocurvature spectrum

$$\beta_{iso} \simeq (\Omega_b/\Omega_c)^2 B_B^2 \simeq 4.8 \times 10^{-4} < 0.0087 \quad (\text{Planck})$$

First principles paradigm:

- significant progress in model building
- interesting implications: dynamics and observations
- $V = \frac{\lambda}{4}\phi^4$ in perfect agreement with Planck
- Prospects: viscosities, BSM, model building...

Thank **you!**

