

High-Precision Mass Measurements of Atomic (Exotic) Nuclei

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Outline

1. Overview

- Importance of the mass
- Physics goals & mass uncertainty
- Evolution in mass uncertainty

2. Mass-spectrometry techniques for exotic nuclei

- Time-of-flight measurements
- Frequency measurements

3. The Penning trap

- Fundamentals
- Techniques for exotic nuclei
- Existing facilities & physics survey
- Mass uncertainty & half-life

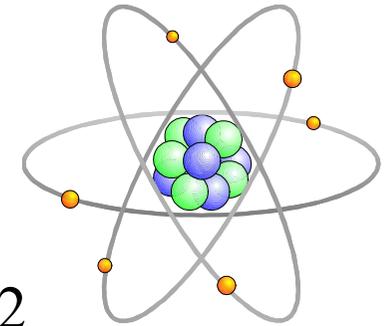
4. Outlook: mass measurements on exotic nuclei in the future

Overview

Importance of the mass

The nuclear binding energy (B)

The nuclear binding energy reflects all forces acting in a nucleus



$$B(N, Z) = \left\{ Zm_p + Nm_n - \underbrace{M(N, Z)}_{\text{Nuclear mass}} \right\} \cdot c^2$$

$$B(N, Z) = \left\{ Zm_p + Nm_n - \left[\underbrace{m(^A X)}_{\text{Atomic mass}} - Zm_e - \sum_i B_i^e \right] \right\} \cdot c^2$$

$$ME(N, Z) = m(^A X) - A \cdot u$$

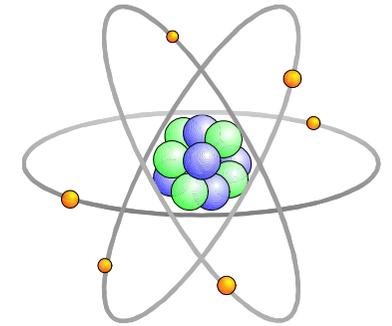
Mass excess (keV)

$$1 \text{ keV} \rightarrow \frac{\delta m}{m} \approx 10^{-8}$$

Importance of the mass

Nucleon separation energies (S_n , S_p , S_{2n} , S_{2p})

Nuclear Structure & astrophysics



$$S_{2n} = B(N, Z) - B(N - 2, Z)$$

$$S_{2p} = B(N, Z) - B(N, Z - 2)$$

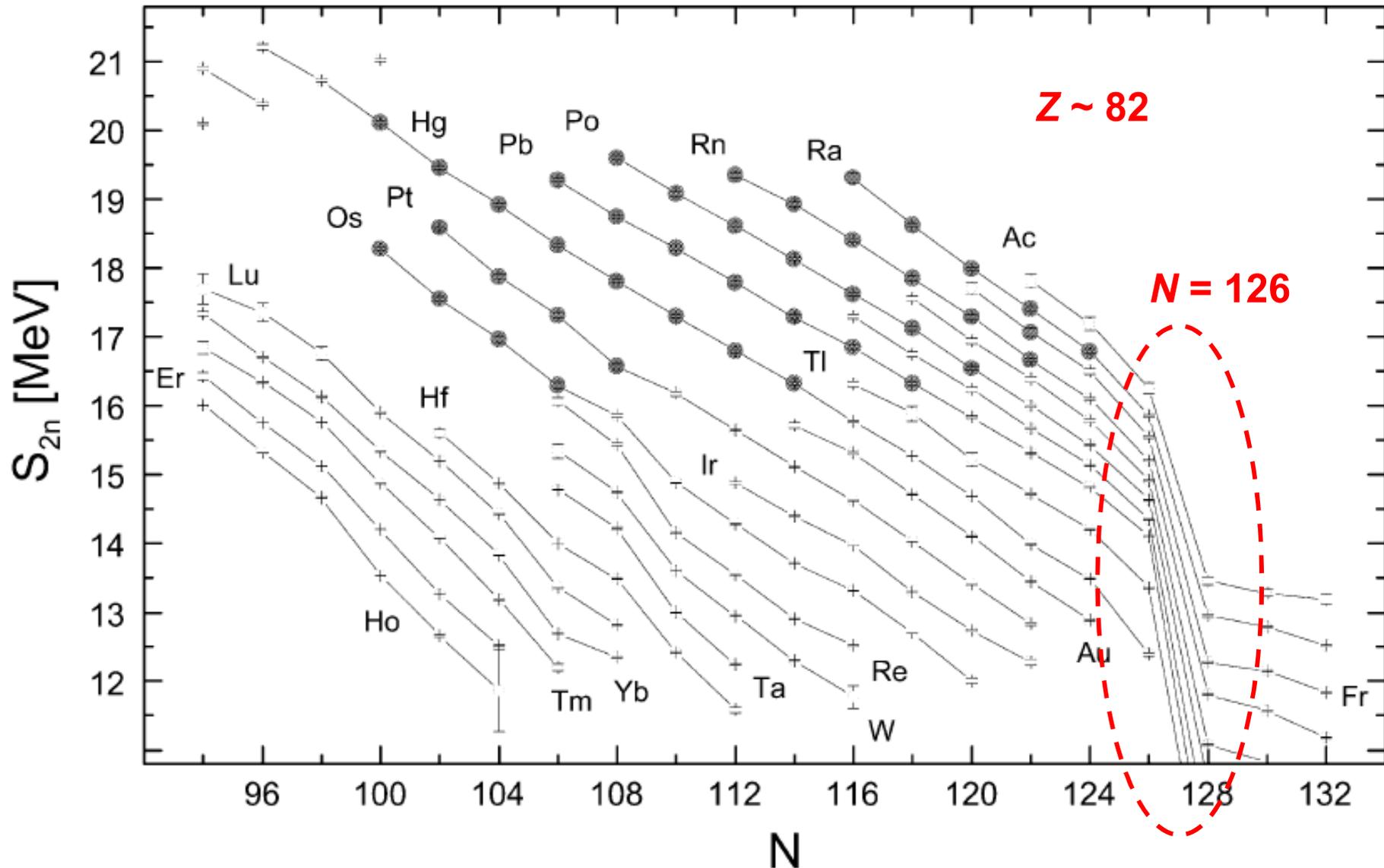
$$S_n = B(N, Z) - B(N - 1, Z)$$

$$S_p = B(N, Z) - B(N, Z - 1)$$

Other quantities are also important for nuclear structure, all of them obtained from nuclear binding energies

Importance of the mass

Nucleon separation energies (S_n , S_p , S_{2n} , S_{2p})



S. Schwarz *et al.*, Nucl. Phys. A 693 (2001) 533

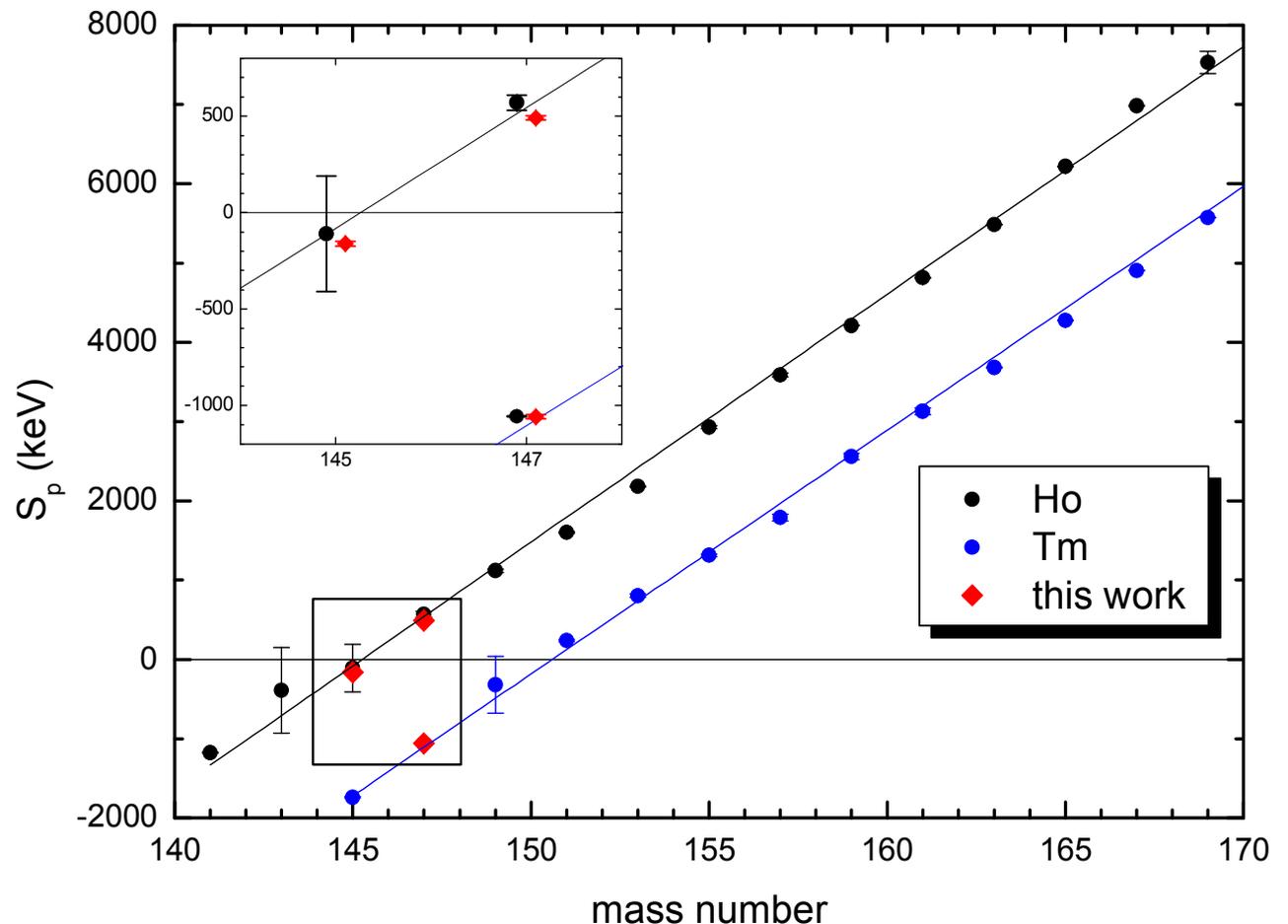
Importance of the mass

Nucleon separation energies (S_n , S_p , S_{2n} , S_{2p})

Location of the proton-drip line

146 69 Tm 77	147 69 Tm 78	148 69 Tm 79	149 69 Tm 80	150 69 Tm 81	151 69 Tm 82
145 68 Er 77	146 68 Er 78	147 68 Er 79	148 68 Er 80	149 68 Er 81	150 68 Er 82
144 67 Ho 77	145 67 Ho 78	146 67 Ho 79	147 67 Ho 80	148 67 Ho 81	149 67 Ho 82
143 66 Dy 77	144 66 Dy 78	145 66 Dy 79	146 66 Dy 80	147 66 Dy 81	148 66 Dy 82
142 65 Tb 77	143 65 Tb 78	144 65 Tb 79	145 65 Tb 80	146 65 Tb 81	147 65 Tb 82

One-proton separation energy for odd-A (even-N) isotopes of Ho (Z=67) and Tm (Z=69):



C. Rauth *et al.*, PRL 100 (2008) 012501

Importance of the mass Astrophysics

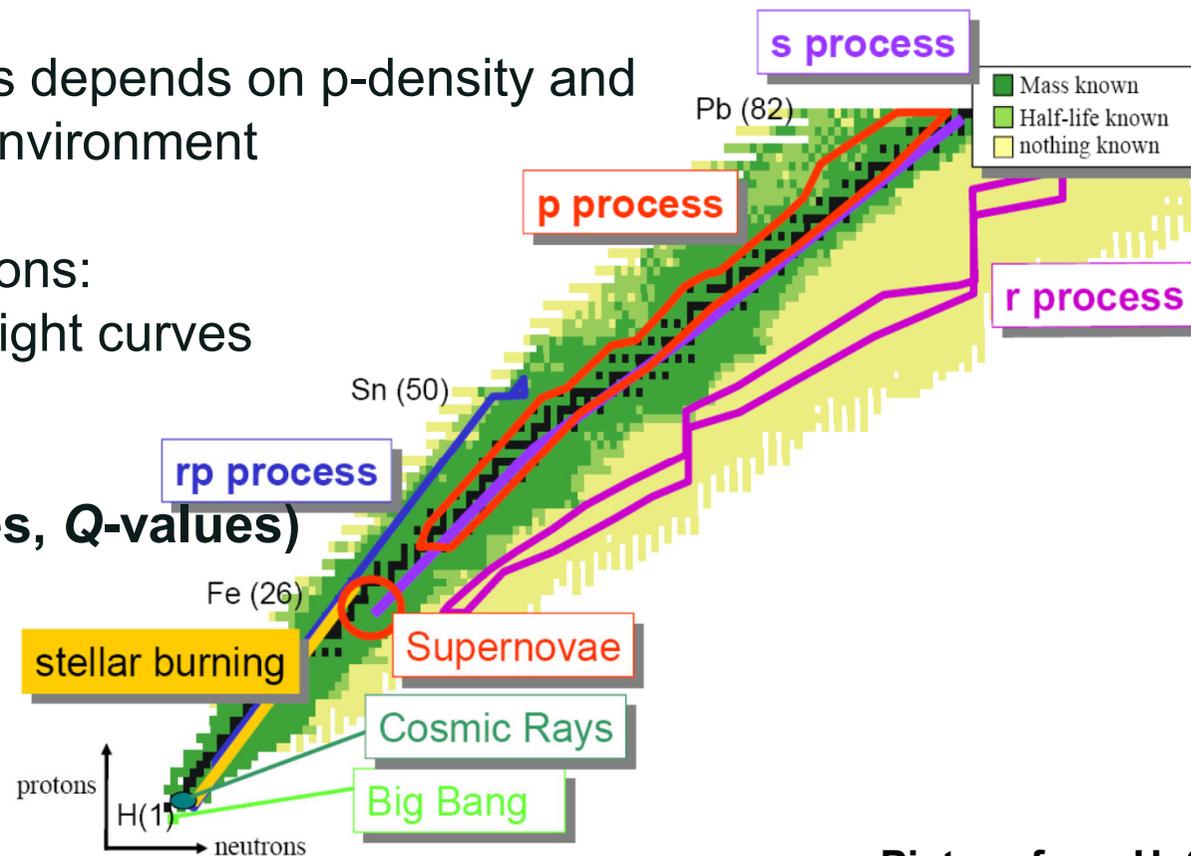
- Nucleosynthesis e.g. on accreting neutron stars
- Explosive hydrogen burning (X-ray bursts)
- Steady-state burning

Evolution of the process depends on p-density and temperature in stellar environment

Astrophysical observations:
elemental abundance, light curves

Required nuclear data:

- **Masses (sep. energies, Q-values)**
- β decay half-lives
- Reaction rates



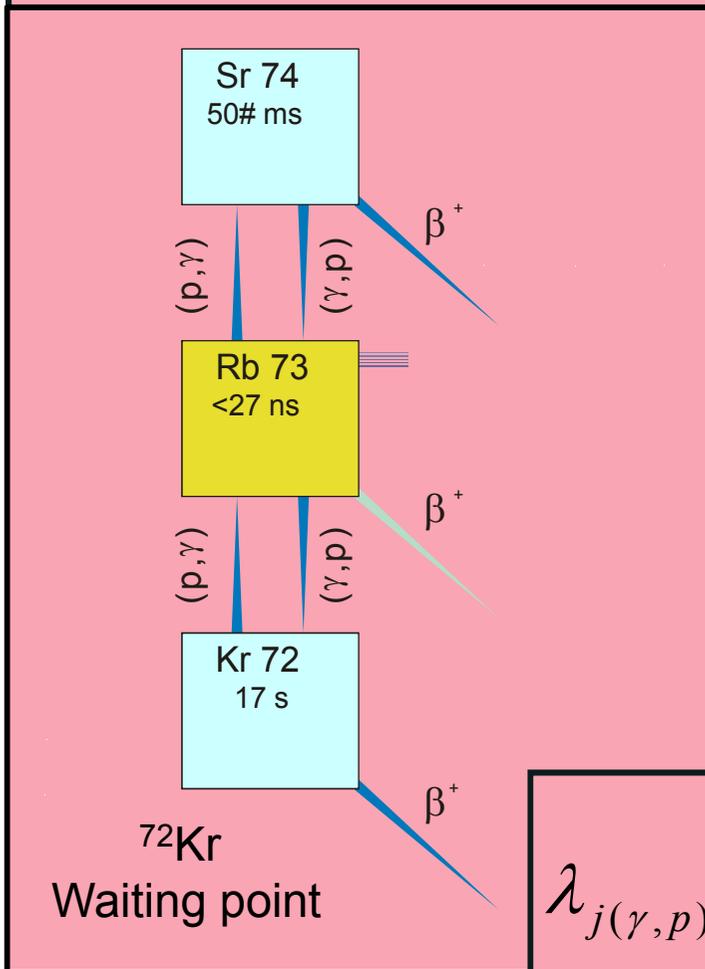
Picture from H. Schatz

Importance of the mass

Nucleon separation energies (S_n , S_p , S_{2n} , S_{2p})

$$Q_{i(p,\gamma)j} = S_j - S_i$$

structure & astrophysics



$$B(N, Z) - B(N - 2, Z)$$

$$B(N, Z) - B(N, Z - 2)$$

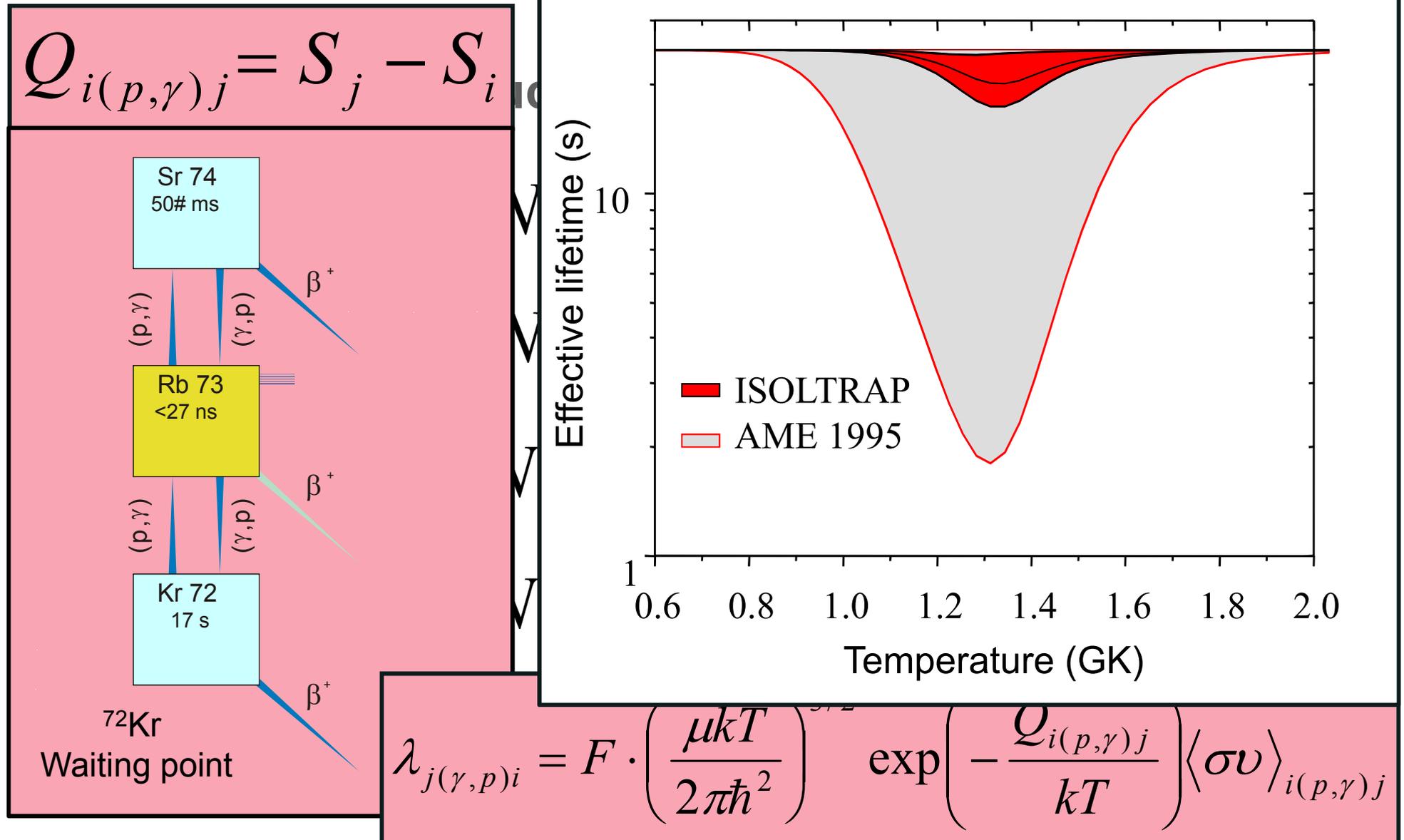
$$B(N, Z) - B(N - 1, Z)$$

$$B(N, Z) - B(N - 1, Z)$$

$$\lambda_{j(\gamma,p)i} = F \cdot \left(\frac{\mu kT}{2\pi\hbar^2} \right)^{3/2} \exp\left(-\frac{Q_{i(p,\gamma)j}}{kT} \right) \langle \sigma v \rangle_{i(p,\gamma)j}$$

Importance of the mass

Nucleon separation energies (S_n , S_p , S_{2n} , S_{2p})



D. Rodríguez et al., PRL 93 (2004) 161104

Importance of the mass

Q-values from β -decay

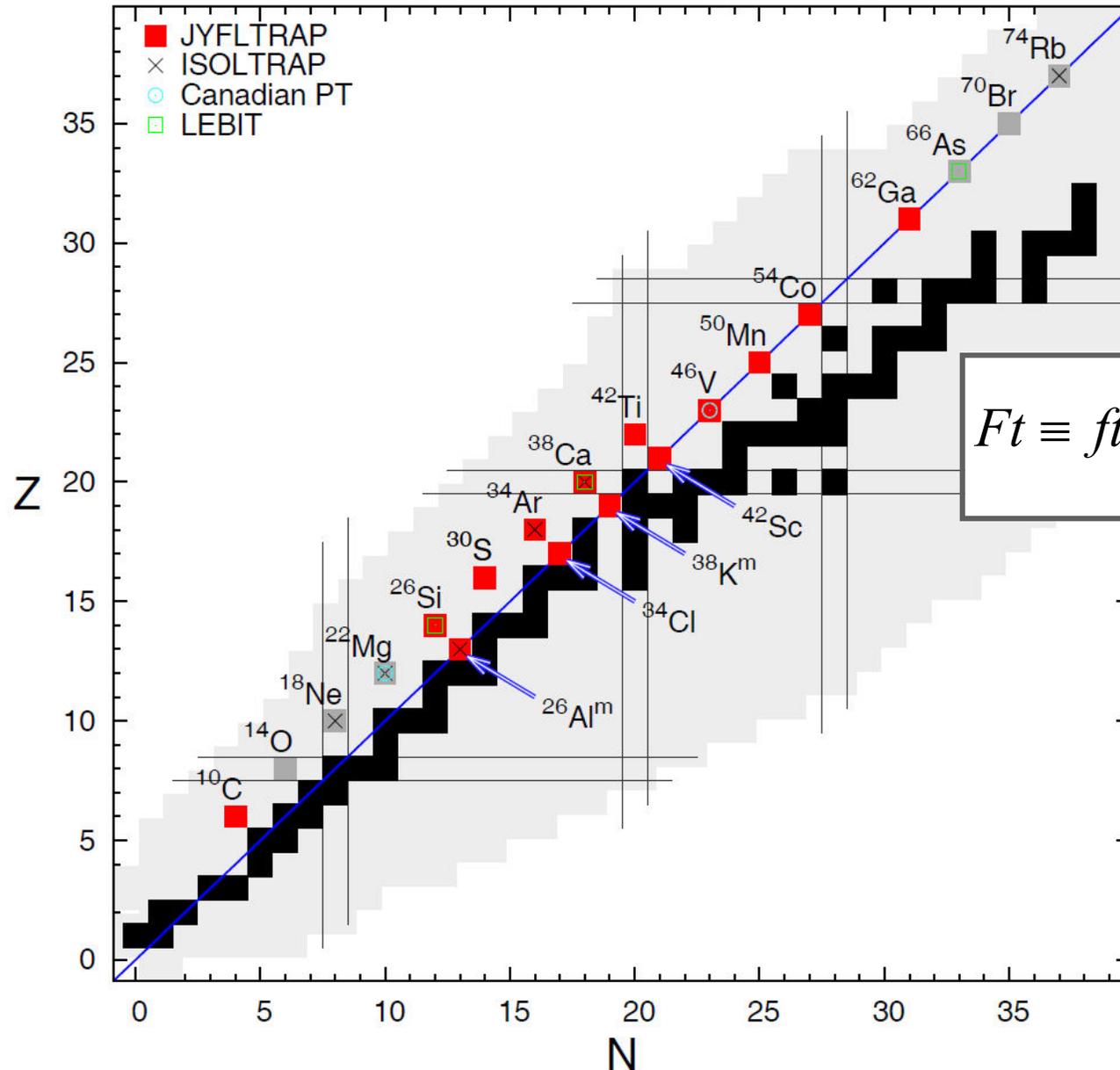
Fundamental symmetries & Neutrino-related physics

$$Q_{EC} = ME(Z, N) - ME(Z - 1, N + 1)$$

$$Q_{\beta^+} = ME(Z, N) - ME(Z - 1, N + 1) - 2m_e$$

$$Q_{\beta^-} = ME(Z, N) - ME(Z + 1, N - 1)$$

Importance of the mass Q-values from β -decay



CVC hypothesis
Superallowed
 β emitters
 $0^+ \rightarrow 0^+$

$$Ft \equiv ft(1 + \delta_R)(1 - \delta_C) = \frac{K}{2G_V^2(1 + \Delta_R^V)}$$

$f(Q^5)$ statistical rate function
 t ($t_{1/2}$, $1/R$) partial half-life

$$V_{ud}^2 = G_V^2 / G_\mu^2$$

$$|V_{ud}^2| + |V_{us}^2| + |V_{ub}^2| = 1 - \Delta$$

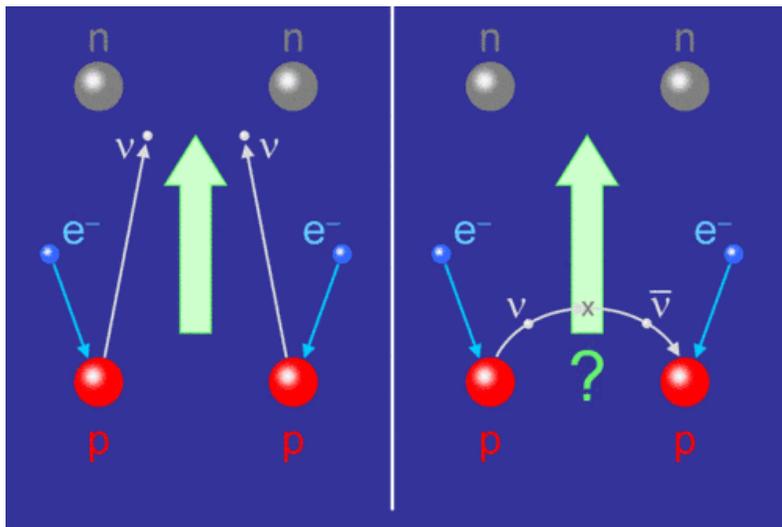
Importance of the mass Q-values from *EC*-decay

Is the neutrino a Majorana or Dirac particle?

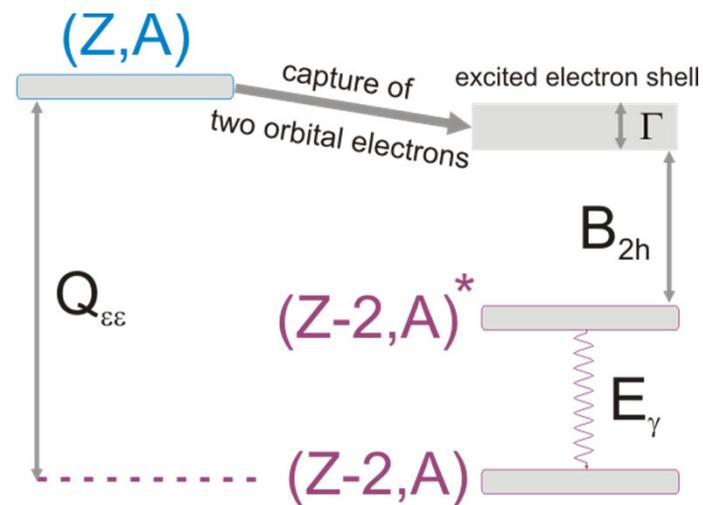
$2\nu 2EC$ ($T_{1/2} > 10^{24}$ y)

$0\nu 2EC$ ($T_{1/2} > 10^{30}$ y)

$$\frac{1}{T_{1/2}} = C \times m_\nu^2 \times |M|^2 \times |\Psi_{1e}|^2 \times |\Psi_{2e}|^2 \times \frac{\Gamma}{(Q - B_{2h} - E_\gamma^2)^2 + \frac{1}{4}\Gamma^2}$$



$0\nu 2EC$ might be resonantly enhanced ($T_{1/2} \sim 10^{25}$ y)



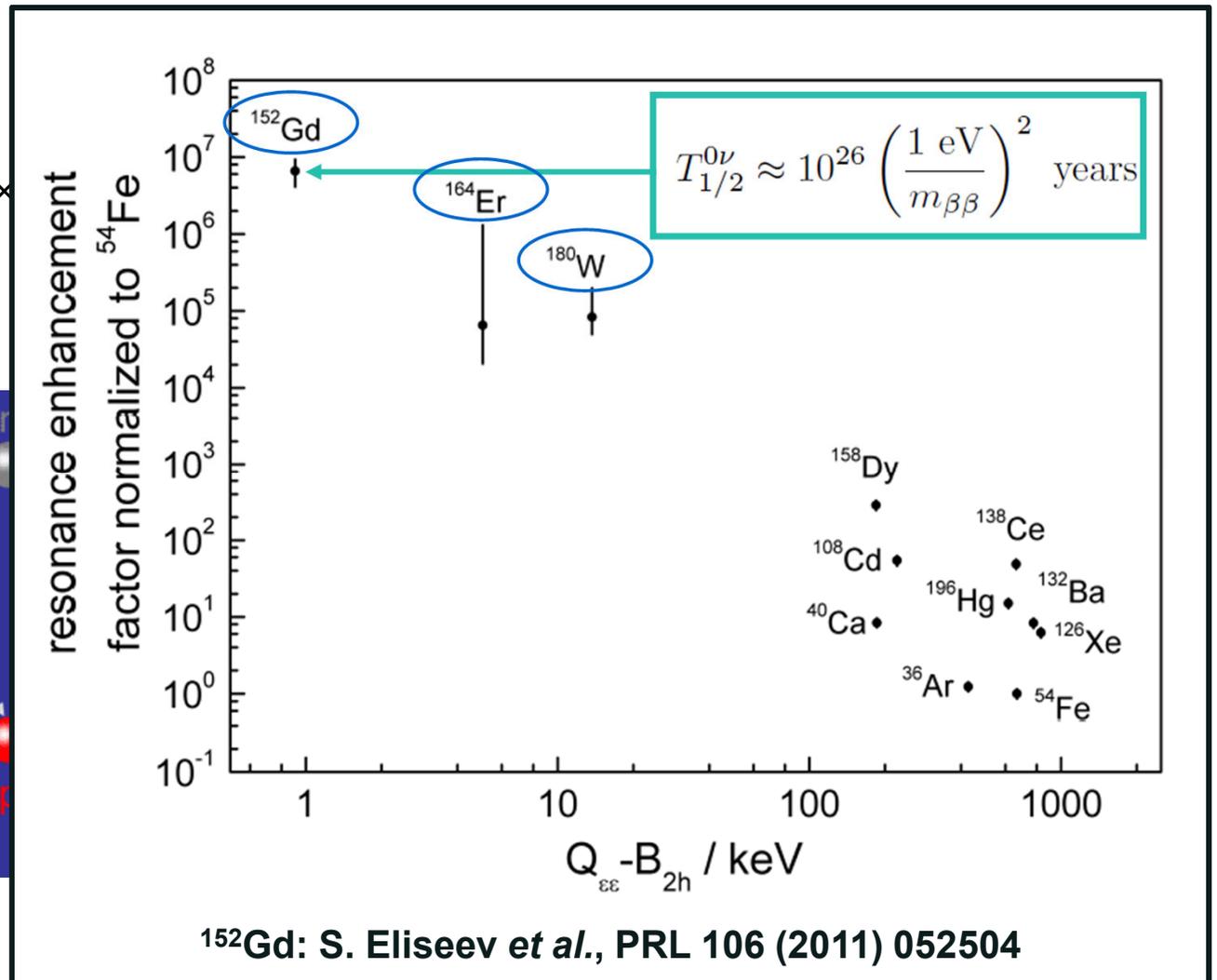
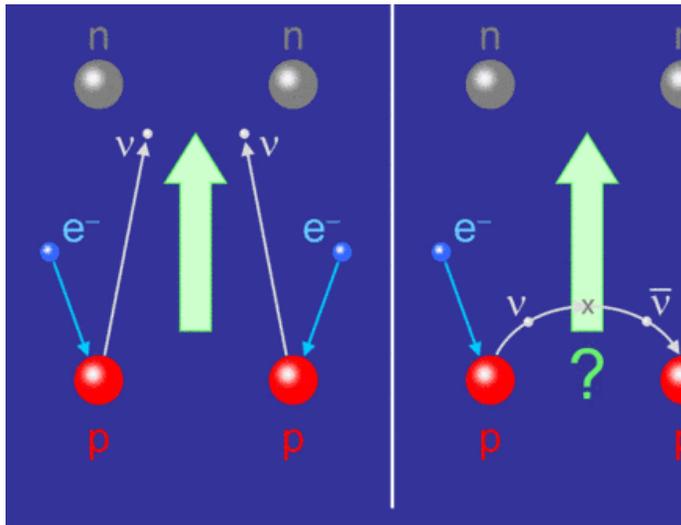
Search for nuclides with $D = (Q_{ee} - B_{2h} - E_g) < 1$ keV

Importance of the mass Q-values from *EC*-decay

Is the neutrino a Majorana or Dirac particle?

$2\nu 2EC$ ($T_{1/2} > 10^{24}$ y)

$$\frac{1}{T_{1/2}} = C \times m_\nu^2 \times |M|^2 \times$$

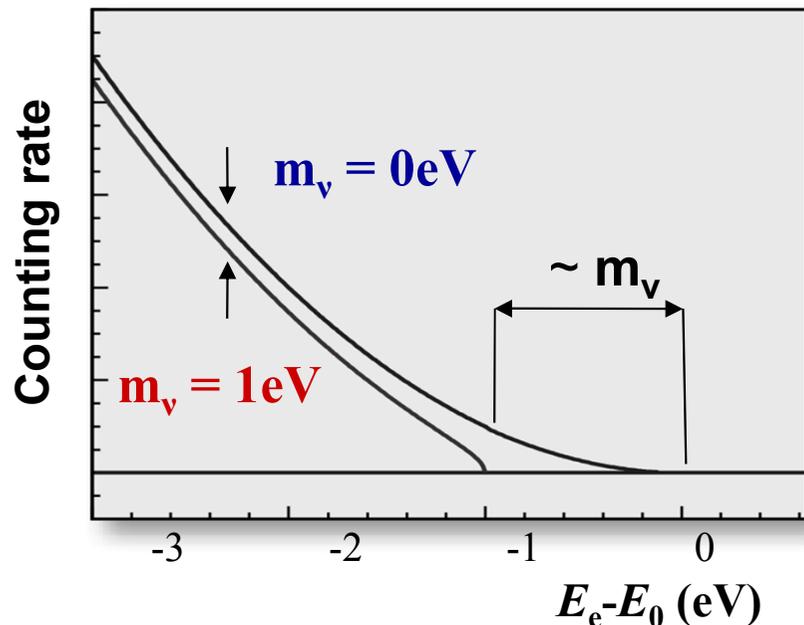


Importance of the mass

Q-values from β -decay

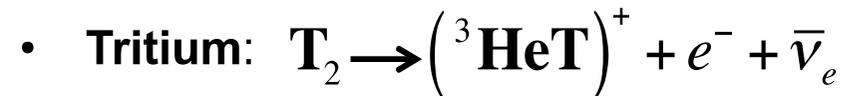
- Direct measurement of the mass by analyzing the end of the β -decay spectrum independent whether the neutrino is a Majorana or a Dirac particle

$$\frac{dN}{dE} = k \times F(E, Z) \times p \times E_{total} \times (E_0 - E_e)^2 \times \left[(E_0 - E_e)^2 - m_\nu^2 \right]^{1/2}$$

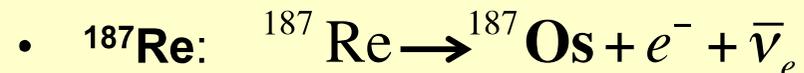


$$F(\delta E) = \int_{-\delta E}^0 N(E_\beta) dE \approx 2 \left(\frac{\delta E}{E_0} \right)^3$$

- The best candidates will be decaying nuclei with low *end-point energies*



$$Q = 18.6 \text{ keV} \quad T_{1/2} = 12.3 \text{ y}$$

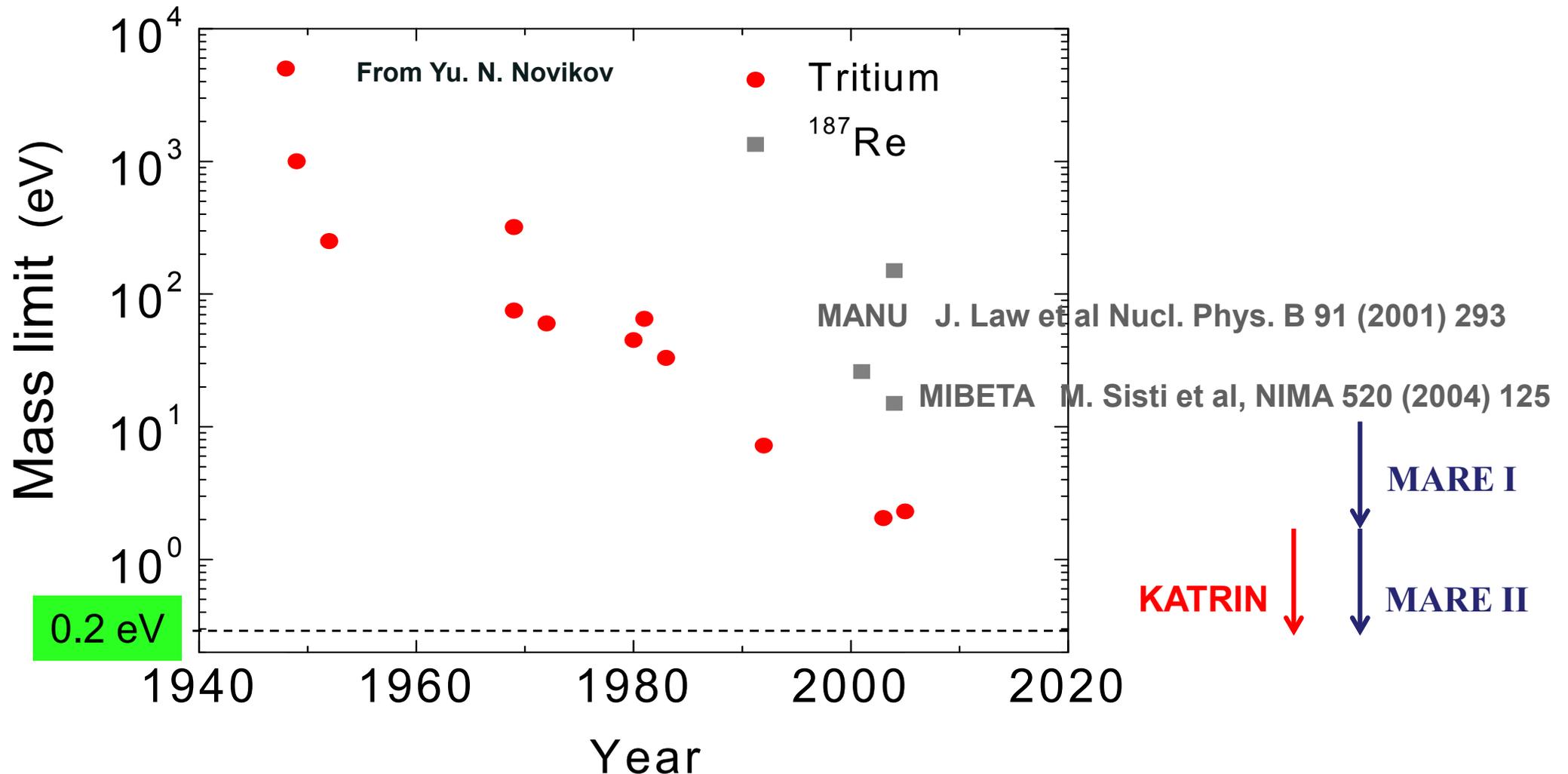


$$Q = 2.5 \text{ keV} \quad T_{1/2} = 43.2 \times 10^9 \text{ y}$$

Q has to be unambiguously determined (with $\delta m/m$ of at least 10^{-11}) from the mass difference using a Penning trap

History of antineutrino mass measurements

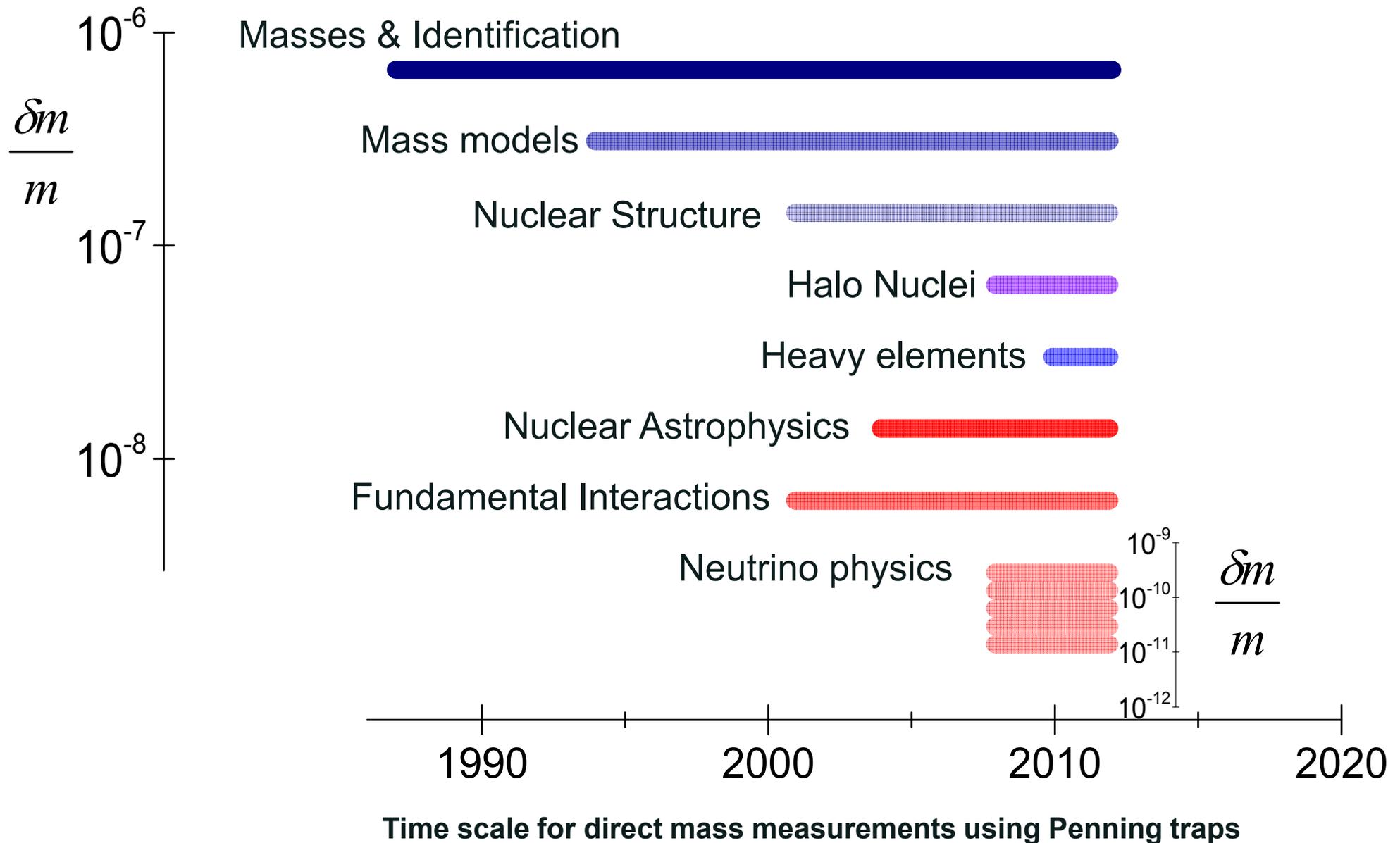
Tritium and ^{187}Re



- There are two International collaborations aiming at measuring with a sensitivity of about $0.2 \text{ eV}/c^2$: KATRIN (KARlsruhe TRItium Neutrino experiment) and MARE (Micro calorimeter Array for a Rhenium Experiment)

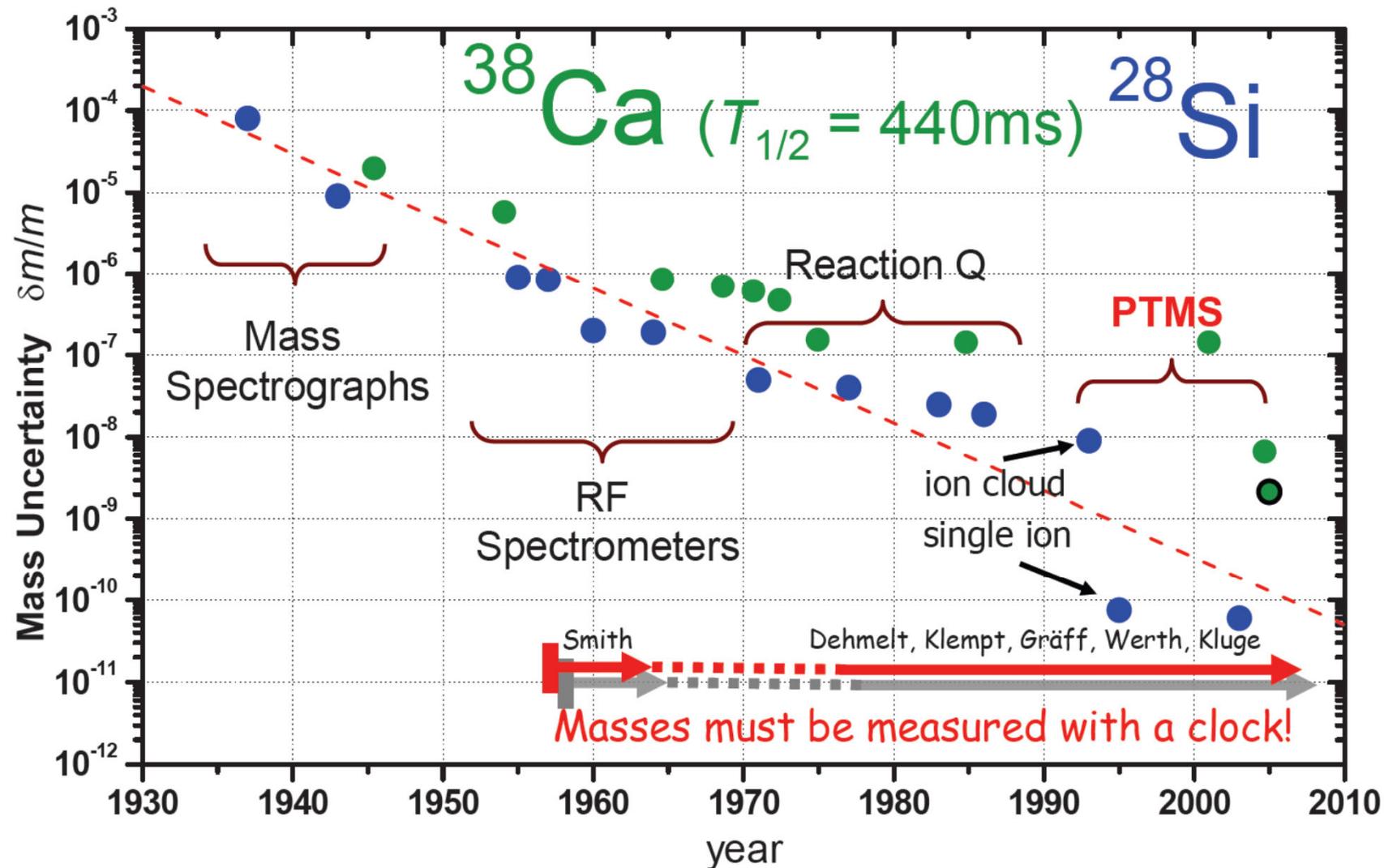
Physics goals

Mass uncertainty required



Mass uncertainty

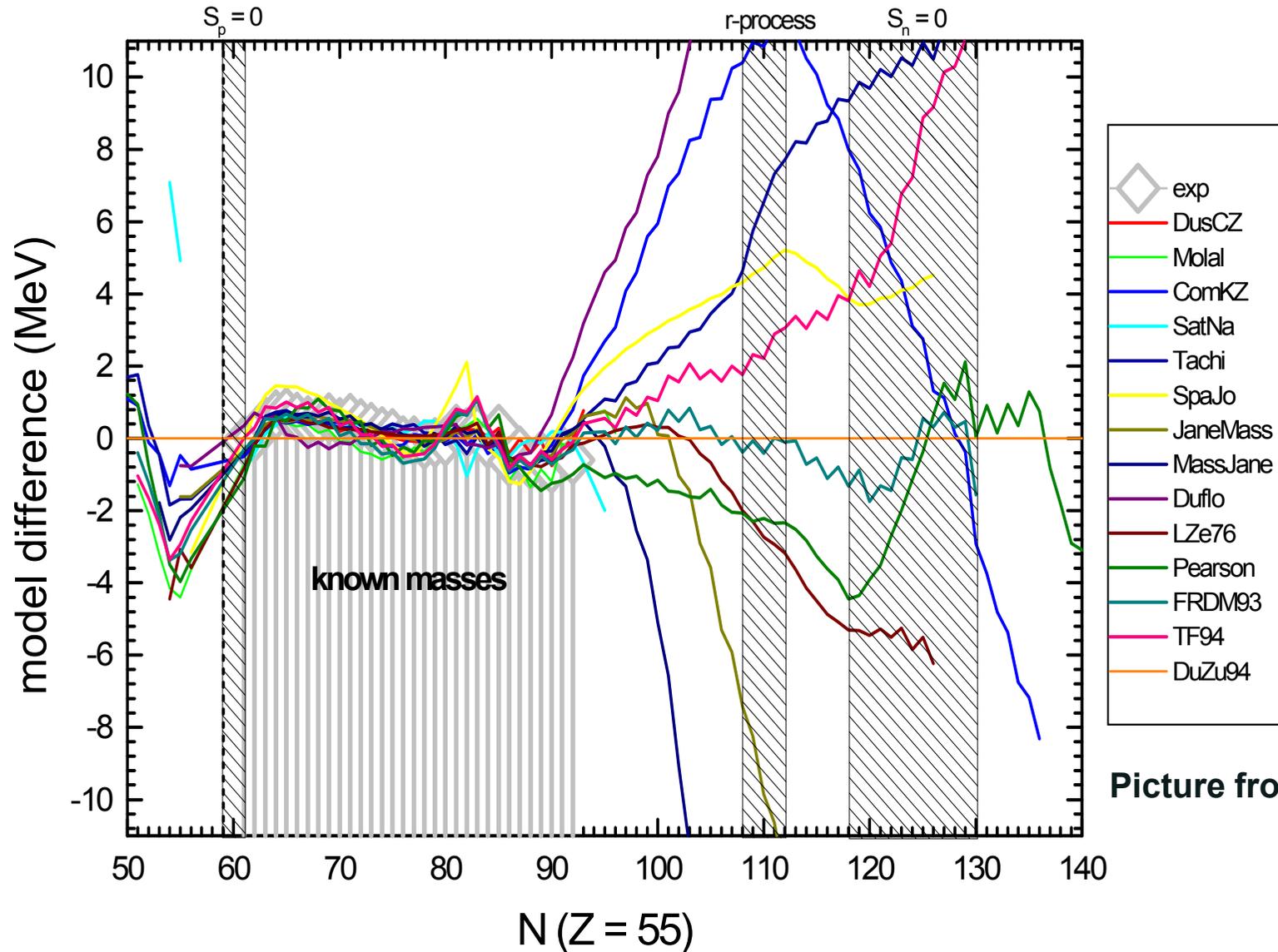
Evolution in history



Picture from H.-J. Kluge

Experimental results and models

Discrepancies between the models

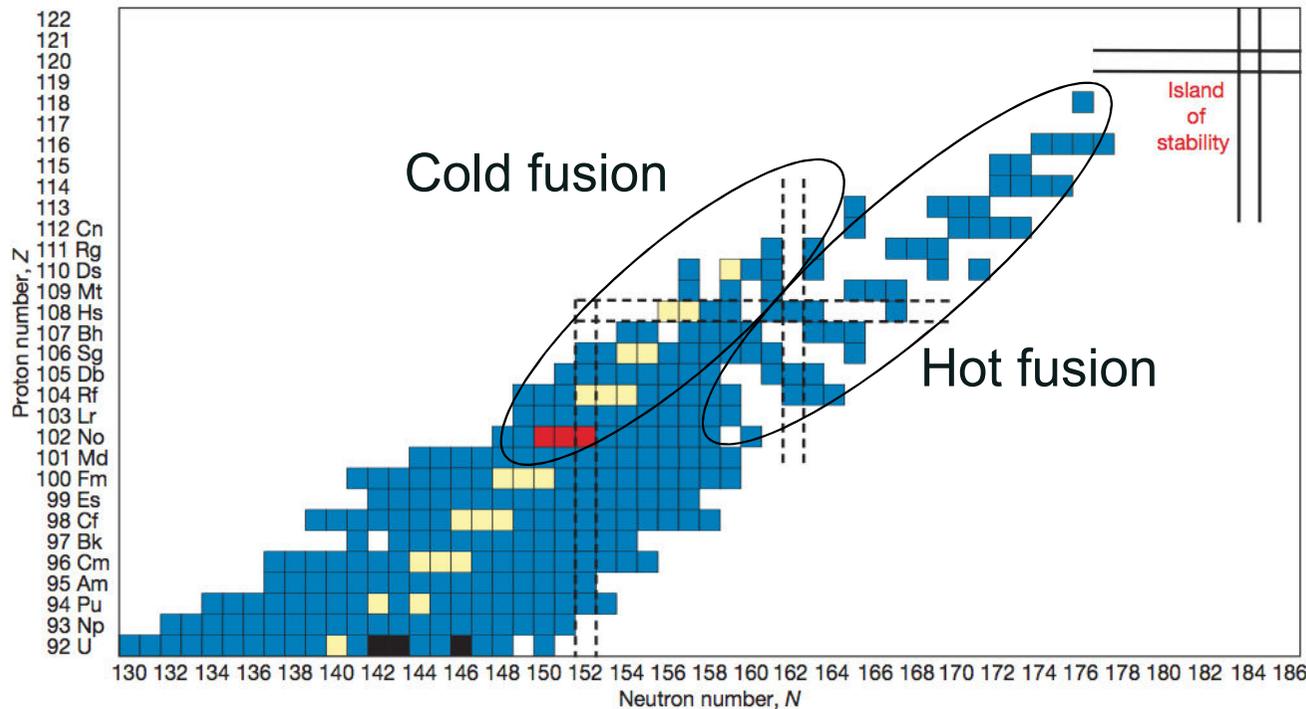


Picture from D. Lunney

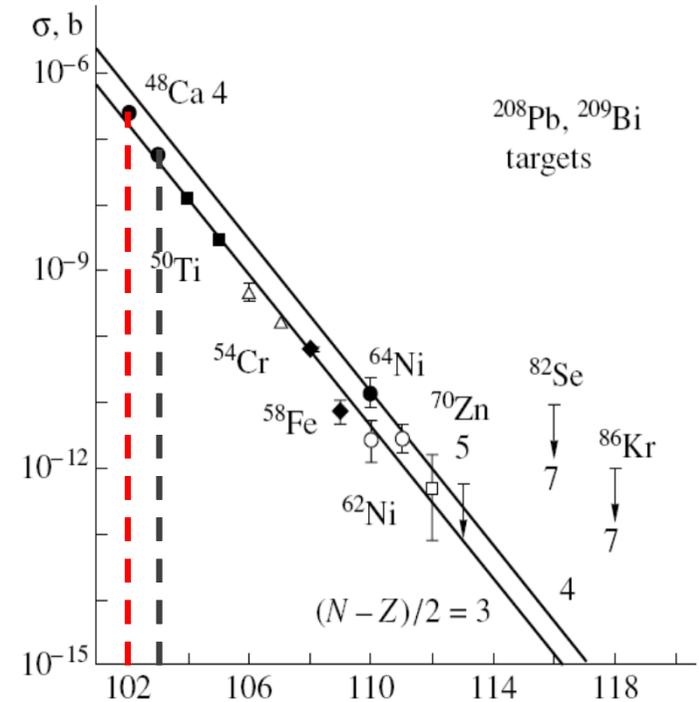
Particularly relevant in the region of Superheavy Elements

Experimental results

Influence in the atomic mass evaluation



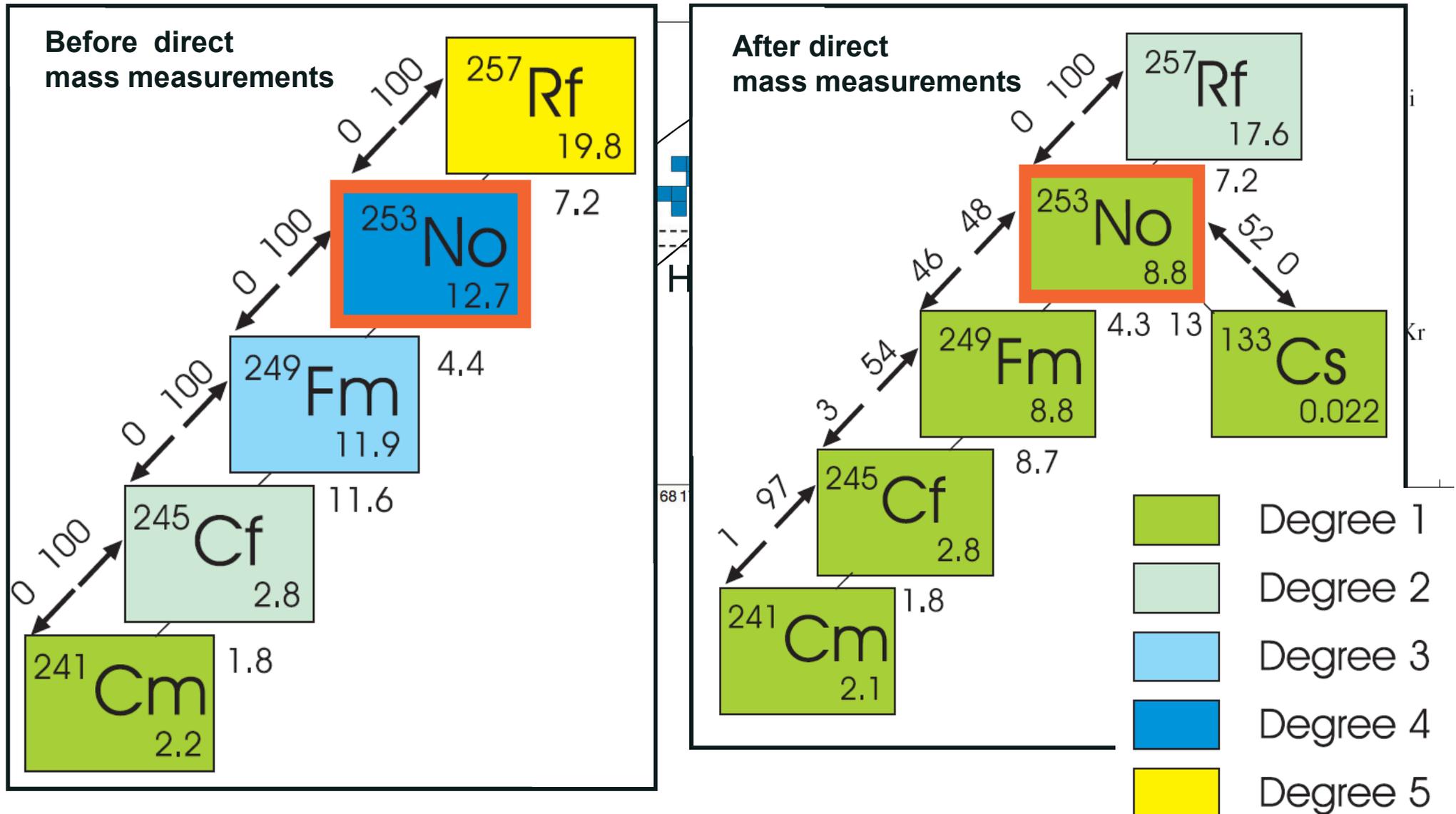
M. Block et al, Nature 463 (2010) 785



Very low production cross section

Experimental results

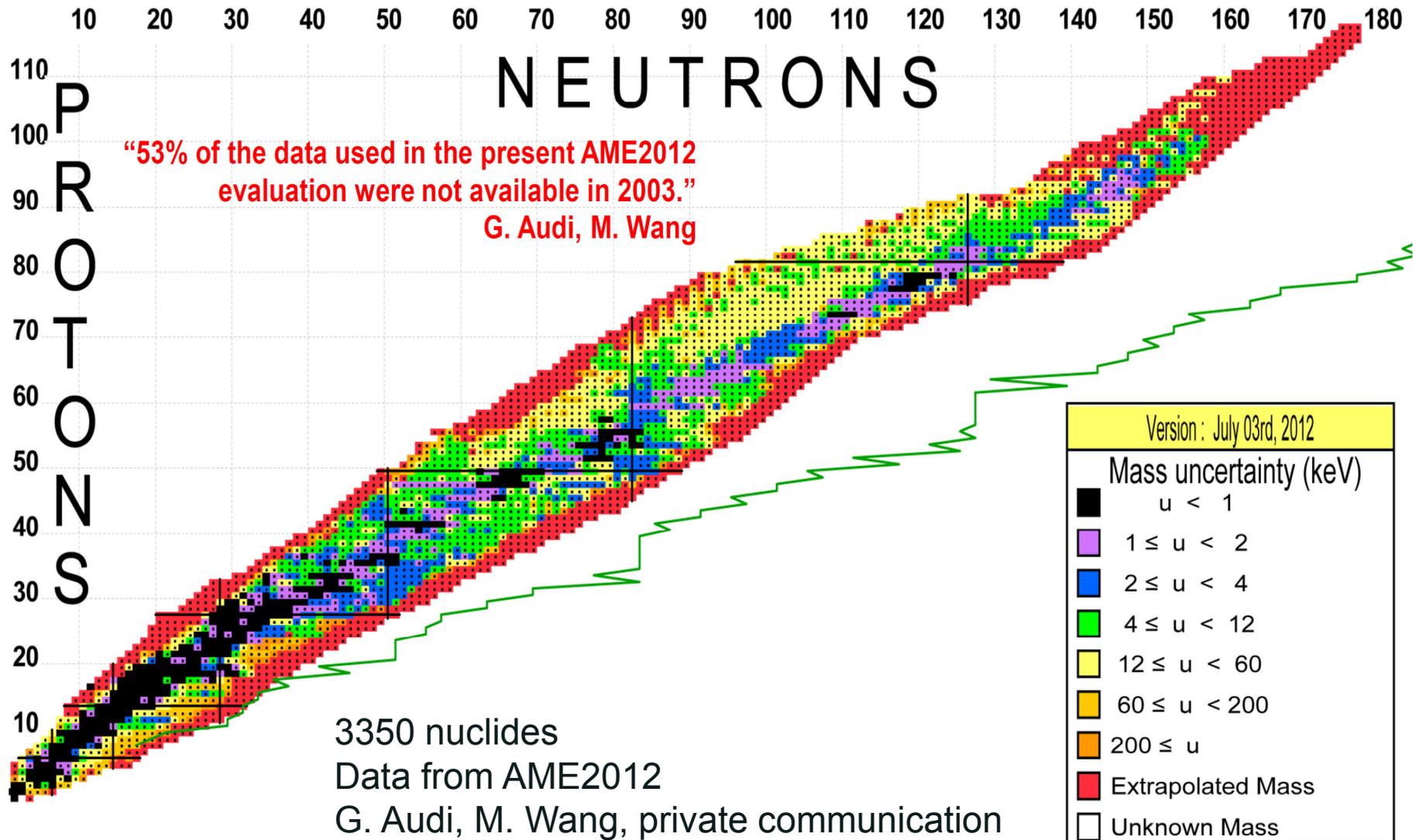
Influence in the atomic mass evaluation



M. Dworschak *et al.* Phys. Rev. C 81 (2010) 064312

Mass measurements in the last decade

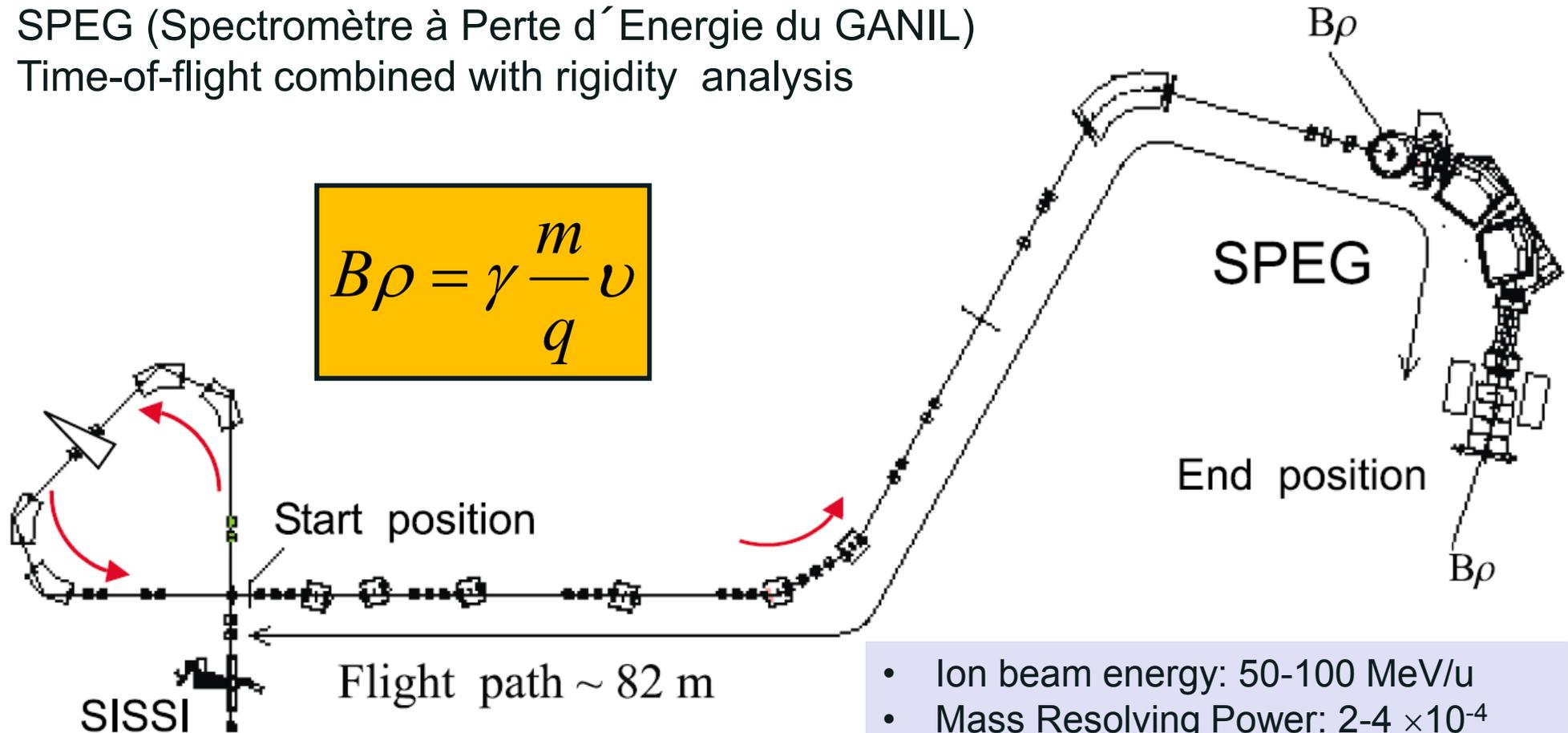
The 2012 Atomic Mass Evaluation



MASS SPECTROMETRY TECHNIQUES FOR EXOTIC NUCLEI

Time-of-flight measurements Large Spectrometers (an example)

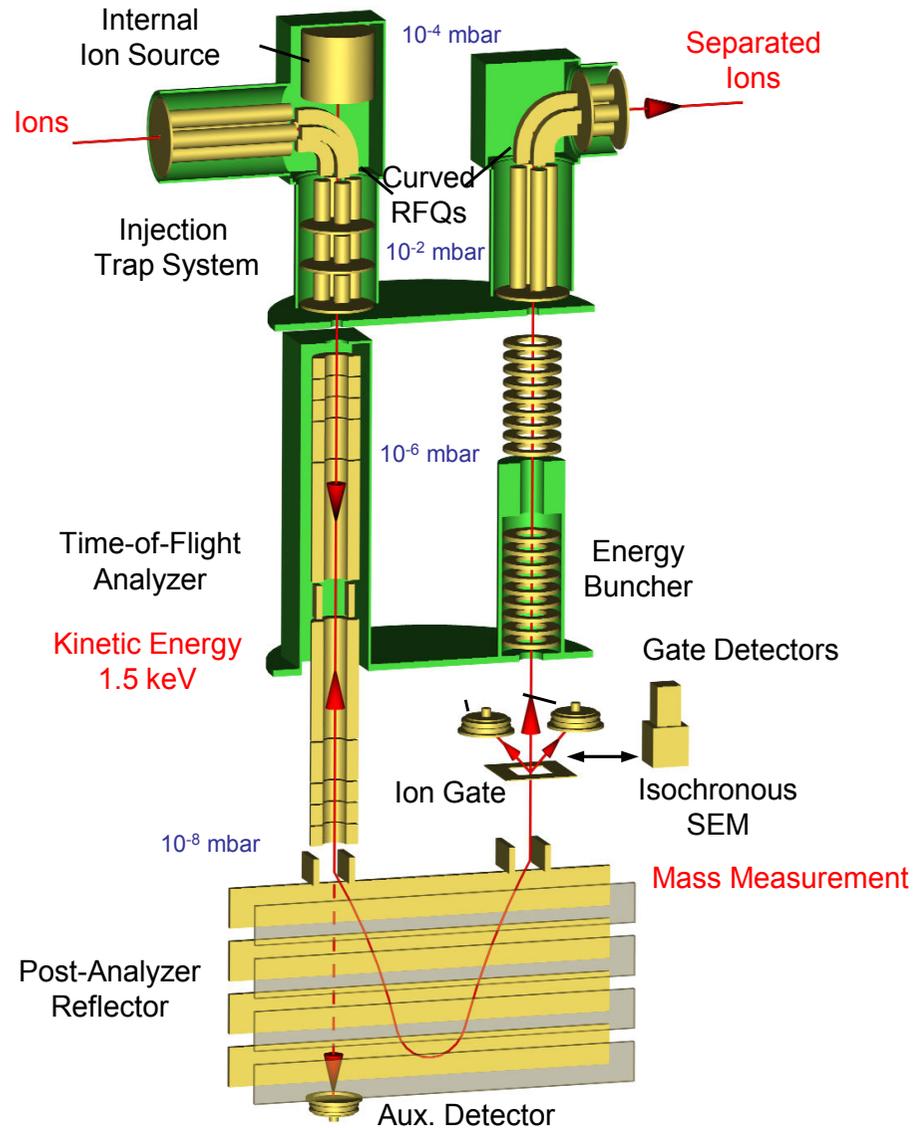
SPEG (Spectromètre à Perte d'Énergie du GANIL)
Time-of-flight combined with rigidity analysis



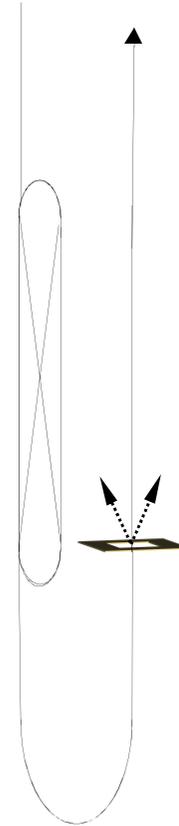
- Ion beam energy: 50-100 MeV/u
- Mass Resolving Power: $2-4 \times 10^4$
- High sensitivity and very short-lived nuclei
- Masses up to $A \sim 70$

Time-of-flight measurements

Compact devices (MR-TOF-MS)

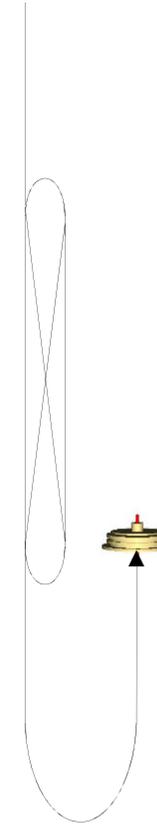


Isobar Separation Mode



$$m/\Delta m > 10^5$$

High Resolution Mode



$$m/\Delta m > 10^5,$$

$$\text{Mass Uncertainty } 10^{-6}-10^{-7}$$

Broadband Mode



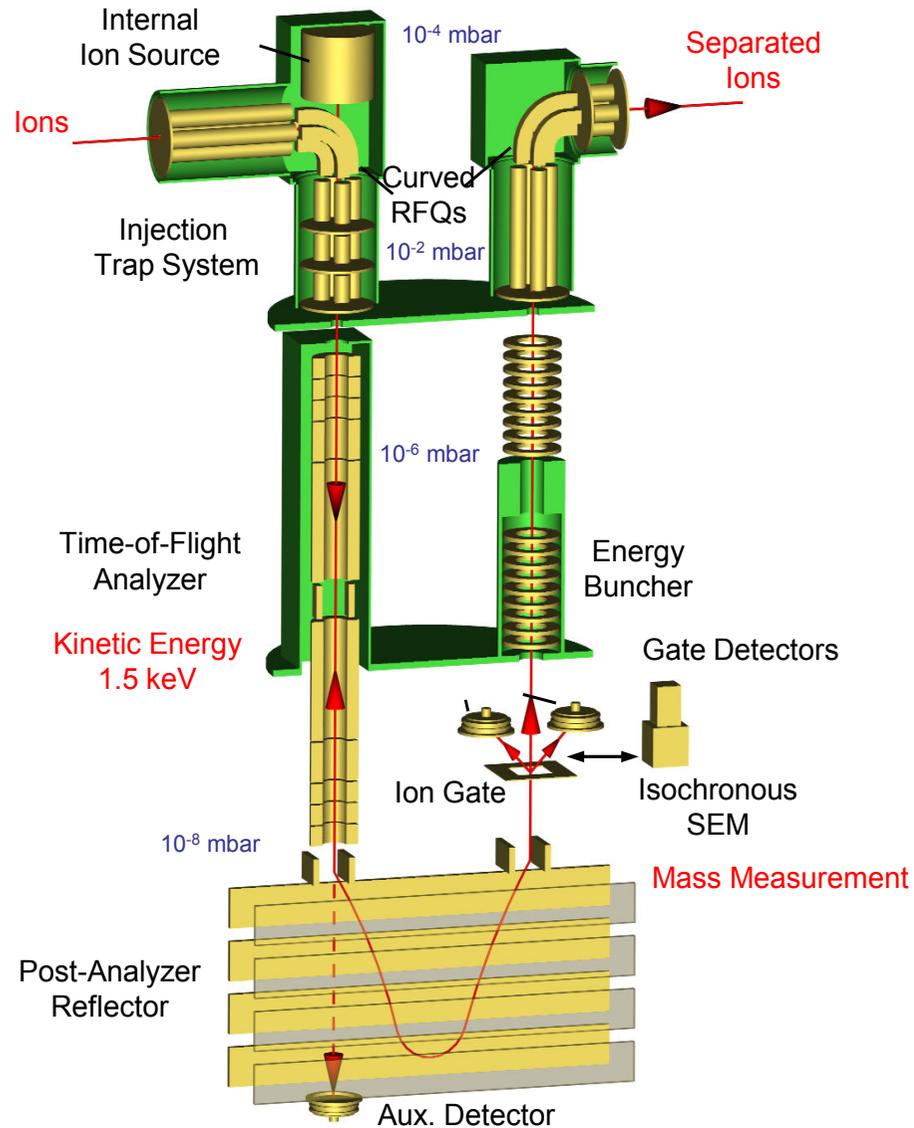
$$\text{Full Mass Range,}$$

$$m/\Delta m \sim 4000$$

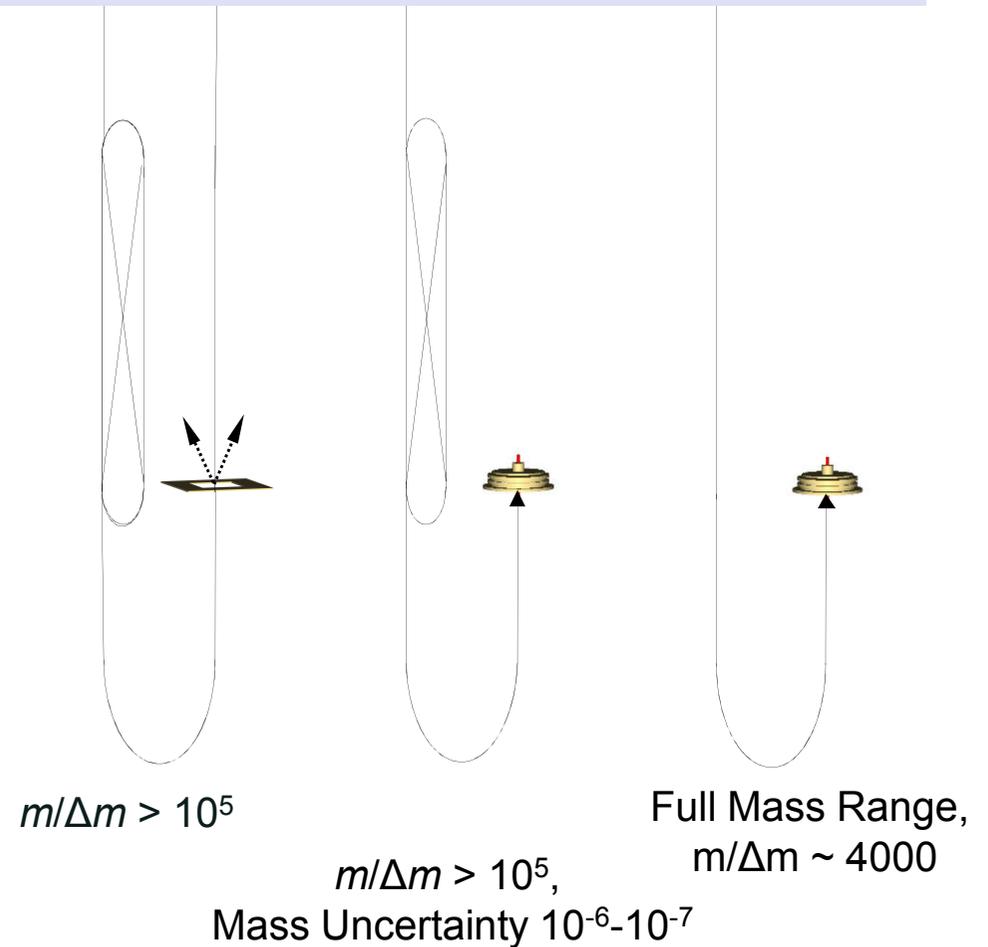
W.R. Plaß *et al.*, NIMB 266 (2008) 4560

Time-of-flight measurements

Compact devices (MR-TOF-MS)



- Ion beam energy: 1-2 keV
- Mass Resolving Power: 600,000
- Mass Measurement Uncertainty: $\sim 10^{-7}$
- Sensitivity: ~ 10 ions



W.R. Plaß *et al.*, NIMB 266 (2008) 4560

MR-TOF-MS at accelerators

FRS at GSI

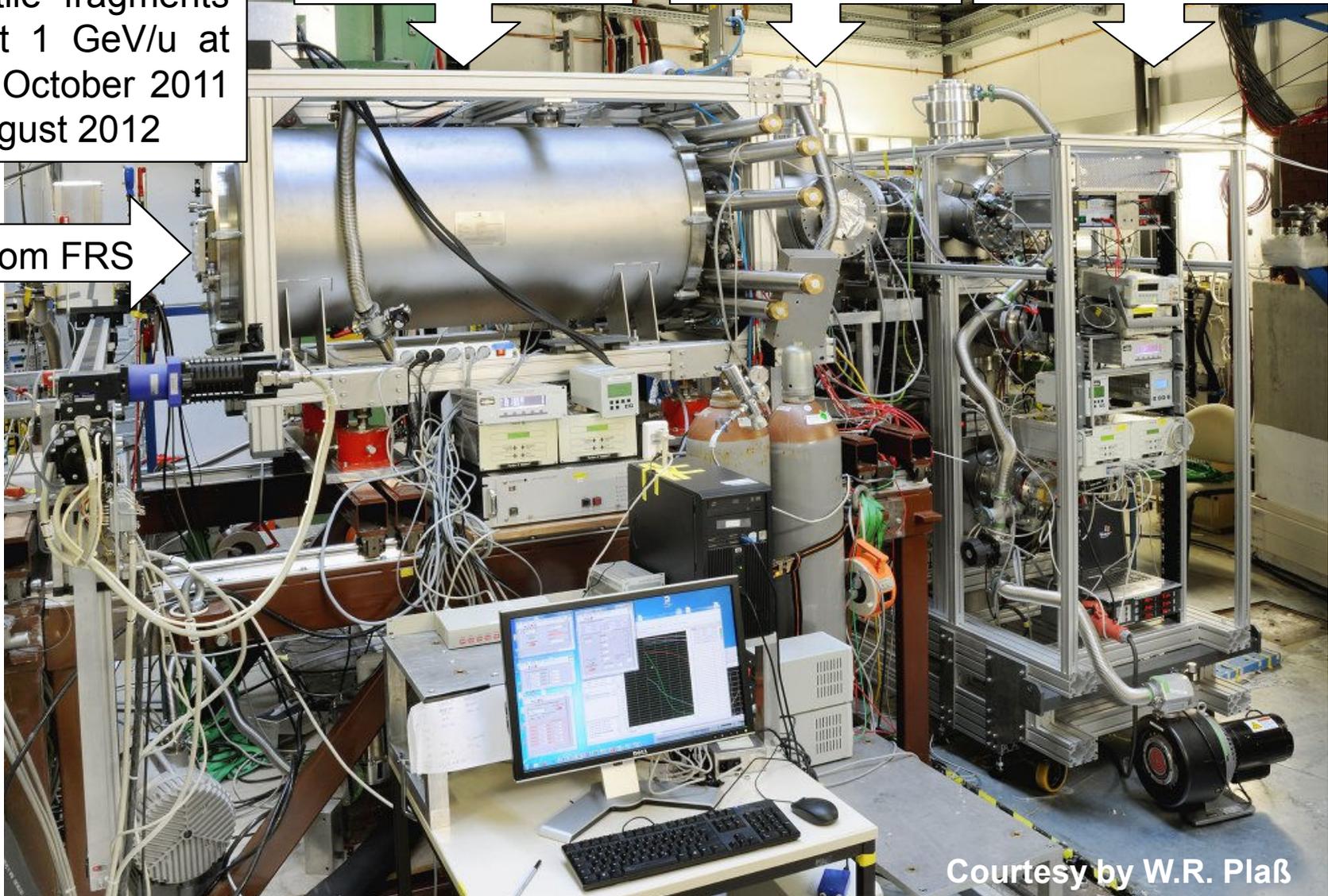
On-line test of the with ^{238}U projectile fragments produced at 1 GeV/u at the FRS in October 2011 and July/August 2012

Beam from FRS

Cryogenic
stopping cell

Diagnostics
unit

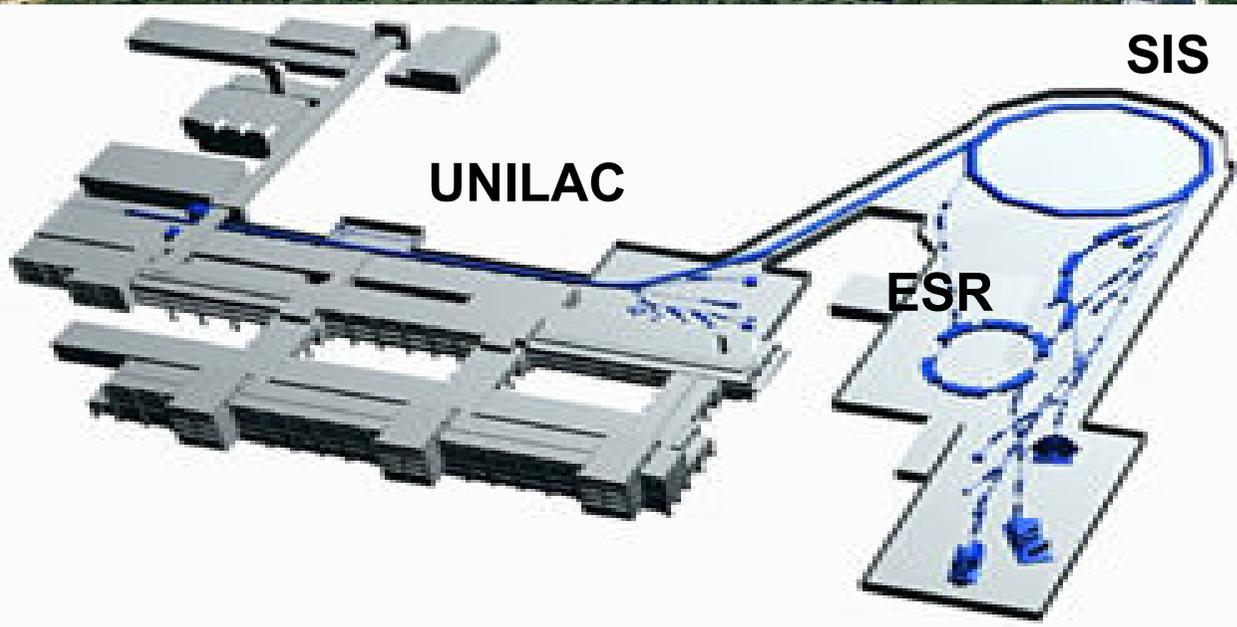
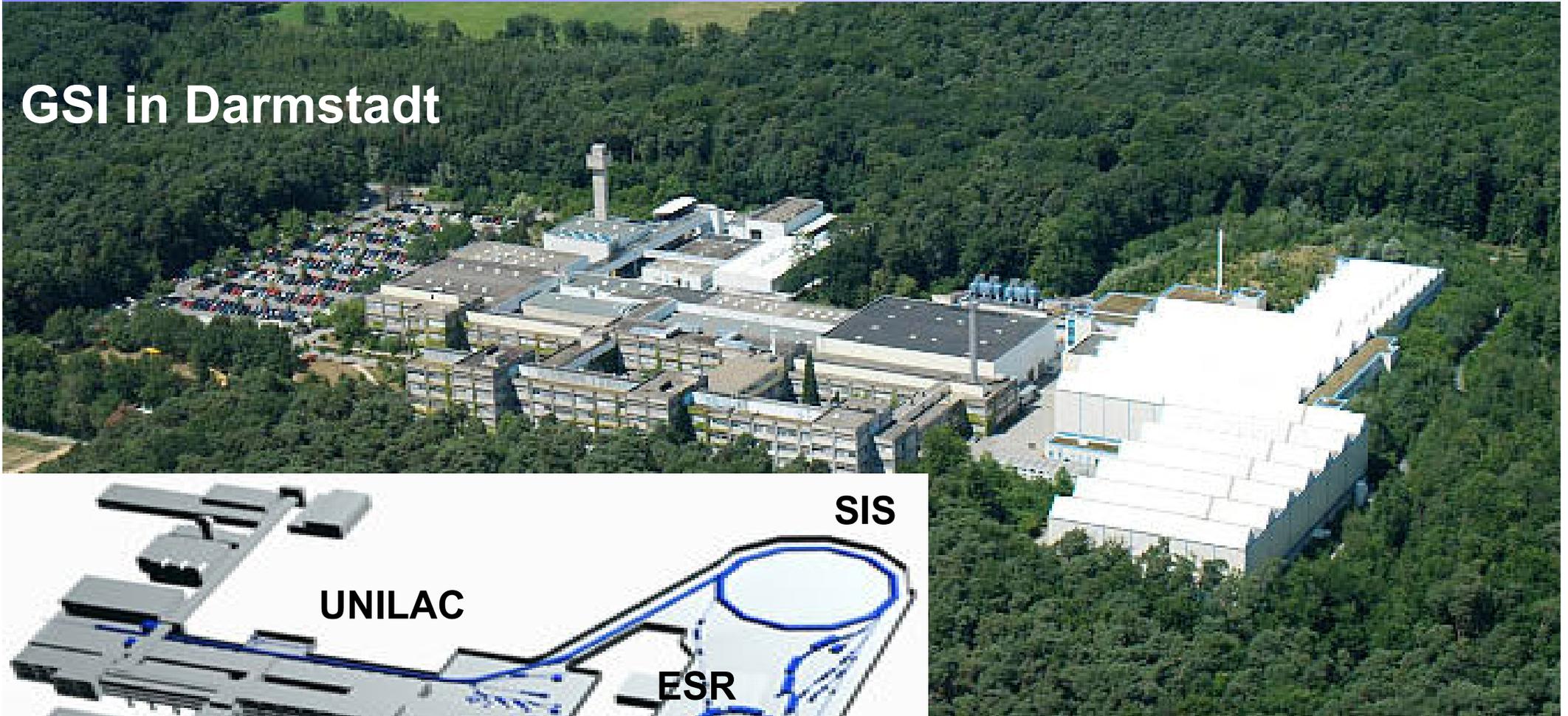
Time-of-flight
mass spectrometer



Courtesy by W.R. Plaß

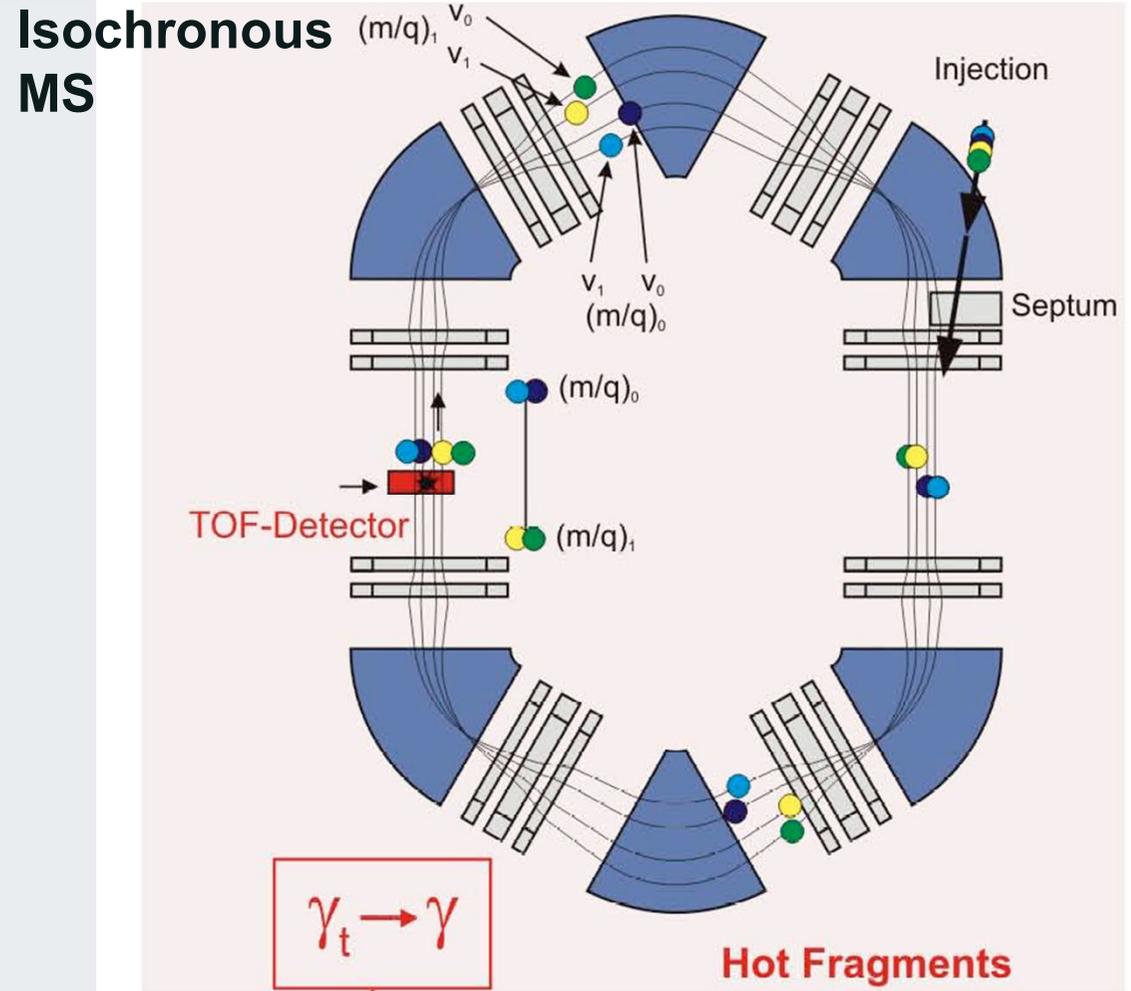
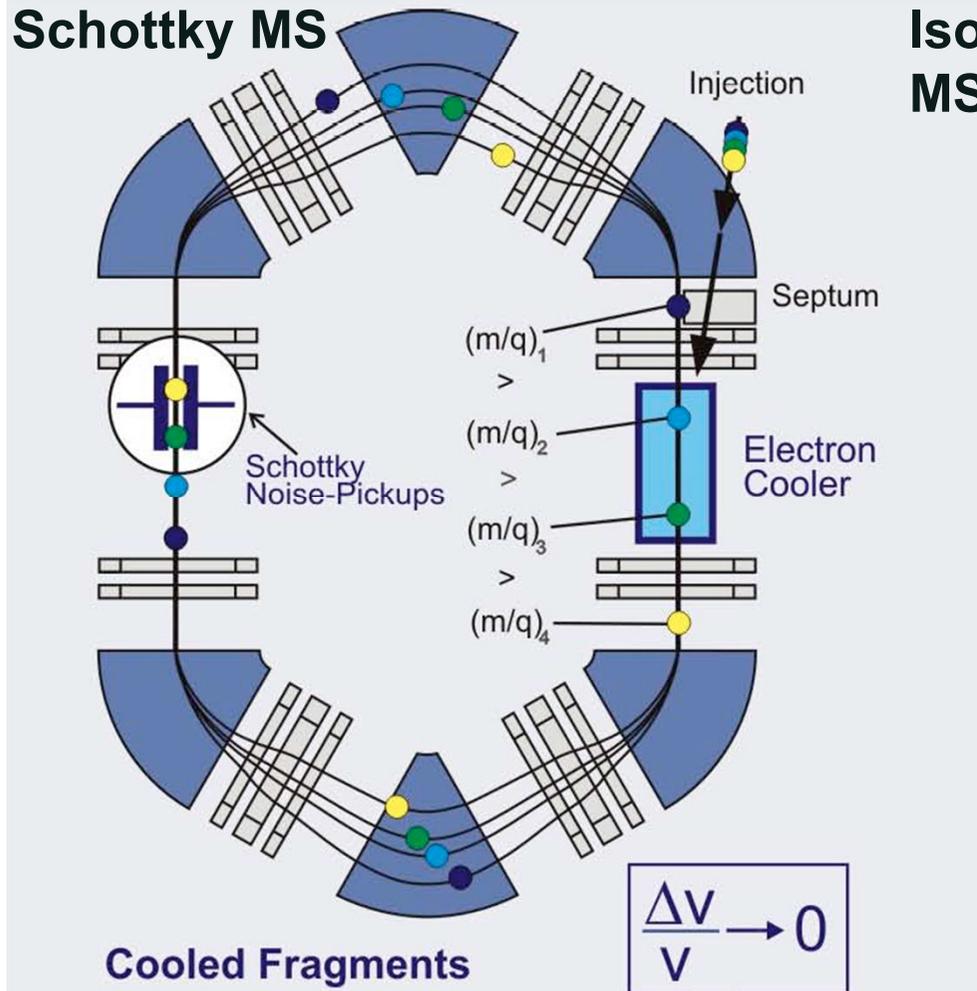
Revolution frequency measurements Experimental storage ring (ESR)

GSI in Darmstadt



Revolution frequency measurements

Experimental storage ring (ESR)



Stochastic &
Electron cooling

$$\frac{\Delta f}{f} = -\frac{1}{\gamma_t^2} \frac{\Delta(m/q)}{m/q} + \frac{\Delta V}{V} \left(1 - \frac{\gamma^2}{\gamma_t^2} \right)$$

Revolution frequency measurements Experimental storage ring (ESR)



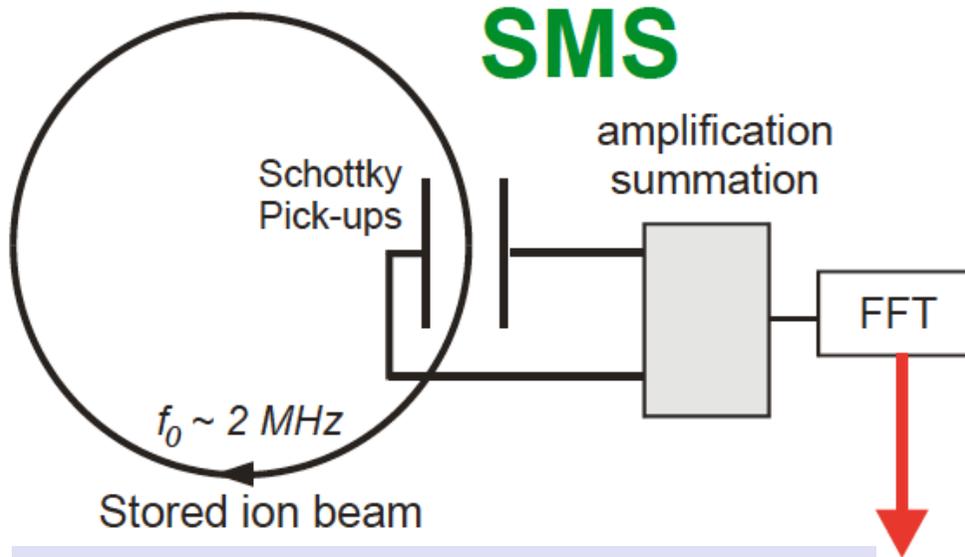
More than 1000 masses measured

$$\frac{\Delta f}{f} = -\frac{1}{\gamma_t^2} \frac{\Delta(m/q)}{m/q} + \frac{\Delta V}{V} \left(1 - \frac{\gamma^2}{\gamma_t^2} \right)$$

Revolution frequency measurements

Detectors and performance

Schottky MS



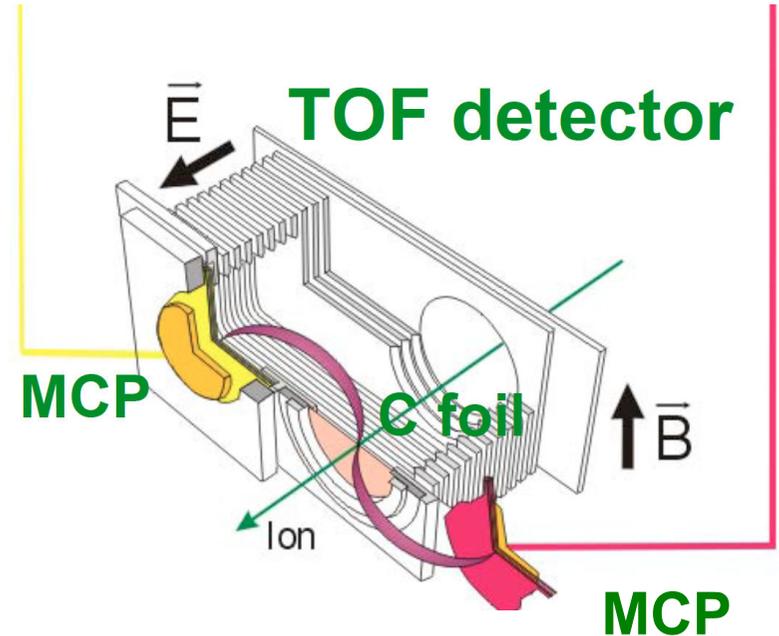
- Ion beam energy: ~ 10 MeV/u
- Mass Resolving Power: 2×10^6
- Measurement Uncertainty: 1.5×10^{-7}
- Sensitivity: ~1 ion
- Half-life: above 10 s

$$\frac{\Delta f}{f} = -\frac{1}{\gamma_t^2} \frac{\Delta(m/q)}{m/q} + \frac{\Delta V}{V} \left(1 - \frac{\gamma^2}{\gamma_t^2} \right)$$

Isochronous MS

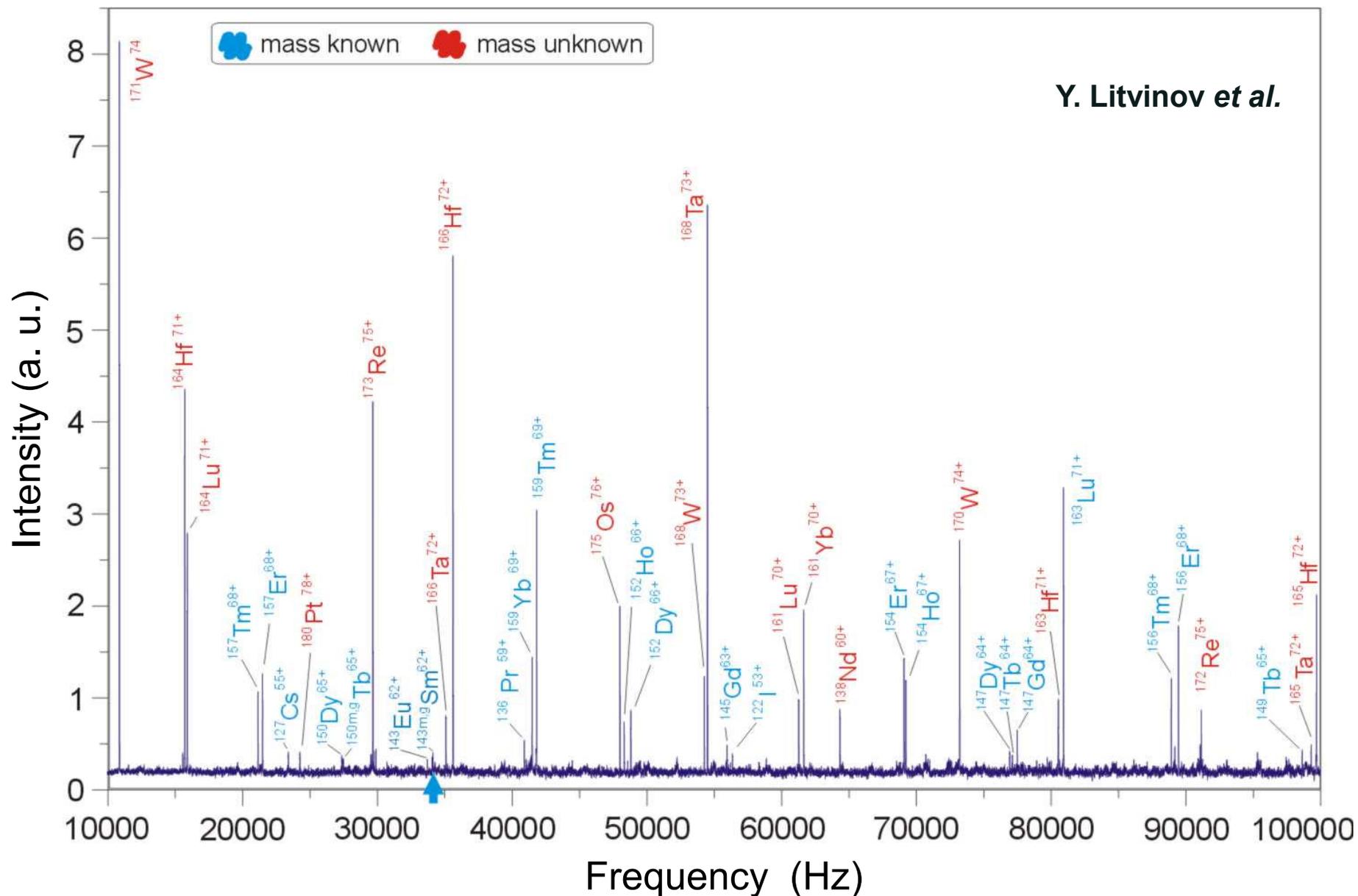
Oscilloscope channel A

Oscilloscope channel B



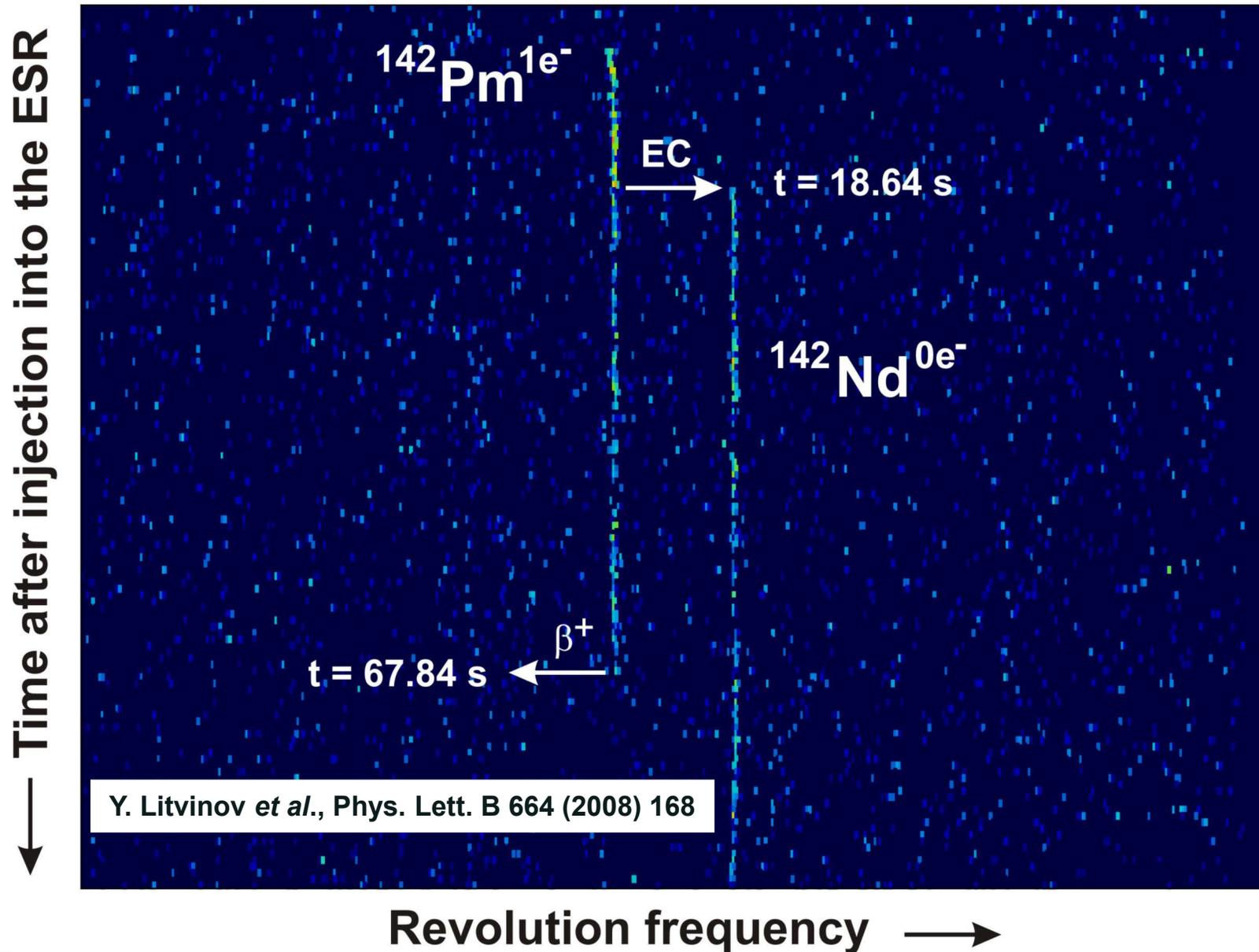
- Ion beam energy: ~ 10 MeV/u
- Mass Resolving Power: 3×10^5
- Measurement Uncertainty: $10^{-5} - 10^{-6}$
- Sensitivity: ~1 ion
- Half-life: above 100 μ s

Revolution frequency measurements at ESR SMS spectrum

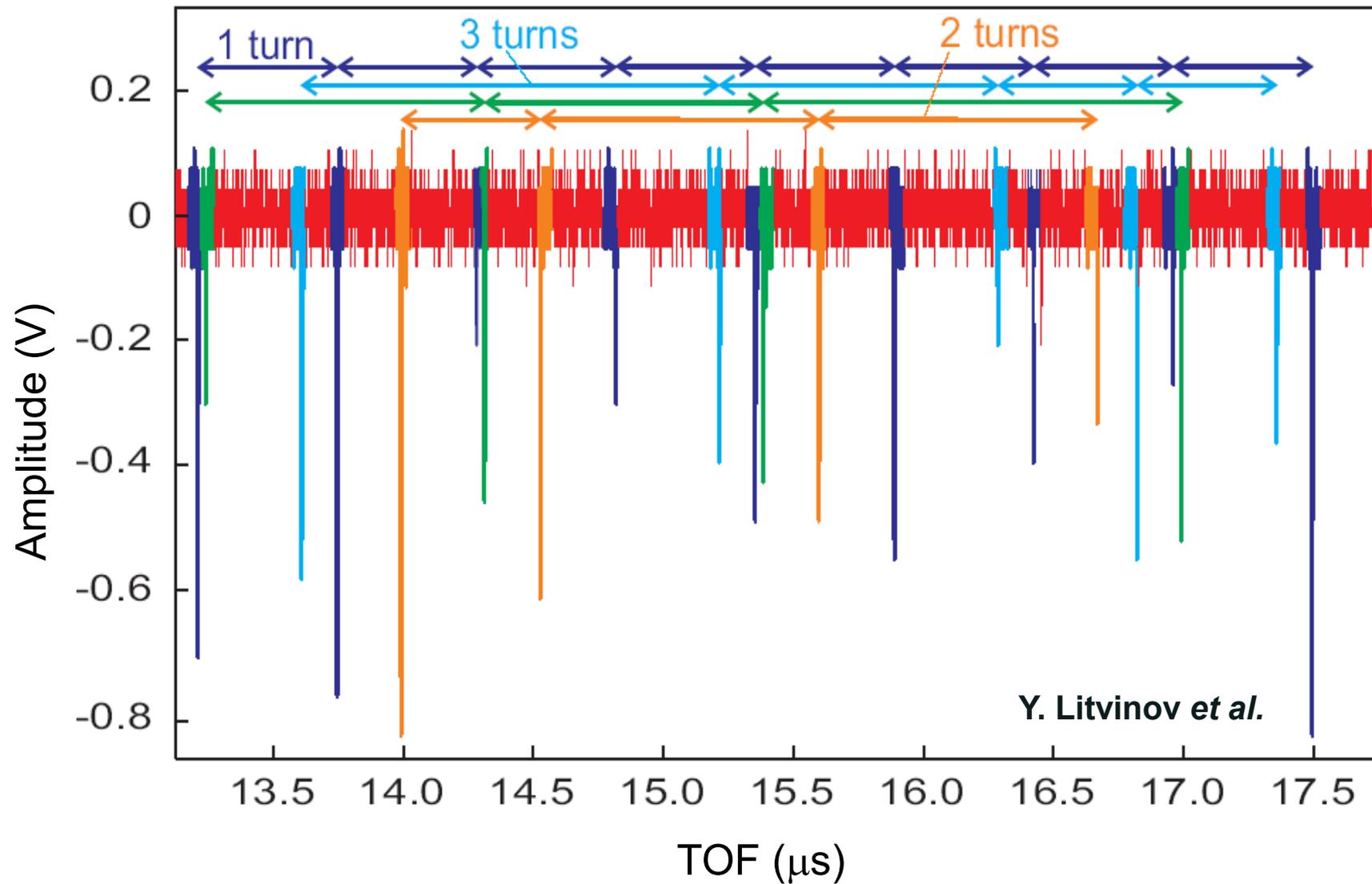


A prominent result of SMS at ESR

Decay of two hydrogen-like promethium ions of $^{142}\text{Pm}^{60+}$



Revolution frequency measurements at ESR IMS spectrum

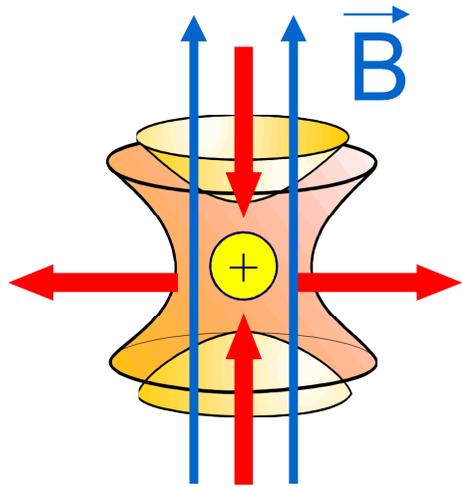


THE PENNING TRAP:
CYCLOTRON FREQUENCY
MEASUREMENTS

Penning traps

Fundamentals

Magnetic + Electrostatic field



Axial motion

$$\omega_z = \sqrt{\frac{q\Phi_0}{md^2}}$$

Reduced cyclotron motion

$$\omega_+ = \frac{1}{2} \left[\omega_c + \sqrt{(\omega_c^2 - 2\omega_z^2)} \right]$$

Magnetron motion

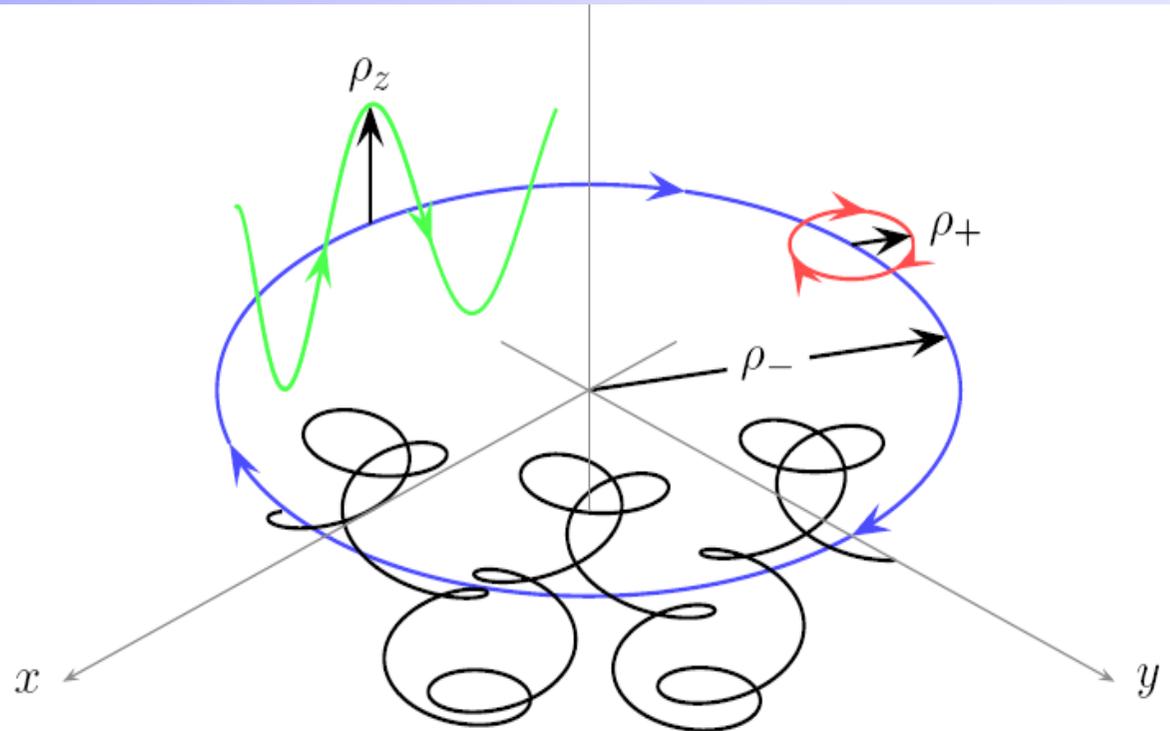
$$\omega_- = \frac{1}{2} \left[\omega_c - \sqrt{(\omega_c^2 - 2\omega_z^2)} \right]$$

$$\omega_c = \frac{q}{m} B$$

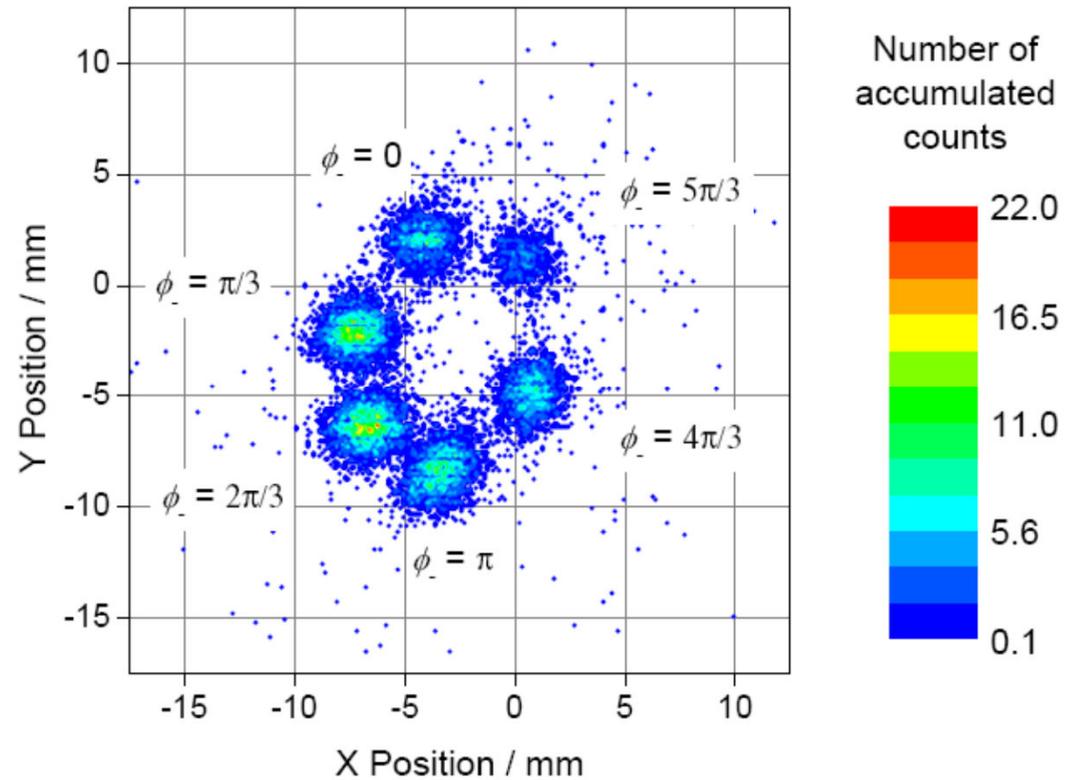
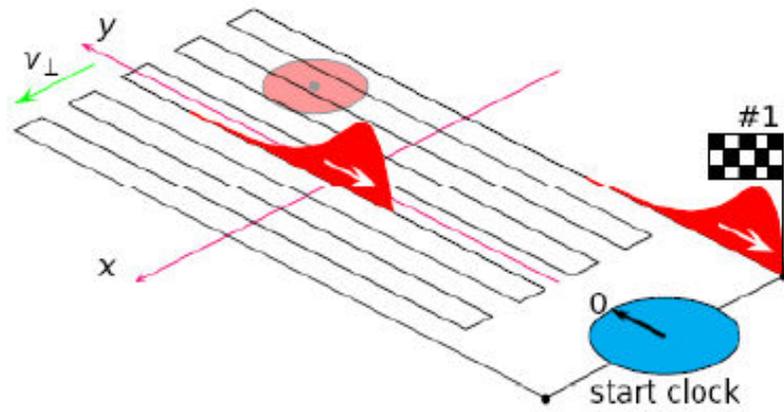
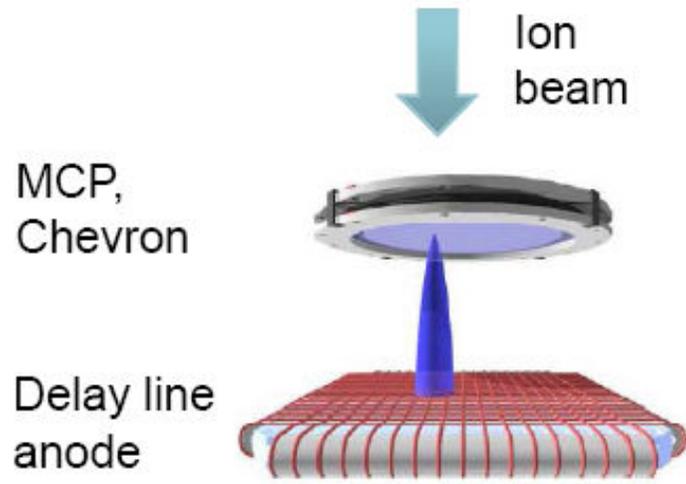
$$\omega_c = \omega_+ + \omega_-$$

$$\omega_+ > \omega_z > \omega_-$$

$$\omega_+^2 + \omega_z^2 + \omega_-^2 = \omega_c^2$$

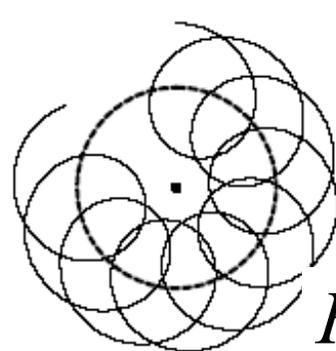


Penning traps Fundamentals

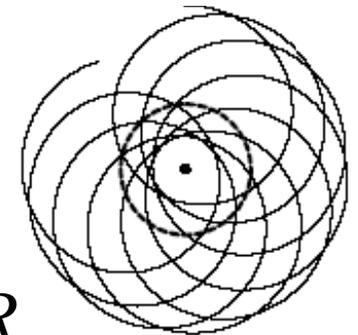


G. Eitel *et al.*, NIM A 606 (2009) 475

$$R_{+} < R_{-}$$

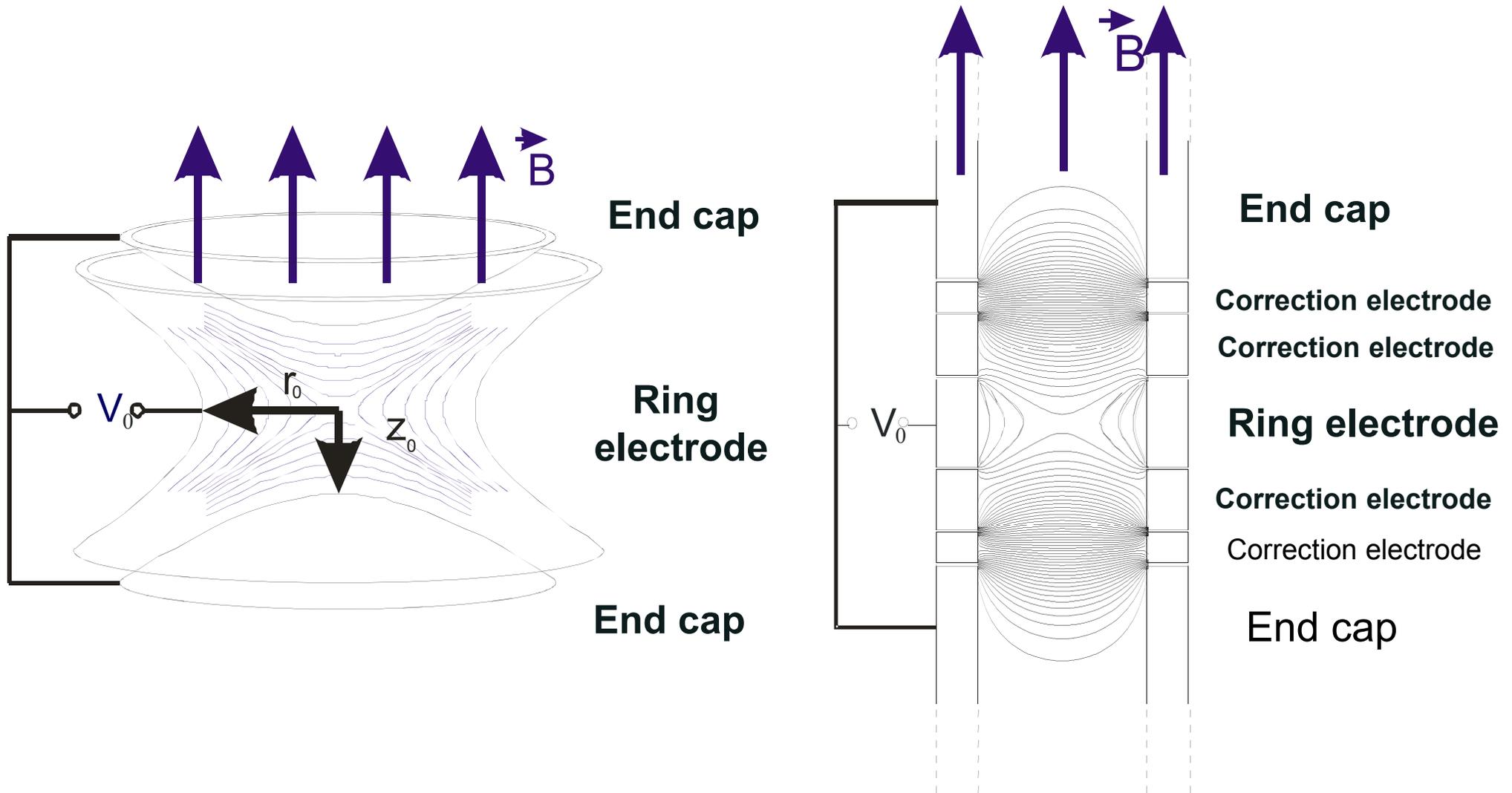


$$R_{+} > R_{-}$$



Penning traps

Different geometries

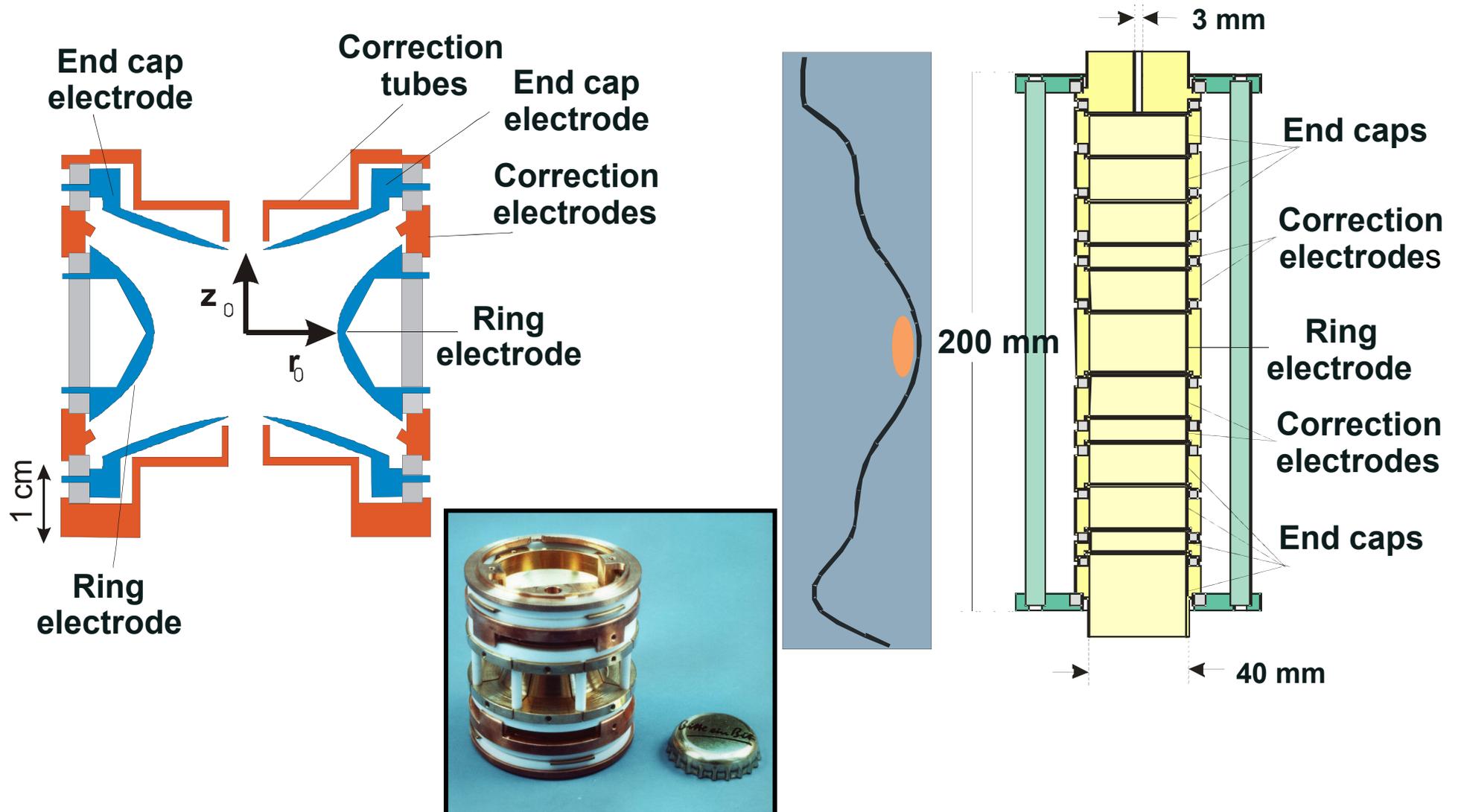


L.S. Brown and G. Gabrielse, Rev. Mod. Phys. 58 (1986) 233

Penning traps

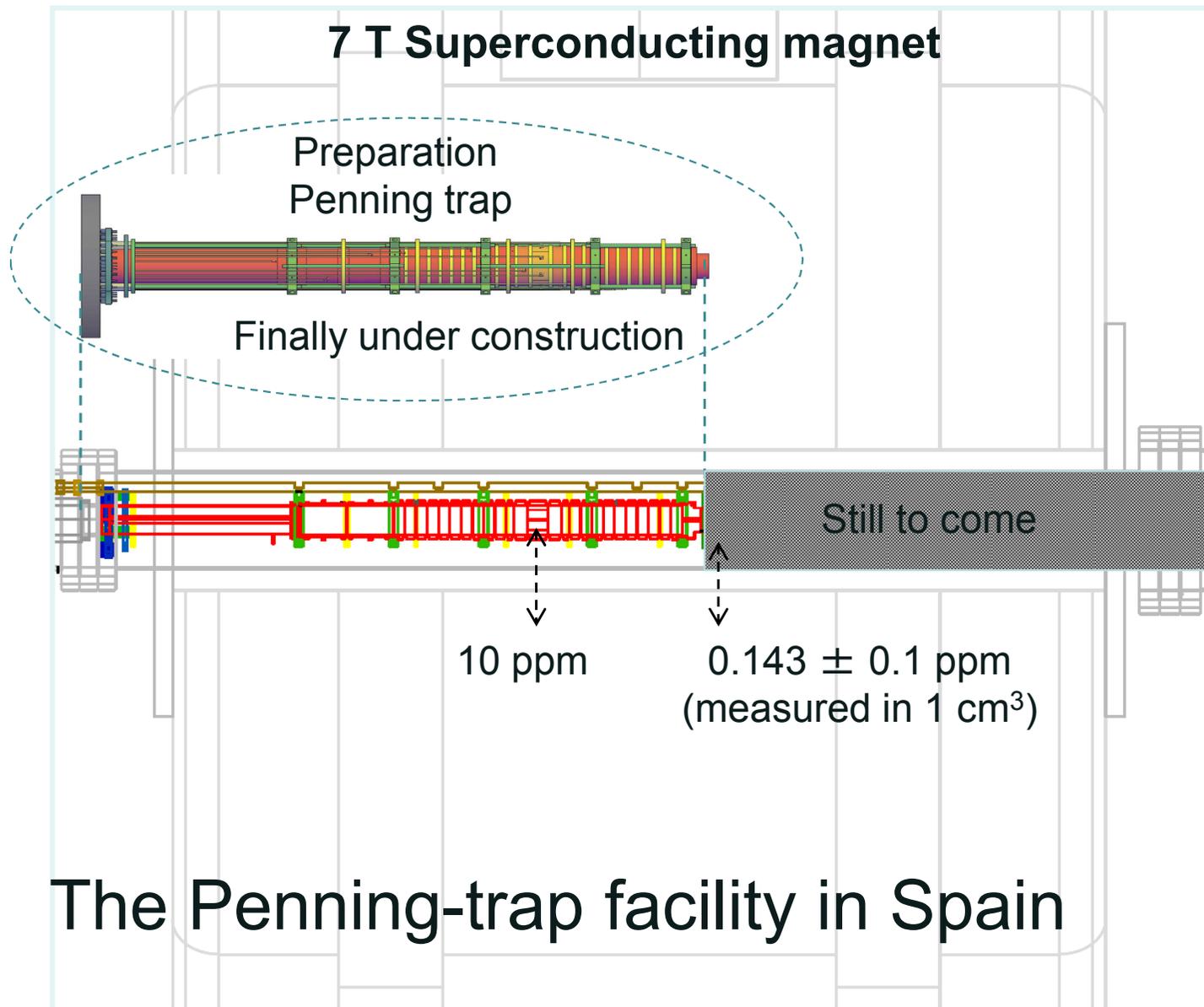
Different geometries

The Pioneering ISOLTRAP at ISOLDE



Penning traps

Different geometries

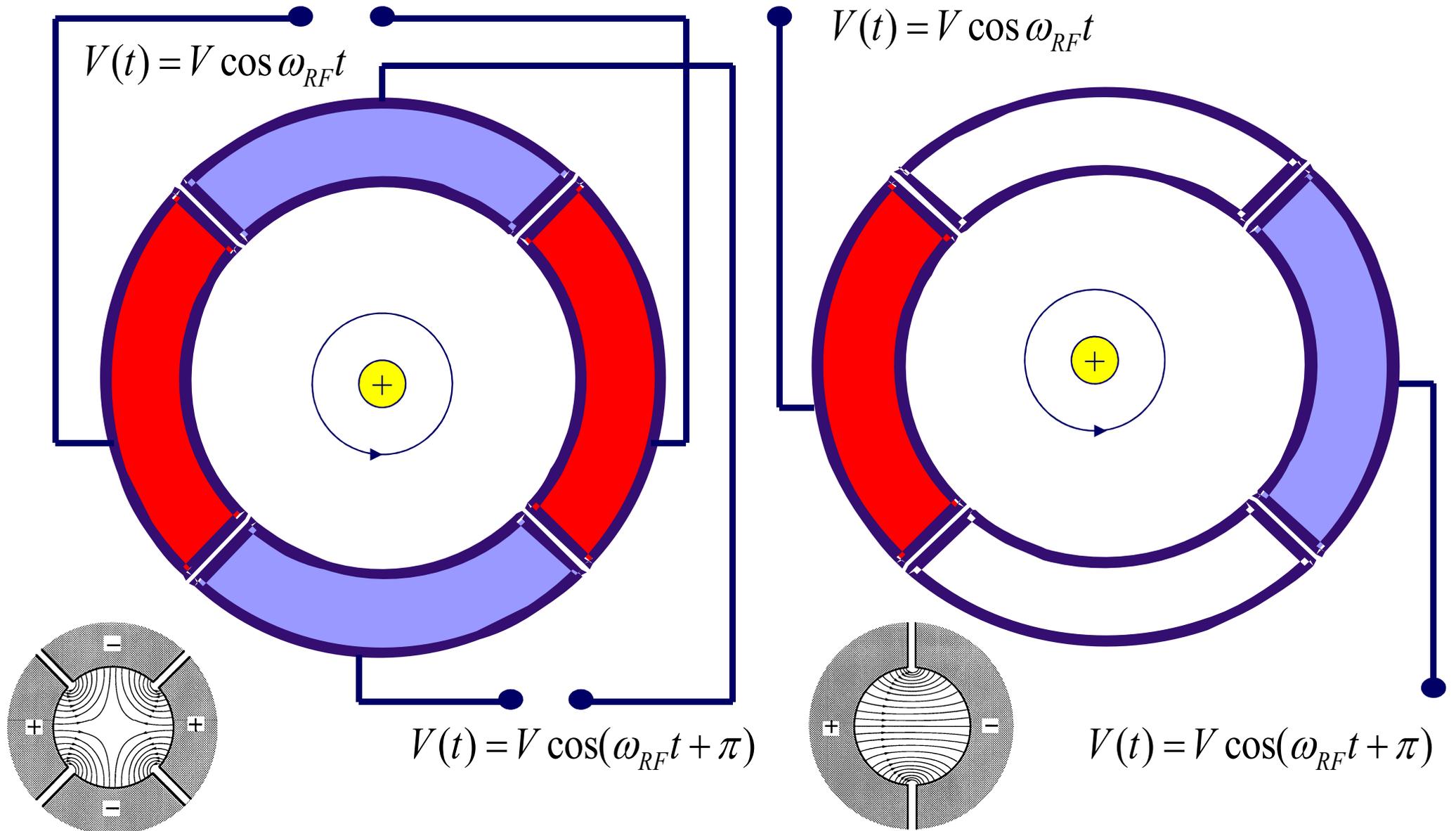


The Penning-trap facility in Spain



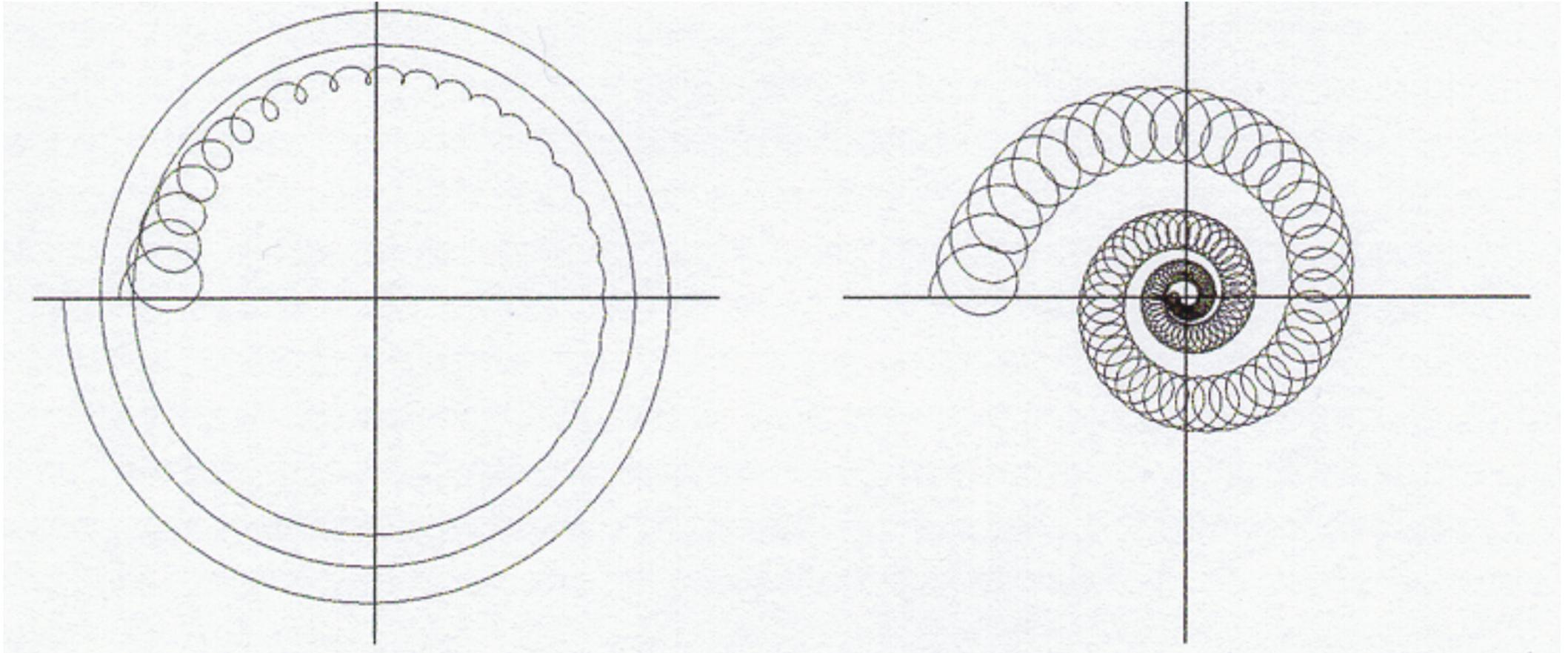
Penning-trap techniques

External Fields



Penning-trap techniques

Isobaric separation



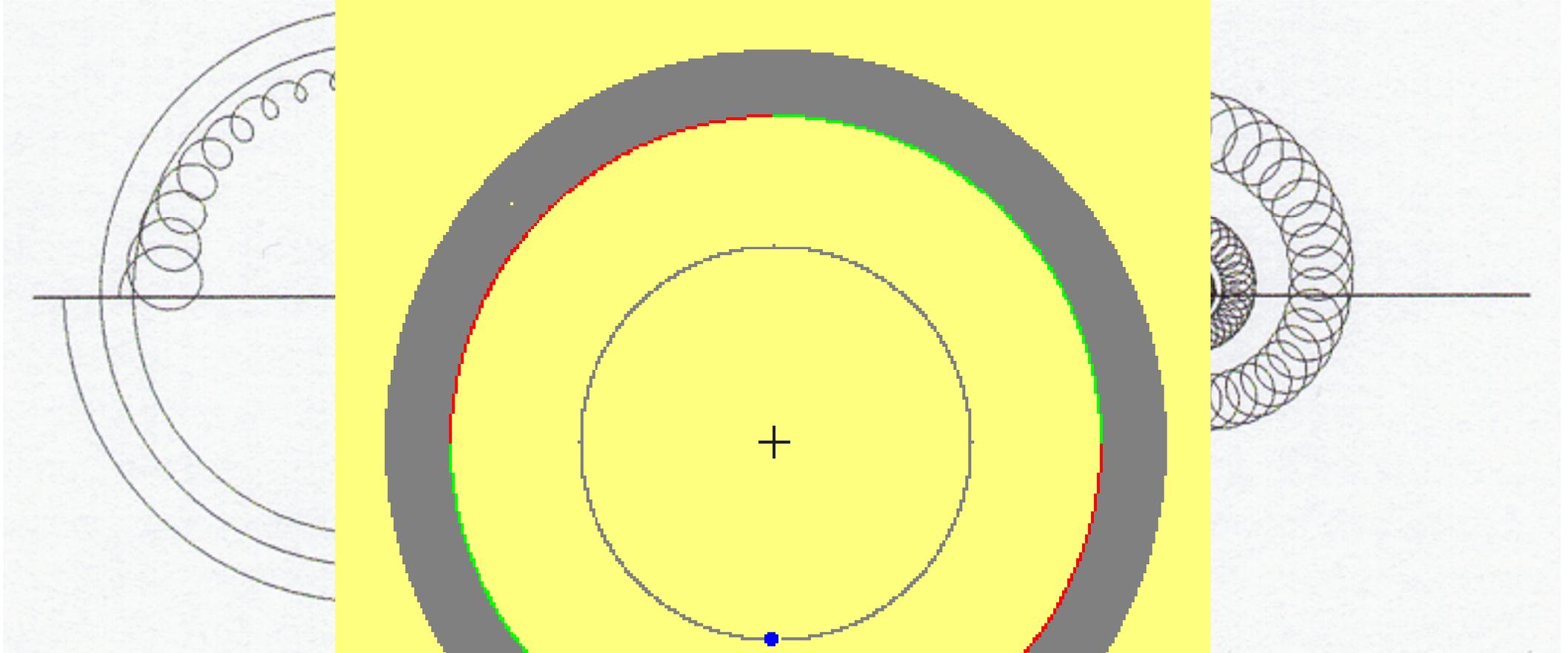
- Increase of the cyclotron radius when no excitation is applied.

- Centering of the ions after excitation at $\omega_{RF} = \omega_c$.

G. Savard *et al.*, Phys. Lett. A, 158 (1991) 247

Penning-trap techniques

Isobaric separation



- Increase of ...
when no exc

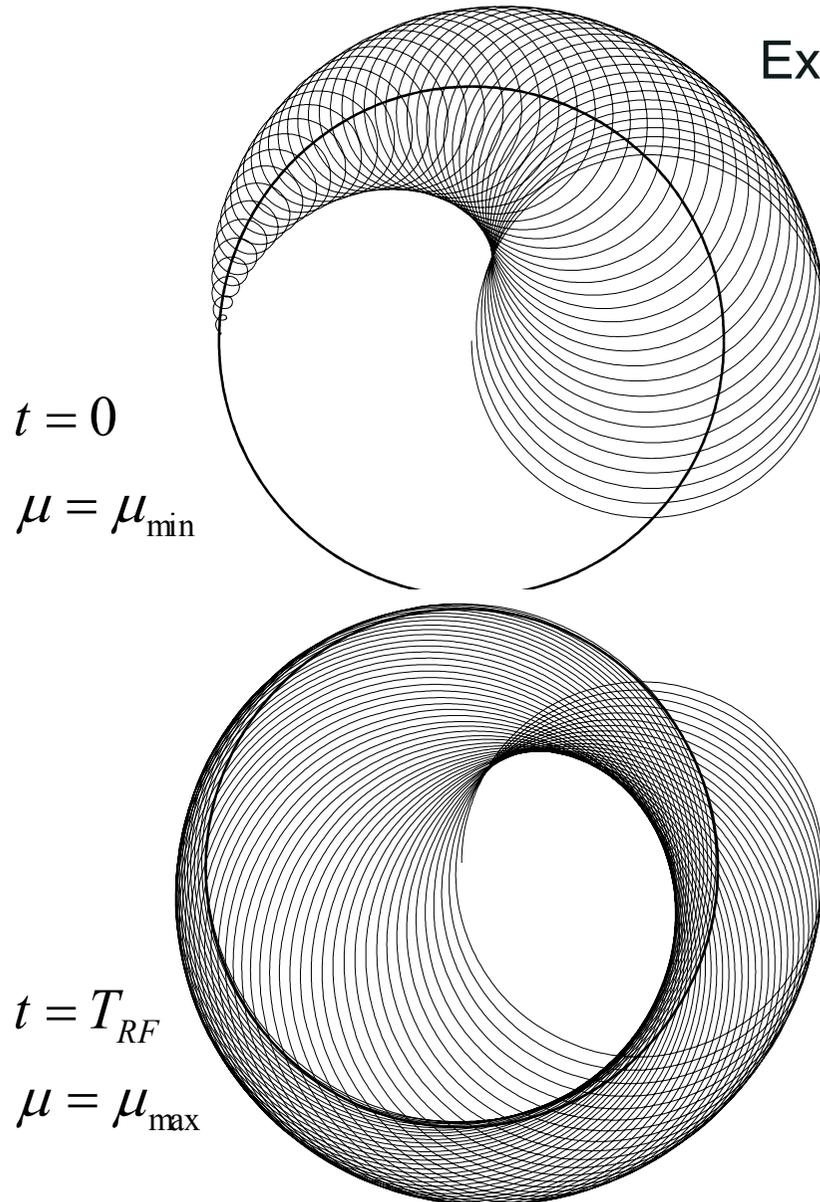
ions after
 ω_c .

G. Savard *et al.*, I

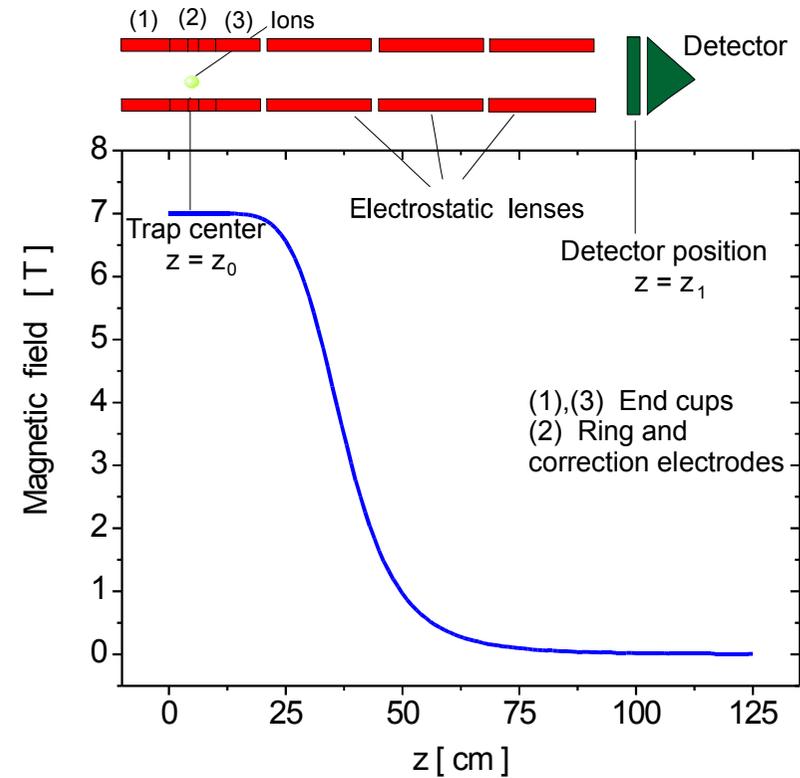
Movie Courtesy of Georg Bollen

Penning-trap techniques

Time-of-Flight Ion Cyclotron Resonance (TOF-ICR)



External RF quadrupole field with $\omega_{RF} = \omega_c (= \omega_+ + \omega_-)$



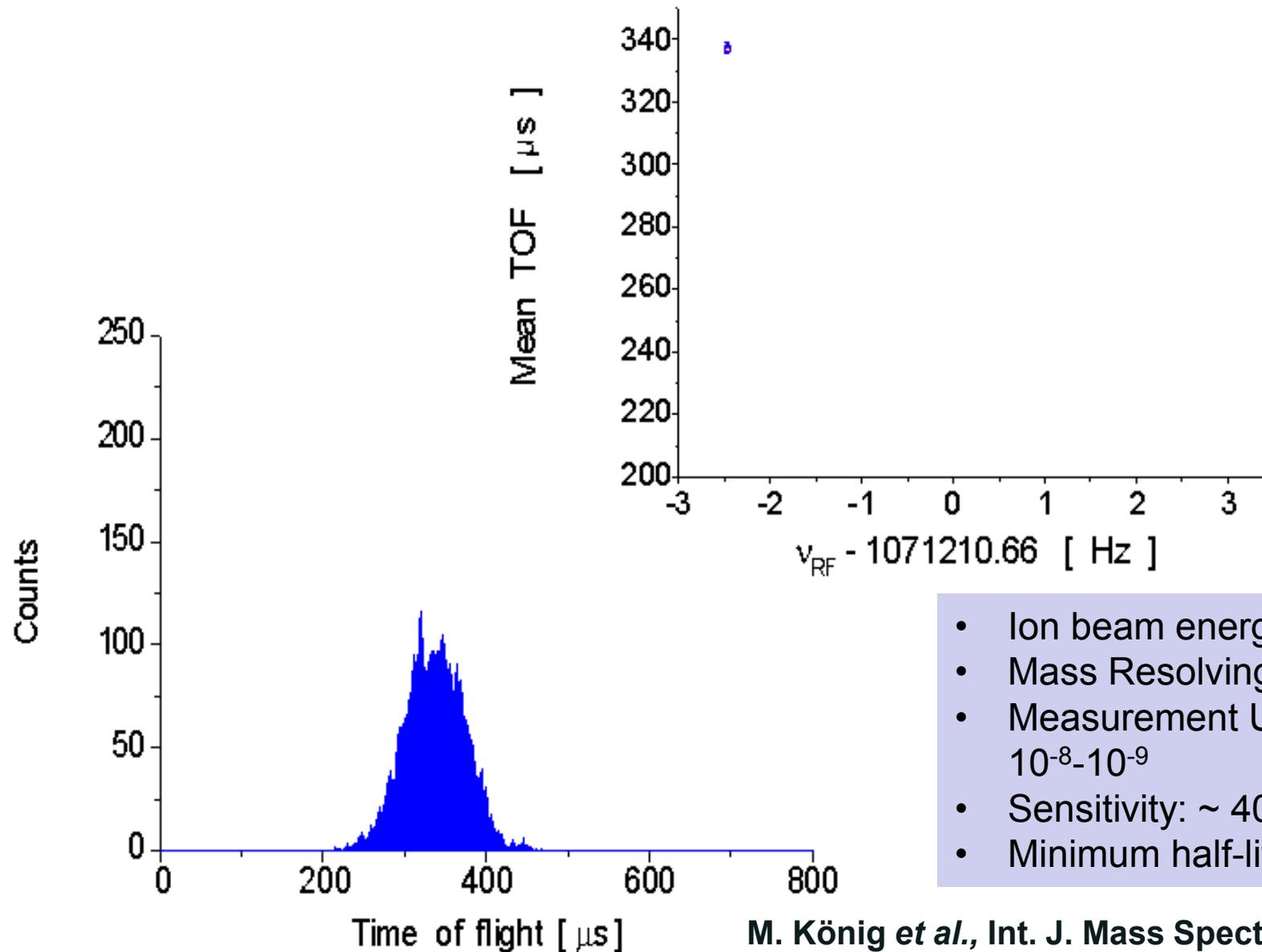
$$F_z = \mu \frac{\partial B}{\partial z}$$

$$TOF_{\min} \Leftrightarrow \mu_{\max} \Leftrightarrow \omega_{RF} = \omega_c$$

G. Gräff et al., Z. Phys. A 297 (1980) 35

Penning-trap techniques

Time-of-Flight Ion Cyclotron Resonance (TOF-ICR)

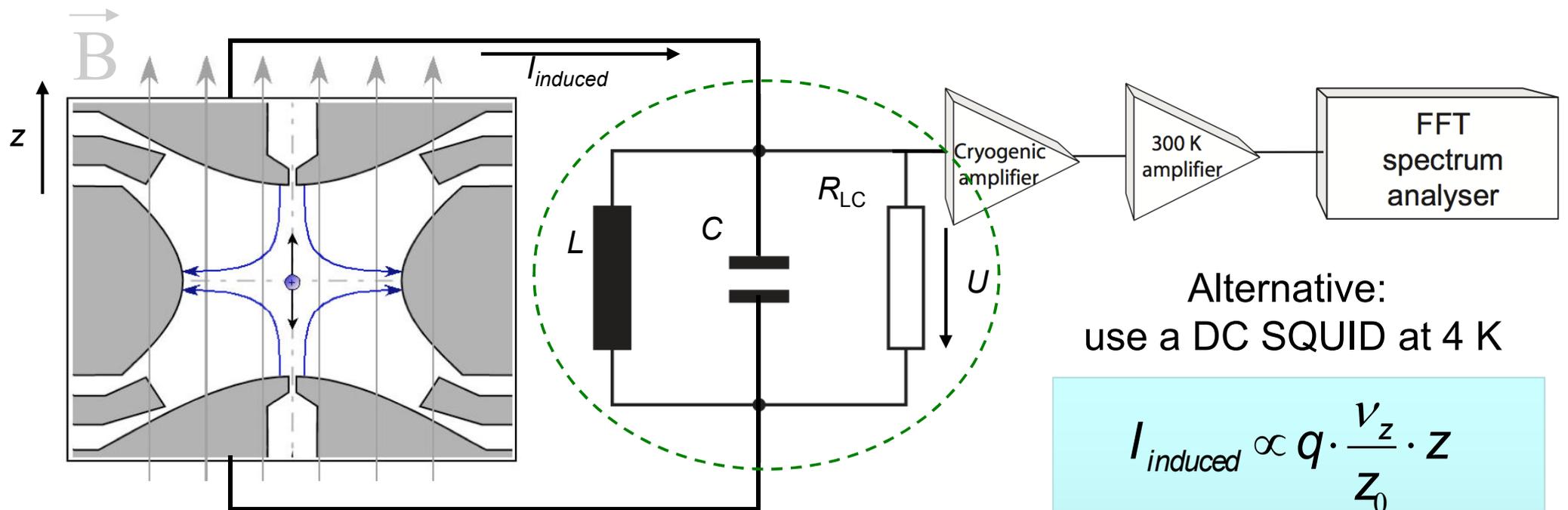


- Ion beam energy: ~ 1 eV
- Mass Resolving Power: 1×10^6
- Measurement Uncertainty: 10^{-8} - 10^{-9}
- Sensitivity: ~ 40 ions
- Minimum half-life: 8 ms

M. König *et al.*, Int. J. Mass Spectrom. 142 (1995) 95

Penning-trap techniques

Induced image current detection

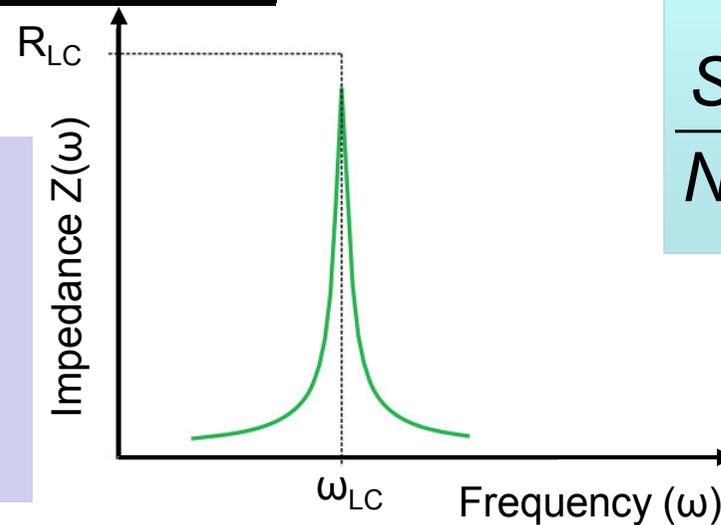


Alternative:
use a DC SQUID at 4 K

$$I_{induced} \propto q \cdot \frac{v_z}{z_0} \cdot z$$

$$\frac{S}{N} \propto q \cdot \frac{z}{z_0} \cdot \sqrt{\frac{v_z}{\Delta v}} \cdot \sqrt{\frac{Q}{TC}}$$

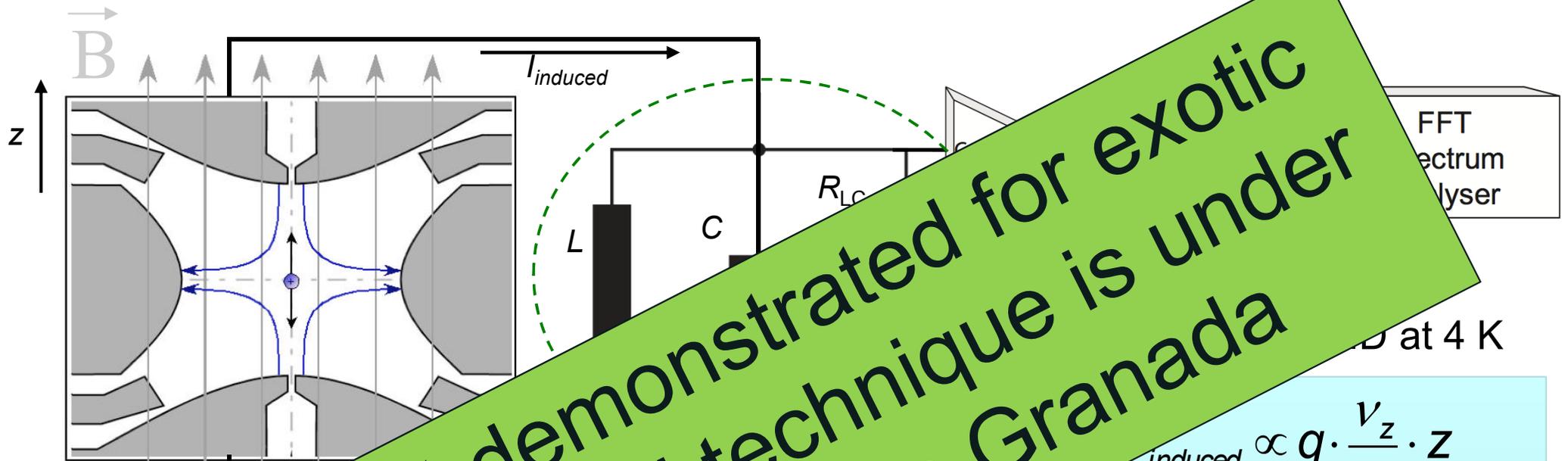
$$Z = \frac{Q}{2\pi\nu C}$$



- Ion beam energy: ~ 1 eV
- Measurement Uncertainty: 10^{-11}
- Sensitivity: ~ 1 ions
- Minimum half-life: (to be determined)

Penning-trap techniques

Induced image current detection

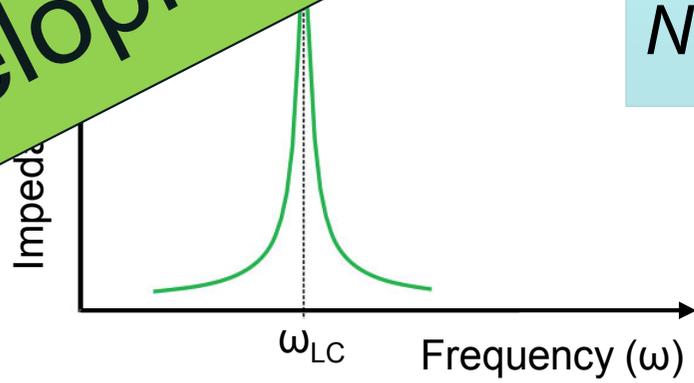


Not yet demonstrated for exotic ions. A novel technique is under development in Granada

- Ion beam
- Measurement up to 10
- Sensitivity
- Minimum height (determined)

$$I_{\text{induced}} \propto q \cdot \frac{v_z}{z_0} \cdot z$$

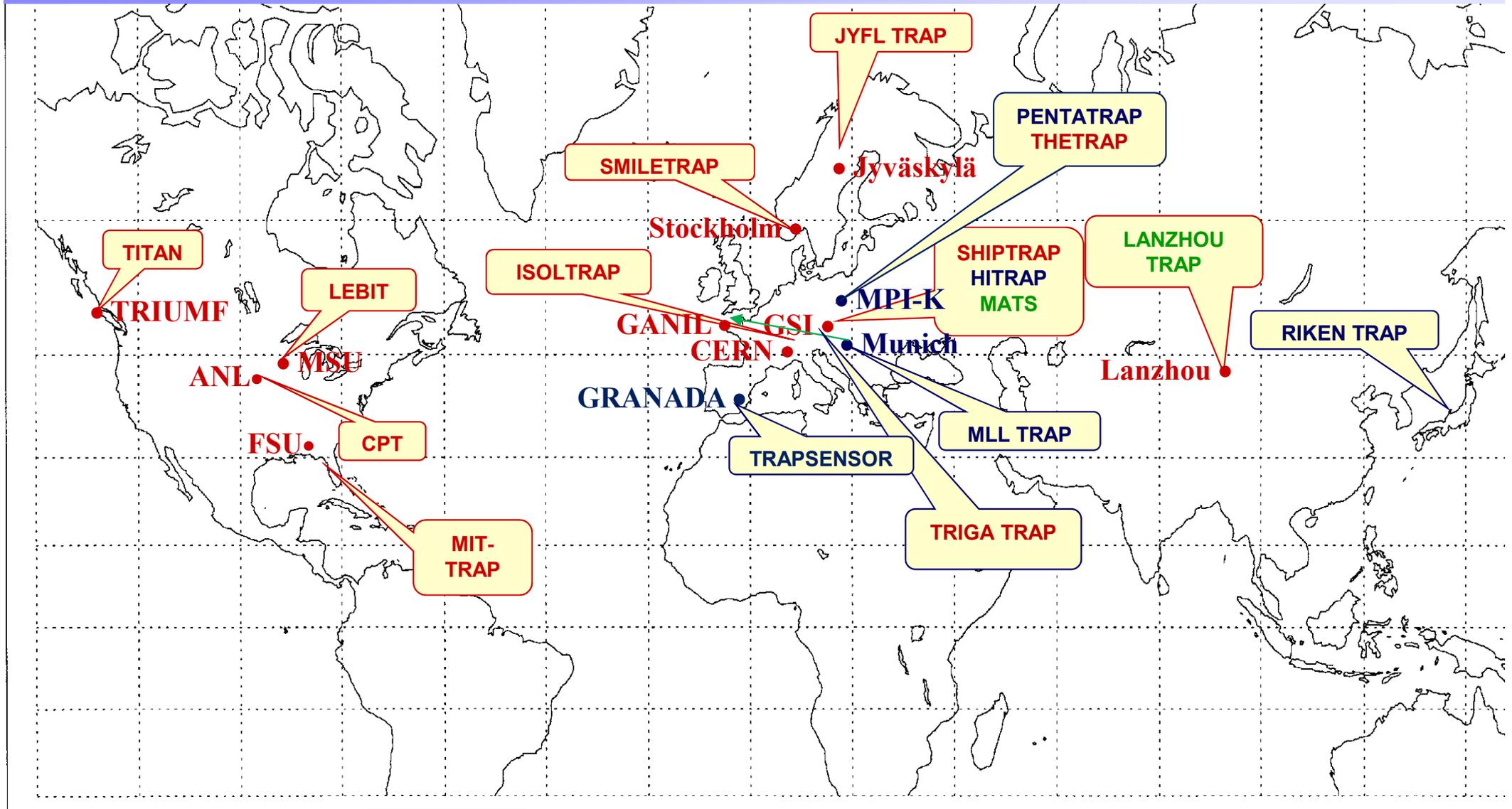
$$\frac{S}{N} \propto q \cdot \frac{z}{z_0} \cdot \sqrt{\frac{v_z}{\Delta v}} \cdot \sqrt{\frac{Q}{TC}}$$



$$Z = \frac{Q}{2\pi\nu C}$$

Existing Penning-trap facilities

Mass spectrometry of atomic nuclei

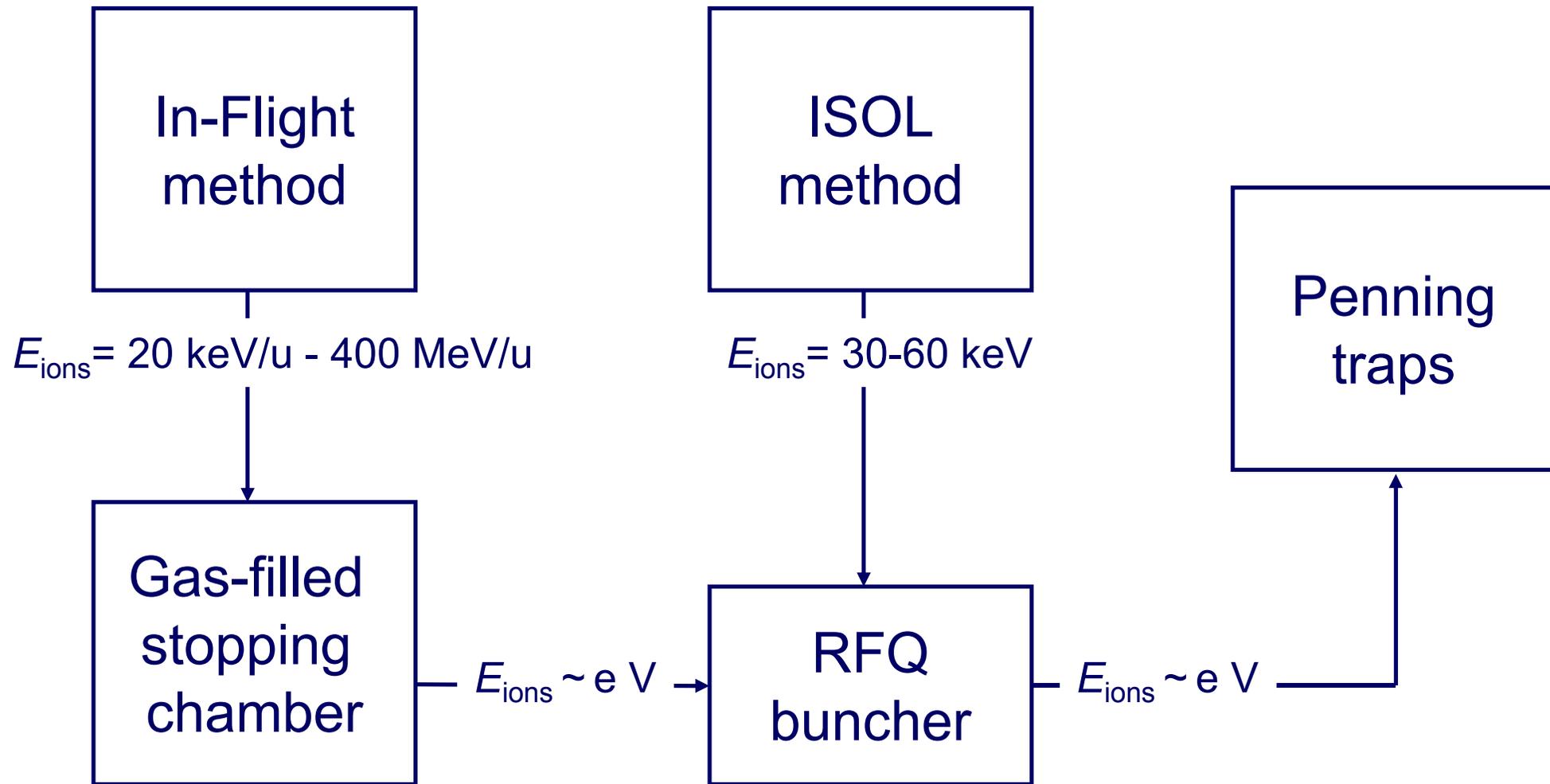


● **Running facilities**

● **Under construction**

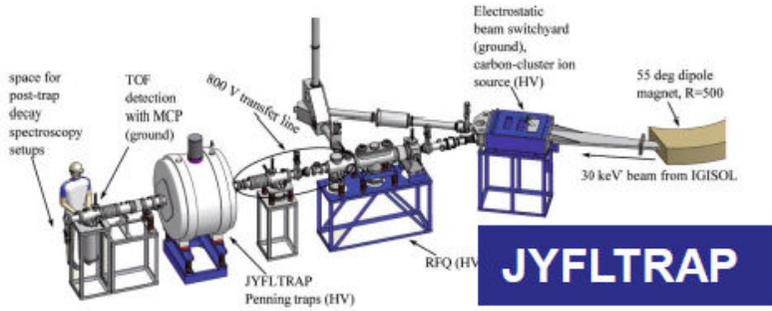
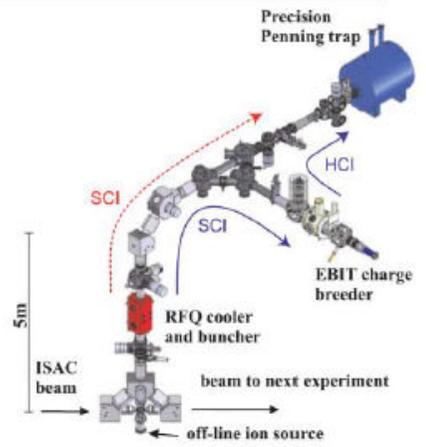
● **Future facilities**

Coupling Penning traps to RIB facilities

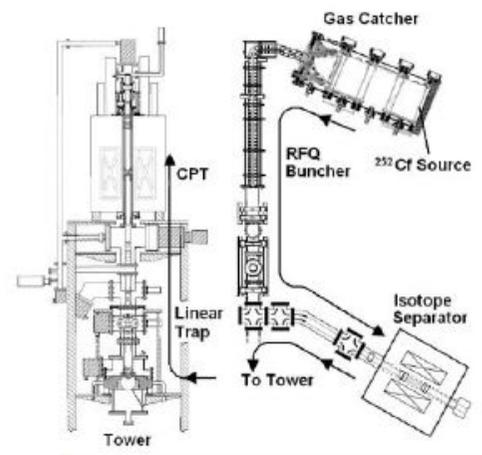


Growing facilities

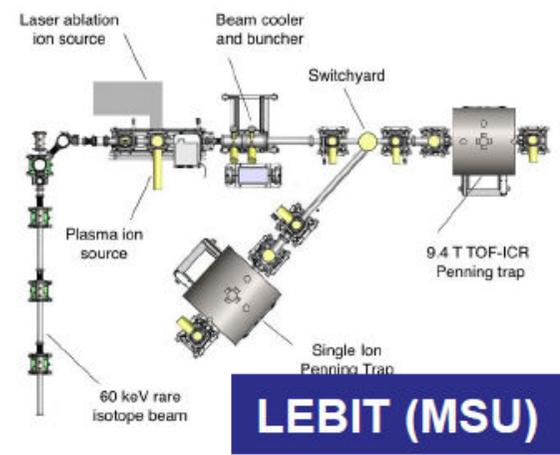
TITAN (TRIUMF)



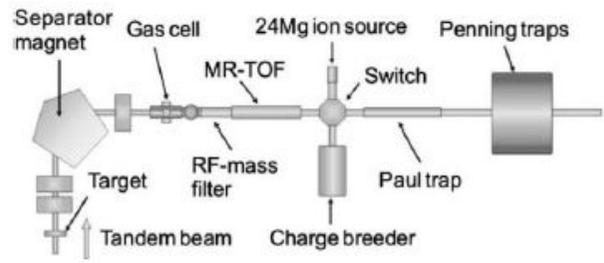
JYFLTRAP



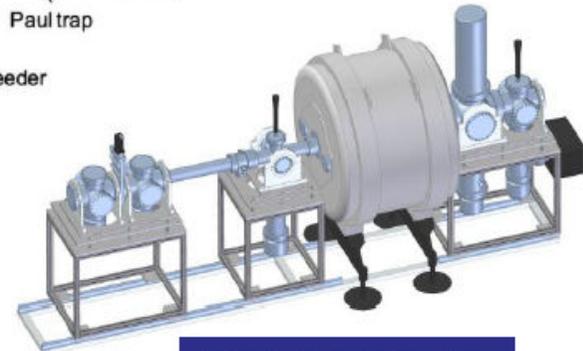
CPT (ARGONNE)



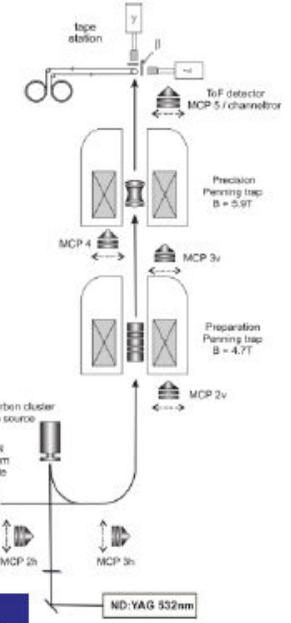
LEBIT (MSU)



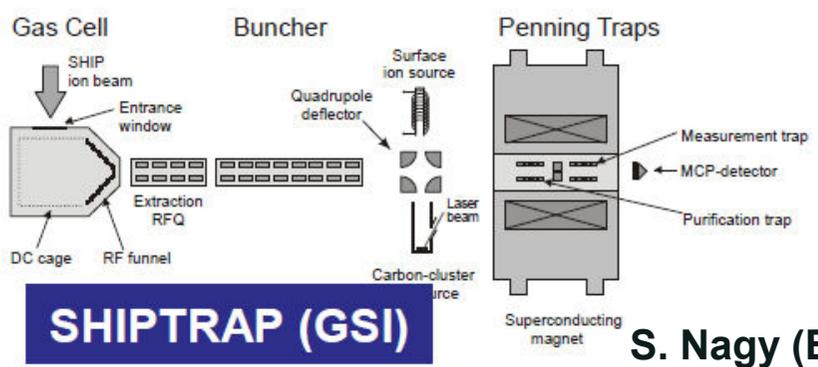
MLLTRAP



TRIGA-TRAP



ISOLTRAP (CERN)

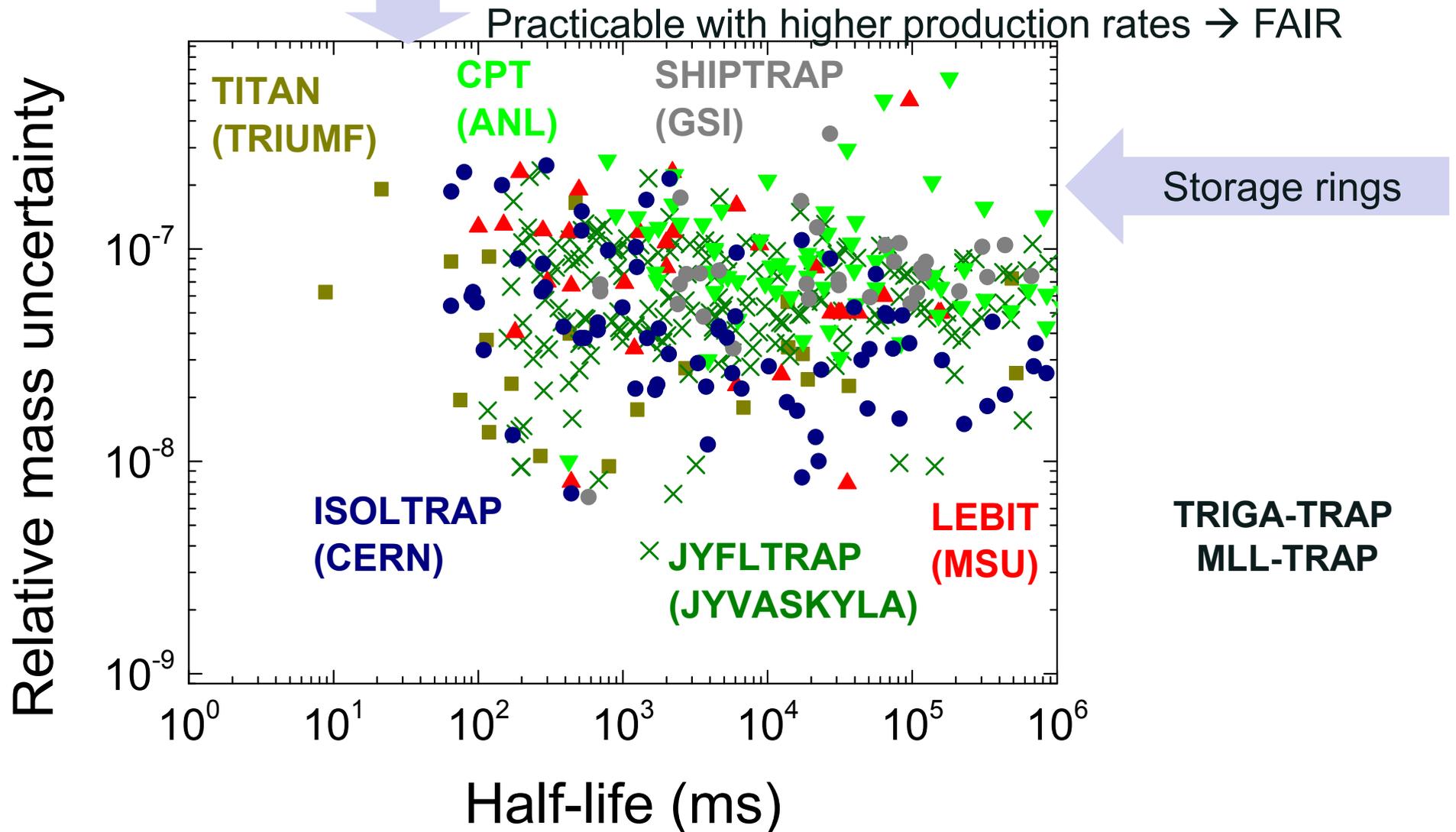


SHIPTRAP (GSI)

S. Nagy (EMIS2012)

Physics survey

Penning-trap mass measurements at RIB

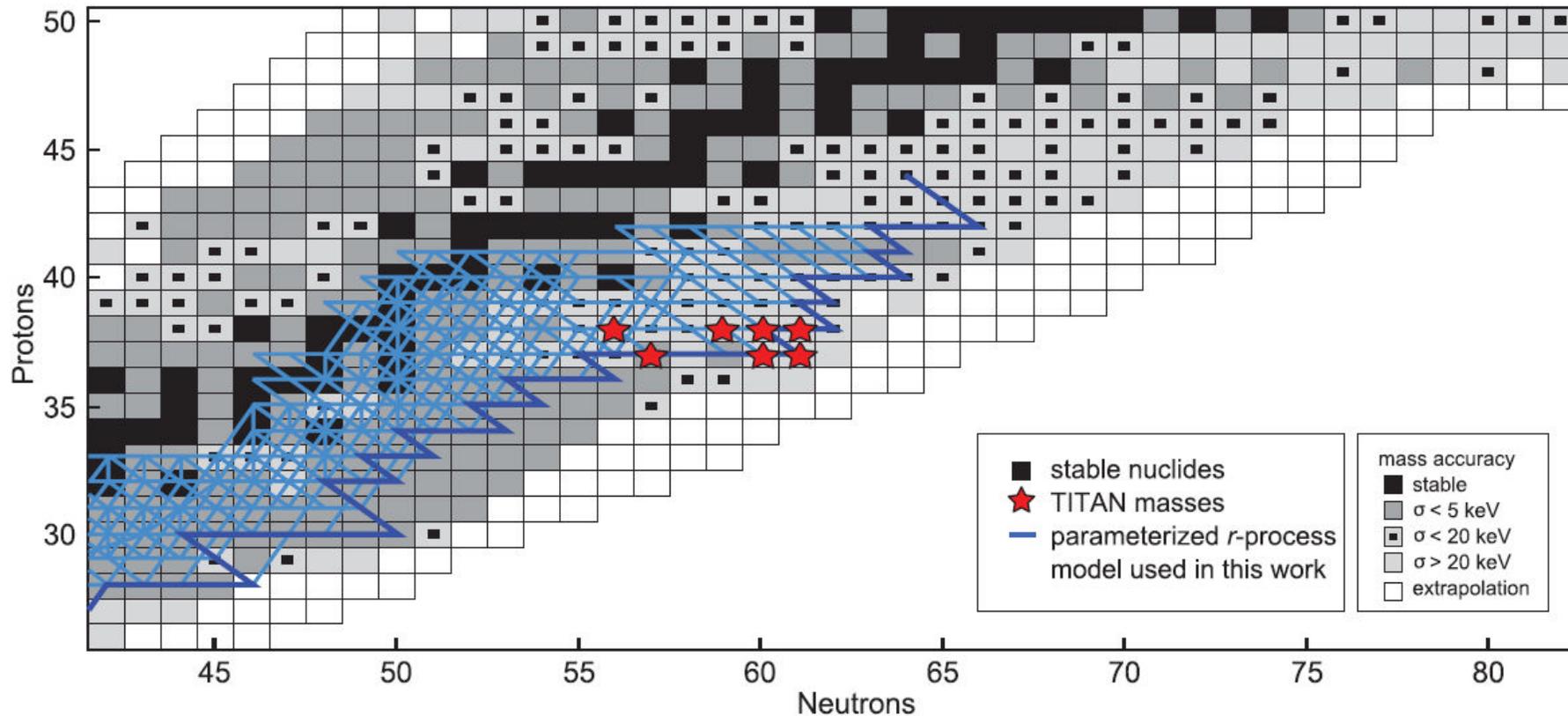


High-lights 2011-2012: S. Eliseev *et al.*, PRL 106 (2011) 052504, E. Haettner *et al.*, PRL 106 (2011) 122501, S. Eliseev *et al.*, PRL 107 (2011) 152501, S. Ettenauer *et al.*, PRL 107 (2011) 272501, M. Brodeur *et al.*, PRL 108 (2012) 052504, D. Fink *et al.*, PRL 108 (2012) 062502, M. Brodeur *et al.*, PRL 108 (2012) 212501, J. Hakala *et al.*, PRL 109 (2012) 032501, A. T. Gallant *et al.*, PRL 109 (2012) 032506, E. Minaya Ramirez *et al.*, Science 337 (2012) 1207.

Recent highlights (TITAN)

Nuclear (Penning-trap) astrophysics at TRIUMF (June 2012)

rp-process (highly-charged ions)

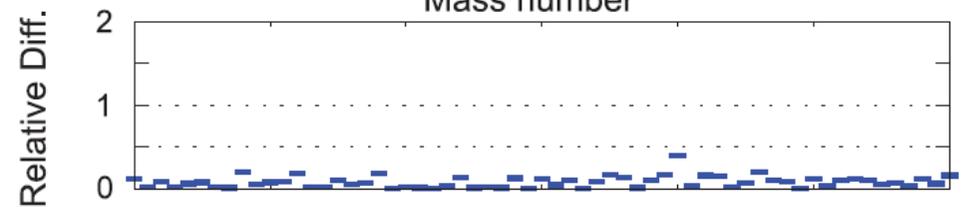
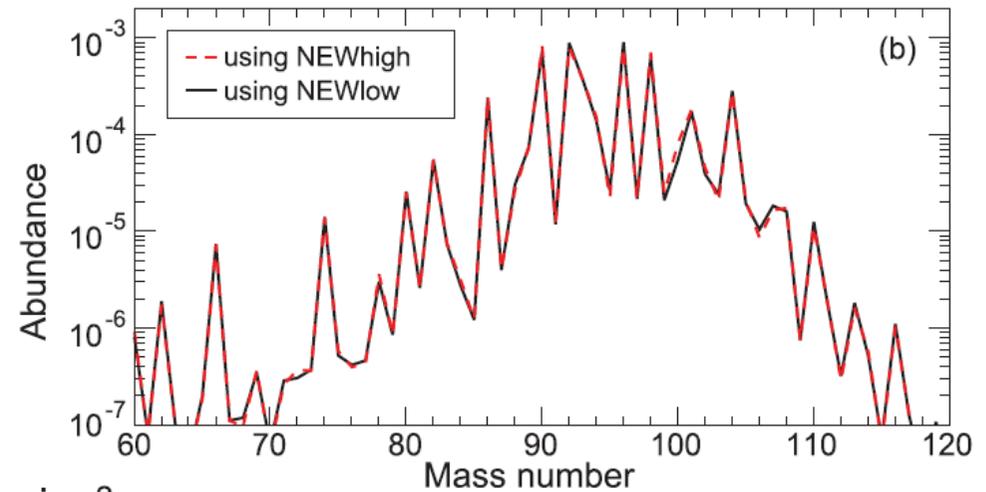
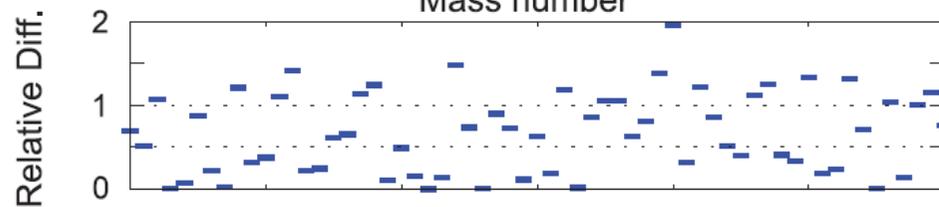
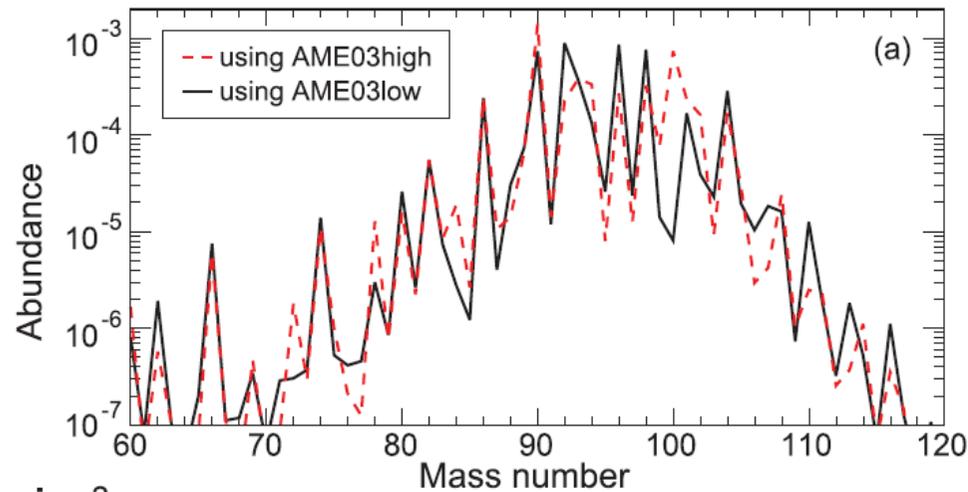
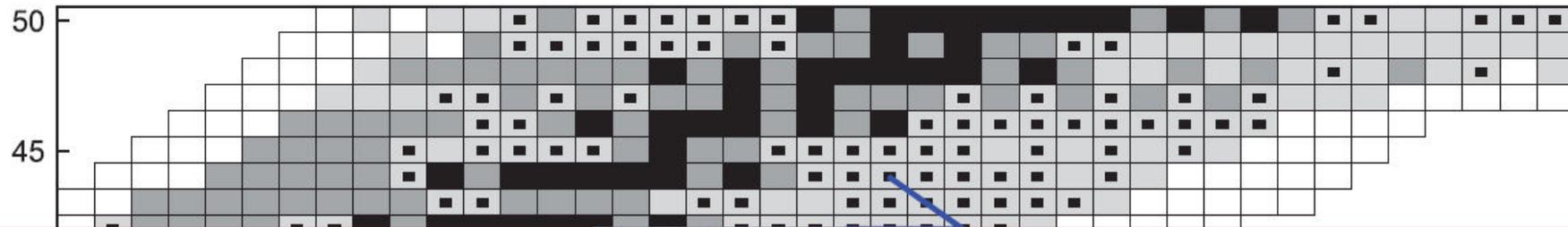


V. Simon *et al.*, Phys. Rev. C 85, 064308 (2012)

Recent highlights (TITAN)

Nuclear (Penning-trap) astrophysics at TRIUMF (June 2012)

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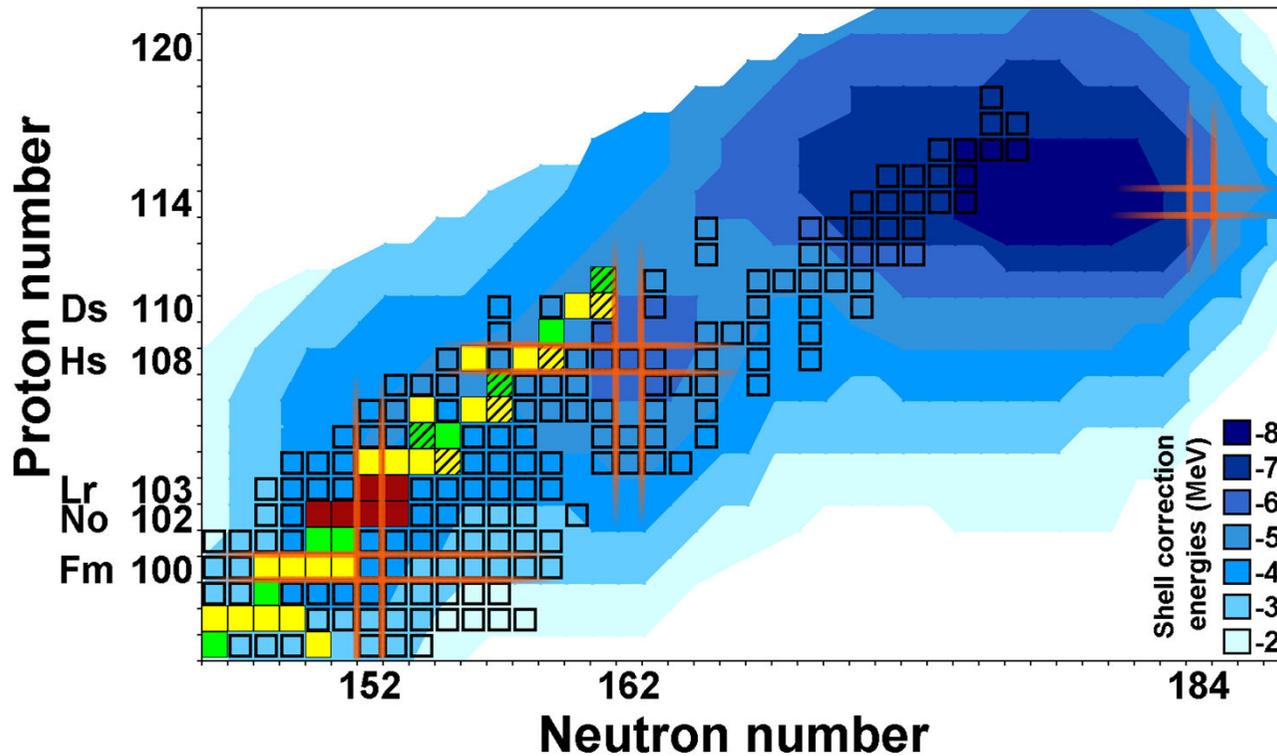


V. Simon *et al.*, Phys. Rev. C 85, 064308 (2012)

Recent highlights (SHIPTRAP)

NUSTAR (Penning-trap) experiment at GSI (September 2012)

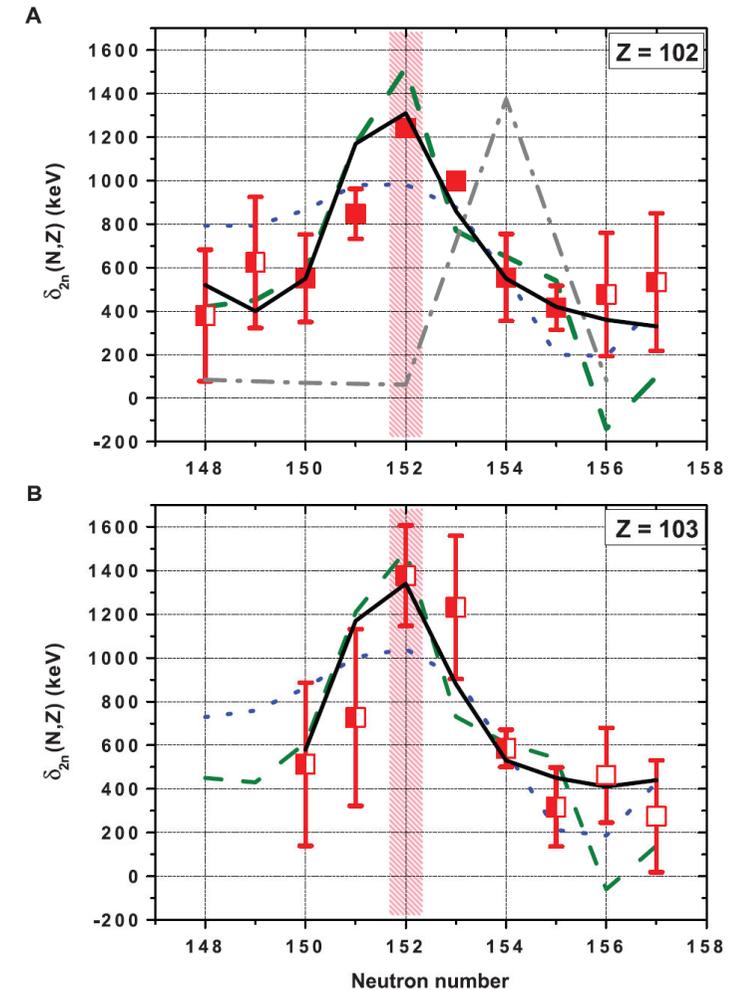
Nuclear Structure (minute production)



- The isotope with lowest production rate ever measured in a Penning trap (^{256}Lr , 60(18) nb) \rightarrow 48 ions detected in 93 hours !!!

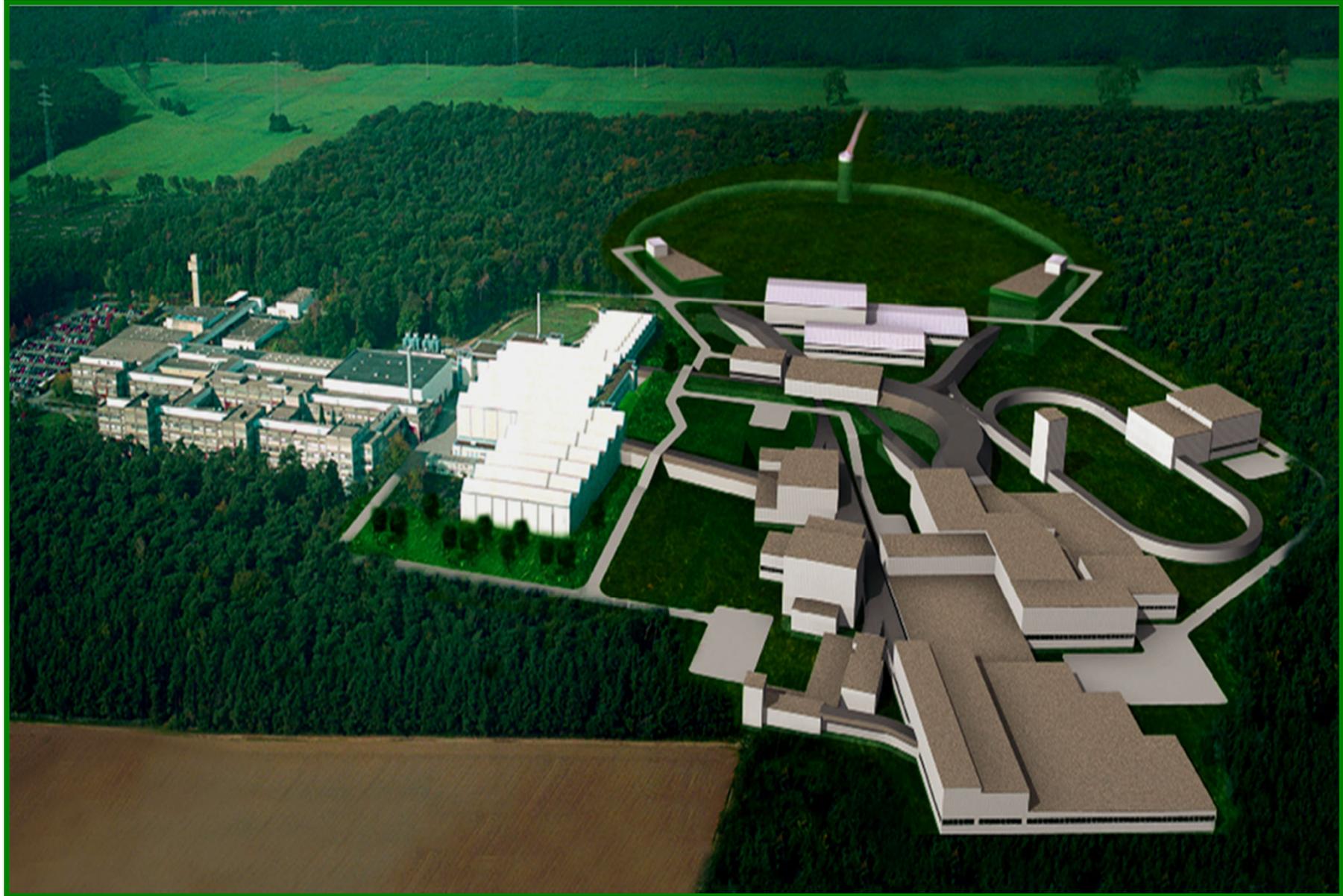
E. Minaya Ramirez *et al.*, Science 337 (2012) 1207

$$\begin{aligned} \delta_{2n}(N,Z) &= S_{2n}(N,Z) - S_{2n}(N+2,Z) \\ &= -2M_{\text{exc}}(N,Z) + M_{\text{exc}}(N-2,Z) \\ &\quad + M_{\text{exc}}(N+2,Z), \end{aligned}$$

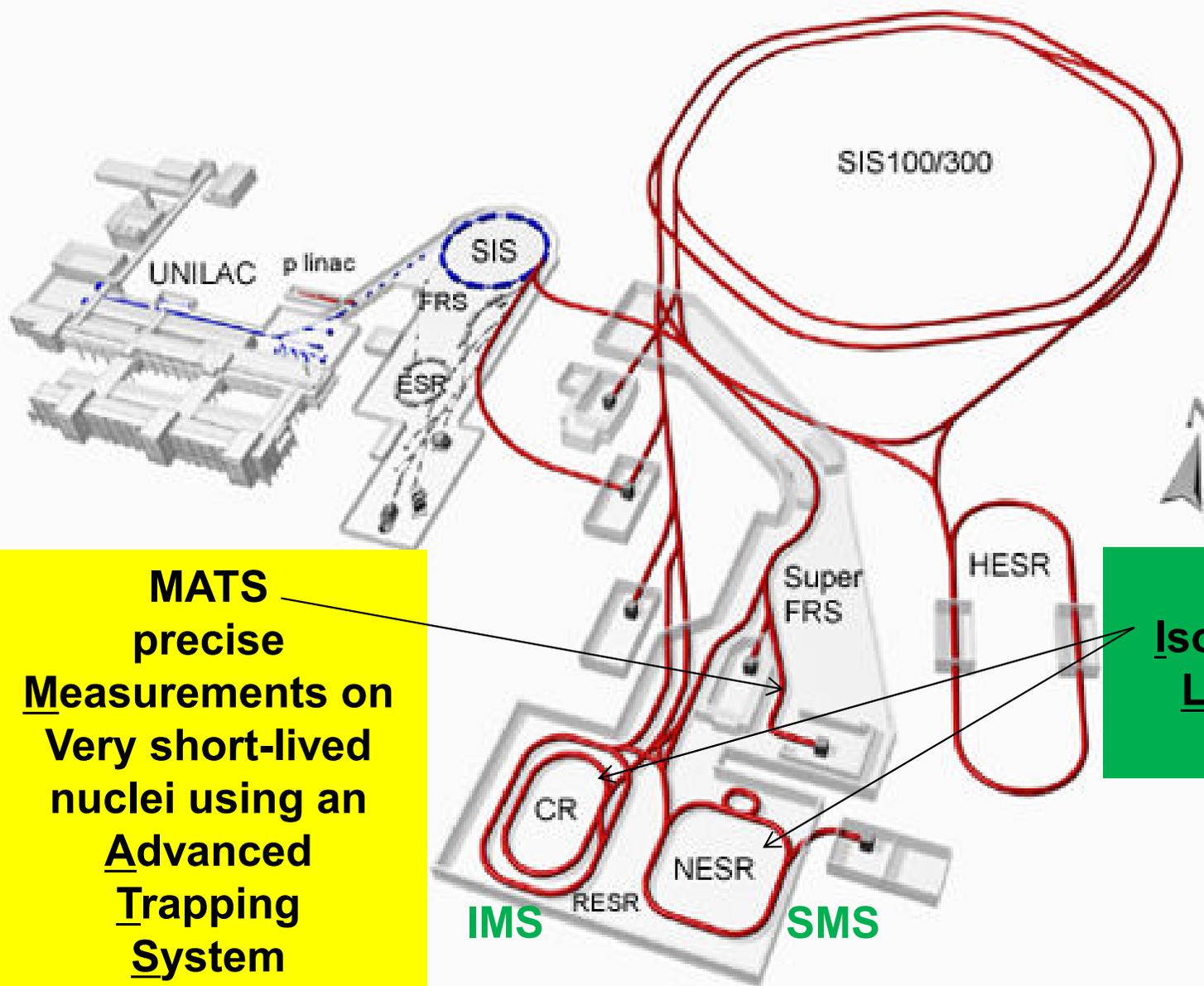


MASS MEASUREMENTS ON EXOTIC NUCLEI IN THE FUTURE

Future Facility for Antiprotons and Ions Research (FAIR) Higher Production yields



Future Facility for Antiprotons and Ions Research (FAIR) Higher Production yields



MATS
precise
Measurements on
Very short-lived
nuclei using an
Advanced
Trapping
System

ILIMA
Isomeric Beams
Lifetimes and
Masses

