

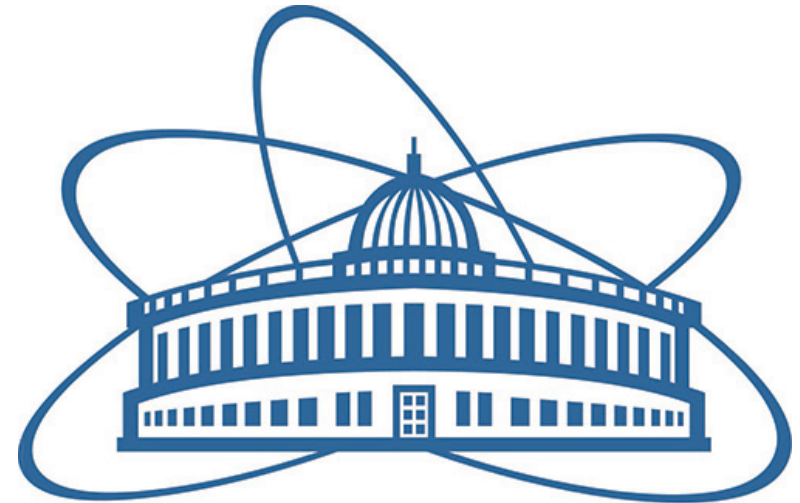
Neutrino magnetic moments in low-energy neutrino scattering on condensed matter systems

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Introduction

- The detectors for searching light particles of dark matter are actively discussed in the literature. To achieve the sensitivity to low-energy signals at a level of ~ 1 meV, condensed matter targets are proposed [1].
- Such detectors can also be used to study the low-energy neutrino scattering in order both to test the Standard Model and to search for the physics beyond the Standard Model [2].

Background

Kinematical regime

$$E_\nu \ll m, \quad T \leq \frac{2E_\nu^2}{m} \ll E_\nu, \quad E_\nu \ll \frac{1}{R_{\text{nuc}}},$$

where T is the energy transfer, m is the atomic mass, and R_{nuc} is the nuclear radius.

The differential cross section for an atom

According to [3]:

$$\frac{d\sigma_{i \rightarrow f}}{dT d\mathbf{q}^2 d\varphi_q} = [\Sigma^{(w)} + \Sigma^{(\mu)}] \delta(T - E_f + E_i).$$

The weak interaction contribution:

$$\Sigma^{(w)} = \frac{G_F^2}{2\pi^2} \left[C_V^2 \left(1 - \frac{q^2}{E_\nu^2}\right) + C_A^2 \left(1 + \frac{q^2}{E_\nu^2}\right) \right],$$

$$C_V = Z \left(\frac{1}{2} - 2 \sin^2 \theta_W\right) - \frac{1}{2}N + Z \left(\mp \frac{1}{2} + 2 \sin^2 \theta_W\right) F_{\text{el}}(q^2),$$

$$C_A^2 = (C_A^{\text{nuc}})^2 + \frac{1}{4} \sum_{n,l} [(L_+^{nl} - L_-^{nl}) F_{\text{el}}^{nl}(q^2)]^2,$$

$$(C_A^{\text{nuc}})^2 = \frac{g_A^2}{4} [(Z_+ - Z_-) - (N_+ - N_-)]^2,$$

where q is the momentum transfer, with $q^2 = 2mT$, the plus (minus) stands for $\nu = \nu_e$ ($\nu = \nu_{\mu,\tau}$), and Z (N) is the number of protons (neutrons). $F_{\text{el}}(q^2)$ is the Fourier transform of the electron density, $g_A = 1.25$, and Z_\pm and N_\pm (L_\pm^{nl}) are the numbers of protons and neutrons (electrons) with spin parallel (+) or antiparallel (-) to the nucleus spin (the total electron spin).

The electromagnetic interaction contribution:

$$\Sigma^{(\mu)} = \frac{\alpha^2 Z^3}{2m_e^2} |\mu_\nu|^2 \left(\frac{2m}{q^2} - \frac{1}{E_\nu} \right) F_{\text{scr}}^2(q^2),$$

where $\mu_\nu \sim 10^{-11} \mu_B$ is the neutrino dipole magnetic moment [4] and $F_{\text{scr}}^2(q^2) = [1 - F_{\text{el}}(q^2)]^2$.

General formalism

The differential cross section for the system of \mathcal{N} atoms:

$$\frac{d\sigma}{dT} = \int_0^\infty dq^2 \int_0^{2\pi} d\varphi_q \left[\Sigma^{(w)} + \Sigma^{(\mu)} \right] S(T, \mathbf{q}),$$

where $S(T, \mathbf{q}) = \sum_{i,f} w_i \left| \langle f | \sum_{j=1}^{\mathcal{N}} e^{i\mathbf{q}\mathbf{R}_j} | i \rangle \right|^2 \delta(T - E_f + E_i)$

is the dynamical structure factor, with \mathbf{R}_j being the position of the j th atom and w_i the probability for the system to be in the initial state $|i\rangle$.

Neutrino scattering on liquid ⁴He

$$\mathcal{N} \text{ free atoms} \quad S(T, \mathbf{q}) = \mathcal{N} \delta\left(T - \frac{q^2}{2m}\right)$$

$$\frac{d\sigma}{dT} = \mathcal{N} \left[C_V^2 \frac{G_F^2 m}{\pi} \left(1 - \frac{mT}{2E_\nu^2}\right) + \frac{\pi \alpha^2 Z^3}{m_e^2} |\mu_\nu|^2 \left(\frac{1}{T} - \frac{1}{E_\nu}\right) F_{\text{scr}}^2 \right],$$

where $C_V^2 = C_V^2(q^2 = 2mT)$ and $F_{\text{scr}}^2 = F_{\text{scr}}^2(q^2 = 2mT)$.

$$\mathcal{N} \text{ bound atoms (liquid)} \quad S(T, \mathbf{q}) = \frac{\mathcal{N}q}{2mu} \delta(T - uq)$$

$$\frac{d\sigma}{dT} = \mathcal{N} \left[C_V^2 \frac{G_F^2 T^2}{2\pi mu^3} \left(1 - \frac{T^2}{4u^2 E_\nu^2}\right) + \frac{\pi \alpha^2 Z^3}{m_e^2} |\mu_\nu|^2 \left(\frac{2mu^2}{T^2} - \frac{1}{E_\nu}\right) F_{\text{scr}}^2 \right],$$

where $C_V^2 = C_V^2(q^2 = T^2/u^2)$, $F_{\text{scr}}^2 = F_{\text{scr}}^2(q^2 = T^2/u^2)$ and u is the phonon speed in liquid ⁴He.

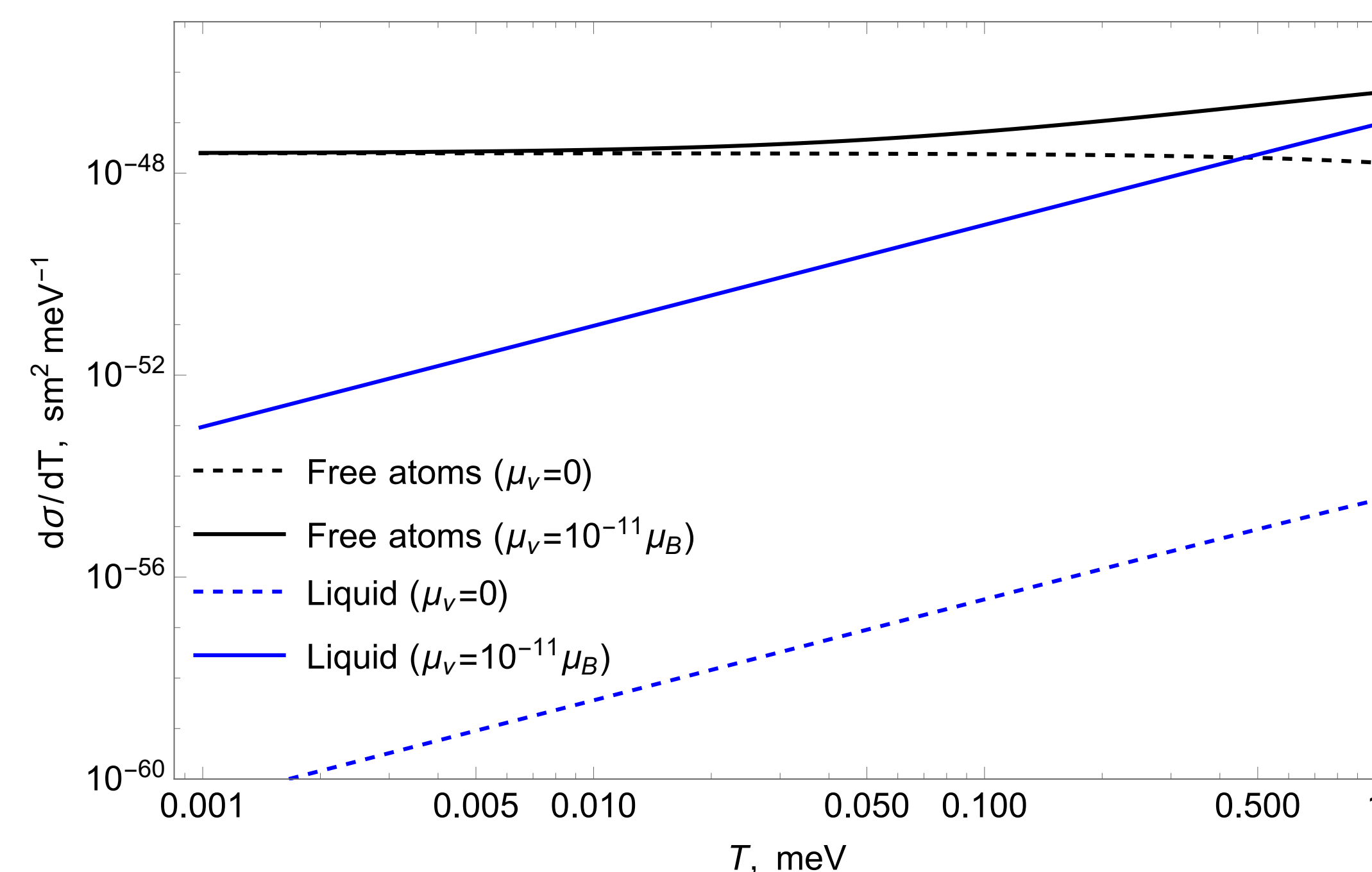


Figure: The differential cross section for tritium antineutrino scattering at $E_\nu = 10$ keV normalized to the number of helium atoms \mathcal{N} .

Conclusions

- The theory of low-energy neutrino scattering by a target in a condensed state with account for the neutrino anomalous magnetic moment is developed.
- It is shown that taking collective effects into account in the neutrino scattering by superfluid He-4 qualitatively changes the dependence of the differential cross section on the energy transfer.
- This fact must be taken into account both in the preparation and in the data analysis of future neutrino experiments with detectors based on liquid He-4 and other materials (e. g., graphene [5]).
- The obtained results will be used in the search for the electromagnetic properties of neutrinos [6].

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