

Higgs boson measurements and their EFT interpretations

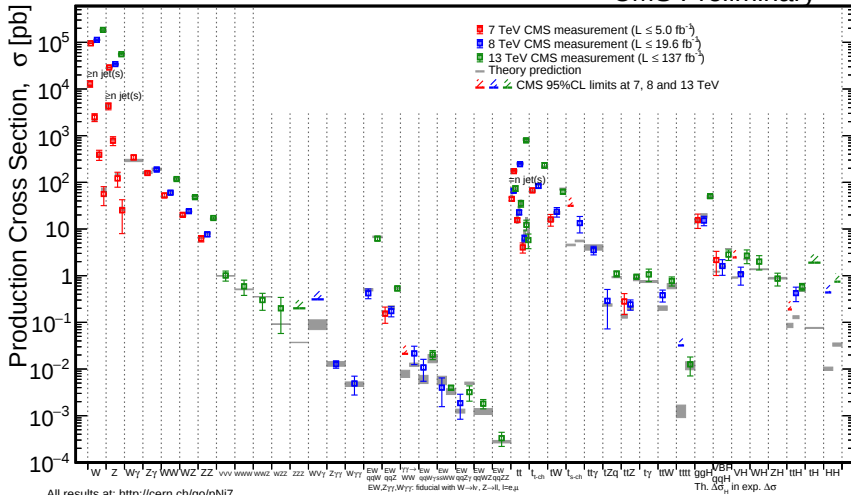
Dermot Moran (CIEMAT)



Triumph of the SM at the LHC

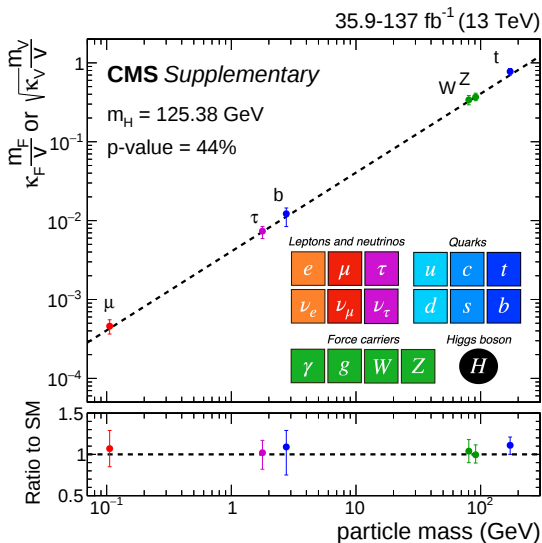
September 2020

CMS Preliminary



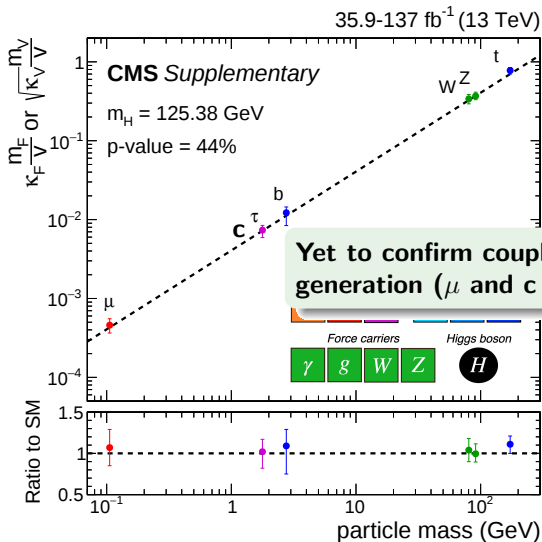
Predictions and measurements agree over **many orders of magnitude**

SM picture of EWSB established



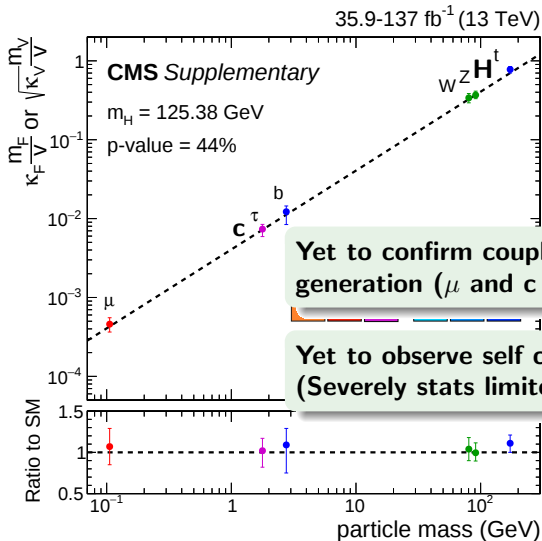
Observed Higgs coupling to **heavy fermions and vector bosons**

SM picture of EWSB established..



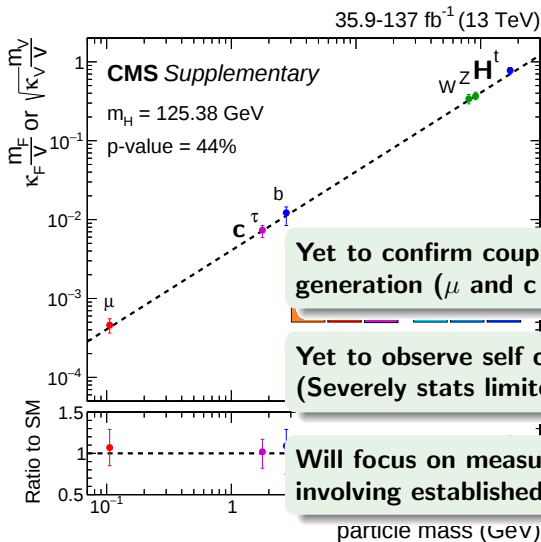
Observed Higgs coupling to **heavy fermions and vector bosons**

SM picture of EWSB established..



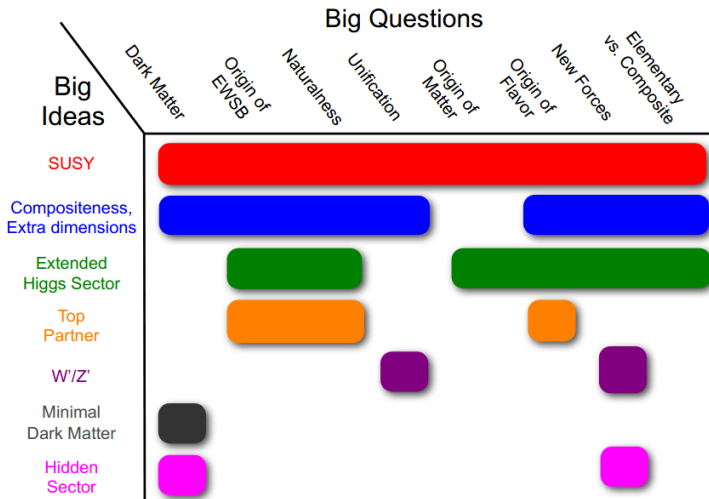
Observed Higgs coupling to **heavy fermions and vector bosons**

SM picture of EWSB established..



Observed Higgs coupling to **heavy fermions and vector bosons**

Big questions remain



Big ideas for what lies beyond the SM

→ Generally predict **new particles at the TeV scale**

Direct searches at LHC find no new particles

ATLAS Exotics Searches* - 95% CL Upper Exclusion Limits

Status: March 2021

ATLAS Preliminary
$$\int \mathcal{L} dt = (3.6 - 139) \text{ fb}^{-1}$$
 $\sqrt{s} = 8, 13 \text{ TeV}$

Model	L, γ	Jets	$E_{\text{miss}}^{\text{min}}$	$[\mathcal{L}(\text{fb}^{-1})]$	Limit	Reference
Extra dimensions	ADD $Q\bar{Q} + g$	$0, \mu, \tau, \gamma$	1-4	Yes	M_{Pl}	2102.10874
	ADD gg -resonant $\gamma\gamma$	2	—	36.7	M_{Pl}	1703.74147
	ADD BH multijet	—	2-3	—	M_{Pl}	1703.08127
	RS1 $Q\bar{Q} + \gamma\gamma$	2	—	37.9	M_{Pl}	1703.15405
	Buk RS $Q\bar{Q} + WW/ZZ$	multi-channel	—	36.1	M_{Pl}	1808.22380
	Buk RS $Q\bar{Q} + W\bar{W} + \nu\bar{\nu}q$	$1, \mu, \tau$	2/1/3	Yes	3.1 TeV	2004.14638
	Buk RS $g\bar{g} + \nu\bar{\nu}$	$1, \mu, \tau$	1.5/3.2/1/2	Yes	1.6 TeV	1804.15822
	ZUED / RPP	$0, \mu, \tau \leq 2$	1-3	Yes	65 MeV	1803.09878
	SSM $Z^* \rightarrow f\bar{f}$	$2, \mu, \tau$	—	139	Z^* mass	1902.06248
	SSM $Z^* \rightarrow \tau\bar{\tau}$	$2, \mu, \tau$	—	36.1	Z^* mass	1703.07242
Gauge bosons	LeptoHiggs $Q\bar{Q} \rightarrow b\bar{b}$	—	2b	—	2.40 TeV	1805.05959
	LeptoHiggs $Q\bar{Q} \rightarrow t\bar{t}$	$0, \mu, \tau$	2.1b, 2.2j	Yes	Z^* mass	2005.10338
	SSM $W^+ \rightarrow \nu\bar{\nu}$	$1, \mu, \tau$	—	139	W mass	1905.05909
	SSM $W^+ \rightarrow \nu\bar{\nu}$	$1, \mu, \tau$	—	Yes	3.1 TeV	1801.06992
	HVT $W^+ \rightarrow WZ \rightarrow f\bar{f}q$ model B	$1, \mu, \tau$	2/1/1j	Yes	W mass	2004.14638
	HVT $Z^* \rightarrow ZH$ model B	$0, \mu, \tau$	1.2b	3	Z^* mass	2007.00293
	HVT $Z^* \rightarrow W^+W^-$ model B	$0, \mu, \tau$	2.1b, 2.2j	Yes	W mass	2007.00293
	LRSM $W^+ \rightarrow \nu\bar{\nu}$	multi-channel	—	36.1	W mass	1807.14475
	LRSM $W^+ \rightarrow \nu\bar{\nu}$	$2, \mu$	1	—	60 MeV	1804.15822
	CI $q\bar{q}q$	$0, \mu, \tau$	2	—	37.0	1703.09127
CI	CI $q\bar{q}q$	$2, \mu, \tau$	—	139	A	2006.12946
	CI $q\bar{q}q$	$2, \mu, \tau$	—	139	A	2006.12946
	CI $q\bar{q}q$	$2, \mu, \tau$	—	139	A	2006.12946
	CI $q\bar{q}q$	$2, \mu, \tau$	—	139	A	2006.12946
	CI $q\bar{q}q$	$2, \mu, \tau$	—	139	A	2006.12946
	CI $q\bar{q}q$	$2, \mu, \tau$	—	139	A	2006.12946
	CI $q\bar{q}q$	$2, \mu, \tau$	—	139	A	2006.12946
	CI $q\bar{q}q$	$2, \mu, \tau$	—	139	A	2006.12946
	CI $q\bar{q}q$	$2, \mu, \tau$	—	139	A	2006.12946
	CI $q\bar{q}q$	$2, \mu, \tau$	—	139	A	2006.12946
DM	Asial-vector med. (Dirac DM)	$0, \mu, \tau, \gamma$	1-4	Yes	$E_{\text{miss}}^{\text{min}}$	2102.10874
	Pseudo-scalar med. (Dirac DM)	$0, \mu, \tau, \gamma$	1-4	Yes	$E_{\text{miss}}^{\text{min}}$	2102.10874
	Vector med. $Z^*, 2HDM$ Dirac	$0, \mu, \tau$	2b	Yes	$E_{\text{miss}}^{\text{min}}$	2102.10874
	Pseudo-scalar med. $2HDM +$	$0, \mu, \tau$	2b	Yes	$E_{\text{miss}}^{\text{min}}$	2102.10874
	Scalar sector, $d = 5$ (Dirac DM)	$0, \mu, \tau, \gamma$	1b, 0.5j	Yes	$E_{\text{miss}}^{\text{min}}$	2102.10874
	Scalar LO 1 st gen	$2, \mu$	2	Yes	LO med.	2006.05872
	Scalar LO 2 nd gen	$2, \mu$	2	Yes	LO med.	2006.05872
	Scalar LO 3 rd gen	$1, \mu$	2b	Yes	LO med.	2006.05872
	Scalar LO 2 nd gen	$0, \mu, \tau$	2.1b, 2.2b	Yes	LO med.	2006.05872
	Scalar LO 3 rd gen	$2, \mu, \tau$	1.2c, 1.1b	Yes	LO med.	2006.05872
LO	Scalar LO 2 nd gen	$2, \mu, \tau$	1.2c, 1.1b	Yes	LO med.	2006.05872
	Scalar LO 3 rd gen	$2, \mu, \tau$	1.2c, 1.1b	Yes	LO med.	2006.05872

*Only a selection of the available mass limits on new states or phenomena is shown.

† Small-radius (large-radius) jets are denoted by the letter i (J).

Suggests **energy gap** between SM and New Physics (NP)

Energy scale hierarchy motivates use of **Effective Field Theory (EFT)**

Different physics at different Energy scales

Fermi Theory : Low energy EFT of the SM (Full underlying theory)

The image shows two Feynman diagrams. The top diagram represents the full theory, with an incoming fermion line f_1 and an outgoing fermion line f_3 on the left, and an incoming fermion line f_2 and an outgoing fermion line f_4 on the right. They are connected by a wavy line representing a W boson propagator with momentum p . The vertices are labeled g . The full amplitude is given as $\mathcal{M}_{\text{full}} \sim \frac{g^2}{p^2 - m_W^2}$. The bottom diagram represents the effective theory, showing a contact interaction where the four fermion lines (f_1, f_2, f_3, f_4) meet at a single point labeled G_F . The effective amplitude is given as $\mathcal{M}_{\text{EFT}} \sim G_F \propto -\frac{g^2}{m_W^2}$.

$$\mathcal{M}_{\text{full}} \sim \frac{g^2}{p^2 - m_W^2}$$
$$\mathcal{M}_{\text{EFT}} \sim G_F \propto -\frac{g^2}{m_W^2}$$

W propagator $\rightarrow$ contact interaction

Correctly describes physics if $p^2 \ll m_W^2$

In turn can **treat SM as low energy EFT** of some unknown underlying theory

Recipe for a SM EFT

Don't know underlying theory so take a “**bottom-up**” approach :

- **1) Particles** : Include **all particles with $m \ll \Lambda$** $\rightarrow$ Operators ($\mathcal{O}_i$) built from combinations of SM Fields and derivatives
- **2) Symmetries** : Require $\mathcal{O}_i$ obey **SM symmetries** (Gauge, Lorentz)
- **3) Dimension Counting** : Split coupling in front of $\mathcal{O}_i$ into constant c_i and some power of $\Lambda \rightarrow$ Organize $\mathcal{O}_i$ with respect to their mass dimension D

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{c_i^{(5)}}{\Lambda} \mathcal{O}_i^{(5)} + \sum_i \frac{c_i^{(6)}}{\Lambda^2} \mathcal{O}_i^{(6)} + \sum_i \frac{c_i^{(7)}}{\Lambda^3} \mathcal{O}_i^{(7)} + \sum_i \frac{c_i^{(8)}}{\Lambda^4} \mathcal{O}_i^{(8)} + \dots,$$

Recipe for a SM EFT

Don't know underlying theory so take a “**bottom-up**” approach :

- **1) Particles** : Include **all particles with $m \ll \Lambda$** $\rightarrow$ Operators ($\mathcal{O}_i$) built from combinations of SM Fields and derivatives
- **2) Symmetries** : Require $\mathcal{O}_i$ obey **SM symmetries** (Gauge, Lorentz)
- **3) Dimension Counting** : Split coupling in front of $\mathcal{O}_i$ into constant c_i and some power of $\Lambda \rightarrow$ Organize $\mathcal{O}_i$ with respect to their mass dimension D

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \cancel{\sum_i \frac{c_i^{(5)}}{\Lambda} \mathcal{O}_i^{(5)}} + \sum_i \frac{c_i^{(6)}}{\Lambda^2} \mathcal{O}_i^{(6)} + \cancel{\sum_i \frac{c_i^{(7)}}{\Lambda^3} \mathcal{O}_i^{(7)}} + \sum_i \frac{c_i^{(8)}}{\Lambda^4} \mathcal{O}_i^{(8)} + \dots,$$

$\mathcal{O}_i^{(5)}$ violate lepton number and $\mathcal{O}_i^{(7)}$ violate Lepton or Baryon number

Recipe for a SM EFT

Don't know underlying theory so take a “**bottom-up**” approach :

- **1) Particles** : Include **all particles with $m \ll \Lambda$** $\rightarrow$ Operators ($\mathcal{O}_i$) built from combinations of SM Fields and derivatives
- **2) Symmetries** : Require $\mathcal{O}_i$ obey **SM symmetries** (Gauge, Lorentz)
- **3) Dimension Counting** : Split coupling in front of $\mathcal{O}_i$ into constant c_i and some power of $\Lambda \rightarrow$ Organize $\mathcal{O}_i$ with respect to their mass dimension D

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{c_i^{(5)}}{\Lambda} \mathcal{O}_i^{(5)} + \sum_i \frac{c_i^{(6)}}{\Lambda^2} \mathcal{O}_i^{(6)} + \sum_i \frac{c_i^{(7)}}{\Lambda^3} \mathcal{O}_i^{(7)} + \sum_i \frac{c_i^{(8)}}{\Lambda^4} \mathcal{O}_i^{(8)} + \dots,$$

$\mathcal{O}_i^{(5)}$ violate lepton number and $\mathcal{O}_i^{(7)}$ violate Lepton or Baryon number

$\mathcal{O}_i^{(8)}$ suppressed by $1/\Lambda^4 \rightarrow$ Typically neglected

A SM EFT

$$\mathcal{L}_{\text{EFT}} = \mathcal{L}_{\text{SM}} + \sum_i \bar{c}_i^{(6)} \mathcal{O}_i^{(6)}.$$

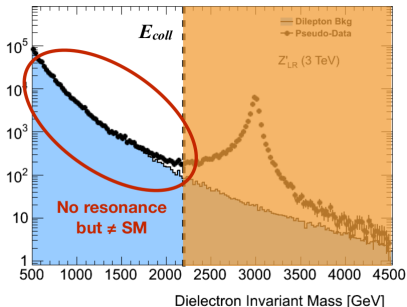
Effects of heavy **NP** mapped onto $\mathcal{O}_i^{(6)}$

c_i (**Wilson coefficient**) specify the strength of the new interactions

Contribution of $\mathcal{O}_i^{(6)}$ scale as $(E/\Lambda)^2 \rightarrow$ Can induce growth with $\sqrt{s}$

EFT only **valid** at $E < \Lambda$

For example : Effect of Z' on dielectron mass when $E < m_{Z'}$



A SM EFT

$$\mathcal{L}_{\text{EFT}} = \mathcal{L}_{\text{SM}} + \sum_i \bar{c}_i^{(6)} \mathcal{O}_i^{(6)}.$$

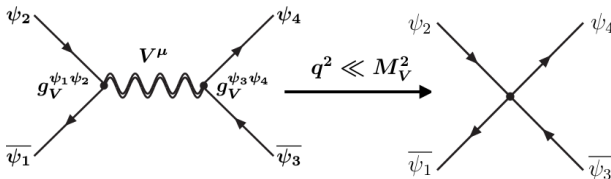
Effects of heavy **NP** mapped onto $\mathcal{O}_i^{(6)}$

c_i (**Wilson coefficient**) specify the strength of the new interactions

Contribution of $\mathcal{O}_i^{(6)}$ scale as $(E/\Lambda)^2 \rightarrow$ Can induce growth with $\sqrt{s}$

EFT only **valid** at $E < \Lambda$

For example : Effect of Z' on dielectron mass when $E < m_{Z'}$



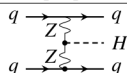
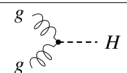
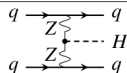
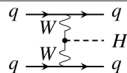
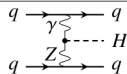
Deformation in dielectron mass described by $\mathcal{O}_i^{(6)}$ effective interaction

Basis of Dimension 6 Operators

Warsaw basis : First complete and non-redundant set of operators (59)

Popular SMEFT basis in the theoretical community

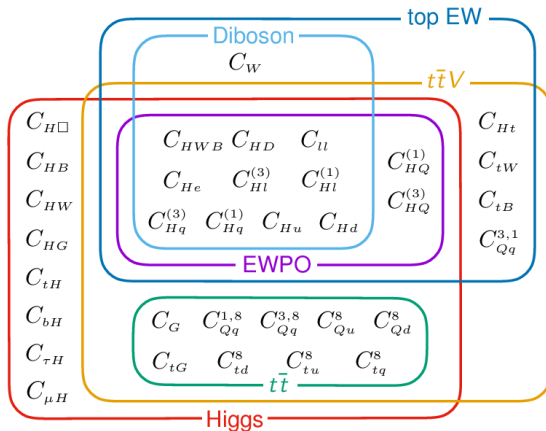
Most relevant for HVV interactions ($V = \text{Gauge boson}$) :

Coefficient	Operator	Example process
c_{HDD}	$(H^\dagger D^\mu H)^* (H^\dagger D_\mu H)$	
c_{HG}	$H^\dagger H G_{\mu\nu}^A G^{A\mu\nu}$	
c_{HB}	$H^\dagger H B_{\mu\nu} B^{\mu\nu}$	
c_{HW}	$H^\dagger H W_{\mu\nu}^I W^{I\mu\nu}$	
c_{HWB}	$H^\dagger \tau^I H W_{\mu\nu}^I B^{\mu\nu}$	

+ corresponding CP–Odd operators $\tilde{O}_{HG}, \tilde{O}_{HB}, \tilde{O}_{HW}, \tilde{O}_{HWB}$

Basis of Dimension 6 Operators

Operators may contribute to **several classes of measurements**



Long term goal : ATLAS and CMS combination (**Global Fit**) of EW+Higgs+Top measurements to detect SM deformations

Basis of Dimension 6 Operators

Higgs basis : Convenient (experimentally) to express EFT in terms of **mass eigenstates after EWSB** :

$$\begin{aligned} \mathcal{L}_{\text{hvv}} = & \frac{h}{v} \left[(1 + \delta c_z) \frac{(g^2 + g'^2)v^2}{4} Z_\mu Z_\mu + c_{zz} \frac{g^2 + g'^2}{4} Z_{\mu\nu} Z_{\mu\nu} + c_{z\Box} g^2 Z_\mu \partial_\nu Z_{\mu\nu} + \tilde{c}_{zz} \frac{g^2 + g'^2}{4} Z_{\mu\nu} \tilde{Z}_{\mu\nu} \right. \\ & + (1 + \delta c_w) \frac{g^2 v^2}{2} W_\mu^+ W_\mu^- + c_{ww} \frac{g^2}{2} W_{\mu\nu}^+ W_{\mu\nu}^- + c_{w\Box} g^2 (W_\mu^- \partial_\nu W_\mu^+ + \text{h.c.}) + \tilde{c}_{ww} \frac{g^2}{2} W_{\mu\nu}^+ \tilde{W}_{\mu\nu}^- \\ & + c_{z\gamma} \frac{e\sqrt{g^2 + g'^2}}{2} Z_{\mu\nu} A_{\mu\nu} + \tilde{c}_{z\gamma} \frac{e\sqrt{g^2 + g'^2}}{2} Z_{\mu\nu} \tilde{A}_{\mu\nu} + c_{\gamma\Box} gg' Z_\mu \partial_\nu A_{\mu\nu} \\ & \left. + c_{\gamma\gamma} \frac{e^2}{4} A_{\mu\nu} A_{\mu\nu} + \tilde{c}_{\gamma\gamma} \frac{e^2}{4} A_{\mu\nu} \tilde{A}_{\mu\nu} + c_{gg} \frac{g_s^2}{4} G_{\mu\nu}^a G_{\mu\nu}^a + \tilde{c}_{gg} \frac{g_s^2}{4} G_{\mu\nu}^a \tilde{G}_{\mu\nu}^a \right], \end{aligned}$$

More transparent connection to measurable quantities

SILH basis : Strongly-Interacting Light Higgs "**top-down**" EFT

More direct mapping to some BSM scenarios (Composite Higgs)

Both related to Warsaw basis through linear transformations

So how are Higgs related EFT coefficients measured?

Dedicated (detector-level) measurements :

Use **dedicated discriminants** and **full simulation** of Signal PDFs

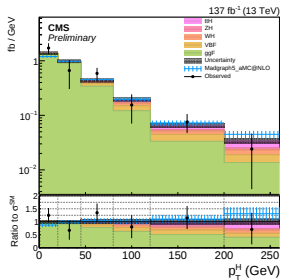
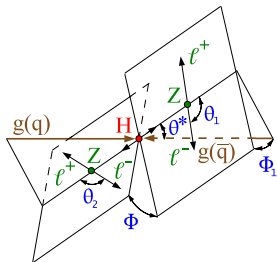
Interpretations of other measurements :

2 types : Differential fiducial σ and simplified template σ (STXS)

Signal **parameterised in gen-level** fiducial bins

Pros and cons for all approaches

Instructive to go through some recent examples from Run 2!



STXS measurements

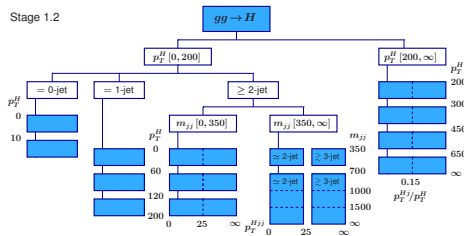
Measure σ in **pre-defined kinematic bins per production mode**

Aim to **minimize theory dependence** while
maximizing sensitivity to BSM effects

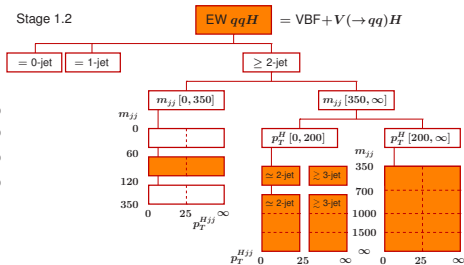
No fiducial phase space (except for $|y_H| < 2.5$)

Can combine decay channels but significant model dependence

Stage 1.2



Stage 1.2



ATLAS STXS interpretations with Warsaw basis

Using $H \rightarrow \gamma\gamma$, $H \rightarrow ZZ$ and $V(\text{lep})H \rightarrow b\bar{b}$

SMEFT σ parameterisation :

$$\sigma \propto |\mathcal{M}_{SM} + \sum_i \bar{c}_i \mathcal{M}_i|^2$$

$$\sigma = \sigma_{SM} + \sigma_{Int} + \sigma_{BSM}$$

$$\frac{\sigma}{\sigma_{SM}} = 1 + \sum_i A_i \bar{c}_i + \sum_{ij} B_{ij} \bar{c}_i \bar{c}_j$$

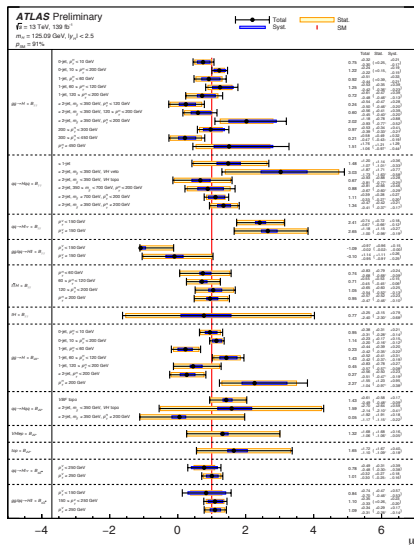
Remember $\bar{c}_i = c_i/\Lambda^2$ so we have :

Linear ($\bar{c}_i$) term suppressed by $1/\Lambda^2$

Quadratic ($\bar{c}_i \bar{c}_j$) term suppressed by $1/\Lambda^4$

A_i and B_{ij} are extracted from simulation

(Same procedure for decay param)



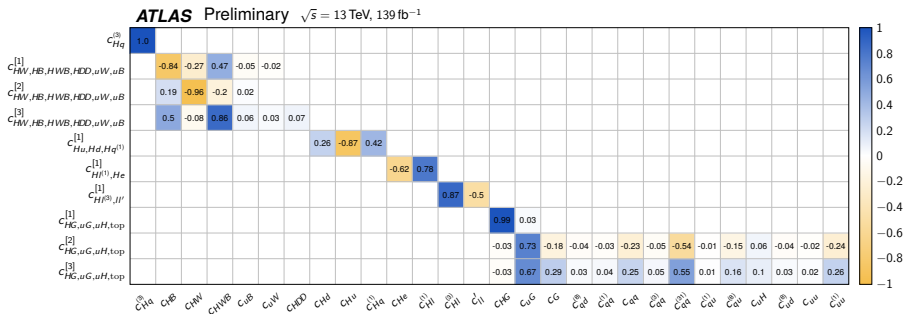
A modified Warsaw basis

Consider **all CP-Even** $\mathcal{O}_i^6$ - Insufficient info to constrain all c_i simultaneously

Use Fisher information and find eigenvectors

→ **Sensitive directions** in $\mathcal{O}_i$ space

Set constraints on linear combinations of c_i (Modified basis)



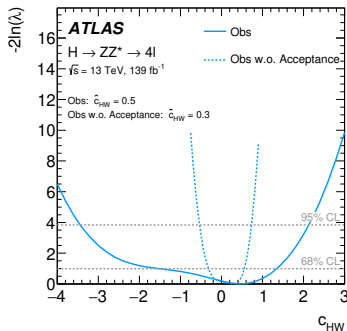
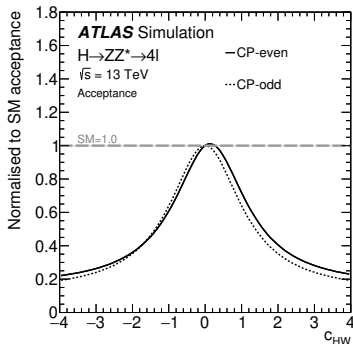
Degeneracy broken with info from other measurements? (VV, top, ...)

Acceptance effects

Possibility of **acceptance changes** induced by operators

$H \rightarrow ZZ \rightarrow 4\ell$ is significantly impacted
(Particularly due to m_{Z_2} requirements in the analysis)

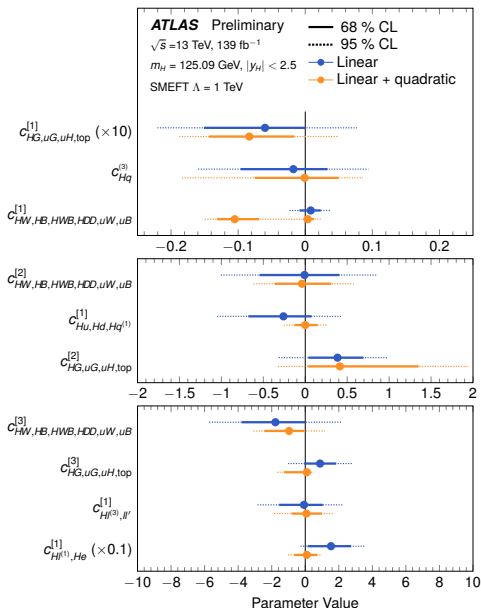
Ad-hoc correction derived and applied to ZZ decay parameterisation



For all other H decays acceptance effects are neglected

***Effect could be significant for HWW also**

ATLAS STXS interpretations with Warsaw basis



CMS STXS interpretations with SILH basis

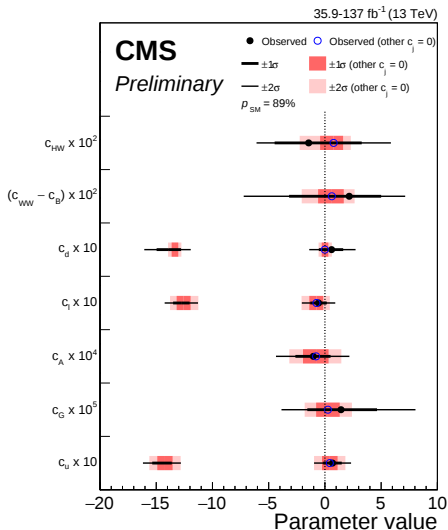
Using H decays to ZZ, $\gamma\gamma$, WW, bb
and $\tau\tau$

Considered **8 leading CP-even**
operators

Linear + Quadratic terms
considered

For full run 2 analysis **moving to**
Warsaw basis

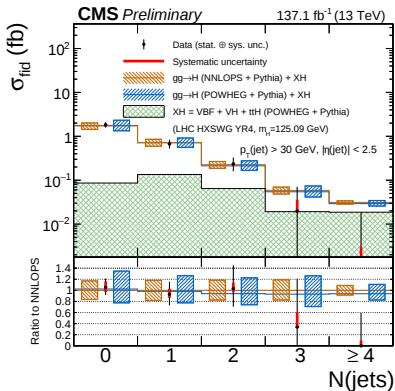
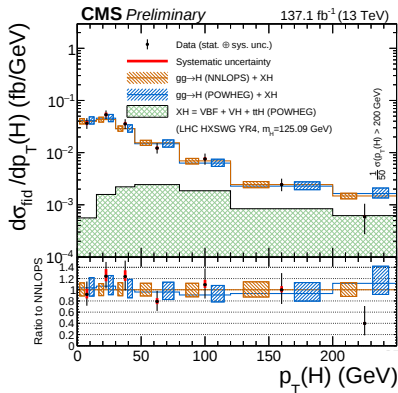
Also **acceptance effects to be**
included!



Fiducial Differential measurements

Fiducial : Target phase-space region matching experimental selection
(Minimize model dependence)

Differential : In bins of a variable (p_T^H , $N_{Jets}, \dots$)
(generally 1 or 2 dimensions)



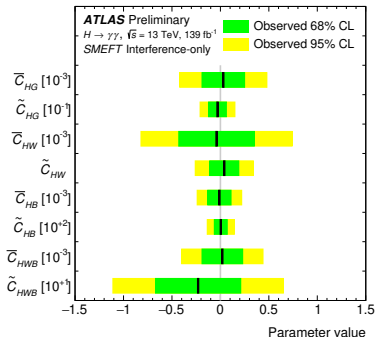
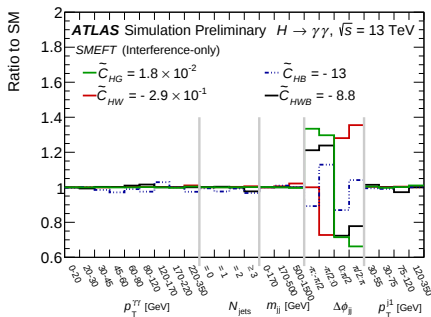
Fiducial depends on final state → Not trivial to combine different modes

ATLAS $H \rightarrow \gamma\gamma$ interpretation with Warsaw basis

Several 1D measurements : $p_T^{\gamma\gamma}$, N_{Jets} , m_{jj} , $\Delta\phi_{jj}$ and p_T^{j1}

Correlations between bins in different distributions calculated using bootstrapping
 $\rightarrow$ Use all variables in interpretation

Target **8 leading CP-Even and Odd** (Sensitivity from $\Delta\phi_{jj}$) $\mathcal{O}_i^6$
 $\rightarrow$ Constraining one at a time



Constraints much greater when including quadratic terms for CP-odd $\mathcal{O}_i^6$

Dedicated HZZ EFT measurements at CMS

General Lorentz invariant form of **HVV** scattering amplitude :

$$A(\text{HVV}) = \frac{1}{v} \left[a_1^{\text{VV}} + \frac{\kappa_1^{\text{VV}} q_{V1}^2 + \kappa_2^{\text{VV}} q_{V2}^2}{(\Lambda_1^{\text{VV}})^2} + \frac{\kappa_3^{\text{VV}} (q_{V1} + q_{V2})^2}{(\Lambda_Q^{\text{VV}})^2} \right] m_{V1}^2 \epsilon_{V1}^* \epsilon_{V2}^* \\ + \frac{1}{v} a_2^{\text{VV}} f_{\mu\nu}^{*(1)} f^{*(2),\mu\nu} + \frac{1}{v} a_3^{\text{VV}} f_{\mu\nu}^{*(1)} \tilde{f}^{*(2),\mu\nu},$$

SU(2)xU(1) enforces relations between ZZ,WW, $\gamma\gamma$ and $Z\gamma$ couplings :

$$\begin{aligned} a_1^{\text{WW}} &= a_1^{\text{ZZ}}, \\ a_2^{\text{WW}} &= c_w^2 a_2^{\text{ZZ}} + s_w^2 a_2^{\gamma\gamma} + 2s_w c_w a_2^{Z\gamma}, \\ a_3^{\text{WW}} &= c_w^2 a_3^{\text{ZZ}} + s_w^2 a_3^{\gamma\gamma} + 2s_w c_w a_3^{Z\gamma}, \\ \frac{\kappa_1^{\text{WW}}}{(\Lambda_1^{\text{WW}})^2} (c_w^2 - s_w^2) &= \frac{\kappa_1^{\text{ZZ}}}{(\Lambda_1^{\text{ZZ}})^2} + 2s_w^2 \frac{a_2^{\gamma\gamma} - a_2^{\text{ZZ}}}{m_Z^2} + 2 \frac{s_w}{c_w} (c_w^2 - s_w^2) \frac{a_2^{Z\gamma}}{m_Z^2}, \\ \frac{\kappa_2^{\text{Z}\gamma}}{(\Lambda_1^{\text{Z}\gamma})^2} (c_w^2 - s_w^2) &= 2s_w c_w \left(\frac{\kappa_1^{\text{ZZ}}}{(\Lambda_1^{\text{ZZ}})^2} + \frac{a_2^{\gamma\gamma} - a_2^{\text{ZZ}}}{m_Z^2} \right) + 2(c_w^2 - s_w^2) \frac{a_2^{Z\gamma}}{m_Z^2} \end{aligned}$$

Dedicated HZZ EFT measurements at CMS

General Lorentz invariant form of **HVV** scattering amplitude :

$$A(\text{HVV}) = \frac{1}{v} \left[a_1^{\text{VV}} + \frac{\kappa_1^{\text{VV}} q_{V1}^2 + \kappa_2^{\text{VV}} q_{V2}^2}{(\Lambda_1^{\text{VV}})^2} + \frac{\kappa_3^{\text{VV}} (q_{V1} + q_{V2})^2}{(\Lambda_Q^{\text{VV}})^2} \right] m_{V1}^2 \epsilon_{V1}^* \epsilon_{V2}^* \\ + \frac{1}{v} a_2^{\text{VV}} f_{\mu\nu}^{*(1)} f^{*(2),\mu\nu} + \frac{1}{v} a_3^{\text{VV}} f_{\mu\nu}^{*(1)} \tilde{f}^{*(2),\mu\nu},$$

SU(2)xU(1) enforces relations between ZZ,WW, $\gamma\gamma$ and $Z\gamma$ couplings :

$$\begin{aligned} a_1^{\text{WW}} &= a_1^{\text{ZZ}}, \\ a_2^{\text{WW}} &= c_w^2 a_2^{\text{ZZ}} + s_w^2 a_2^{\gamma\gamma} + 2s_w c_w a_2^{Z\gamma}, \\ a_3^{\text{WW}} &= c_w^2 a_3^{\text{ZZ}} + s_w^2 a_3^{\gamma\gamma} + 2s_w c_w a_3^{Z\gamma}, \\ \frac{\kappa_1^{\text{WW}}}{(\Lambda_1^{\text{WW}})^2} (c_w^2 - s_w^2) &= \frac{\kappa_1^{\text{ZZ}}}{(\Lambda_1^{\text{ZZ}})^2} + 2s_w^2 \frac{a_2^{\gamma\gamma} - a_2^{\text{ZZ}}}{m_Z^2} + 2 \frac{s_w}{c_w} (c_w^2 - s_w^2) \frac{a_2^{Z\gamma}}{m_Z^2}, \\ \frac{\kappa_2^{\text{Z}\gamma}}{(\Lambda_1^{\text{Z}\gamma})^2} (c_w^2 - s_w^2) &= 2s_w c_w \left(\frac{\kappa_1^{\text{ZZ}}}{(\Lambda_1^{\text{ZZ}})^2} + \frac{a_2^{\gamma\gamma} - a_2^{\text{ZZ}}}{m_Z^2} \right) + 2(c_w^2 - s_w^2) \frac{a_2^{Z\gamma}}{m_Z^2} \end{aligned}$$

Assume $a^{\gamma\gamma}$ and $a^{Z\gamma}$ constrained by $\text{H} \rightarrow \gamma\gamma$ and $\text{H} \rightarrow \text{Z}\gamma$ measurements

4 independent couplings : a_1 (SM), a_2 , a_3 (CP-Odd), k_1

Dedicated HZZ EFT measurements at CMS

Amplitude a_i map directly to **Higgs basis EFT couplings** :

$$\begin{aligned}\delta c_z &= \frac{1}{2}a_1 - 1, \\ c_{z\Box} &= \frac{m_Z^2 s_w^2}{e^2} \frac{\kappa_1}{(\Lambda_1)^2}, \\ c_{zz} &= -\frac{2s_w^2 c_w^2}{e^2} a_2, \\ \tilde{c}_{zz} &= -\frac{2s_w^2 c_w^2}{e^2} a_3.\end{aligned}$$

Full detector simulation of corresponding HZZ states

For 1 HVV vertex $\sigma \propto |a_1 \mathcal{M}_1 + a_2 \mathcal{M}_2 + a_3 \mathcal{M}_3 + k_1 \mathcal{M}_{k1}|^2$

Many interference terms!

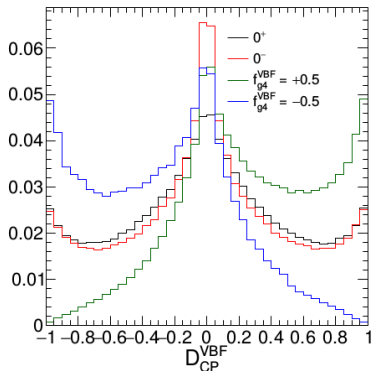
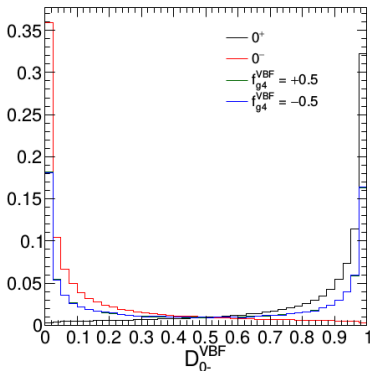
Use reweighting rather than simulate all required terms

Dedicated Observables

Exploit full production (ggF,VBF,VH) and decay info ($ZZ \rightarrow 4\ell$)
with **ME based discriminants**

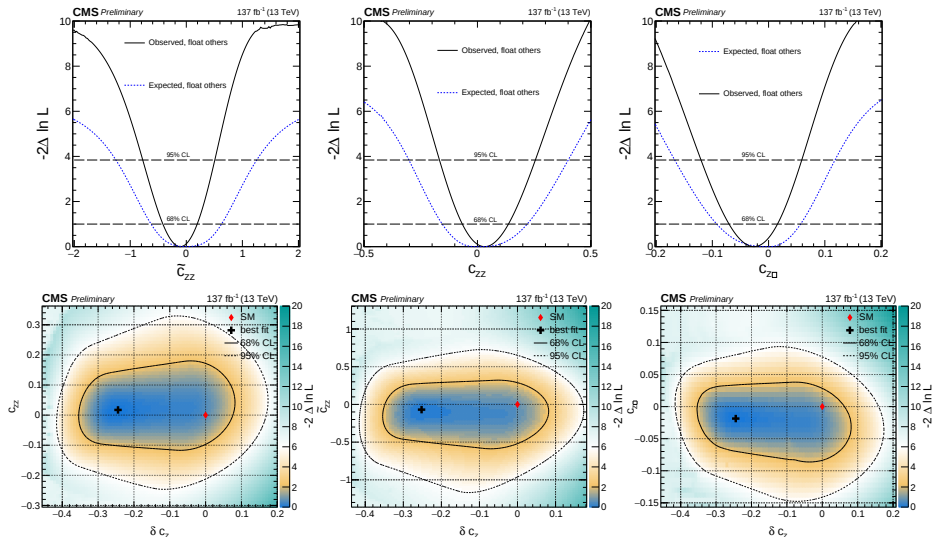
Can target **production mode**, **Higgs coupling model** and **interference**

For example to study a_3 (CP-Odd) in VBF HZZ can use :



Dedicated studies of **other Higgs vertices** include ggH, ttH, $H\tau\tau$

Scans of Higgs basis EFT couplings



Can rotate to warsaw basis → In principle could include in **global fits**

Dedicated HWW analysis also under development

Some Lessons going forward

STXS/Differential approach promising for EFT global fits but.....

1) Carefull when **extrapolating to "fiducial regions"**

Often relying on SM templates to determine efficiencies/acceptances

2) **Higgs decay** information not included

Angular info of decay (e.g HZZ 4l) sensitive to BSM effects

3) STXS not sensitive to **CP-Odd** operators

Some Lessons going forward

STXS/Differential approach promising for EFT global fits but.....

1) Carefull when **extrapolating to "fiducial regions"**

Often relying on SM templates to determine efficiencies/acceptances

2) **Higgs decay** information not included

Angular info of decay (e.g HZZ 4l) sensitive to BSM effects

3) STXS not sensitive to **CP-Odd** operators

Dedicated analysis also promising and avoids these issues but....

1) Lacks **re-interpretation flexibility** of STXS/Differential

Necessary to redo analysis for any change in theory

2) Limited to a **smaller set of operators**

3) **Time-consuming and computer-intensive** to implement

Some Lessons going forward

STXS/Differential approach promising for EFT global fits but.....

1) Carefull when **extrapolating to "fiducial regions"**

Often relying on SM templates to determine efficiencies/acceptances

2) **Higgs decay** information not included

Angular info of decay (e.g HZZ 4l) sensitive to BSM effects

3) STXS not sensitive to **CP-Odd** operators

Dedicated analysis also promising and avoids these issues but....

1) Lacks **re-interpretation flexibility** of STXS/Differential

Necessary to redo analysis for any change in theory

2) Limited to a **smaller set of operators**

3) **Time-consuming and computer-intensive** to implement

And finally for all approaches we are neglecting **effect of operators on backgrounds....**

Global fits from Theorists

Several theorists already use LHC measurements in their **global EFT fits**

Several fitting frameworks available
(different statistical interpretations, Linear/Quadratic, LO/NLO...)

EFT*fitter*

Fitmaker

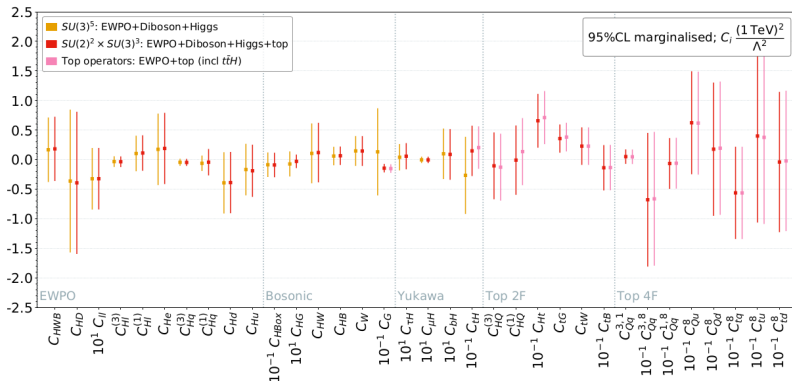


Combining **statistically independent measurements**
(include correlation info using published covariance matrices)

Combination with non-LHC constraints (LEP, Tevatron, etc)

Global fit with FitMaker

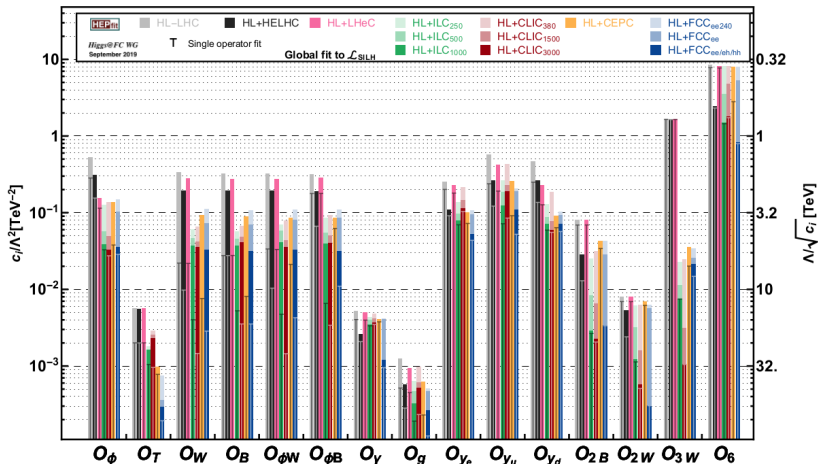
Top, Higgs, diboson and EWK fit with warsaw basis
(Neglecting quadratic $\mathcal{O}_i^6$ contributions)



Input from Higgs : Run 1 μ + Run 2 STXS + (CMS) HWW Differential

Global fit projections with HEPFit

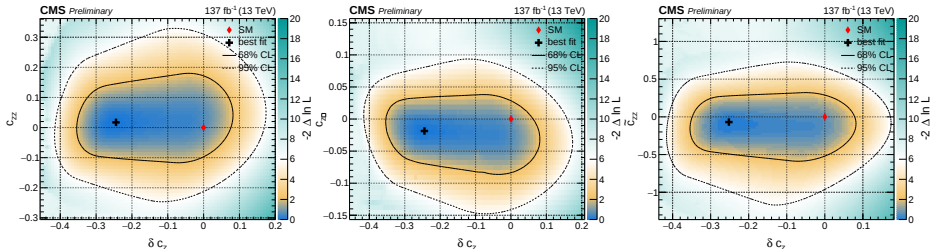
Projected Higgs (μ), diboson and EWK fit with SILH basis
(Neglecting quadratic $\mathcal{O}_i^6$ contributions)



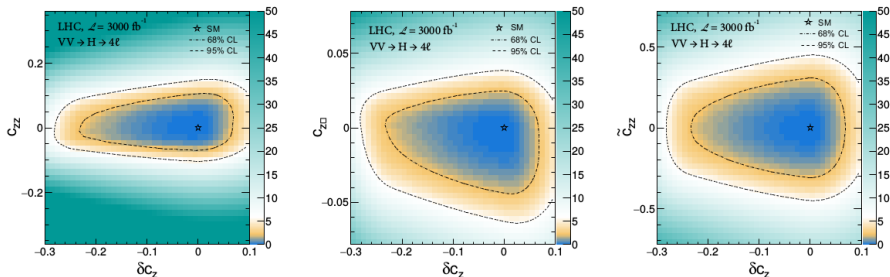
Large difference between global and single operator fit $\rightarrow$ large correlations

Projections for dedicated HZZ studies at CMS

Run 2 :



HL-LHC :



SUMMARY

- Given apparent energy gap between SM and new physics an EFT approach is well motivated at the LHC
- Long term goal would be ATLAS and CMS global fit of EW+Higgs+Top measurements to detect SM deformations
- Within Higgs groups there are a number of types of measurements that could be exploited in a global fit :

STXS

Fiducial Differential

Dedicated (detector-level) analysis

- Pros and cons of each and many recent developments discussed
- Exciting and fast moving area so watch this space!!