

(New) Tau Physics measurements for present and future facilities

Pablo Roig (IFAE)

Third Workshop on Flavour Physics in the LHC era: Theoretical and experimental views

4-6 February 2013

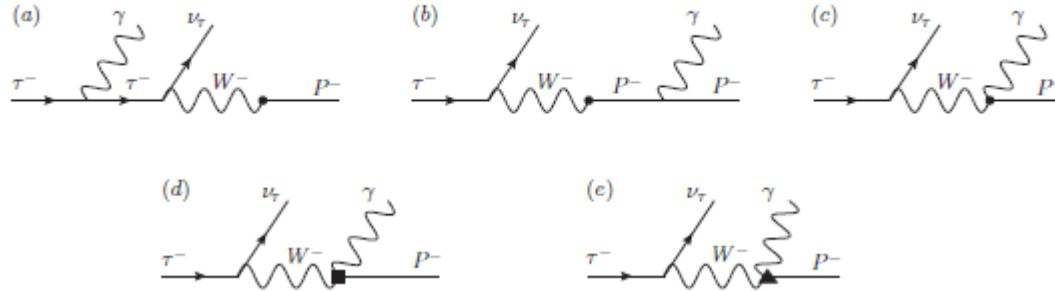
IFIC

Valencia

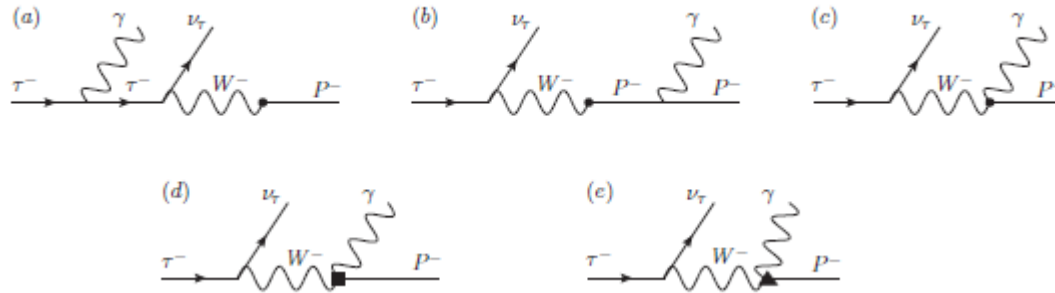
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- $\tau^- \longrightarrow \eta \pi^- \nu_\tau$
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- $\tau^- \longrightarrow \eta K^- \nu_\tau$

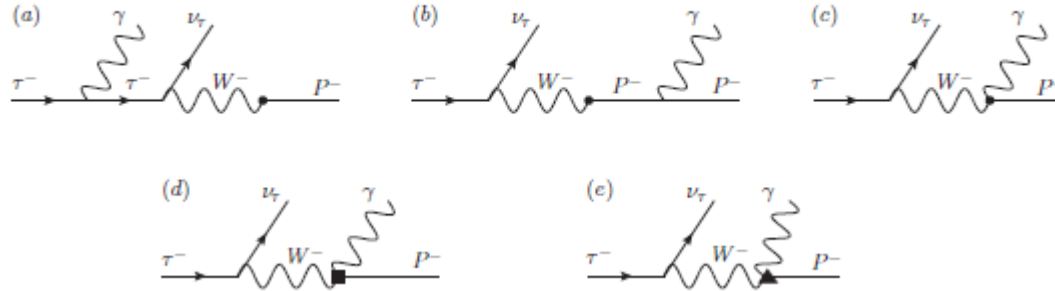
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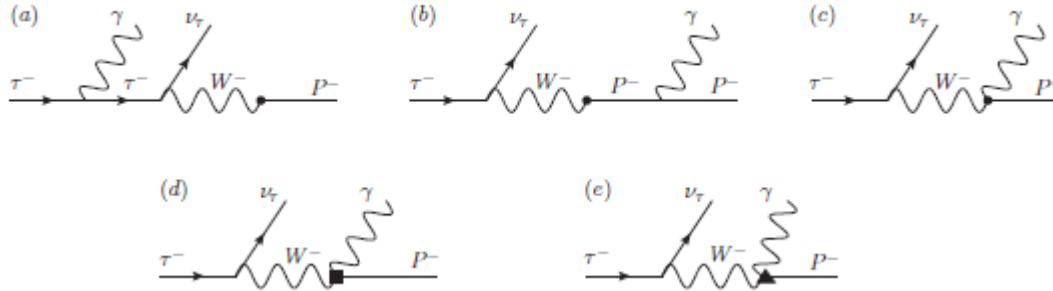
Guo and Roig '10
~1.5% for $E_{\gamma_0} \sim 50$ MeV

$$R_{\tau/\pi} \equiv \frac{\Gamma_{\tau^- \rightarrow \pi^- \nu_\tau}}{\Gamma_{\pi^- \rightarrow \mu^- \bar{\nu}_\mu}}$$

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Decker and Finkemeier '94

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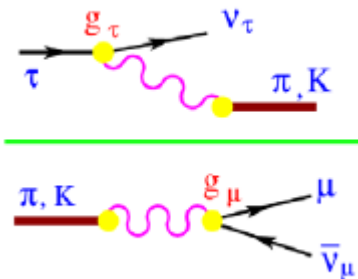
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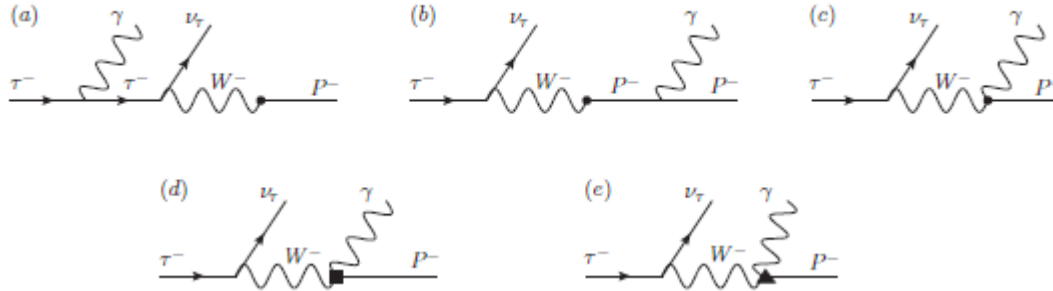
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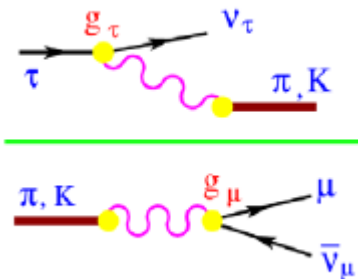
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Not important for charged current universality tests

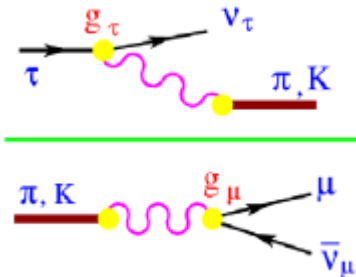
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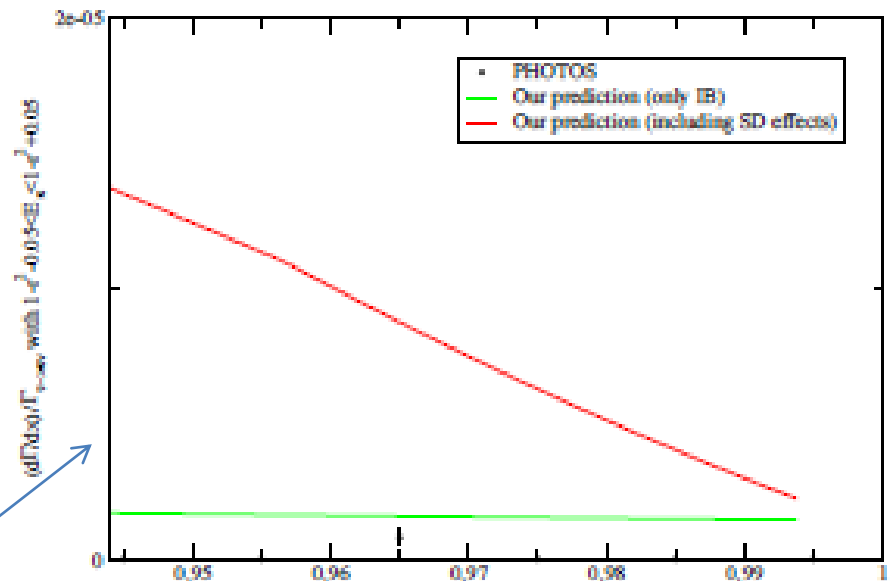
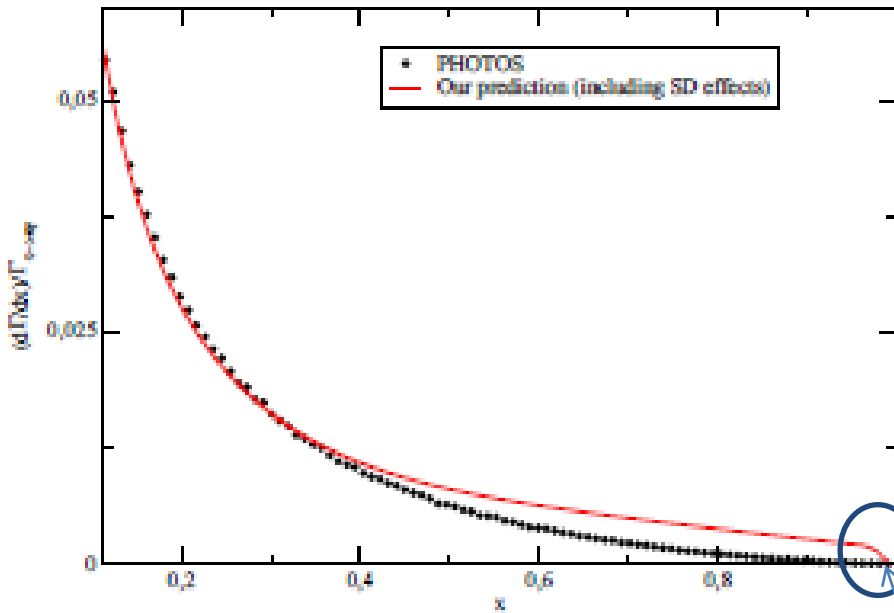
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Main source of background for BaBar and Belle Guo and Roig '10

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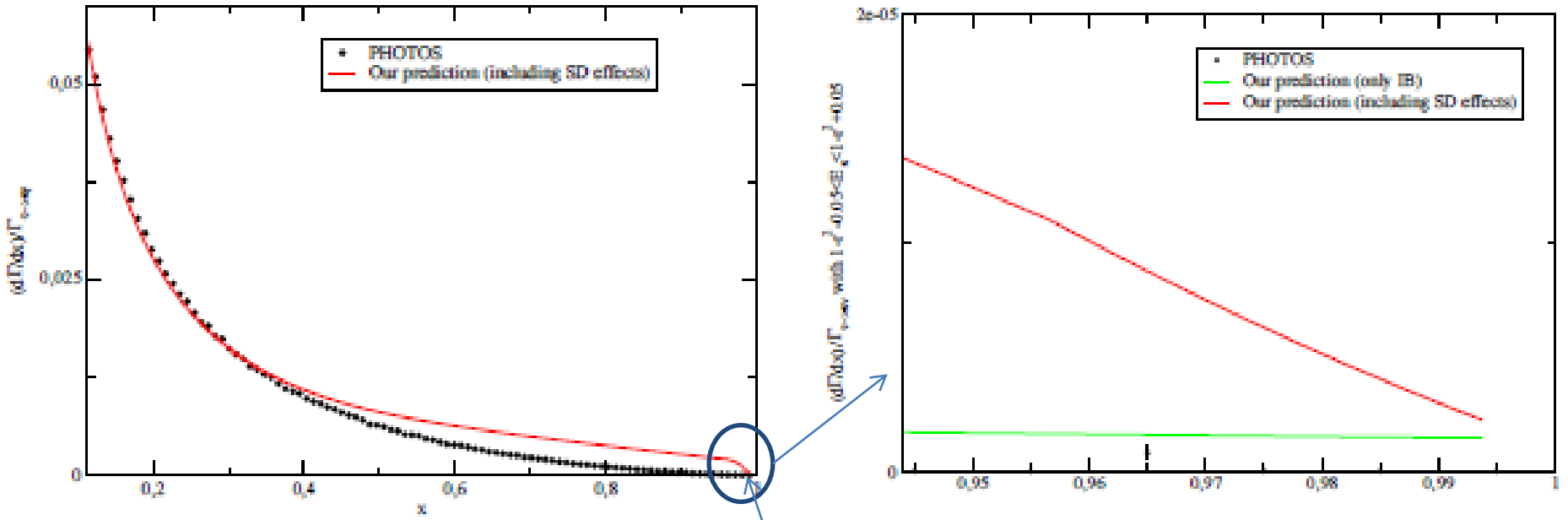


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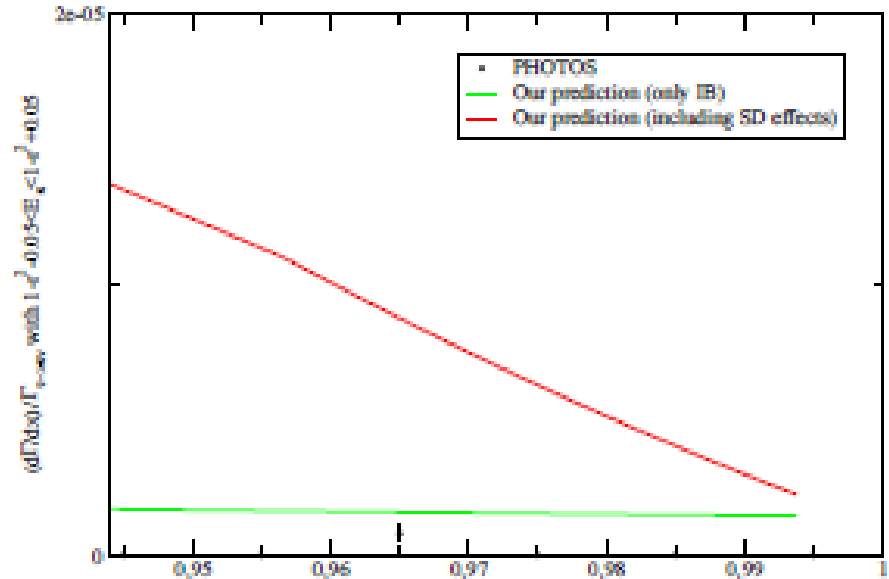
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Guevara, López Castro and Roig

Work in progress on $\tau^- \longrightarrow \pi^- l^+ l^- \nu_\tau$, which can mimic $\tau^- \longrightarrow \mu^- l^+ l^-$

Another interesting measurement!



~~CP~~ in $\tau^- \longrightarrow (K\pi)^- \nu_\tau$

$$\frac{d\Gamma_{K\pi}}{d\sqrt{s}} = \frac{G_F^2 |V_{us}|^2 M_\tau^3}{32\pi^3 s} \left(1 - \frac{s}{M_\tau^2}\right)^2 \left[\left(1 + 2 \frac{s}{M_\tau^2}\right) q_{K\pi}^3 |F_+^{K\pi}(s)|^2 + \frac{3\Delta_{K\pi}^2}{4s} q_{K\pi} |F_0^{K\pi}(s)|^2 \right]$$

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Nice agreement between data and theory

Belle '07

Jamin, Pich, Portolés '06, '08; Boito, Escribano, Jamin '08, '10

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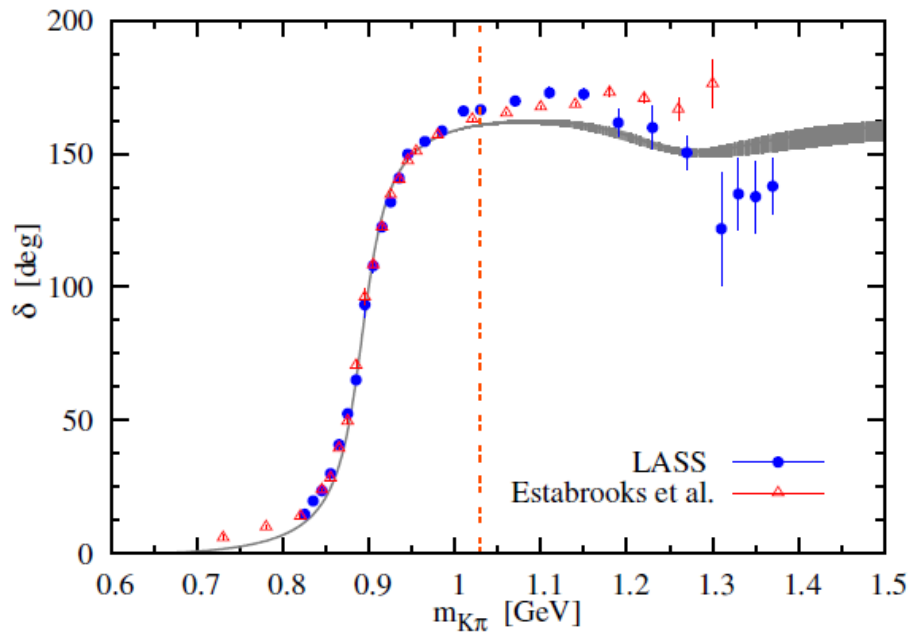
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From coupled channel equations

Jamin, Oller, Pich '06

$\text{Tan}^{-1} [\text{Im}F_+/\text{Re}F_+]$

Figure 2: Phase of the form factor $F_+(s)$ together with experimental results from LASS [46] and Estabrooks *et al.* [47]. The opening of the first inelastic channel, $K^*\pi$, is indicated by the dashed vertical line. The gray band represents the extrema from the fits of Tab. 3.

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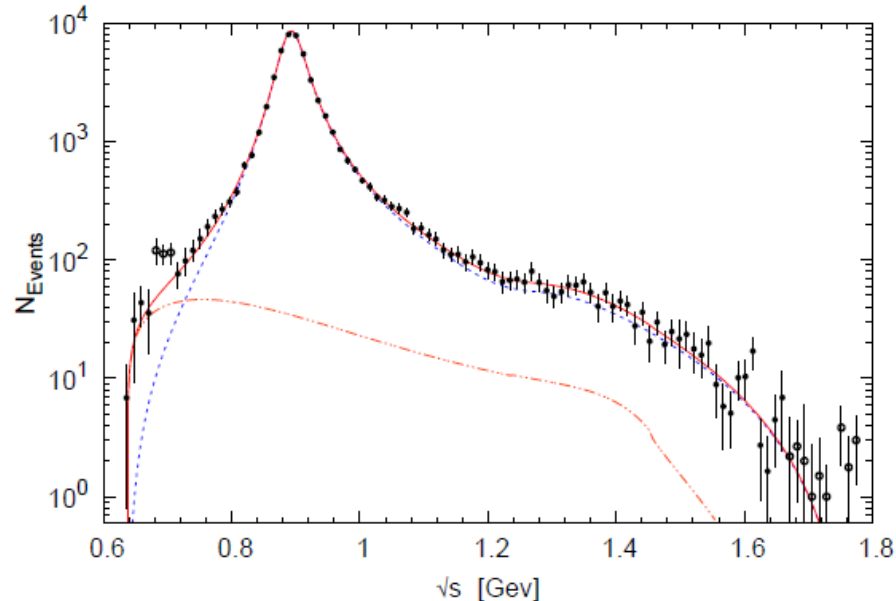


Figure 1: Fit result for the spectrum of $\tau \rightarrow K\pi\nu_\tau$ with $s_{\text{cut}} = 4 \text{ GeV}^2$, third column of Tab. 3. The data are from the Belle collaboration [18]. Points represented with unfilled circles are excluded from the fit (see text in Section 3). The solid red line represents the full fit including contributions from $F_+(s)$ and $F_0(s)$. The scalar contribution alone is represented by the dot-dashed orange line whereas the dashed blue line gives the vector contribution.

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→ However, not all experimental information is condensed in these two 1-D plots. It would be desirable to confront them to **data on** the two spectral functions proportional to **Re[F₊F₀]** and **Im[F₊F₀]**

Kuhn, Mirkes '97

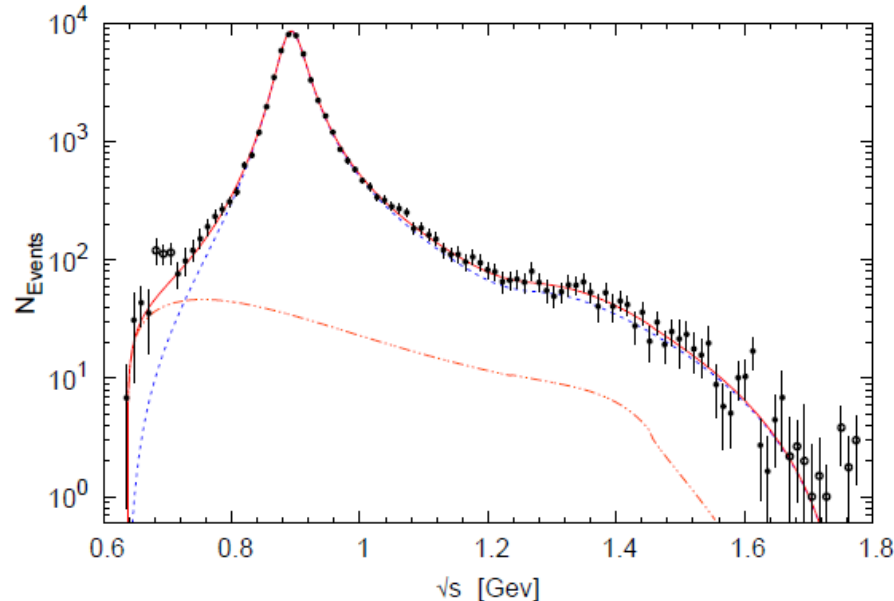


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→ In the mode containing a K^0 there is associated ~~CP~~, but is that small that can't be measured. Besides that, one is interested in measuring ~~CP~~ due to a scalar boson exchange.

BaBar, Belle '11

The CP violating quantity is $\sim \text{Im}[F_+ F_0 |_{\text{H}}]$. This reinforces the case for **measuring $\text{Re}[F_+ F_0]$ and $\text{Im}[F_+ F_0]$** separately to check the SM predictions and reduce the error in searches for CPV in these decays.

and some angular dependence

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TAUOLA (with new hadronic currents Shekhovtsova, Przedzinski, Roig and Was '12 and work in progress with I. Nugent) has all the machinery **to make this kind of analysis**.

$$\tau^- \longrightarrow \eta \pi^- \nu_\tau$$

$G(\eta\pi^-) \neq G(V_\mu)$: Second-class current weak decays. Vanishing in the isospin limit.

Non-zero width because $m_u \neq m_d$ and $e^2 \neq 0$

$$\langle \eta \pi^+ | \bar{u} \gamma^\mu d | 0 \rangle = -\sqrt{2} \left[f_+^{\eta\pi}(s) (p_\eta - p_\pi)^\mu + f_-^{\eta\pi}(s) (p_\eta + p_\pi)^\mu \right]$$

$$f_0^{\eta\pi}(s) = f_+^{\eta\pi}(s) + \frac{s}{\Delta_{\eta\pi}} f_-^{\eta\pi}(s), \quad \Delta_{\eta\pi} = m_\eta^2 - m_\pi^2$$

$$f_+^{\eta\pi}(s) = f_0^{\eta\pi}(s)|_{LO} = \epsilon = \frac{\sqrt{3}(m_d - m_u)}{4(m_s - m_{ud})} \simeq 0.99 \times 10^{-2}$$

(1.21 ± 0.15) × 10⁻² PDG'12

(m_u + m_d)/2

Strong suppression, BR ~ 10⁻⁵

Experimentally, BR ≤ 9.9 × 10⁻⁵ BaBar'11

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Strong suppression, BR $\sim 10^{-5}$ \longrightarrow Sensitivity to New Physics (Bramon, Narison, Pich '87)
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Theoretically, **ongoing works** (Descotes-Genon, Kou, Mousallam; Escribano, González-Solís, Roig) **predict BR $< 10^{-5}$** , roughly a factor of two smaller than earlier estimates

μ anomalous magnetic moment

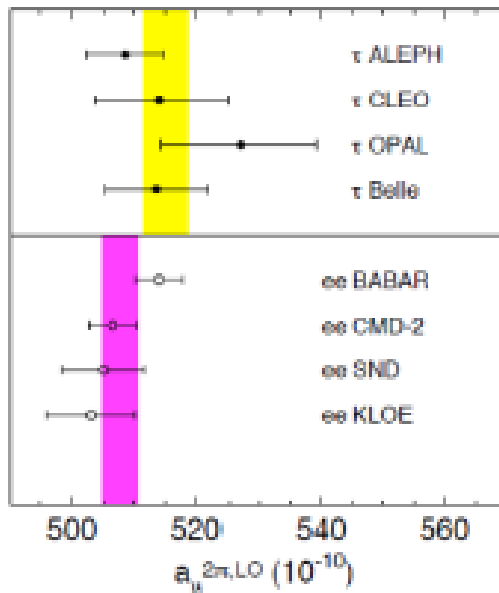
(addressed in previous talk)

$$a_{\mu}^{\text{exp}} - a_{\mu}^{\text{th}} = 701.5 \pm 4.7 (\tau) \quad 692.4 \pm 4.1 (e^+e^-)$$

$(\times 10^{-10})$ 2.3σ 3.6σ

Main contribution comes from $\pi\pi$

Integral from threshold to 1.8 GeV



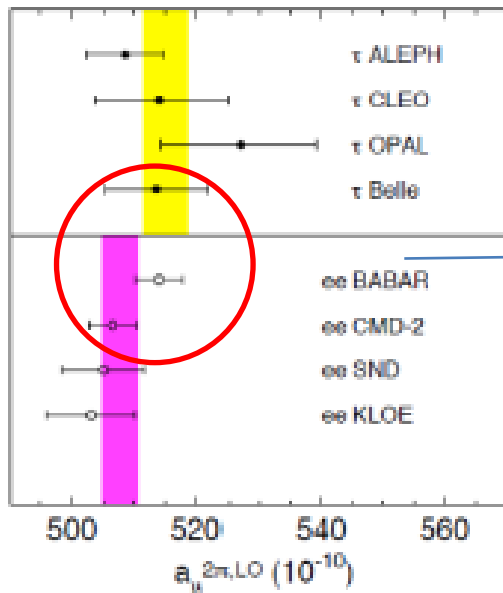
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Last measurements tend to reduce the tension between τ and e^+e^- extractions.

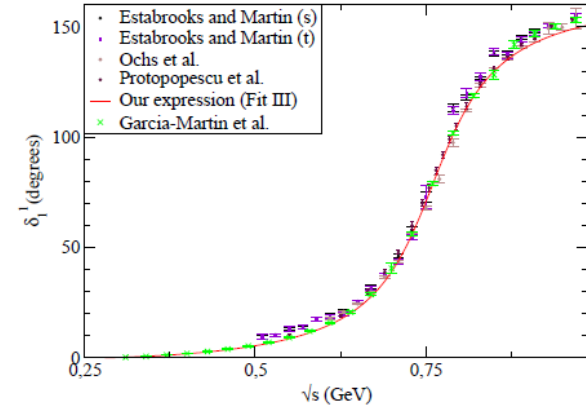
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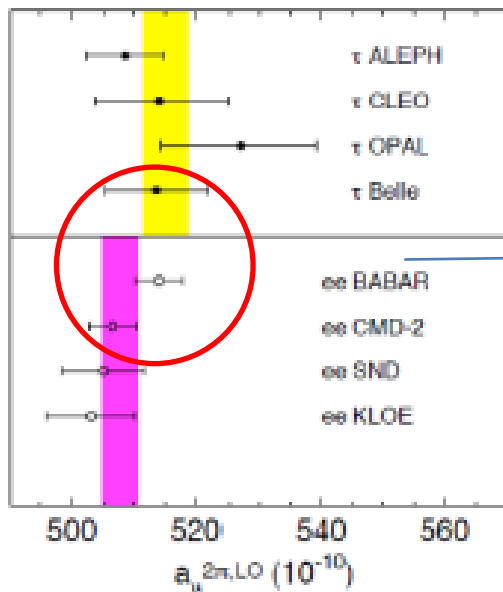
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Dumm and Roig '13
 τ analysis

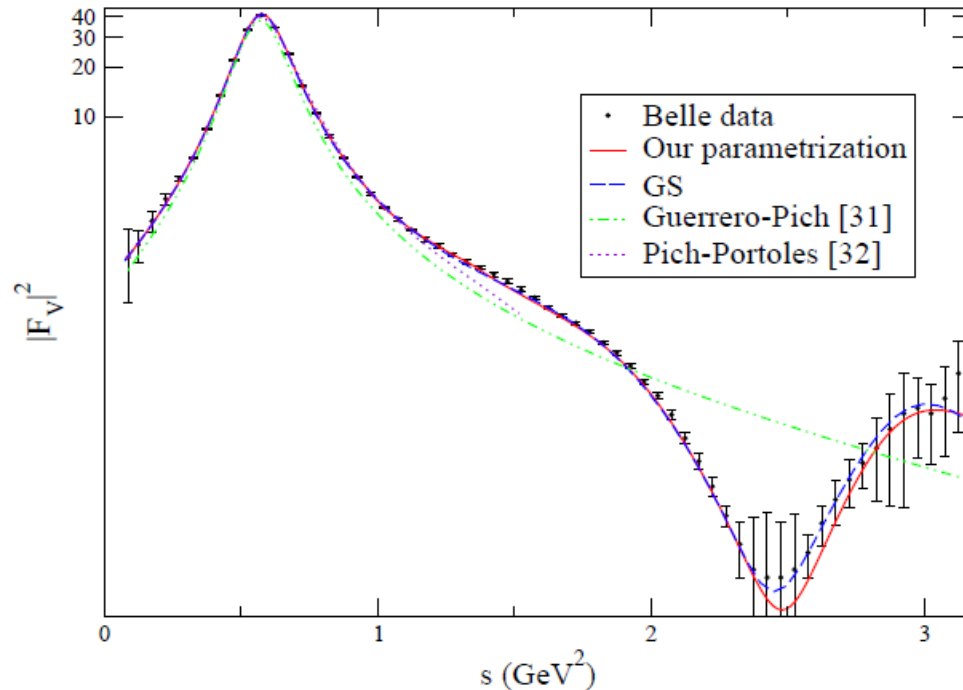


However, isospin corrections need to be understood properly.

Integral from threshold to 1.8 GeV



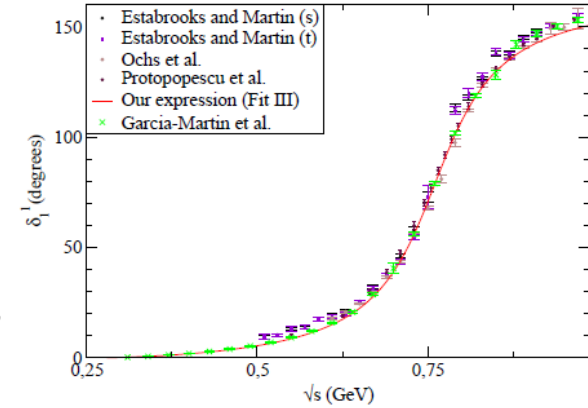
Last measurements tend to reduce the tension between τ and e^+e^- extractions.



μ anomalous magnetic moment

$$a_{\mu}^{\text{exp}} - a_{\mu}^{\text{th}} = 701.5 \pm 4.7 (\tau) \quad 692.4 \pm 4.1 (e^+e^-)$$

$(\times 10^{-10})$
 2.3σ
 3.6σ



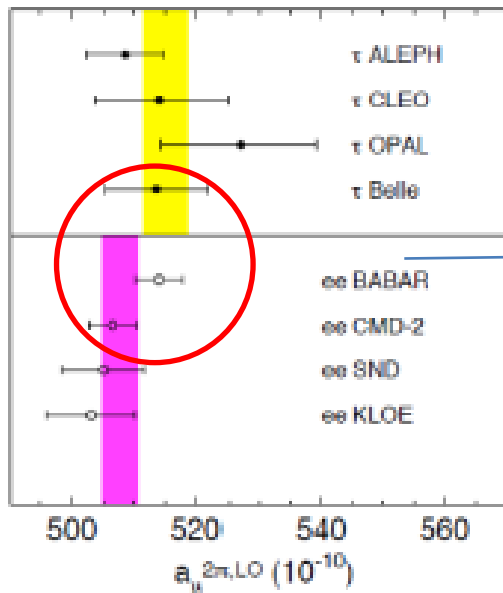
Dumm and Roig '13
 τ analysis

Main contribution comes from $\pi\pi$

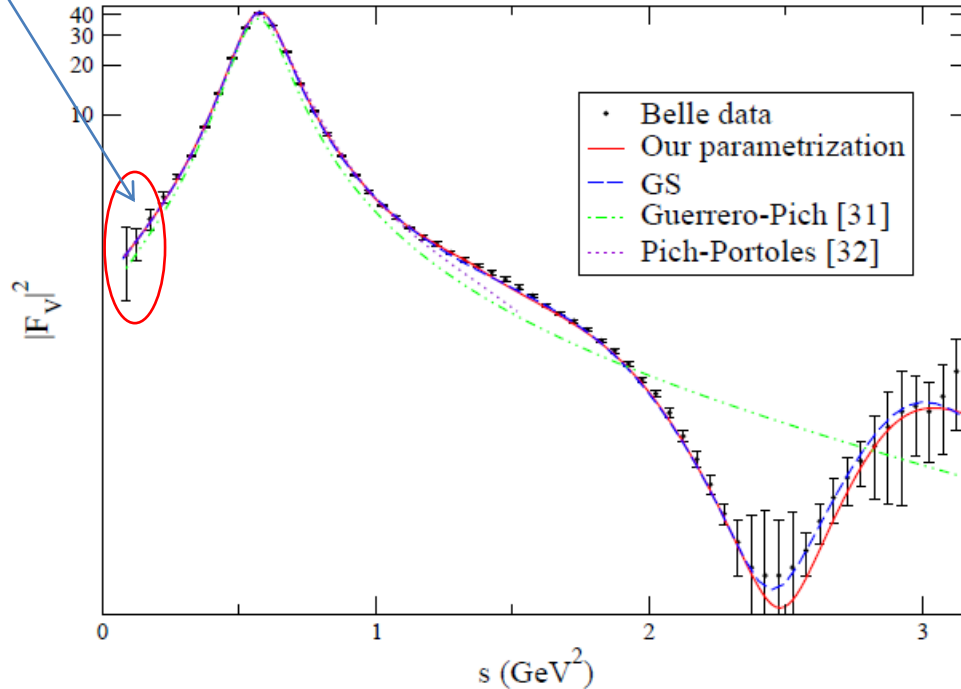
Is it possible to further improve these data?

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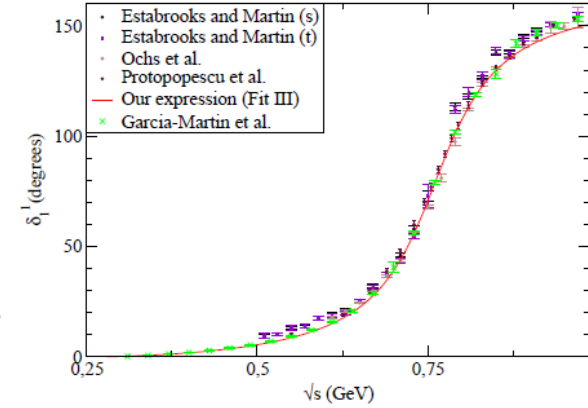
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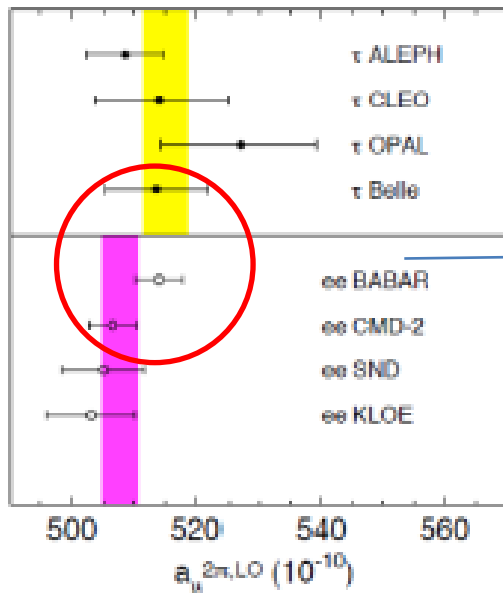
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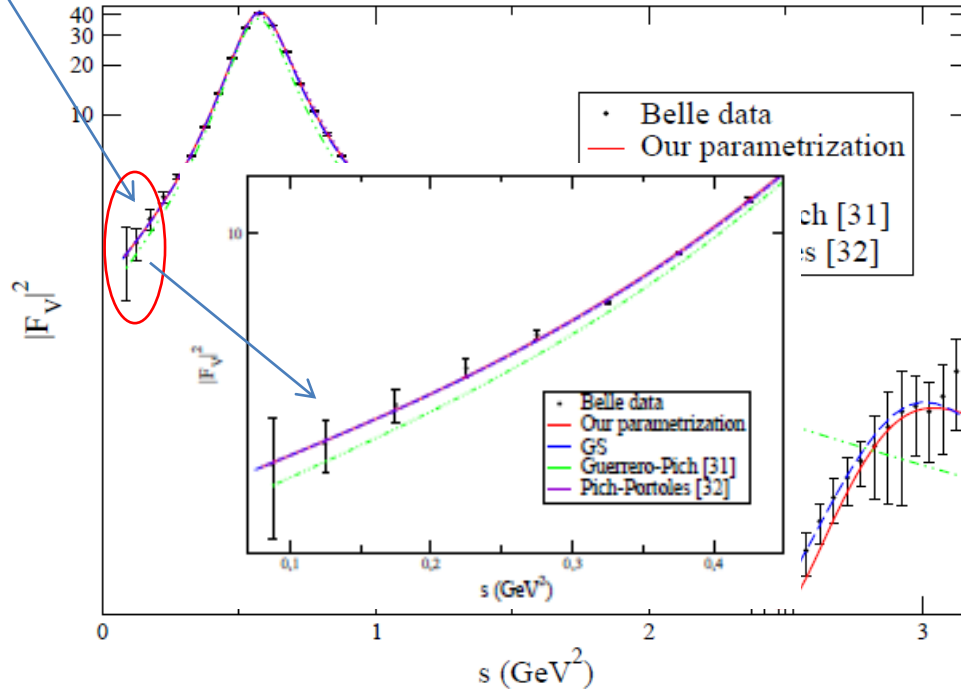
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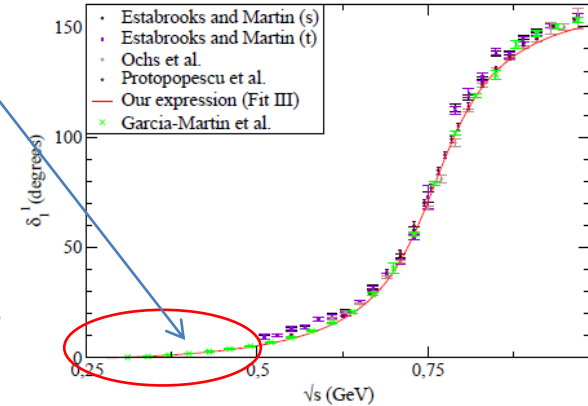
$$a_{\mu}^{\text{exp}} - a_{\mu}^{\text{th}} = 701.5 \pm 4.7 (\tau) \quad (x10^{-10})$$

2.3 σ

$$692.4 \pm 4.1 (e^+e^-)$$

3.6 σ

Is it possible to measure here?



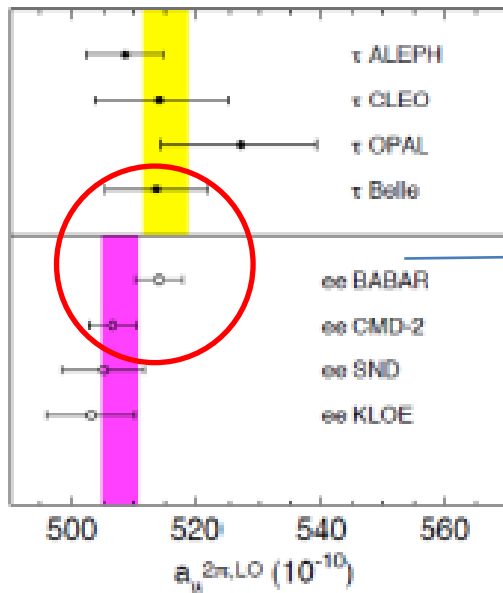
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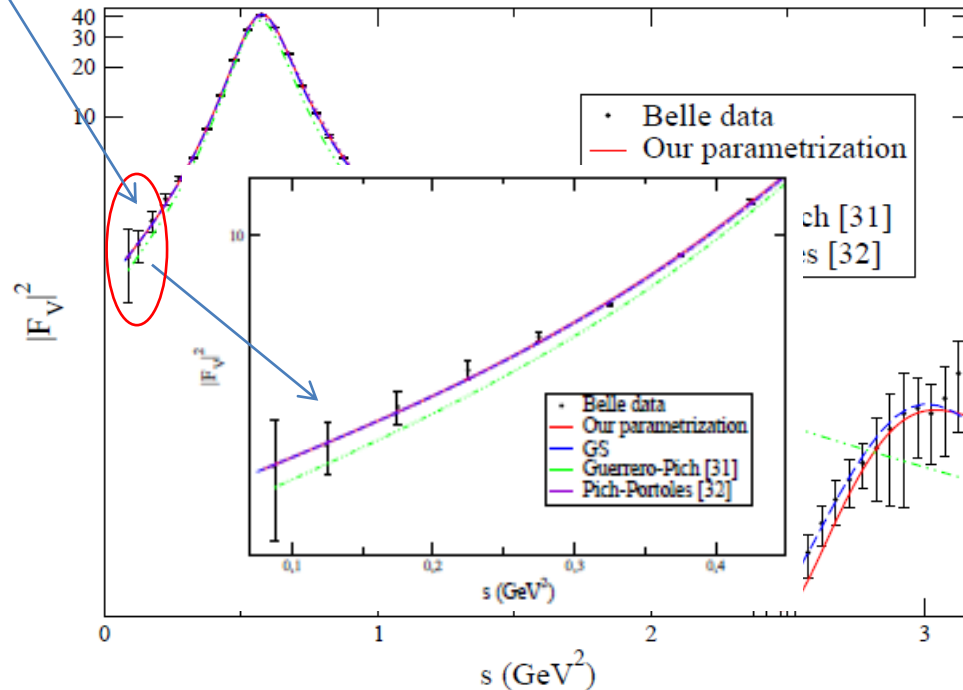
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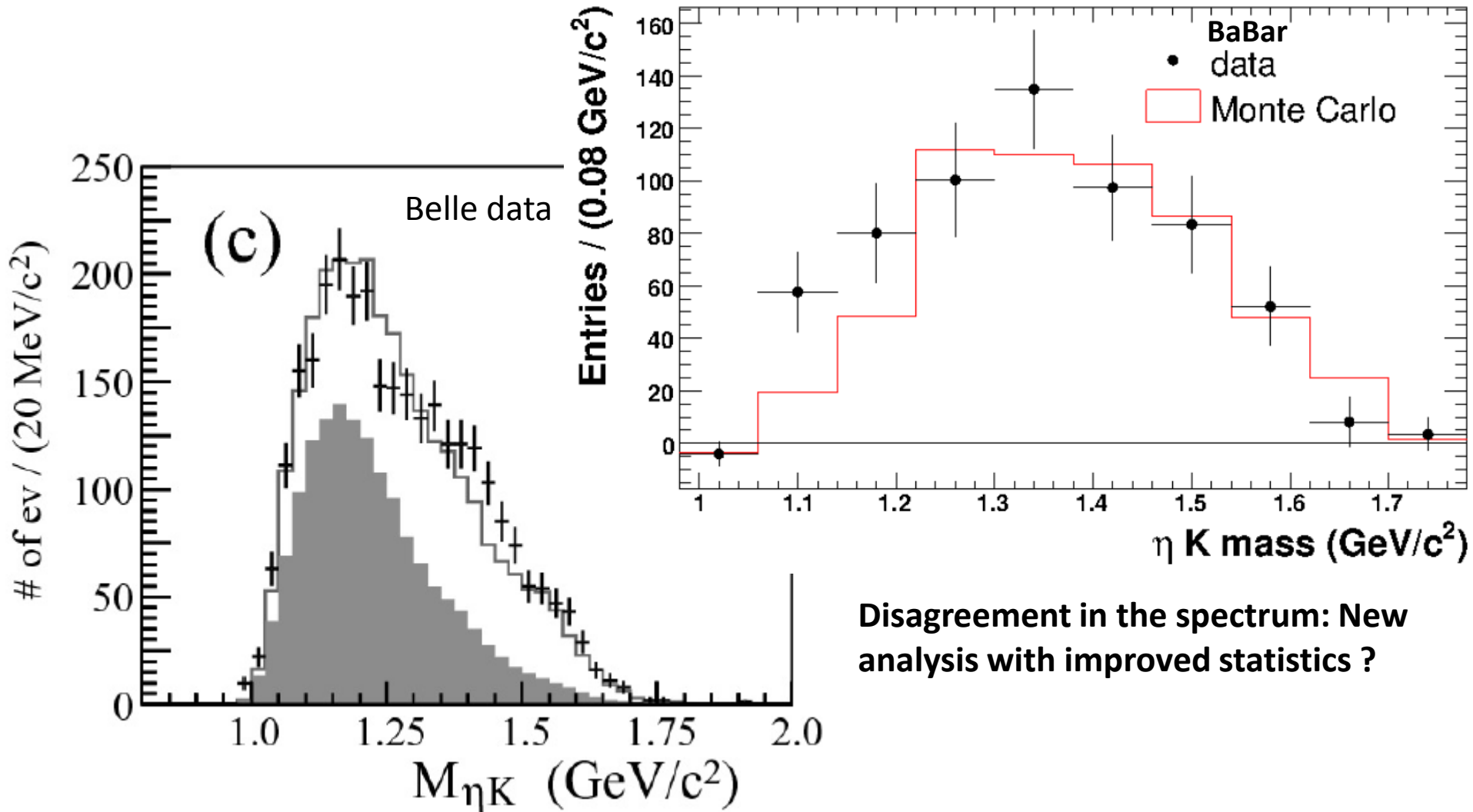
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SUMMARY: INTERESTING MEASUREMENTS

*In addition to what has
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previous talks...*

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- $\tau^- \longrightarrow \eta \text{K}^- \nu_\tau$ **F₀** + Scalar resonance dynamics /nature

SKIPPED SLIDES

Belle analysis:

$$\begin{aligned}
 A_i^{\text{CP}} &= \frac{\iiint_{Q_{1,t}^2}^{Q_{2,t}^2} \cos \beta \cos \psi \left(\frac{d\Gamma_{\tau-}}{d\omega} - \frac{d\Gamma_{\tau+}}{d\omega} \right) d\omega}{\frac{1}{2} \iiint_{Q_{1,t}^2}^{Q_{2,t}^2} \left(\frac{d\Gamma_{\tau-}}{d\omega} + \frac{d\Gamma_{\tau+}}{d\omega} \right) d\omega} \\
 &\simeq \langle \cos \beta \cos \psi \rangle_{\tau-}^i - \langle \cos \beta \cos \psi \rangle_{\tau+}^i \qquad d\omega = dQ^2 d\cos \theta d\cos \beta
 \end{aligned}$$

(see [8]). The angle β is defined by $\cos \beta = \vec{n}_L \cdot \hat{q}_1$ where $\hat{q}_1 = \vec{q}_1/|\vec{q}_1|$ is the direction of the K_S^0 and \vec{n}_L is the direction of the e^+e^- center of mass (CM) system, both observed in the hadronic rest frame. The azimuthal angle α is not observable in this experiment and has to be integrated over. The variable θ is the angle between the direction opposite to the direction of the CM system and the direction of the hadronic system in the τ rest frame.

where ψ denotes the angle between the direction of the CM frame and the direction of the τ as seen from the hadronic rest frame and can be calculated as

$$\cos \psi = \frac{x(m_\tau^2 + Q^2) - 2Q^2}{(m_\tau^2 - Q^2)\sqrt{x^2 - 4Q^2/s}}. \qquad (5)$$

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 \end{aligned}$$

$K_S^0 \pi^\pm \nu_\tau$ decays. It should be noted that CPV in K^0 decays leads to a small SM CP asymmetry of $O(10^{-3})$ in the rates of this τ decay mode [6, 7]. This asymmetry is just below our experimental sensitivity. Here the focus will be on CPV that could arise from a charged scalar boson exchange [8], e.g., a charged Higgs boson. This type of CPV cannot be observed from measurement of τ^\pm decay rates. However, it can be detected as a difference in the τ^\pm decay angular distributions and is accessible without requiring information about the τ polarization or the determination of the τ rest frame. Limits for the

BaBar analysis: $(-0.36 \pm 0.23 \pm 0.11)\%$

$(0.36 \pm 0.01)\%$

SM

$$A_Q = \frac{\Gamma(\tau^+ \rightarrow \pi^+ K_S^0 \bar{\nu}_\tau) - \Gamma(\tau^- \rightarrow \pi^- K_S^0 \nu_\tau)}{\Gamma(\tau^+ \rightarrow \pi^+ K_S^0 \bar{\nu}_\tau) + \Gamma(\tau^- \rightarrow \pi^- K_S^0 \nu_\tau)}$$

2.8 standard deviations