

# QCD and applications to nuclear physics and matter in extreme conditions



**IFIC**  
INSTITUT DE FÍSICA  
CORPUSCULAR



Miguel Albaladejo (IFIC)

Workshop CSIC Strategic Theme  
“Understanding the basic components  
of the Universe, its structure and evolution”  
Madrid Feb. 22-23, 2023



## SOLVING QUANTUM CHROMO- DYNAMICS

### CHAPTER 3

#### Coordinators

M. Pilar Hernández  
(IFIC, CSIC - UV, Coordinator)

Laura Tolos  
(ICE, CSIC - IEEC, Adjunct Coordinator)

#### Researchers and Centers

(in alphabetical order)

José F. Barbón (IFT, CSIC - UAM)

Margarita García-Pérez  
(IFT, CSIC - UAM)

Oscar García Prada  
(ICMAT, CSIC - UAM - UCM - UC3M)

Antonio González-Arroyo  
(IFT, CSIC - UAM)

Esperanza López (IFT, CSIC - UAM)

Cristina Manuel (ICE, CSIC - IEEC)

Juan Nieves (IFIC, CSIC - UV)

Carlos Pena (IFT, CSIC - UAM)

Alberto Ramos (IFIC, CSIC - UV)

Germán Rodrigo (IFIC, CSIC - UV)

Agustín Sabio Vera (IFT, CSIC - UAM)

Lacking an analytical solution of QCD, a multi-prong strategy has been implemented that includes a variety of tools:

- 1 the **perturbative treatment** whenever high-momentum scales are involved,
- 2 the construction of **effective field theories** that exploit symmetries in a maximal way,
- 3 and a **first-principles** formulation, based on the introduction of a space- time lattice, that allows, in particular, for a **numerical solution** [LQCD].

Notwithstanding, a formal solution to the theory continues to be a key challenge for the future both in physics and mathematics.

### In this talk...

In Alberto's very nice previous talk we have already seen quite a lot about points 1 and 3. In this talk, I will focus more on point **EFTs** (point 2), and its application to **hadron spectroscopy** and **extreme conditions**, without forgetting their **mutual connection**.

# Effective Field Theories

- QCD lagrangian built in terms of (degrees of freedom) quarks and gluons:  $\mathcal{L}[q, \bar{q}, g]$
- Quarks are **confined** into colorless hadrons, which are ultimately the objects detected in experiments.
- EFTs: build QFT with hadrons ( $\pi, K, \dots, D, B, \dots$ ) as degrees of freedom

## Example 1: Chiral Perturbation Theory

- Light quarks are... light:  $m_u \simeq m_d \simeq m_s \simeq 0$
- 8 Goldstone bosons SU(3):  $\phi_i \rightarrow \pi, K, \eta$
- [Weinberg, PRL,17,616('66); Phys. A96,327 ('79)]  
[Gasser, Leutwyler, AP,158,142('84); NP,B250,465('85)]
- LO lagrangians:

$$\mathcal{L}_2 = \frac{f^2}{4} \langle D_\mu U^\dagger D^\mu U \rangle + \frac{f^2}{4} \langle \chi^\dagger U + \chi U^\dagger \rangle$$

$$U[\phi] = \exp \left( i \sum_{i=1}^8 \phi_i \lambda_i \right)$$

$$\chi = 2B_0 \text{diag} (m_u, m_d, m_s)$$

$$\simeq \text{diag} (m_\pi^2, m_\pi^2, 2m_K^2 - m_\pi^2)$$

*Symmetry* **2022**, *14*, 1884. <https://doi.org/10.3390/sym14091884>

Article

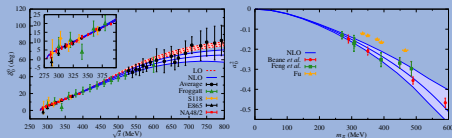
### Weinberg's Compositeness<sup>†</sup>

Ubirajara van Kolck<sup>1,2</sup>

<sup>1</sup> Université Paris-Saclay, CNRS/IN2P3, IJCLab, 91405 Orsay, France; vankolck@ijclab.in2p3.fr

<sup>2</sup> Department of Physics, University of Arizona, Tucson, AZ 85721, USA

<sup>†</sup> Dedicated to the memory of Steven Weinberg, who always chose the right degrees of freedom.



# Effective Field Theories

- QCD lagrangian built in terms of (degrees of freedom) quarks and gluons:  $\mathcal{L}[q, \bar{q}, g]$
- Quarks are **confined** into colorless hadrons, which are ultimately the objects detected in experiments.
- EFTs: build QFT with hadrons ( $\pi, K, \dots, D, B, \dots$ ) as degrees of freedom

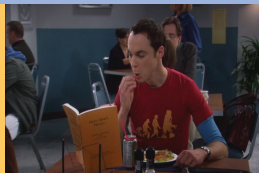
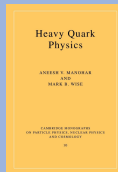
## Example 2: HQSS + SU(3) flavour

- Heavy quarks are... heavy:  $m_c, m_b \rightarrow \infty$
- $H_a^{(Q)} = \frac{1 + \not{v}}{2} \left( P_{a\mu}^{*(Q)} \gamma^\mu - P_a^{(Q)} \gamma_5 \right)$
- HQSS and SU(3) light flavour symmetry:  $H_a^{(Q)} \sim (Q\bar{u}, Q\bar{d}, Q\bar{s}) \sim (D^{(*)0}, D^{(*)+}, D_s^{(*)+})$

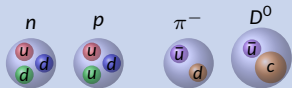
[Grinstein *et al.*, NP,B380('92); Alfiky *et al.*, PL,B640('06), ...]

$H\bar{H} \rightarrow H\bar{H}$  LO lagrangian:

$$\begin{aligned} \mathcal{L}_{4H} &= \frac{1}{4} \text{Tr} \left[ \bar{H}^{(Q)a} H_b^{(Q)} \gamma_\mu \right] \text{Tr} \left[ H^{(\bar{Q})c} \bar{H}_d^{(\bar{Q})} \gamma^\mu \right] \left( F_A \delta_a^b \delta_c^d + F_A^\lambda \vec{\lambda}_a^b \cdot \vec{\lambda}_c^d \right) \\ &+ \frac{1}{4} \text{Tr} \left[ \bar{H}^{(Q)a} H_b^{(Q)} \gamma_\mu \gamma_5 \right] \text{Tr} \left[ H^{(\bar{Q})c} \bar{H}_d^{(\bar{Q})} \gamma^\mu \gamma_5 \right] \left( F_B \delta_a^b \delta_c^d + F_B^\lambda \vec{\lambda}_a^b \cdot \vec{\lambda}_c^d \right), \end{aligned}$$



## (Conventional and exotic) Spectroscopy. The recent LHCb $T_{cc}^+$ “tetraquark”

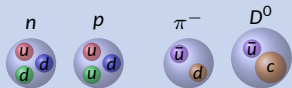


### Conventional hadrons:

- Mesons:  $q\bar{q}'$ :  $\pi^+ = u\bar{d}$ ,  $D^0 = c\bar{u}$ , ...
- Baryons:  $q_1q_2q_3$ :  $p = uud$ ,  $n = udd$ , ...

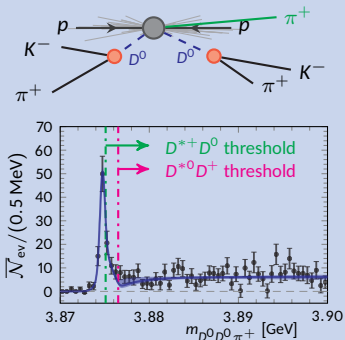
- Constituent **quark models** have successfully described most of (but not all!) the hadrons discovered so far.
- Only possibilities? No, the only requirement is to be **color singlets**. There can be tetraquarks ( $q_1q_2q_3q_4$ ), pentaquarks ( $\bar{q}_1q_2q_3q_4$ ), hybrids ( $\bar{q}_1q_2g$ ), glueballs ( $gg$ ), **hadronic molecules** ( $MM'$ ,  $MB$ ,  $BB'$ ),...

# (Conventional and exotic) Spectroscopy. The recent LHCb $T_{cc}^+$ “tetraquark”



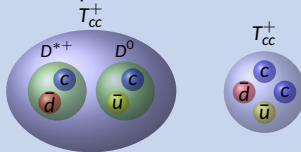
- Conventional hadrons:
  - Mesons:  $q\bar{q}'$ :  $\pi^+ = u\bar{d}$ ,  $D^0 = c\bar{u}$ , ...
  - Baryons:  $q_1q_2q_3$ :  $p = uud$ ,  $n = udd$ , ...

[LHCb: Nature Phys. 18, 751('22);  
Nature Commun., 13, 3351('22)]



- Constituent **quark models** have successfully described most of (but not all!) the hadrons discovered so far.
- Only possibilities? No, the only requirement is to be **color singlets**. There can be tetraquarks ( $q_1q_2q_3q_4$ ), pentaquarks ( $\bar{q}_1q_2q_3q_4$ ), hybrids ( $\bar{q}_1q_2g$ ), **hadronic molecules** ( $MM'$ ,  $MB$ ,  $BB'$ ),...

- In 2021, the LHCb collaboration discovers  $T_{cc}^+$ , with quark content  $cc\bar{u}\bar{d}$ , and very close to  $D^{*+}D^0$  threshold
- Even if something is **explicitly exotic**, you still have to understand how are quarks distributed inside the hadron.



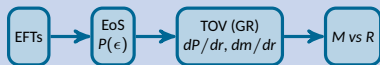
- [MA, PL,B829,137052('21); MA, Nieves, EPJ,C82,724('22)]
- Our analysis favours the **molecular** picture.
- Also, relevant to understand QCD confinement, and color combinations.

# Nuclear and Astroparticle Physics Group @ ICE-CSIC (Barcelona)

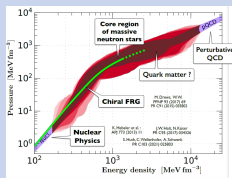
Physics of **compact/dense stellar** objects from its hadron/nuclear constituents:

- Equation of state  $P(\epsilon)$  of dense phases, including the interplay of dense matter with dark matter, to determine  $M$  and  $R$ :

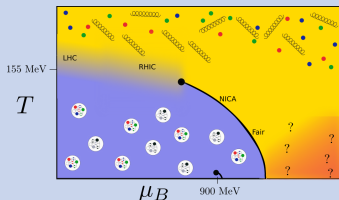
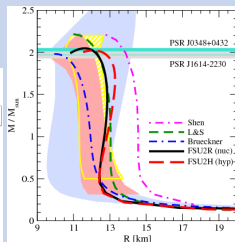
[Tolos et al., Ast.J., 834,3('17); Ast.Soc.Aus., 34,65('17)]



- Transport properties inside NS (hydrodynamical modes, rotational properties, and cooling)



[W. Weise@HYP2022]

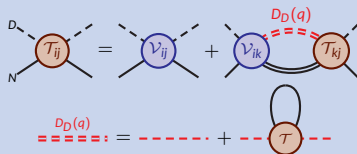


Matter at **extreme large temperatures**, as those occurring few microseconds after the Big Bang;

- Exploration of QCD phase diagram
- Properties of quarks and hadrons in hot and dense matter
- Effective descriptions of the dynamics of the hot quark-gluon plasma
- Studied in present particle accelerators, such as the LHC/CERN or RHIC/BNL.

# $D^{(*)}$ and $\bar{D}^{(*)}$ in nuclear matter

- Suppose you make your experiments “inside” of a nuclear medium, with nucleon density  $\rho \lesssim \rho_0 \approx 0.17 \text{ fm}^{-3}$

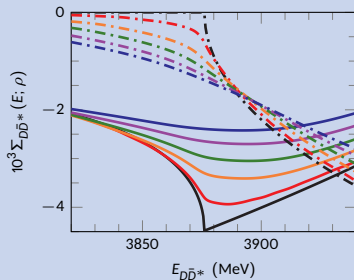


- $D^{(*)}, \bar{D}^{(*)}$  propagators in nuclear matter:

$$D_V^{-1}(q) = q_0^2 - \vec{q}^2 - m_V^2 - \Pi_V(q^0, \vec{q})$$

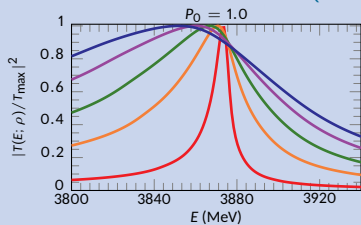
- Self-consistent calculation:

- $D^{(*)}$ : Tolos *et al.*, PR,C80,065202('09)
- $\bar{D}^{(*)}$ : Garcia-Recio *et al.*, PR,C85,025203('12)

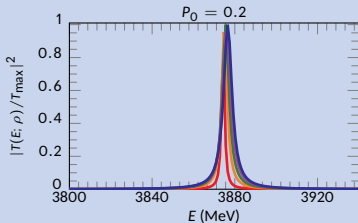


- [MA, Nieves, Tolos, PR,C104,035203('21)]
- With these modified propagators, one can compute  $D\bar{D}^*$  loop functions in nuclear medium.
- Threshold effect** washed out as  $\rho$  increases.  $D^{(*)}, \bar{D}^{(*)}$  acquire in-medium widths.

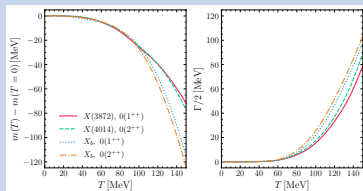
## $D\bar{D}^*$ amplitudes and the $X(3872)$ peak (high and low $P_0$ )



- For  $P_0$  high ( $\gtrsim 0.7$ ) large modifications induced by the nuclear medium
- Very large broadening (larger for larger  $P_0$ )
- $P_0 = 1$ ,  $|T|^2$  peak shifted to the left
- $P_0 = 0.8$ ,  $|T|^2$  peak shifted to the right

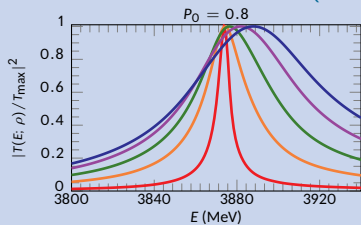


- For  $P_0$  low, medium-induced modifications are small
- Broadening is small (smaller for smaller  $P_0$ )
- Some “pre-existence” of  $X(3872)$ , not directly associated with  $D\bar{D}^*$ , makes that the state does not wash out with increasing  $\rho$



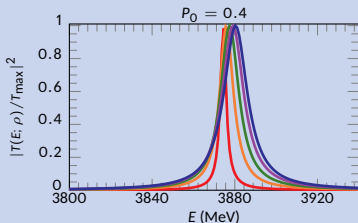
- Together with **density** modifications, one can also study the modifications due to **temperature** of  $X(3872)$  [Montaña *et al.*, 2211.01896]

## $D\bar{D}^*$ amplitudes and the $X(3872)$ peak (high and low $P_0$ )

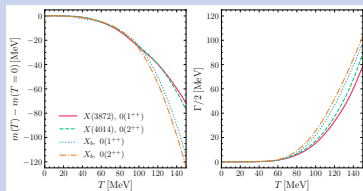


- For  $P_0$  high ( $\gtrsim 0.7$ ) large modifications induced by the nuclear medium
- Very large broadening (larger for larger  $P_0$ )
- $P_0 = 1$ ,  $|T|^2$  peak shifted to the left
- $P_0 = 0.8$ ,  $|T|^2$  peak shifted to the right

- For  $P_0$  low, medium-induced modifications are small
- Broadening is small (smaller for smaller  $P_0$ )
- Some “pre-existence” of  $X(3872)$ , not directly associated with  $D\bar{D}^*$ , makes that the state does not wash out with increasing  $\rho$



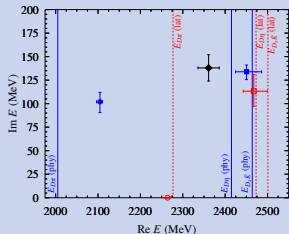
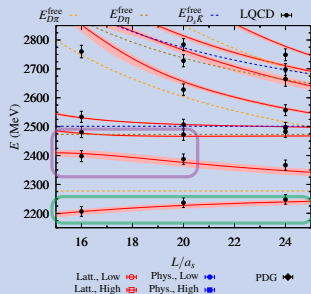
- Together with **density** modifications, one can also study the modifications due to **temperature** of  $X(3872)$  [Montaña *et al.*, 2211.01896]



# $D\pi$ , $D\eta$ , $D_s\bar{K}$ and $D^*(2300)$ : Comparison with LQCD energy levels

## Important connections between LQCD and EFTs!

[MA et al., Phys. Lett. B 767, 465 (2017)]

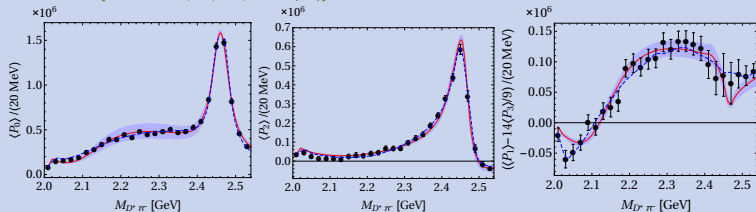


- $E_n(L)$  are provided for  $D\pi$ ,  $D\eta$ ,  $D_s\bar{K}$  in a recent LQCD simulation. [G. Moir et al., JHEP 1610, 011 (2016)]
- **Red Bands:** Our amplitude in a finite volume. [MA et al., Phys. Lett. B 767, 465 (2017)]
- **No fit** is performed (LECs previously determined)
- Level **below threshold**, associated with a **bound state**.
- **Second level** has large shifts w. r. t. thresholds, non-interacting energy levels.
- For lattice masses, we find a **bound state** and a **resonance**
- For physical masses:
  - The bound state evolves into a **resonance** above  $D\pi$  threshold.
  - The resonance varies very little.
- PDG: only one resonance, “suspiciously” lying between both.
- We also study  $DK$ ,  $D_s\eta$ ,  $(S, I) = (1, 0)$ ,  $D_{s0}^*(2317)$  bound state:  $M = 2315_{-28}^{+18}$  MeV.

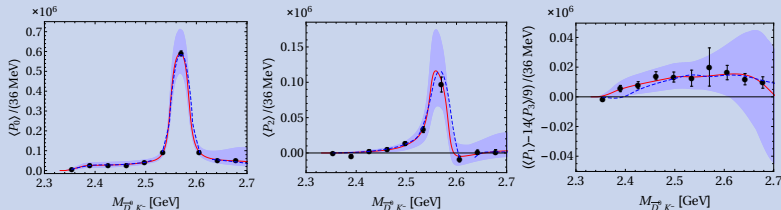
# $D\pi$ , $D\eta$ , $D_s\bar{K}$ and $D^*(2300)$ : Comparison with experimental data from LHCb

Du, MA, Fernández-Soler, Guo, Hanhart, Meißner, Nieves, Yao, PR,D98,094018('18)

- $B^- \rightarrow D^+ \pi^- \pi^-$  [LHCb Collab., PR,D94,072001('16)]



- $B_s^0 \rightarrow \bar{D}^0 K^- \pi^+$  [LHCb Collab., PR,D90,072003('14)]



- Rapid movement in  $\langle P_{13} \rangle$  [no  $D_2(2460)$ ] between 2.4 and 2.5 GeV. Related to  $D\eta$  and  $D_s\bar{K}$  openings.
- Recall: these are the amplitudes with **two states** in the  $D_0^*(2300)$  region, and no fit of the  $T$ -matrix parameters is done.

# QCD and Hadronic Physics Thematic Network



MINISTERIO  
DE CIENCIA  
E INNOVACIÓN



AGENCIA ESTATAL DE INVESTIGACIÓN - Convocatoria 2022  
«Redes de Investigación»

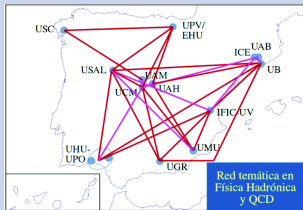
## 1. DATOS DE LA ACTUACIÓN

**COORDINADORA DE LA RED** (Nombre y apellidos): **Laura Tolós Rigueiro**

**TÍTULO DE LA ACTUACIÓN EN ESPAÑOL:** **Red Temática de Física Hadrónica y QCD**

**TÍTULO EN INGLÉS:** **QCD and Hadronic Physics Thematic Network**

hadQCDnet



- Application submitted to 2022 call for “Redes de Investigación”
- PI (coordinator): Laura Tolós (ICE)
- Several coPIs (representatives) are also from CSIC
- Many of the people in the proposal from CSIC too
- Joint application between more QCD- and hadron-oriented researchers
- Can be an important, strategic tool for this chapter of the challenge

## Summary / prospectives

- Solving QCD is certainly a fundamental challenge
- Currently, in the non-perturbative regimes, we can provide approximate numerical calculations of the exact solution through LQCD simulations.
- Perturbative calculations with (rather complicated) loops, which are currently one of the bottlenecks in precision calculations, are being improved with new methods developed.
- EFTs, supplemented with non-perturbative (unitarity, crossing, ...) methods, are also living great advances, which are playing an important role in the field of hadron spectroscopy, where new resonances/particles/states are being constantly discovered in experiments (LHCb, BES,...)
- Many of these aspects can be applied together to increase our knowledge of strong interactions

### Summary of summary

In the context of **“Understanding the basic components of the Universe, its structure and evolution”**, the problem of **“Solving Quantum Chromodynamics”** its a lively and exciting challenge, with many developments in several branches, which can also interact together to achieve this goal