

Falsifying the Universe with Martin since 2014

Frank Deppisch f.deppisch@ucl.ac.uk

University College London

Baryon Asymmetry Generation and Washout

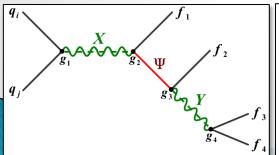
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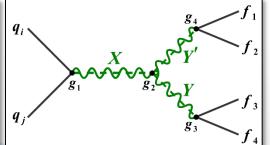


- Generation via heavy neutrino decays
- Competition with LNV washout processes
- Conversion to baryon asymmetry
 - EW sphaleron processes at $T \approx 100 \text{ GeV}$
 - Observed asymmetry

$$\eta_B \equiv \frac{n_B - n_{\bar{B}}}{n_{\gamma}} = (6.20 \pm 0.15) \times 10^{-10}$$

• What if we observe lepton number violating processes at the LHC or in $0\nu\beta\beta$?









Induced Washout

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Compare LHC cross section with lepton number asymmetry washout

$$\frac{\Gamma_W}{H} > 3 \times 10^{-3} \frac{M_P M_X^3}{T^4} \frac{K_1(M_X/T)}{f_{q_1 q_2}(M_X/\sqrt{s})} \times (s \sigma_{\text{LHC}})$$

- Lower limit on total washout rate
 - Neglecting other washout processes

$$\log_{10} \frac{\Gamma_W}{H} > 7 + 0.6 \left(\frac{M_X}{\text{TeV}} - 1\right) + \log_{10} \frac{\sigma_{\text{LHC}}}{\text{fb}}$$

- Observation of LNV @ LHC corresponds to highly effective washout $\Gamma_W/H \gg 1$
 - Excludes Leptogenesis models that generate asymmetry above M_X



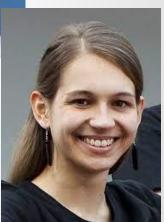


Caveats

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- Cannot exclude scenarios that generate a lepton number asymmetry below observed scale M_x
 - But strong limits still apply
- Asymmetry can be present in one lepton generation only
 - Unambiguous falsification requires observation of LNV in all flavours (or observation of low energy LFV such as $\tau \to e\gamma$)
- Sphalerons only affect LH leptons...What if LNV is observed for RH leptons only?
 - Not an issue as all LH and RH charged fermions are in thermal equilibrium $\approx M_{EW}$
- Symmetry in new sector coupled via hypercharge induces (B L) chemical potential
 (Antaramian, Hall, Rašin '93)





Caveats

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- Cannot exclude scenarios that generate a lepton number asymmetry below observed scale M_x
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Frank F. Deppisch (University Coll. London), Julia Harz (University Coll. London), Martin Hirsch (Valencia U., IFIC) (Dec 16, 2013)

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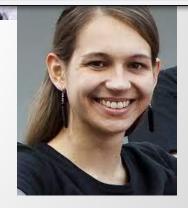
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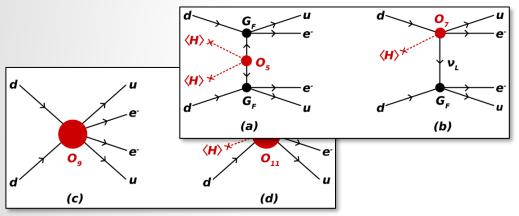
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Washout via $0\nu\beta\beta$ operators

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- Analogous analysis using LNV effective operators of mass dimensions 5, 7, 9, 11
 - 129 Operators



• Matching to $0v\beta\beta$ operators

$$m_e \epsilon_5 = \frac{g^2 v^2}{\Lambda_5}, \ \frac{G_F \epsilon_7}{\sqrt{2}} = \frac{g^3 v}{2\Lambda_7^3}, \ \frac{G_F^2 \epsilon_{\{9,11\}}}{2m_p} = \{\frac{g^4}{\Lambda_9^5}, \frac{g^6 v^2}{\Lambda_{11}^7}\}.$$

$$T_{1/2} = 2.1 \times 10^{25} \text{ y} \cdot (\Lambda_D / \Lambda_D^0)^{2d-8}$$

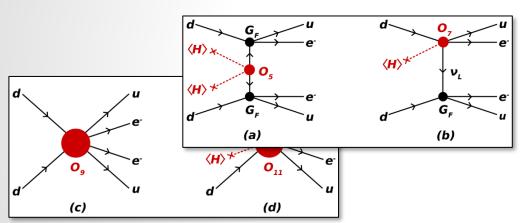


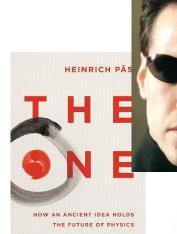


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Matching to 0νββ operators

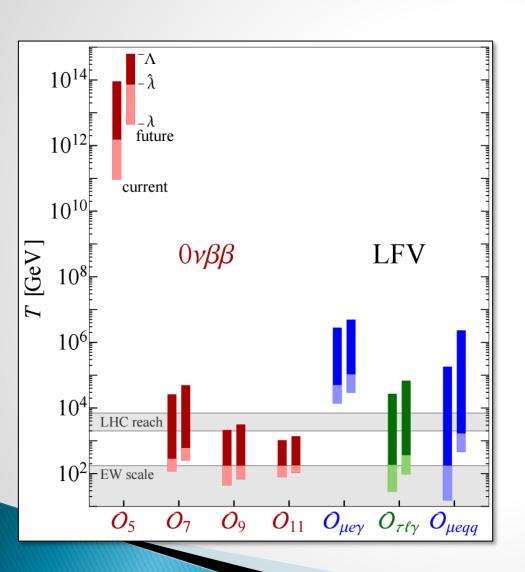
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Observation of 0νββ / LFV



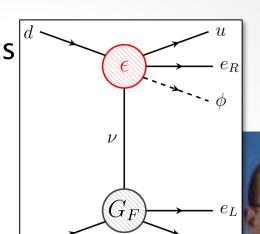




Majorons and RH Currents



Effective RH lepton currents with massless scalar ϕ

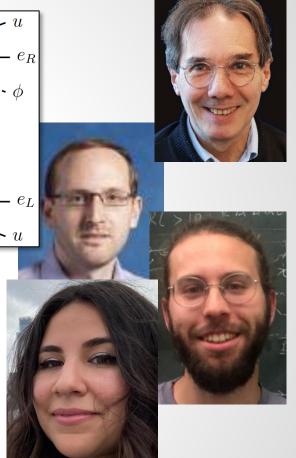


• Giving rise to long-range contribution to $0\nu\beta\beta\phi$ decay

$$\mathcal{M} = \epsilon_{RX}^{\phi} \frac{(G_F \cos \theta_C)^2}{\sqrt{2}m_p} \sum_{N} \int d^3x d^3y \int \frac{d^3q}{2\pi^2 \omega} \phi(\mathbf{y}) e^{i\mathbf{q}(\mathbf{x}-\mathbf{y})}$$

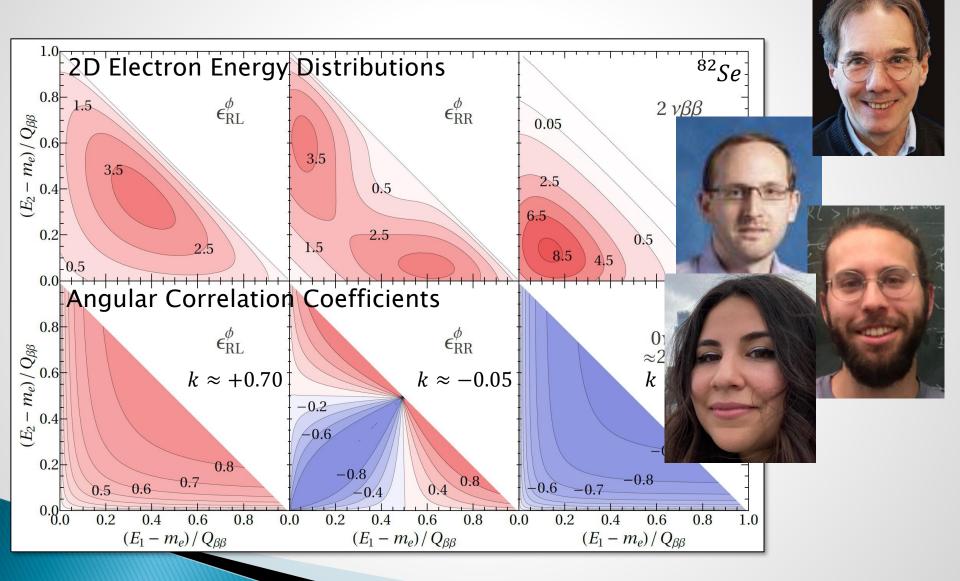
$$\times \left\{ \left[\frac{J_{LX}^{\rho\sigma}(\mathbf{x}, \mathbf{y}) u_{\rho\sigma}^L(E_1\mathbf{x}, E_2\mathbf{y})}{\omega + \mu_N - \frac{1}{2}(E_1 - E_2 - E_{\phi})} - \frac{J_{XL}^{\rho\sigma}(\mathbf{x}, \mathbf{y}) u_{\rho\sigma}^R(E_1\mathbf{x}, E_2\mathbf{y})}{\omega + \mu_N - \frac{1}{2}(E_1 - E_2 + E_{\phi})} \right] - \left[E_1 \leftrightarrow E_2 \right] \right\}$$

- No suppression with ν mass
- Calculation follows long-range η and λ 0νββ modes



Majorons and RH Currents





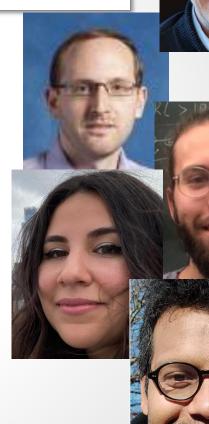
UV Model: LR



Extended Gauge Symmetry

 $G_{LR} = SU(3)_C \times SU(2)_L \times SU(2)_R \times U(1)_X \rightarrow SU(3)_C \times SU(2)_L \times U(1)_Y$

Field	$SU(2)_L$	$SU(2)_R$	B-L	ζ	X	$SU(3)_C$
q_L	2	1	1/3	0	1/3	3
q_R	1	2	1/3	0	1/3	3
ℓ_L	2	1	-1	0	-1	1
ℓ_R	1	2	-1	0	-1	1
$U_{L,R}$	1	1	1/3	+1	4/3	3
$D_{L,R}$	1	1	1/3	-1	-2/3	3
$E_{L,R}$	1	1	-1	-1	-2	1
$N_{L,R}$	1	1	-1	+1	0	1
χ_L	2	1	0	+1	1	1
χ_R	1	2	0	+1	1	1
ϕ	1	1	2	-2	0	1



UV Model: LR or LQ



