

Flavor and Dark Matter

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Standard Model: A story of success

- $SU(3) \times SU(2) \times U(1)$ gauge symmetry describes all observed interactions in Nature.
- Electroweak symmetry breaking established after Higgs discovery in 2012.
- LHC results agree with the model predictions.
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But... \Rightarrow Several open question remain in SM

We still need new physics!!!

Flavor Problem

All Observed *Flavour transitions* can be accommodated in Yukawa couplings:

$$\mathcal{L}_Y = H \bar{Q}_i Y_{ij}^d d_j + H^* \bar{Q}_i Y_{ij}^u u_j$$

But... \Rightarrow

- a) what is the origin of the Yukawa structures??
- b) why is there a CP-violating phase in CKM??
- c) why couplings are so different???

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Dark Matter

Required by rotational velocity of galaxies, Cosmic Microwave Background, etc.

But... \Rightarrow No massive, stable and weakly interacting candidate present in SM.

Dark matter

- Dark matter masses from sub-eV to TeV (or micro black holes).
- Observed DM abundance requires **weak couplings** through freeze-out, with masses $\mathcal{O}(\Lambda_{EW})$: **WIMPs**.
- Weaker couplings and lighter masses possible through freeze-in mechanism: **FIMPs**
- Must be stable or long-lived at cosmological scales.

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Additional unbroken symmetry
needed for DM stability.

Flavour symmetries

- Very different elements in Yukawa matrices: $y_t \simeq 1$, $y_u \simeq 10^{-5}$
- Expect couplings in a “fundamental” theory $\mathcal{O}(1)$
- Small couplings generated as function of small vevs or loops.
- Froggatt-Nielsen mechanism and flavour symmetry to understand small Yukawa elements.

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Can a single symmetry explain flavor and stabilize DM ??

Example $U(1)_{\text{fl}}$

- SM fermion with flavour charges:

$$\begin{pmatrix} a_{Q_1} & a_{Q_2} & a_{Q_3} \\ a_{u_1} & a_{u_2} & a_{u_3} \\ a_{d_1} & a_{d_2} & a_{d_3} \\ a_{L_1} & a_{L_2} & a_{L_3} \\ a_{e_1} & a_{e_2} & a_{e_3} \end{pmatrix} = \begin{pmatrix} 4 & 2 & 0 \\ 4 & 2 & 0 \\ 4 & 3 & 3 \\ 4 & 3 & 3 \\ 4 & 2 & 0 \end{pmatrix}.$$

flavon S of charge -1

Dark sector fermion charged under $U(1)_{\text{fl}}$ with charge $15/2$

$$\mathcal{L}_{\text{DM}} = \frac{1}{2} \bar{\chi} \left(i \gamma^\mu \partial_\mu \right) \chi - y_\chi \left(\frac{S}{\Lambda} \right)^{2n-1} S \bar{\chi}^c \chi + h.c.$$

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\Rightarrow Stable with small flavor interactions

Coincidence problem

- Dark matter (WIMP) produced via freeze-out of a massive, $O(\text{TeV})$, stable particle with weak couplings.
- Baryon density produced via baryo/leptogenesis in presence of CP violation, baryon number violation and out of equilibrium processes.
- Nevertheless, measured values for baryon and dark matter densities surprisingly similar:

$$\frac{\Omega_{DM} h_0^2}{\Omega_B h_0^2} = 4.99 \pm 0.20$$

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Can baryonic and dark-matter have a common origin???

⇒ Asymmetric Dark Matter

Asymmetric Dark Matter

- New quantum number of Dark Matter shared by visible matter.
- Asymmetry in this number shared between the two sectors through new (higher dimension op.) interactions.
- After efficient annihilation of symmetric part, asymmetries of comparable size in visible and dark sectors.
- For total baryons asymmetry, B , and dark matter X .

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B and X related through chemical potentials

$$\text{If } B \simeq X \Rightarrow m_{\text{DM}} \simeq 5 \text{ GeV}$$

Asymmetric Dark Matter model building

- New dark sector particles with non-vanishing charges to generate observed dark matter.
- New interactions connecting visible and dark sectors.
- Generation of related asymmetries in baryonic and dark sectors.
- Efficient **annihilation** of the symmetric components.
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But...

Can we do this in a complete model without new *ad hoc* particles and interactions??

Super Asymmetric Dark Matter

- **Sfermions** share quantum numbers with standard particles and decay to stable dark matter at low temperatures. ✓
- **Gaugino** interactions connect fermion and sfermions. ✓
- **Baryon/lepton** asymmetry shared by the two sectors. ✓
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Asymmetry in supersymmetric sector washed away by gaugino interactions...

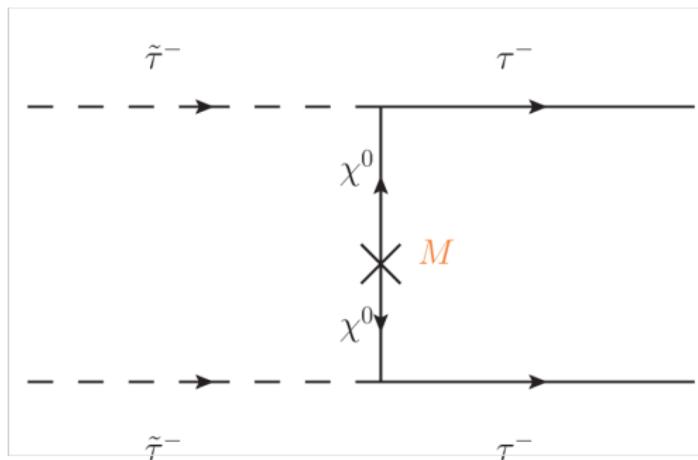
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We must require Dirac gaugino masses to avoid washout of the asymmetry in staus.



Can we decouple this process with large M_χ ??

$$\Gamma(\tilde{\tau}_R + \tilde{\tau}_R \rightarrow \tau\tau) \sim \frac{\tilde{m}^3}{z^{3/2}\pi^2} e^{-z} \frac{4\pi\alpha^2}{M^2} \approx H \sim \frac{(100)^{1/2} m_{\tilde{\tau}}^2}{z^2 M_{Pl}}$$

with $z = m/T$, so that lepton number transfer from the SM to $\tilde{\tau}_R$ s stops at T_b

$$z = \frac{\tilde{m}}{T_b} \approx \ln \left[10^9 \frac{\tilde{m}}{\text{TeV}} \frac{(10\text{TeV})^2}{M^2} \right] \sim 21 + \ln \left[\frac{\tilde{m}}{\text{TeV}} \frac{(10\text{TeV})^2}{M^2} \right]$$

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⇒ No Majorana masses allowed!!

Dirac Asymmetric Dark Matter: to be done...

- Stau decay before nucleosynthesis.
- Low Reheating temperature to avoid gravitino overproduction.
- Allowed Supersymmetric spectrum with $\tilde{\tau}$ NLSP and satisfy LHC bounds.
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keep tuned...