Detection of gravitational waves with axion haloscopes

Camilo A. Garcia Cely

Ramón y Cajal Fellow







Gravitational Waves

• Predicted by Poincaré (1905)



• Einstein provided a firm theoretical background for them (1916) $\Box h_{\mu\nu} = -16\pi G T_{\mu\nu}$

Gravitational Waves

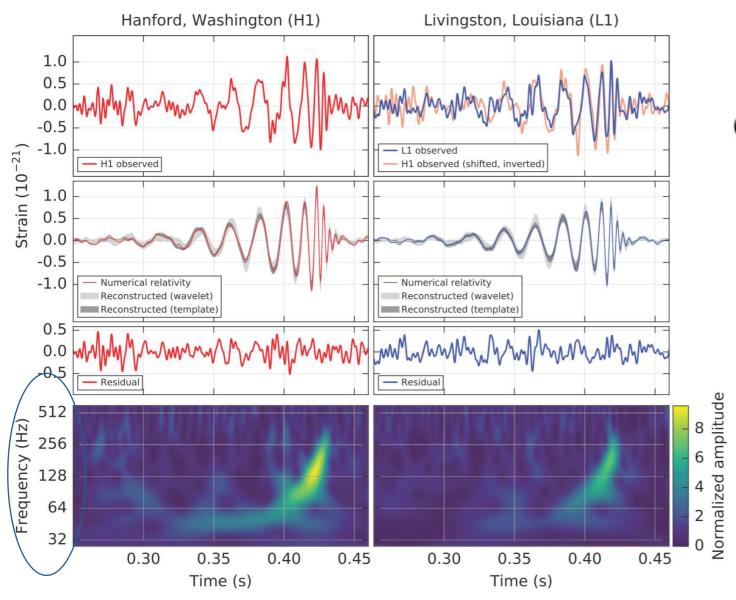
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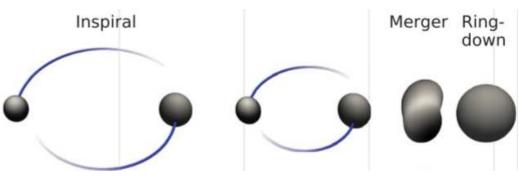
PRL **116,** 061102 (2016)

PHYSICAL REVIEW LETTERS

week ending 12 FEBRUARY 2016

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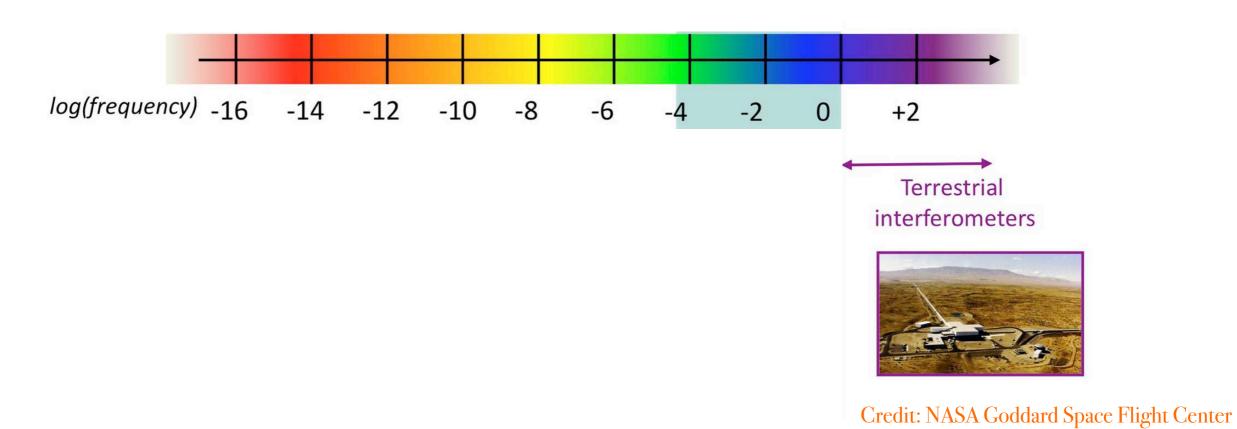




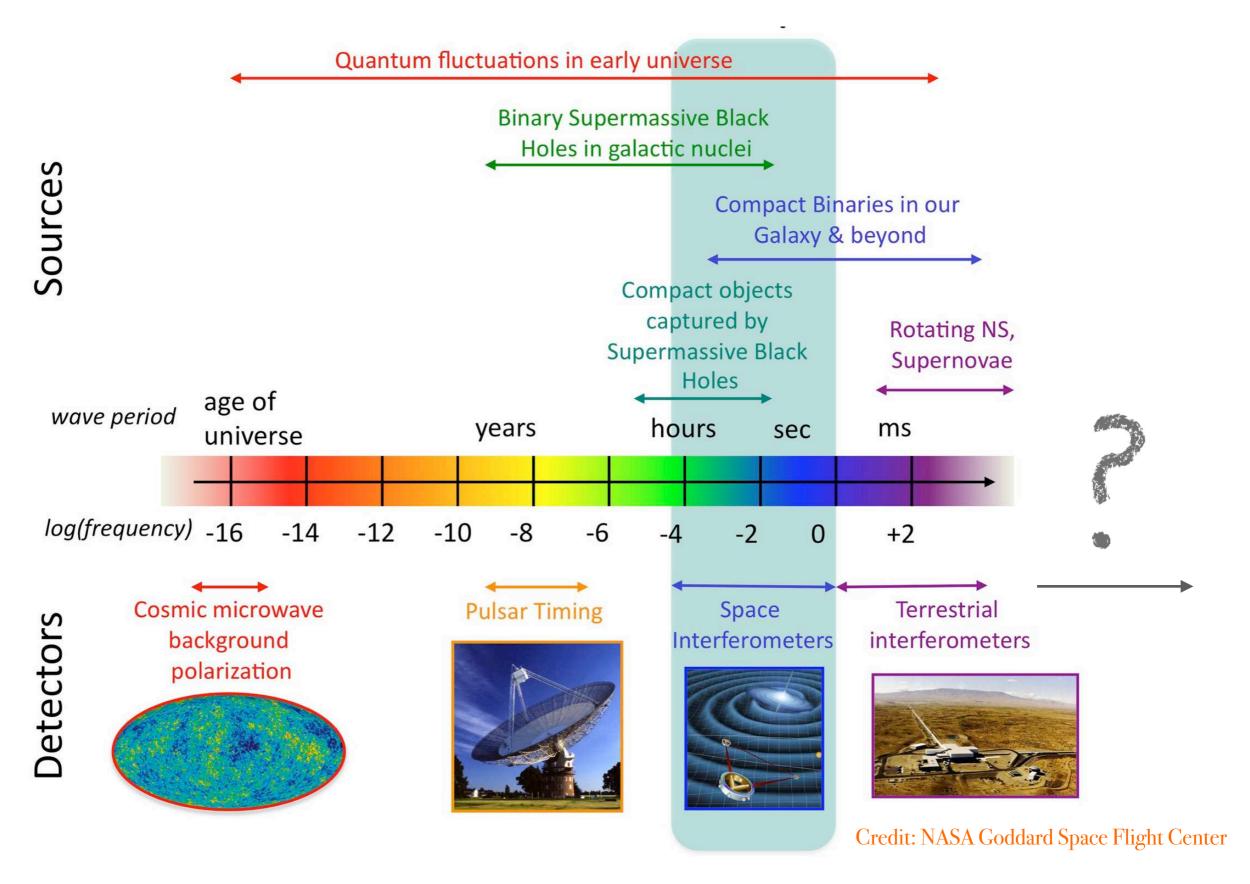
Terrestrial interferometers



Gravitational Wave Spectrum



Gravitational Wave Spectrum



Revisiting Gertsenhstein's ideas

SOVIET PHYSICS JETP

VOLUME 16, NUMBER 2

FEBRUARY, 1963

ON THE DETECTION OF LOW FREQUENCY GRAVITATIONAL WAVES

M. E. GERTSENSHTEĬN and V. I. PUSTOVOĬT

Submitted to JETP editor March 3, 1962

J. Exptl. Theoret: Phys. (U.S.S.R.) 43, 605-607 (August, 1962)

It is shown that the sensitivity of the electromechanical experiments for detecting gravitational waves by means of piezocrystals is ten orders of magnitude worse than that estimated by Weber. [1] In the low frequency range it should be possible to detect gravitational waves by the shift of the bands in an optical interferometer. The sensitivity of this method is investigated.

Terrestrial interferometers

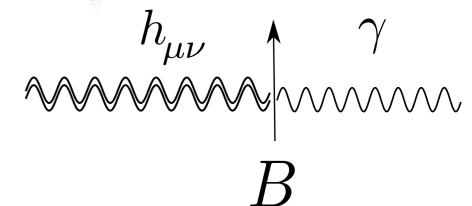


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SOVIET PHYSICS JETP

VOLUME 14, NUMBER 1

JANUARY, 1962



WAVE RESONANCE OF LIGHT AND GRAVITIONAL WAVES

M. E. GERTSENSHTEĬN

Submitted to JETP editor July 29, 1960

J. Exptl. Theoret. Phys. (U.S.S.R.) 41, 113-114 (July, 1961)

The energy of gravitational waves excited during the propagation of light in a constant magnetic or electric field is estimated.

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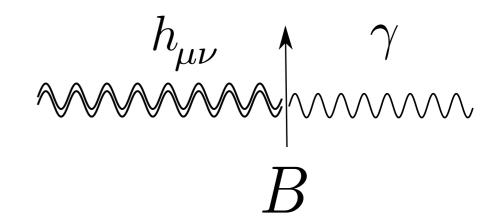
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The (inverse) Gertsenhstein Effect

- The conversion of gravitational waves into electromagnetic waves is a classical process. Its rate does not involve \hbar
- The process is strictly analogous to axionphoton conversion. Raffelt, Stodolski'89

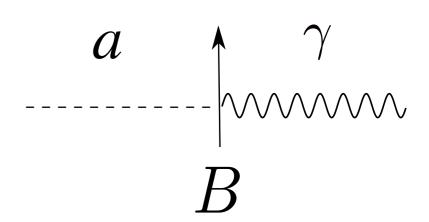


Axions

• The strong CP problem: experiments put a strong upper bound on the electric dipole of the neutron

$$\mathcal{L} = \frac{g_s^2 \theta}{32\pi^2} F_{\mu\nu}^a \tilde{F}^{a\mu\nu} \longrightarrow \frac{g_s^2 a}{32\pi^2 f} F_{\mu\nu}^a \tilde{F}^{a\mu\nu}_{\text{Peccei, Quinn 1977}}$$

• Excellent dark matter candidate Weinberg, Wilczek 1978

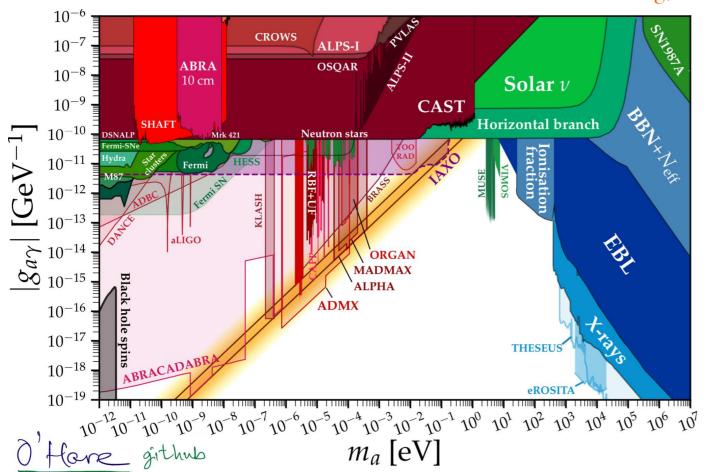


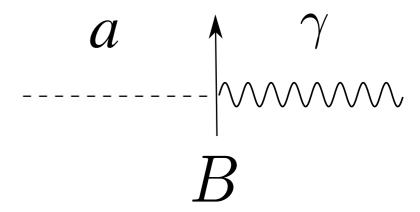
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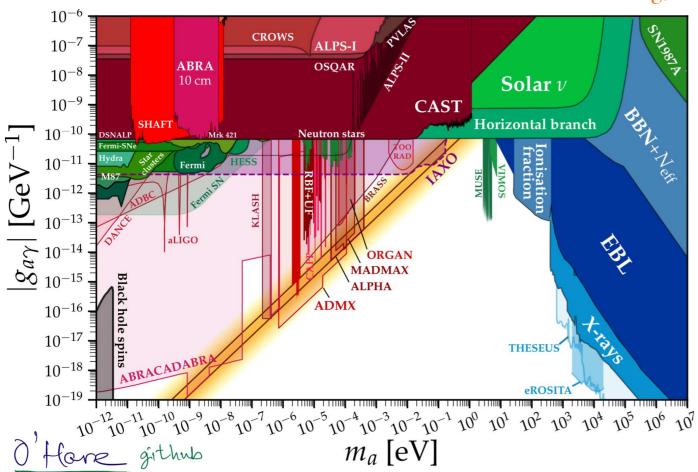


Axions

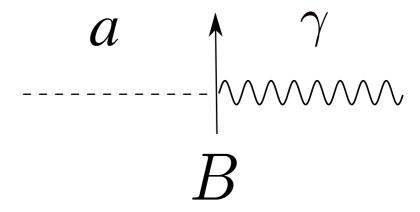
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Ideas and techniques developed for axions can be adapted to gravitational waves



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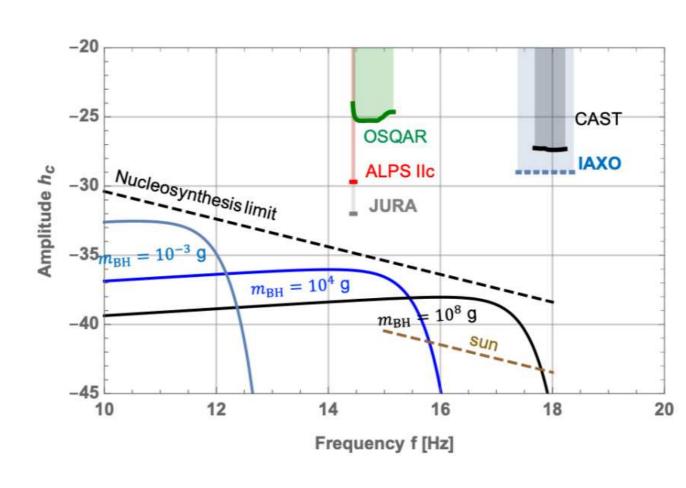
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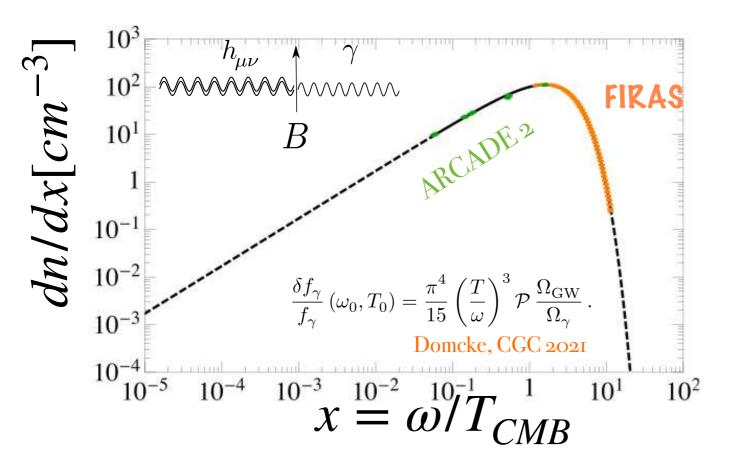
Upper limits on the amplitude of ultra-high-frequency gravitational waves from graviton to photon conversion

A. Ejlli , D. Ejlli, A. M. Cruise, G. Pisano & H. Grote

The European Physical Journal C 79, Article number: 1032 (2019) | Cite this article



Cosmological Detectors



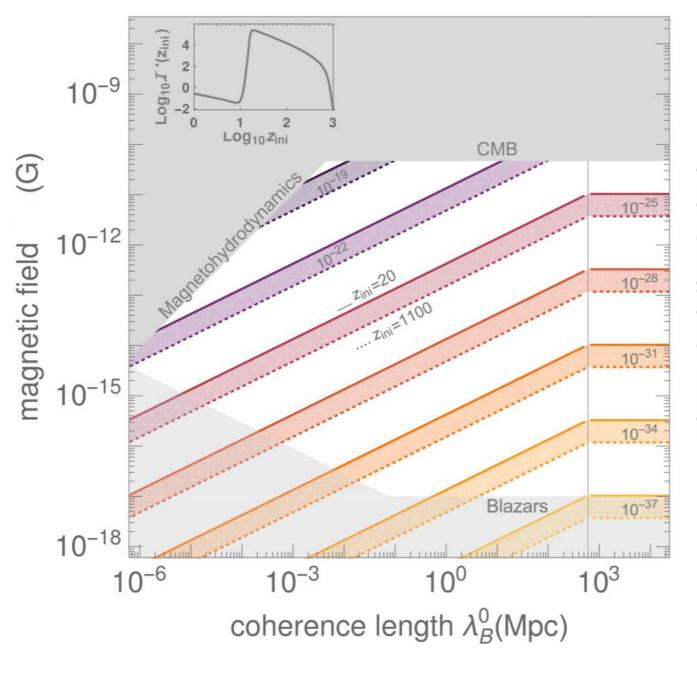
- Involving gravity the conversion probabilities are extremely small. It may be compensated by a 'detector' of cosmological size
- RJ tail remains unexplored with upcoming advances in radio astronomy probing it in the near future.

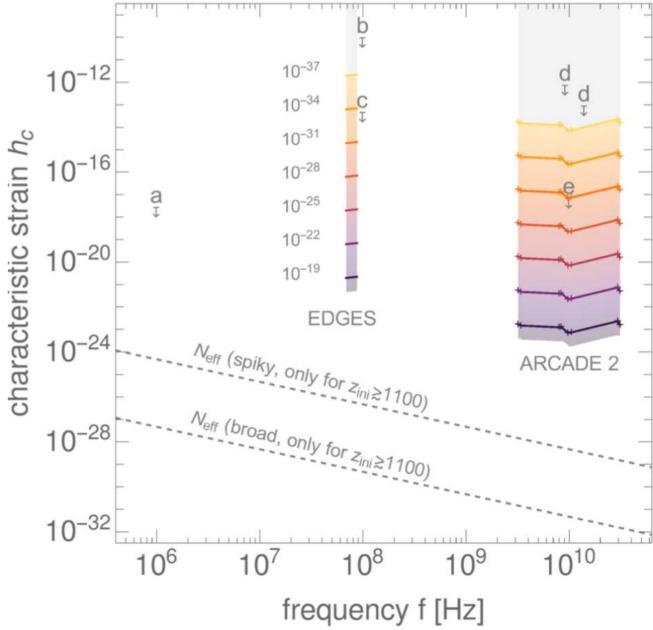
Cosmological Detectors

PHYSICAL REVIEW LETTERS 126, 021104 (2021)

Potential of Radio Telescopes as High-Frequency Gravitational Wave Detectors

Valerie Domcke^{1,2,3,*} and Camilo Garcia-Cely^{1,†}





existing laboratory bounds from

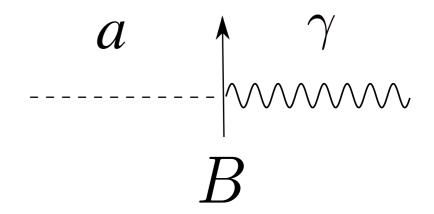
- a) superconducting parametric converter Reece et al '84
- b) waveguide Cruise Ingley '06
- c) 0.75 m interferometer Akutsu '08
- d) magnon detector Ito, Soda '04
- e) magnetic conversion detector Cruise et al '12

Not everything is axion photon conversion

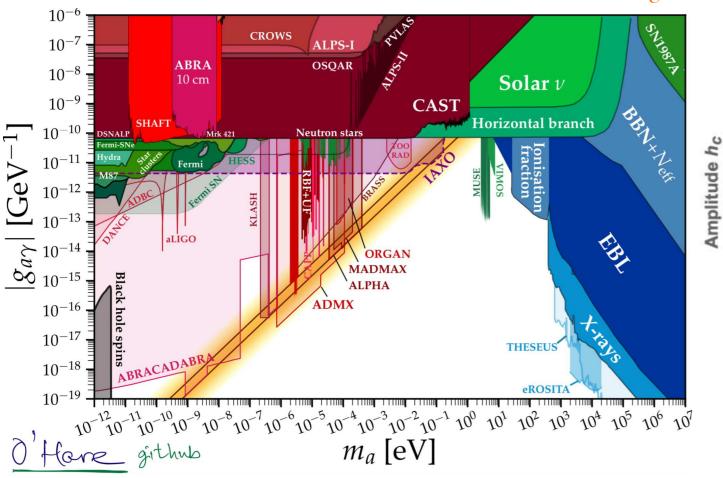
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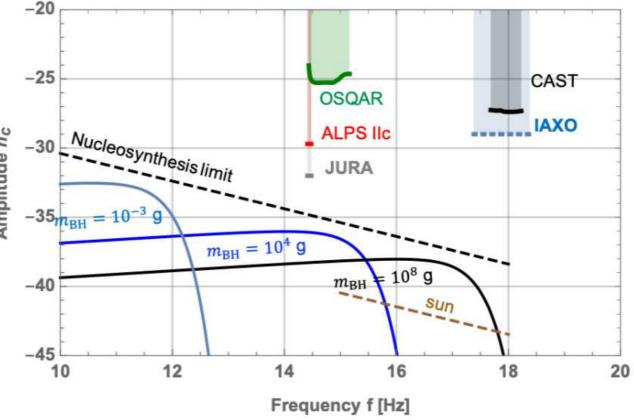
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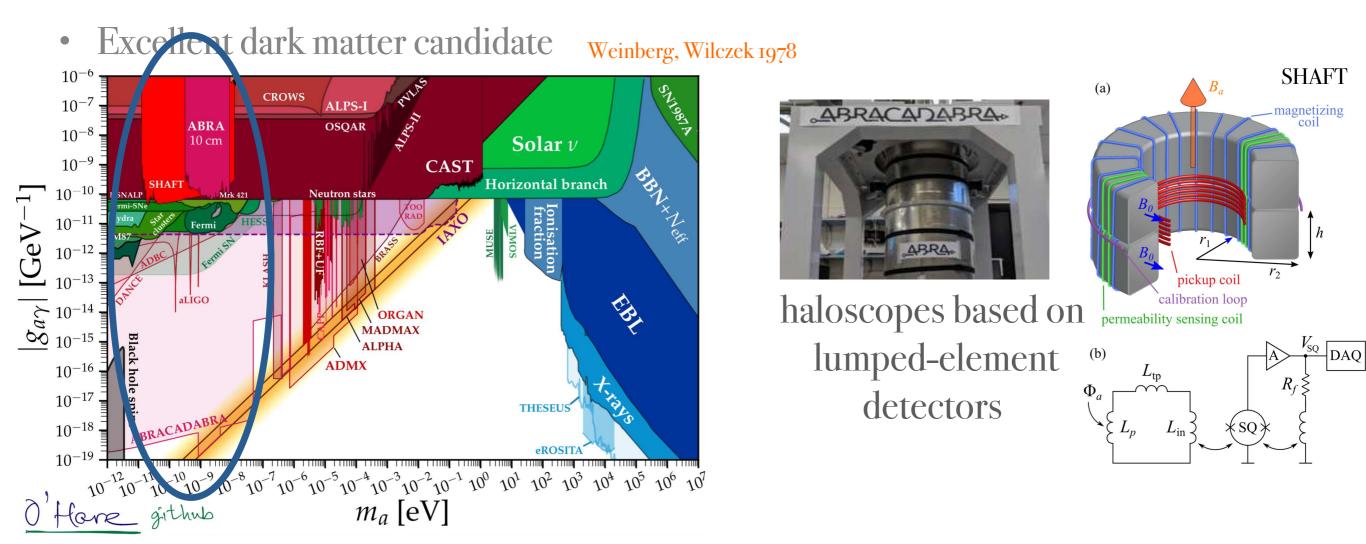
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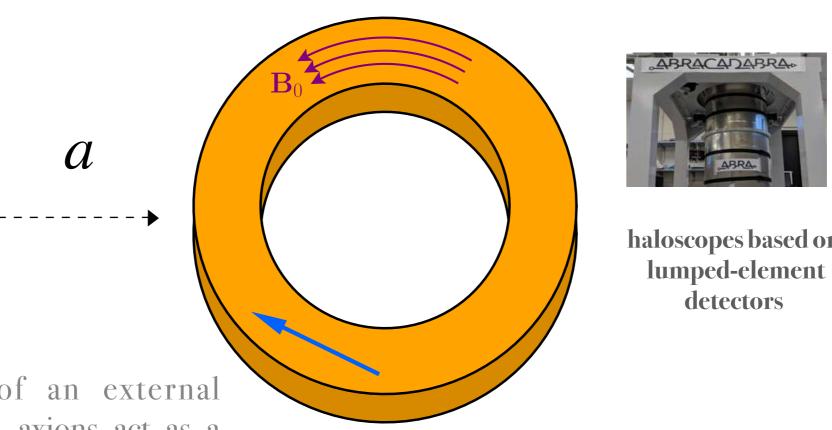
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In the presence of an external electromagnetic field, axions act as a source term to Maxwell's equations, effectively inducing an electromagnetic current.

$$\partial_{\nu}F^{\mu\nu} = j_{\text{eff}}^{\mu}$$

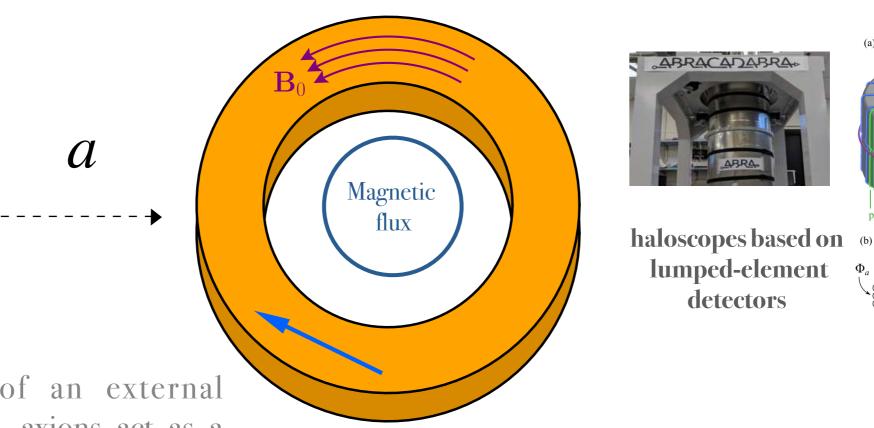


permeability sensing coil haloscopes based on (b)

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SHAFT

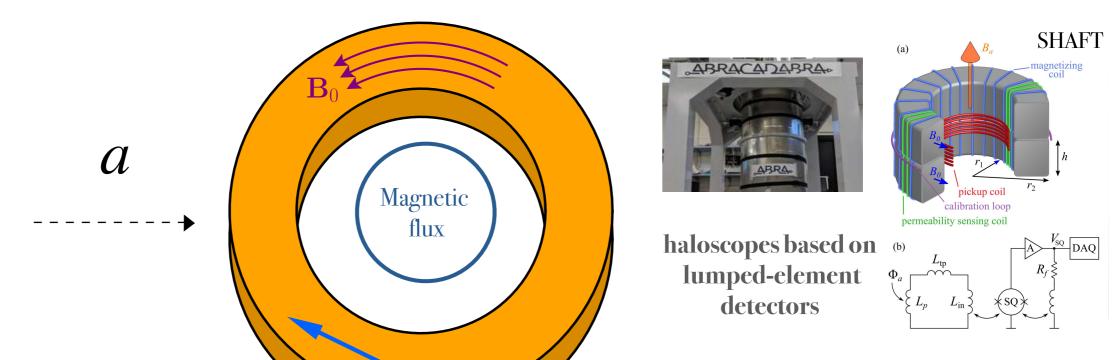


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$$P_i = g_{a\gamma\gamma} a B_i$$

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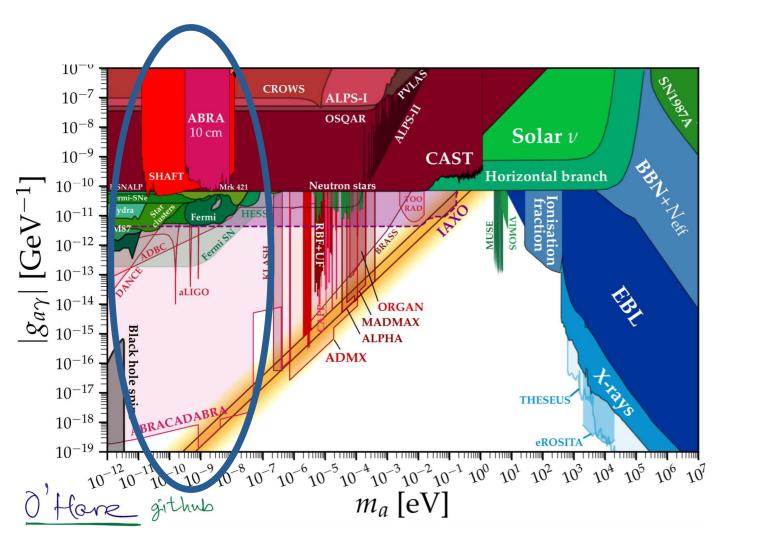
McAllister et al, 1803.07755 Tobar et al, 1809.01654 Ouellet et al, 1809.10709

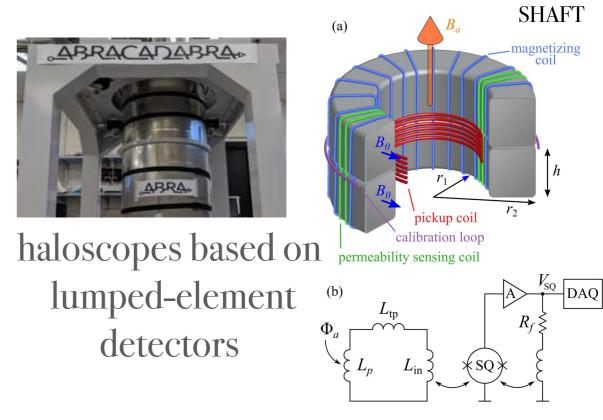
A novel search for high-frequency gravitational waves with low-mass axion haloscopes

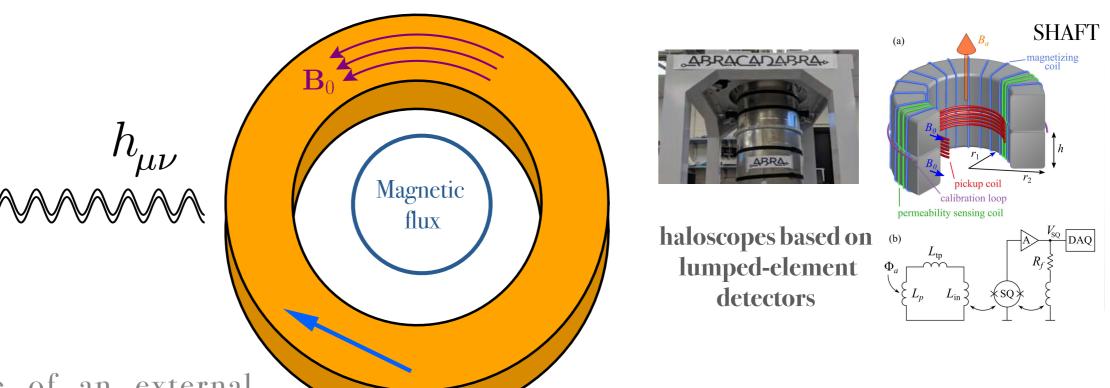
Valerie Domcke, ^{1, 2} Camilo Garcia-Cely, ³ and Nicholas L. Rodd ¹

¹Theoretical Physics Department, CERN, 1 Esplanade des Particules, CH-1211 Geneva 23, Switzerland ²Institute of Physics, École Polytechnique Fédérale de Lausanne (EPFL), CH-1015 Lausanne, Switzerland ³Deutsches Elektronen-Synchrotron DESY, Notkestr. 85, 22607 Hamburg, Germany

Gravitational waves (GWs) generate oscillating electromagnetic effects in the vicinity of external electric and magnetic fields. We discuss this phenomenon with a particular focus on reinterpreting the results of axion haloscopes based on lumped-element detectors, which probe GWs in the 100 kHz-100 MHz range. Measurements from ABRACADABRA and SHAFT already place bounds on GWs, although the present strain sensitivity is weak. However, we demonstrate that the sensitivity scaling with the volume of such instruments is significant – faster than for axions – and so rapid progress will be made in the future. With no modifications, DMRadio-m³ will have a GW strain sensitivity of $h \sim 10^{-20}$ at 200 MHz. A simple modification of the pickup loop used to readout the induced magnetic flux can parametrically enhance the GW sensitivity, particularly at lower frequencies.







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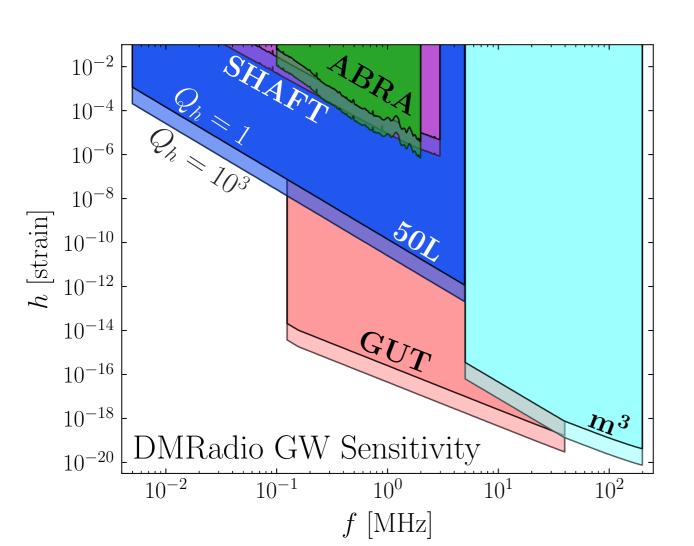
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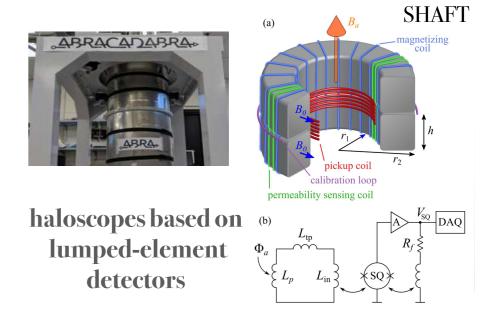
$$P_{i} = -h_{ij}E_{j} + \frac{1}{2}hE_{i} + h_{00}E_{i} - \epsilon_{ijk}h_{0j}B_{k}$$

$$M_i = -h_{ij}B_j - \frac{1}{2}hB_i + h_{jj}B_i + \epsilon_{ijk}h_{0j}E_k$$

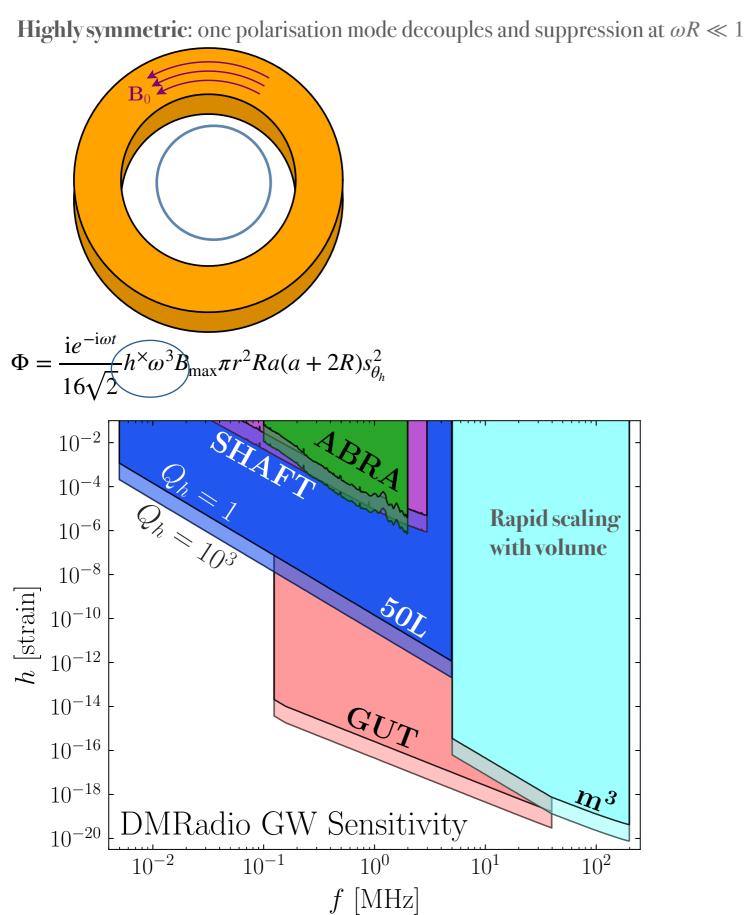
Domcke, CGC, Rodd, 2202.00695

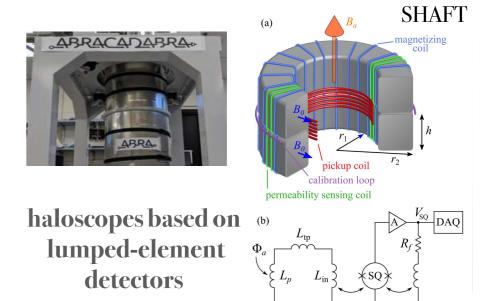
Valerie Domcke, $^{1,\,2}$ Camilo Garcia-Cely, 3 and Nicholas L. Rodd^1





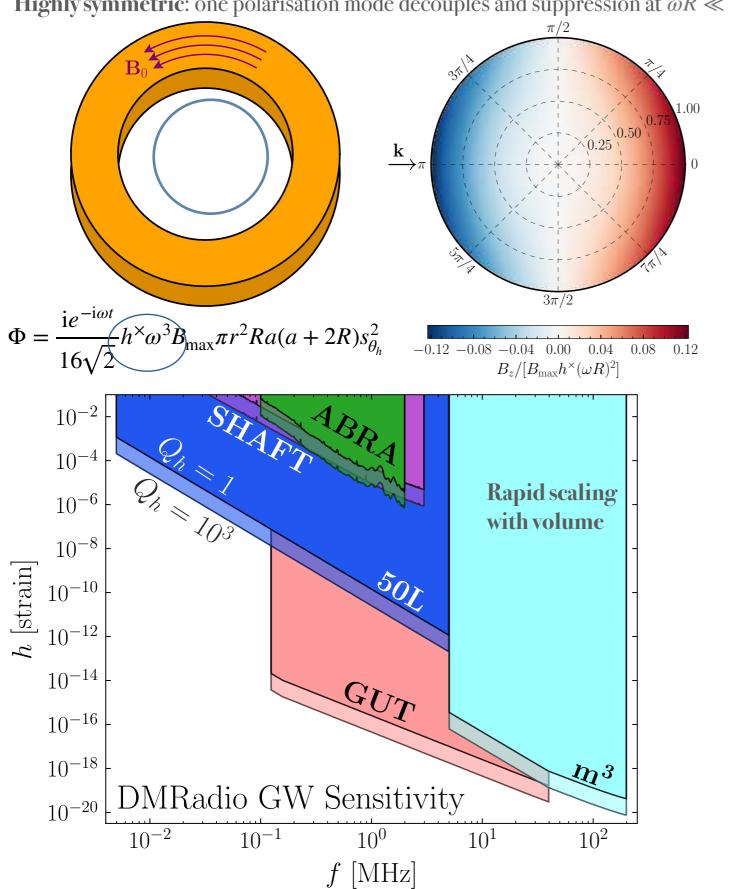
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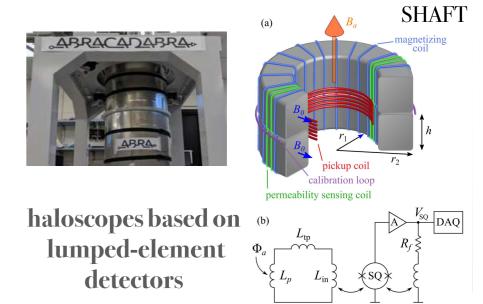




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Highly symmetric: one polarisation mode decouples and suppression at $\omega R \ll 1$





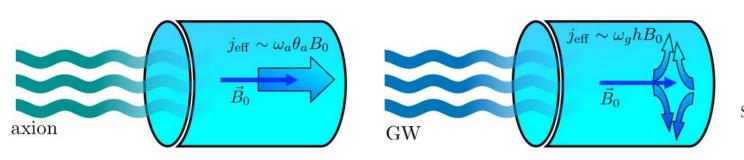
Valerie Domcke, ^{1, 2} Camilo Garcia-Cely, ³ and Nicholas L. Rodd ¹ **Highly symmetric**: one polarisation mode decouples and suppression at $\omega R \ll 1$ 1.00 0.25 0.50 pickup loop $\xrightarrow{\mathbf{k}}_{\pi}$ $\Phi_8 = \frac{e^{-i\omega t}}{3\sqrt{2}} \omega^2 B_{\text{max}} r^3 R \ln(1 + a/R) s_{\theta_h}$ $\times \left(h^{\times} s_{\phi_h} - h^{+} c_{\theta_h} c_{\phi_h} \right).$ $\frac{1}{16}h^{\times}\omega^{3}B_{\max}\pi r^{2}Ra(a+2R)s_{\theta_{h}}^{2}$ $B_z/[B_{\rm max}h^{\times}(\omega R)^2]$ 10^{-8} 10^{-2} 10^{-10} 10^{-4} levitated sensors 10^{-12} 10^{-6} ARCADE 10^{-8} 10^{-14} $\frac{10^{-16}}{2}$ strain $\frac{10^{-18}}{2}$ h [strain] 5011 10^{-10} $DMRadio_8$ 10^{-12} DMR₈-100 10^{-14} 10^{-20} 10^{-16} 10^{-22} 10^{-18} M^3 DMRadio GW Sensitivity UHF-GW Landscape 10^{-24} 10^{-2} 10^{-1} 10^{0} 10^{1} 10^{2} 10^{-2} 10^{-1} 10^{0} 10^{1} 10^{2} 10^{3} 10^{4} f [MHz]

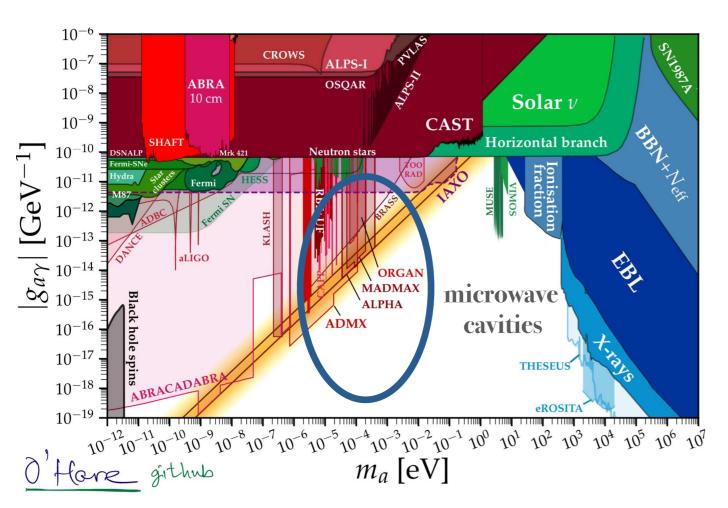
f [MHz]

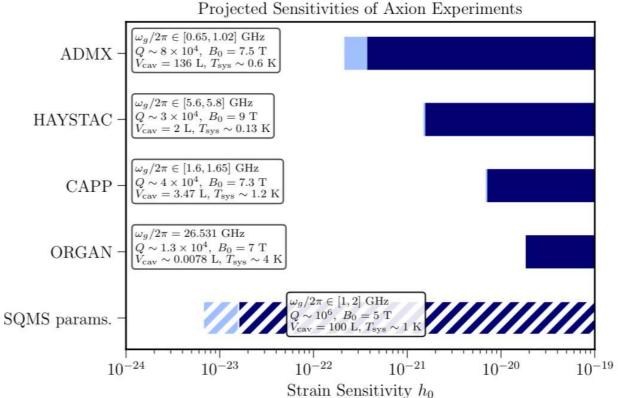
Detecting High-Frequency Gravitational Waves with Microwave Cavities

Asher Berlin, ^{1, 2, 3} Diego Blas, ^{4, 5} Raffaele Tito D'Agnolo, ⁶ Sebastian A. R. Ellis, ^{7, 6} Roni Harnik, ^{2, 3} Yonatan Kahn, ^{8, 9, 3} and Jan Schütte-Engel^{8, 9, 3}

See also Herman, Füzfa, Lehoucq, Clesse, 2012.12189







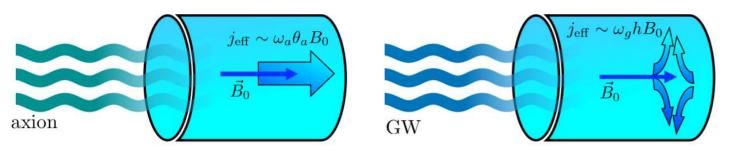
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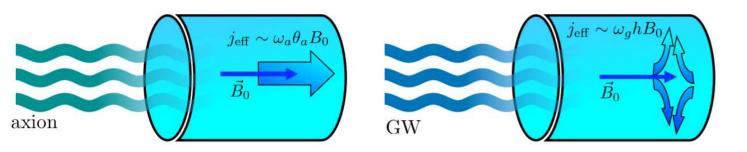
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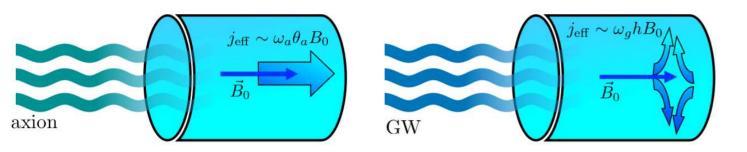
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Eigenmodes
$$\mathbf{E}(\mathbf{x}, t) = \sum_{n} e_n(t) \mathbf{E}_n(\mathbf{x})$$

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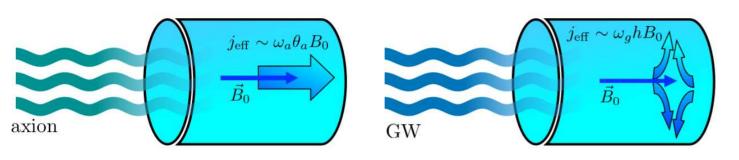
$$\left(\partial_t^2 + \frac{\omega_n}{Q_n}\partial_t + \omega_n^2\right)e_n(t) = -\frac{\int_{V_{\text{cav}}} d^3\mathbf{x} \mathbf{E}_n^* \cdot \partial_t \mathbf{j}_{\text{eff}}}{\int_{V_{\text{cav}}} d^3\mathbf{x} \left|\mathbf{E}_n\right|^2}$$

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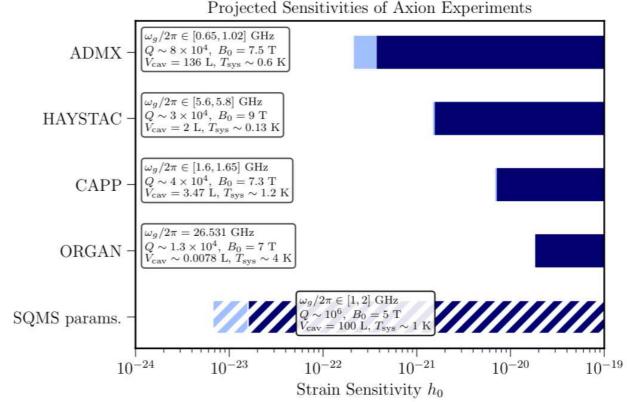
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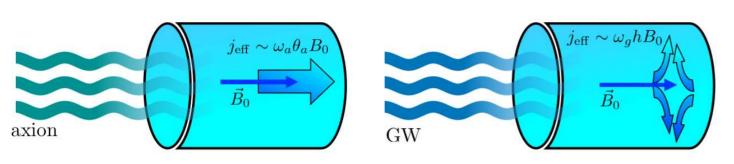
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It resonates when the GW frequency matches one of the eigenmode frequencies

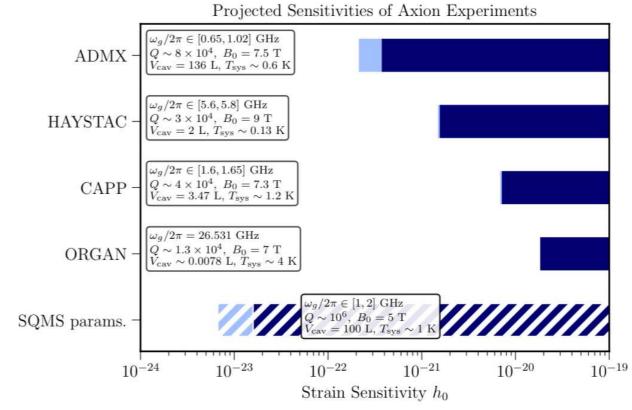
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See also Herman, Füzfa, Lehoucq, Clesse, 2012.12189



In the presence of an external electromagnetic field, GWs act as a source term to Maxwell's equations, effectively inducing an electromagnetic current.



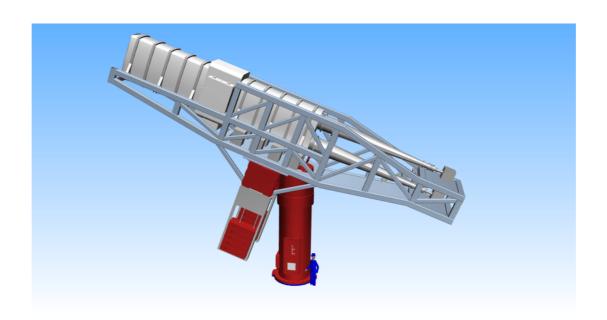
$$\left(\partial_t^2 + \frac{\omega_n}{Q_n}\partial_t + \omega_n^2\right)e_n(t) = -\frac{\int_{V_{\text{cav}}} d^3\mathbf{x} \mathbf{E}_n^* \cdot \partial_t \mathbf{j}_{\text{eff}}}{\int_{V_{\text{cav}}} d^3\mathbf{x} \left|\mathbf{E}_n\right|^2}$$

Eigenmodes
$$\mathbf{E}(\mathbf{x}, t) = \sum_{n} e_n(t) \mathbf{E}_n(\mathbf{x})$$

It resonates when the GW frequency matches one of the eigenmode frequencies

Subtleties due to gauge fixing (TT vs detector frame gauge)

(Baby)IAXO

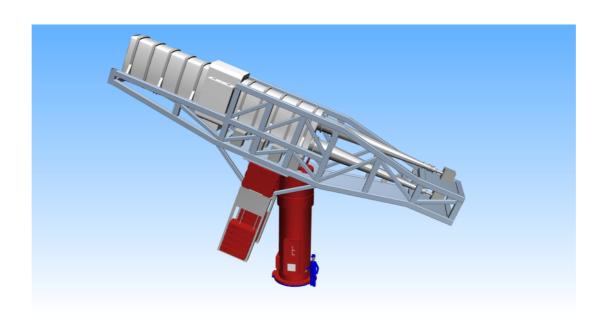


Design of New Resonant Haloscopes in the Search for the Dark Matter Axion: A Review of the First Steps in the RADES Collaboration

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by Alejandro Díaz-Morcillo 1,* © 0, Q José María García Barceló 1 © 0, Q Antonio José Lozano Guerrero 1 © 0, Q Pablo Navarro 1 © 0, Q Benito Gimeno 2 ©, Q Sergio Arguedas Cuendis 3 ©, Q Alejandro Álvarez Melcón 1 © 0, Q Cristian Cogollos 4 ©, Q Sergio Calatroni 3 © 0, Q Babette Döbrich 3 © 0, Q Juan Daniel Gallego-Puyol 5 ©, Q Jessica Golm 3,6 ©, Q Igor García Irastorza 7 ©, Q Chloe Malbrunot 3 © 0, Q Jordi Miralda-Escudé 4,8 ©, Q Carlos Peña Garay 9,10 ©, Q Javier Redondo 7,11 © and Q Walter Wuensch 3 ©
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The future BabyIAXO magnet, a first step of the IAXO experiment... has enough volume to even host <u>single cavities</u> for searching in the 1-2 μ eV mass range (240-480MHz, the lower limit of UHF band). These masses are just below the region explored by ADMX currently.

(Baby)IAXO



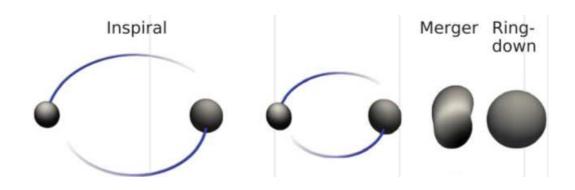
Design of New Resonant Haloscopes in the Search for the Dark Matter Axion: A Review of the First Steps in the RADES Collaboration

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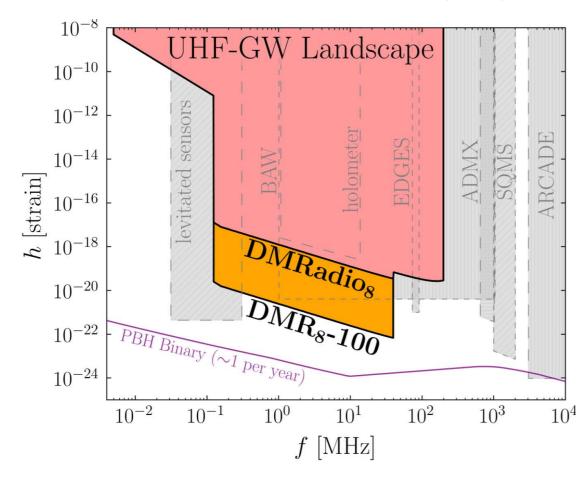
Gravitational waves?

Work in progress

Potential sources

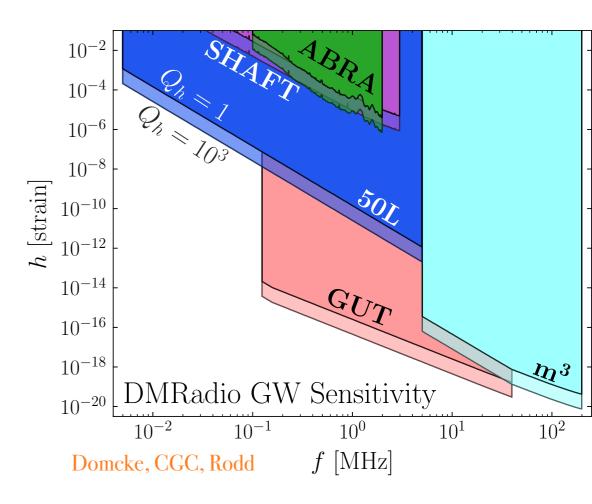


Domcke, CGC, Rodd



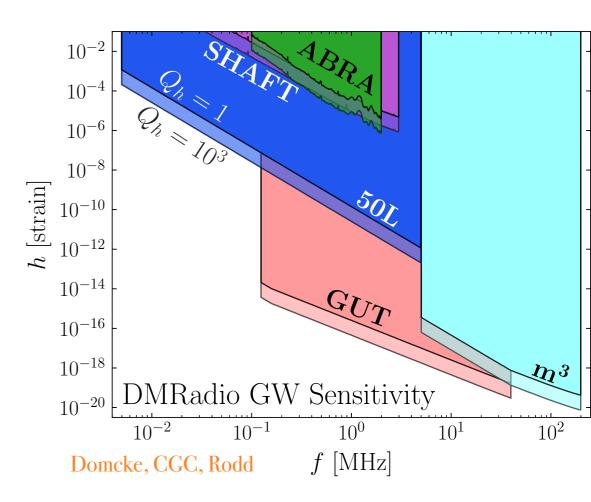
Up-to-date estimates of PBH in binaries and their expected merger rate accounting for the local overdensity in the Milky Way.

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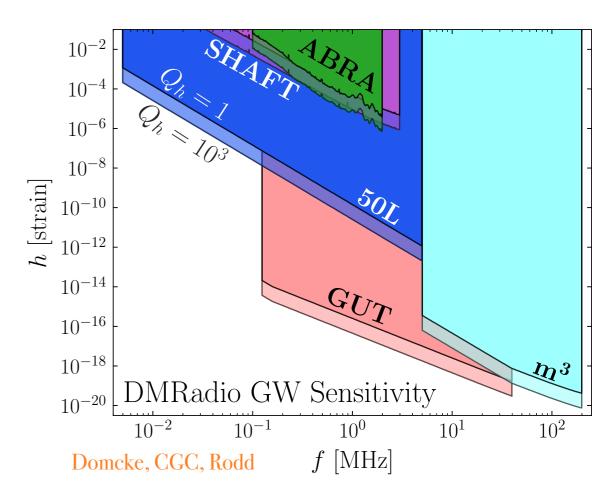
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Through the use of an optimized figure-8 pickup loop geometry, the DMRadio program may discover exotic sources of GWs



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A number of distinct experimental proposals have coalesced on a strain sensitivity of 10⁻²² for MHz GWs, a level that is still orders of magnitude away from any signal of the early Universe. Whether we can hope to probe such strain sensitivities remains to be determined.

