

RESULTS FROM BaBar

P. Villanueva-Pérez
– Universitat de València-CSIC



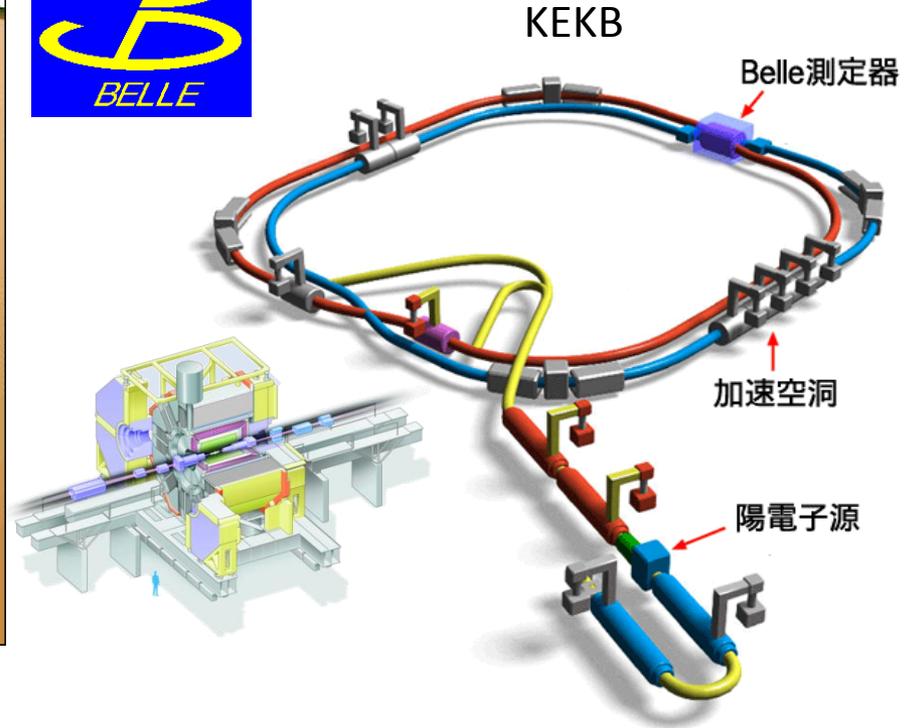
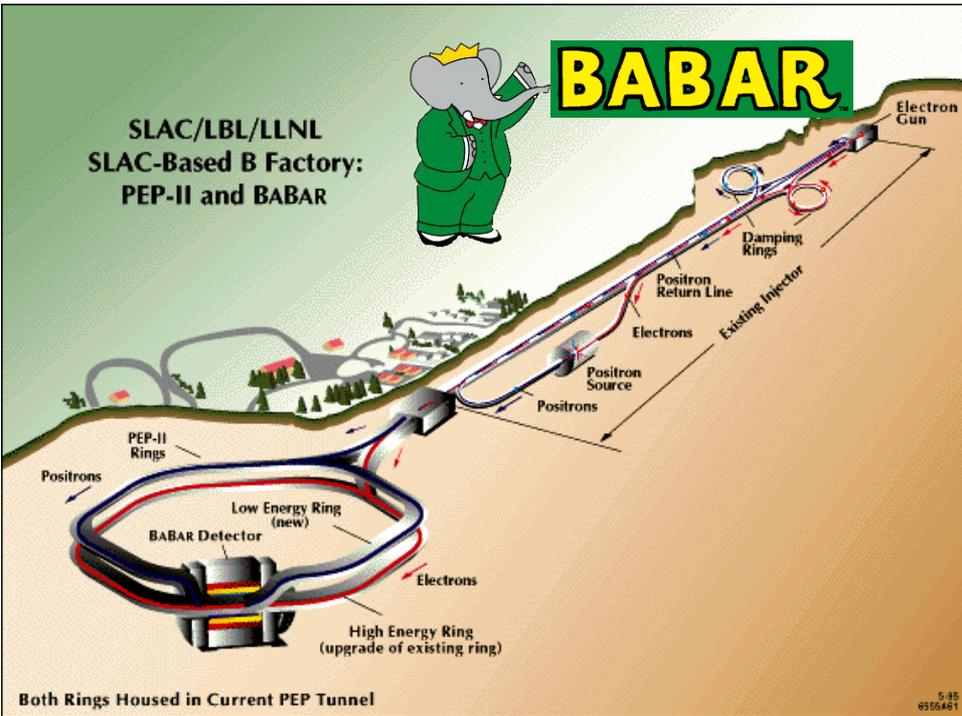
IV Jornadas del CPAN

B FACTORIES

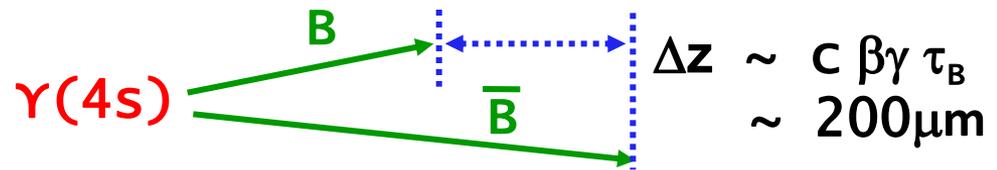
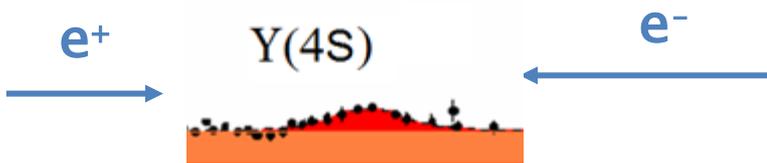
AND

BaBar EXPERIMENT

B Factories



$$\sqrt{s} = 10.58 \text{ GeV}$$



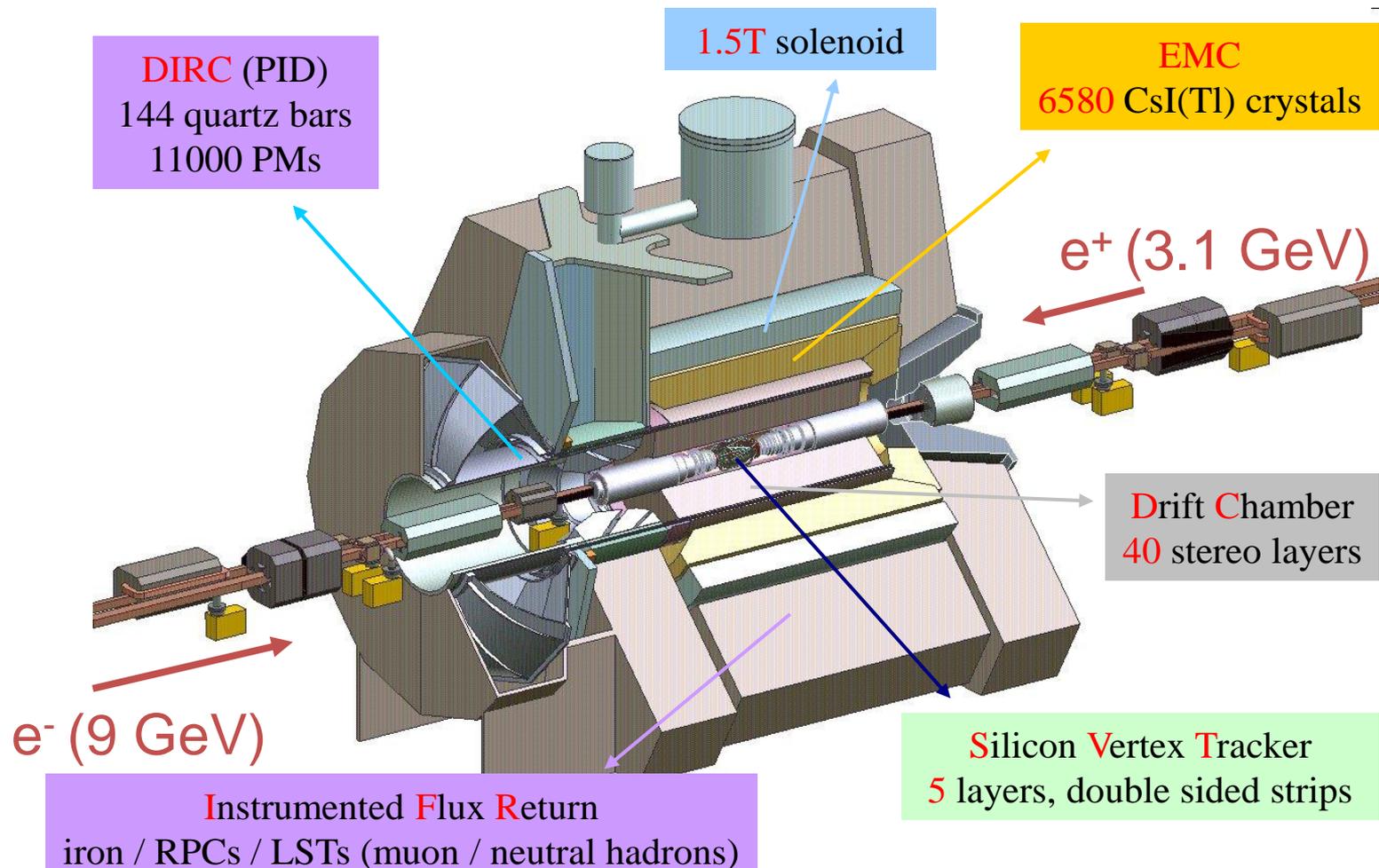
BaBar $p(e^-) = 9 \text{ GeV}$ $p(e^+) = 3.1 \text{ GeV}$

Belle $p(e^-) = 8 \text{ GeV}$ $p(e^+) = 3.5 \text{ GeV}$

$\beta\gamma = 0.56$

$\beta\gamma = 0.42$

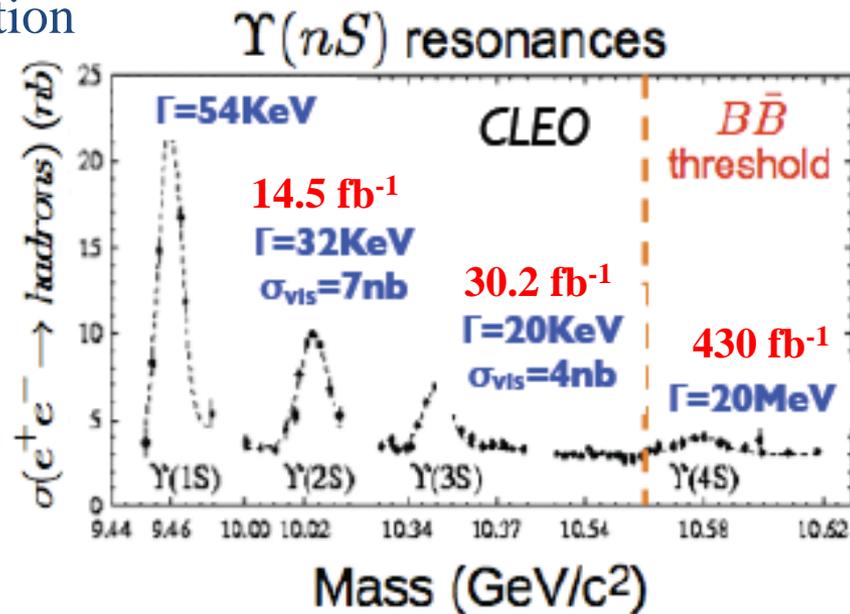
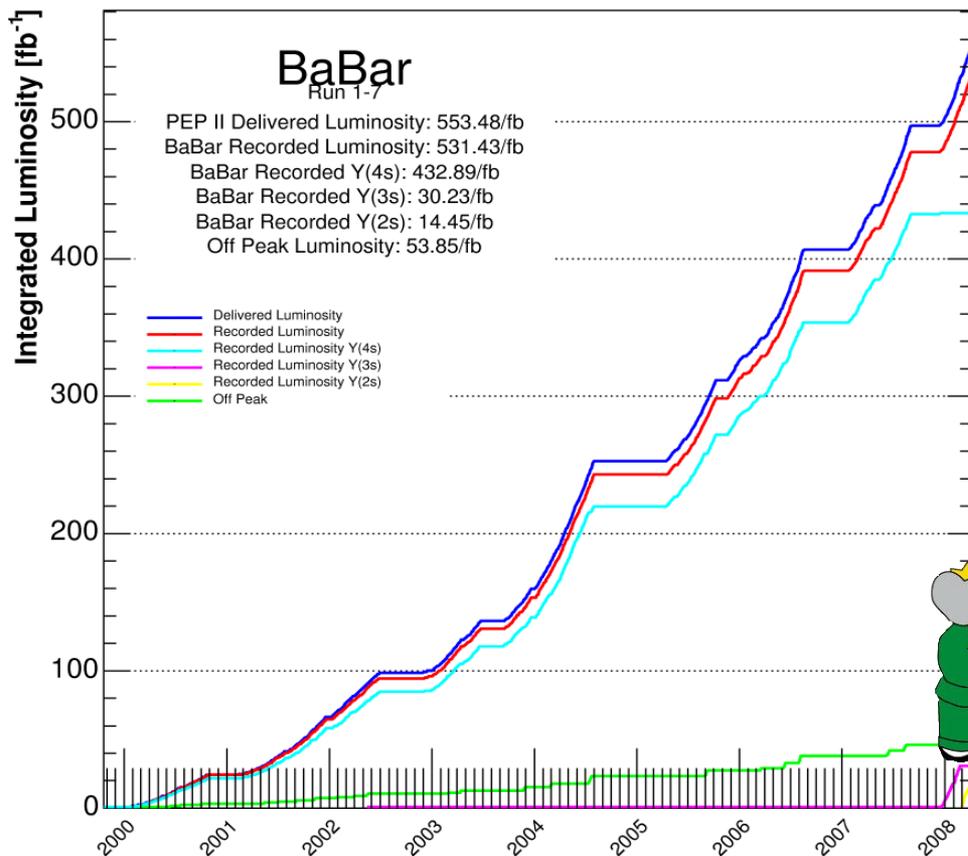
BaBar detector



- Asymmetric B-factory: $E_{\text{cms}} = 10.58 \text{ GeV}$ $e^+e^- \rightarrow \Upsilon(4S) \rightarrow B\bar{B}$
- Performed a wide range of flavor physics results in B, Charm and τ sectors
- General purpose detector in e^+e^- environment: precision tracking, photon/electron detection, particle ID, muon/ K_L identification. Very stable over the 9 years of operation

BaBar data

➤ 530 fb⁻¹ recorded in the 9 years of operation



54 fb⁻¹ Off-Y(nS)
4 fb⁻¹ above Y(4S)

- ≈ 470 × 10⁶ BB (0.5 × Belle)
- ≈ 690 × 10⁶ cc
- ≈ 500 × 10⁶ τ⁺τ⁻
- ≈ 1.2 × 10⁸ Y(3S) (7 × Belle + Cleo)
- ≈ 1.0 × 10⁸ Y(2S) (0.5 × Belle + Cleo)

Outline

- Time Reversal Violation
 - Theoretical motivation
 - Analysis Procedure
 - Results and interpretation

- $B \rightarrow D^{(*)} \tau \nu$
 - Motivation
 - Analysis Procedure
 - Results
 - 2HDM discussion

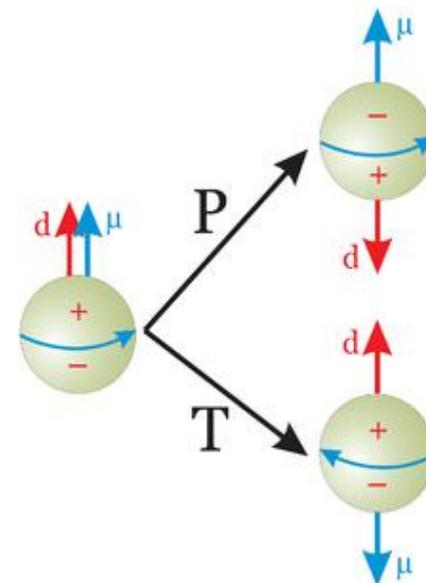
TIME REVERSAL VIOLATION

J. P. Lees et al. Phys. Rev. Lett. 109, 211801 (2012)

Motivation

• Time Reversal in stable systems

- A non-zero value of a T-odd observable in a stationary state, e.g, electric dipole moment of an elementary particle or an atom.
- In an oscillation a difference in the probability of $a \rightarrow b$ from $b \rightarrow a$ at a given time, e.g., $\nu_e \rightarrow \nu_\mu$ vs. $\nu_\mu \rightarrow \nu_e$ experiment proposed for the neutrino factories with muon storage ring.



• Time Reversal in unstable systems

1. Reversal of motion ($t \rightarrow -t$). ~~discard~~
2. $|\text{in}\rangle \leftrightarrow |\text{out}\rangle$ exchange. ~~Odd effects $t \rightarrow -t$~~

→ Experimentally tricky!



CP violation mechanisms

- Decay
- Mixing
- Mixing \times Decay

T violation mechanisms

- Decay
- Mixing
- Mixing \times Decay

CPT



TRV in unstable systems

• Decay TRV searches

$$\text{CP} \left\{ \begin{array}{l} B^0 \rightarrow K^+ \pi^-, R_1 \\ \bar{B}^0 \rightarrow K^- \pi^+, R_2 \end{array} \right\} \xleftrightarrow{\text{CPT}} \left\{ \begin{array}{l} K^- \pi^+ \rightarrow \bar{B}^0, R_1 \\ K^+ \pi^- \rightarrow B^0, R_2 \end{array} \right.$$

Unable to perform the T test:

- Preparation of the initial state.
- The strong processes will swamp the feeble weak processes.

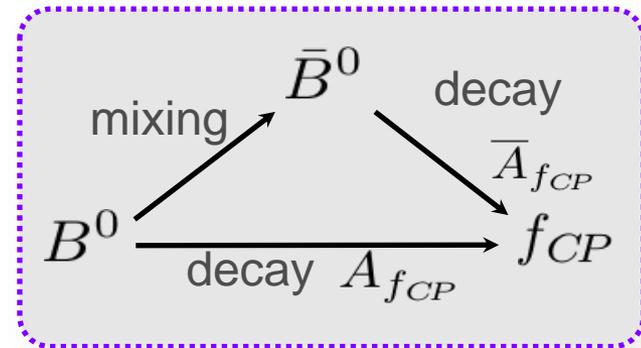
• Mixing TRV searches

$$\begin{array}{ccc} \begin{array}{l} K^0 \rightarrow \bar{K}^0 \\ B^0 \rightarrow \bar{B}^0 \end{array} & \xrightarrow{\text{CPT}} & \begin{array}{l} K^0 \rightarrow \bar{K}^0 \\ B^0 \rightarrow \bar{B}^0 \end{array} \\ \text{CP} \downarrow & & \text{T} \searrow \\ \begin{array}{l} \bar{K}^0 \rightarrow K^0 \\ \bar{B}^0 \rightarrow B^0 \end{array} & & \begin{array}{l} \bar{K}^0 \rightarrow K^0 \\ \bar{B}^0 \rightarrow B^0 \end{array} \end{array} \quad \text{CPLEAR}$$

We cannot distinguish CP and T.

Not a DIRECT observation of TRV

• Interference TRV searches



CPV time dependent (TD) studies:

- There are no exchanges $t \leftrightarrow -t$ and $|in \rangle \leftrightarrow |out \rangle$.
- Assumes CPT invariance and $\Delta\Gamma = 0$.

Foundations of the analysis

- **Ingredients:** M. C. Bañuls and J. Bernabeu, Phys. Lett. B 464, 117 (1999); Nucl. Phys. B 590, 19 (2000).
 - **EPR entanglement** produced by the decay of the $\Upsilon(4S)$.

$$|i\rangle = \frac{1}{\sqrt{2}}[B^0(t_1)\bar{B}^0(t_2) - \bar{B}^0(t_1)B^0(t_2)] \\ = \frac{1}{\sqrt{2}}[B_+(t_1)B_-(t_2) - B_-(t_1)B_+(t_2)]$$



- **Quantum Mechanics.**

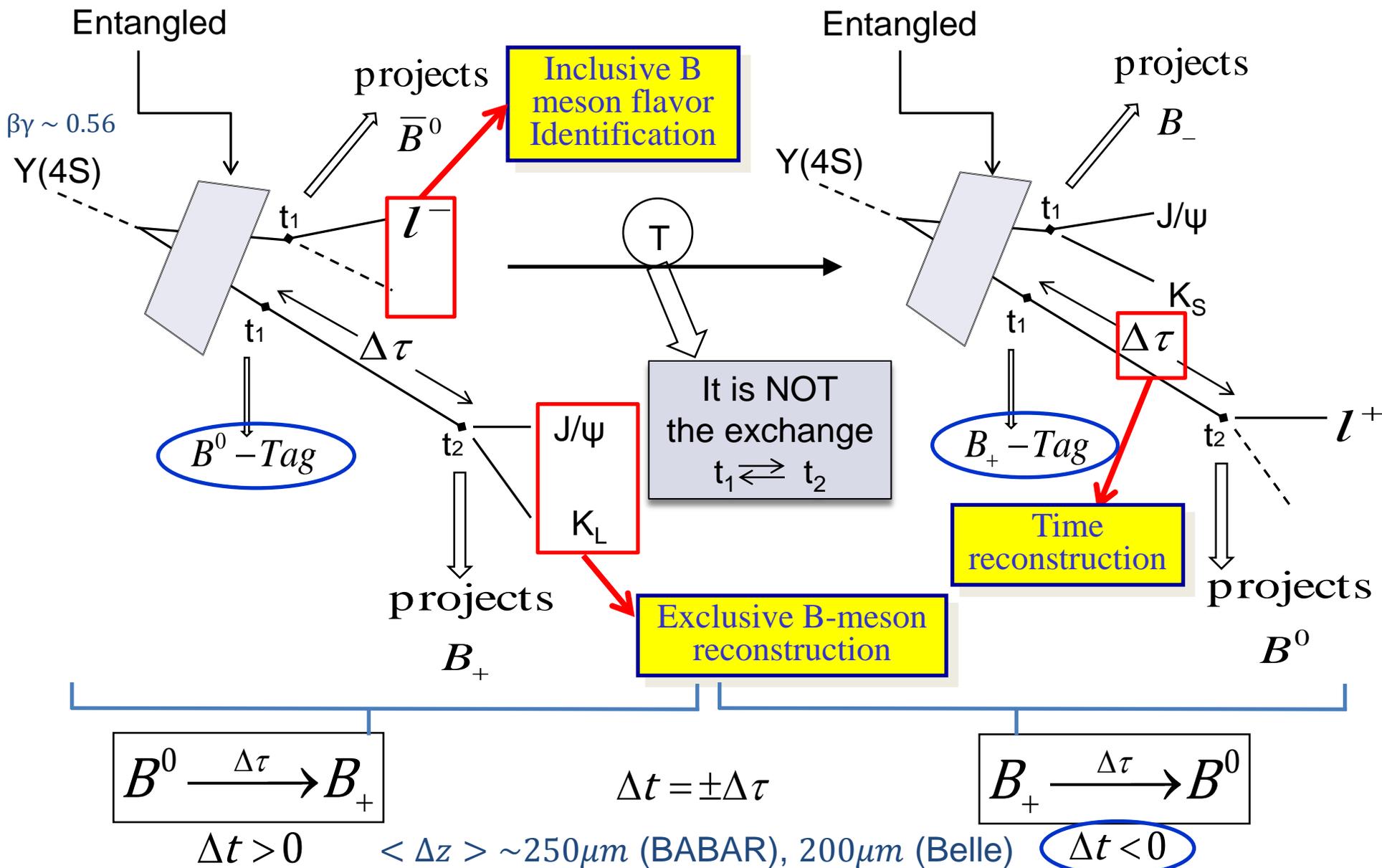
$$\Delta\tau = t_Y - t_X > 0$$

Reference: Physical Process
(X,Y): Reconstructed Final States

Reference (X, Y)	T-Transformed (X, Y)
$B^0 \rightarrow B_+ \quad (l^-, J/\psi K_L)$	$B_+ \rightarrow B^0 \quad (J/\psi K_S, l^+)$
$B^0 \rightarrow B_- \quad (l^-, J/\psi K_S)$	$B_- \rightarrow B^0 \quad (J/\psi K_L, l^+)$
$\bar{B}^0 \rightarrow B_+ \quad (l^+, J/\psi K_L)$	$B_+ \rightarrow \bar{B}^0 \quad (J/\psi K_S, l^-)$
$\bar{B}^0 \rightarrow B_- \quad (l^+, J/\psi K_S)$	$B_- \rightarrow \bar{B}^0 \quad (J/\psi K_L, l^-)$

l^+ and l^- project over the B flavor, i.e., B^0 and \bar{B}^0 respectively

Foundations of the analysis



T-transformed processes

JHEP08 (2012) 064

Reference (X,Y)	T -Transformed
$B^0 \rightarrow B_+$ ($\ell^-, J/\psi K_L^0$)	$B_+ \rightarrow B^0$ ($J/\psi K_S^0, \ell^+$)
$B^0 \rightarrow B_-$ ($\ell^-, J/\psi K_S^0$)	$B_- \rightarrow B^0$ ($J/\psi K_L^0, \ell^+$)
$\bar{B}^0 \rightarrow B_+$ ($\ell^+, J/\psi K_L^0$)	$B_+ \rightarrow \bar{B}^0$ ($J/\psi K_S^0, \ell^-$)
$\bar{B}^0 \rightarrow B_-$ ($\ell^+, J/\psi K_S^0$)	$B_- \rightarrow \bar{B}^0$ ($J/\psi K_L^0, \ell^-$)

(X,Y) is the reconstructed final states (tag, reco.)

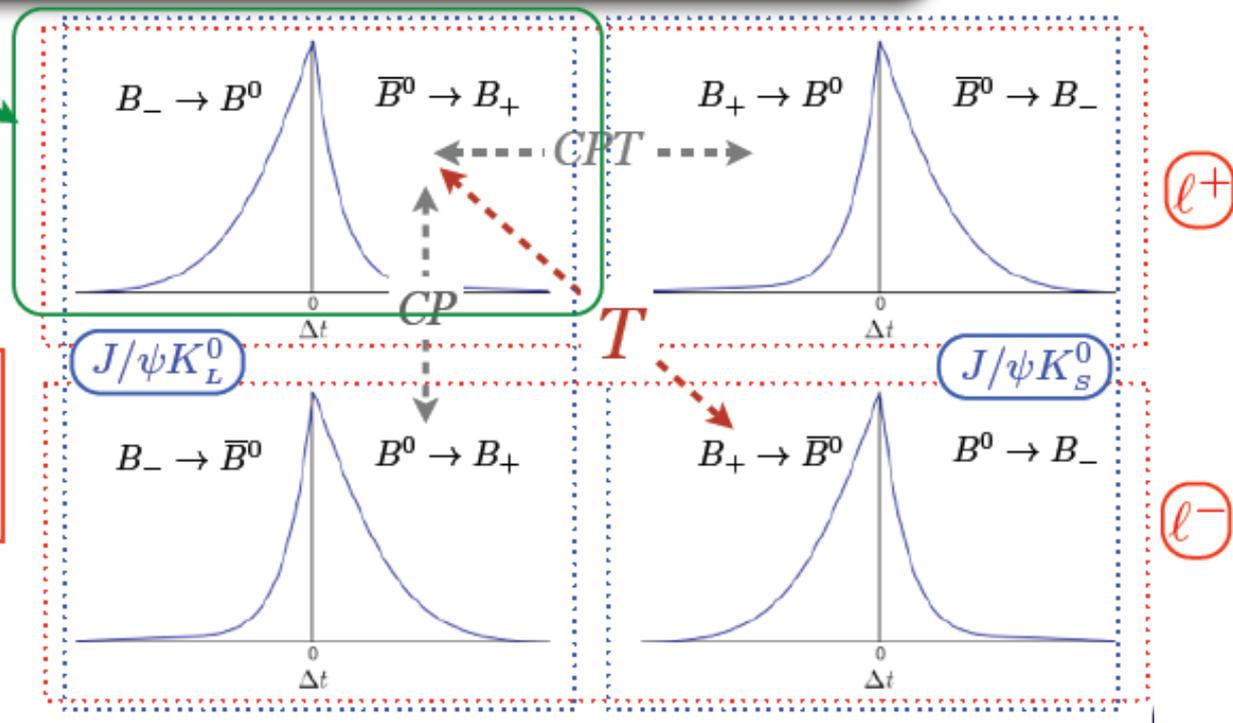
...and similar for CP, CPT

In total we can build:

- 4 independent T comparisons
- 4 independent CP comparisons
- 4 independent CPT comparisons

T implies comparison of:

- 1) Opposite Δt sign
- 2) Different reco states (ψK_S v. ψK_L)
- 3) Opposite flavor states (B^0 v. \bar{B}^0)



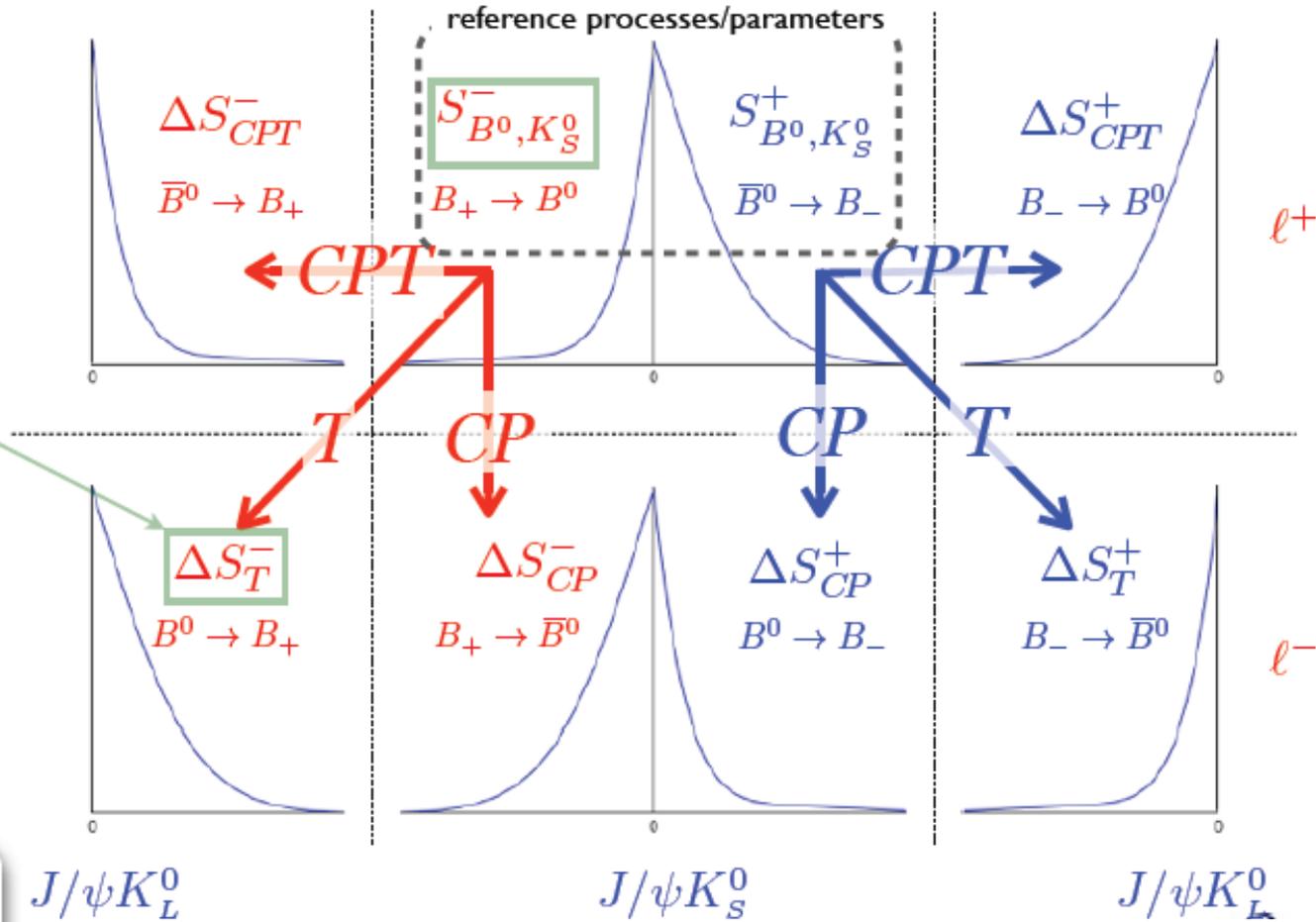
Signal parameters

8 Signal PDFs: $g_{\alpha,\beta}^{\pm}(\Delta\tau) \propto e^{-\Gamma\Delta\tau} \{1 + S_{\alpha,\beta}^{\pm} \sin(\Delta m_d \Delta\tau) + C_{\alpha,\beta}^{\pm} \cos(\Delta m_d \Delta\tau)\}$

$$\Delta t = t_{CP} - t_{flav} = \begin{cases} +\Delta\tau & \text{for "flavor tag"} \\ -\Delta\tau & \text{for "CP tag"} \end{cases} \quad \alpha \in \{B^0, \bar{B}^0\}; \quad \beta \in \{K_S^0, K_L^0\}$$

Prediction from CPV

Parameter	Value
$S_{B^0, K_S^0}^+$	0.7
$\Delta S_T^+ = S_{B^0, K_L^0}^- - S_{B^0, K_S^0}^+$	-1.4
$\Delta S_{CP}^+ = S_{B^0, K_S^0}^- - S_{B^0, K_S^0}^+$	-1.4
$\Delta S_{CPT}^+ = S_{B^0, K_L^0}^- - S_{B^0, K_S^0}^+$	0.0
$S_{B^0, K_S^0}^-$	-0.7
$\Delta S_T^- = S_{B^0, K_L^0}^+ - S_{B^0, K_S^0}^-$	1.4
$\Delta S_{CP}^- = S_{B^0, K_S^0}^+ - S_{B^0, K_S^0}^-$	1.4
$\Delta S_{CPT}^- = S_{B^0, K_L^0}^+ - S_{B^0, K_S^0}^-$	0.0
$C_{B^0, K_S^0}^+$	0.0
$\Delta C_T^+ = C_{B^0, K_L^0}^- - C_{B^0, K_S^0}^+$	0.0
$\Delta C_{CP}^+ = C_{B^0, K_S^0}^- - C_{B^0, K_S^0}^+$	0.0
$\Delta C_{CPT}^+ = C_{B^0, K_L^0}^- - C_{B^0, K_S^0}^+$	0.0
$C_{B^0, K_S^0}^-$	0.0
$\Delta C_T^- = C_{B^0, K_L^0}^+ - C_{B^0, K_S^0}^-$	0.0
$\Delta C_{CP}^- = C_{B^0, K_S^0}^+ - C_{B^0, K_S^0}^-$	0.0
$\Delta C_{CPT}^- = C_{B^0, K_L^0}^+ - C_{B^0, K_S^0}^-$	0.0



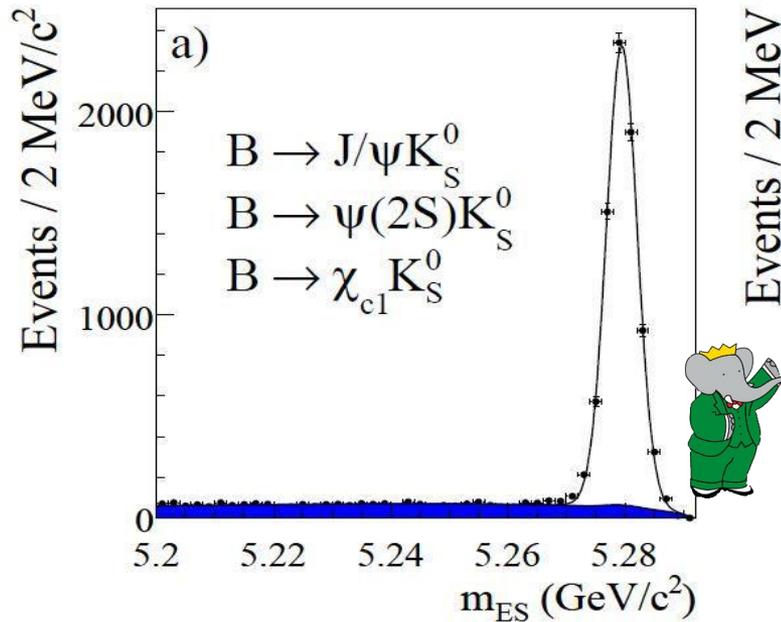
For T Violation
in the interference $\Delta S_T^+ \neq 0, \Delta S_T^- \neq 0$
in the decay $\Delta C_T^+ \neq 0, \Delta C_T^- \neq 0$

Dataset

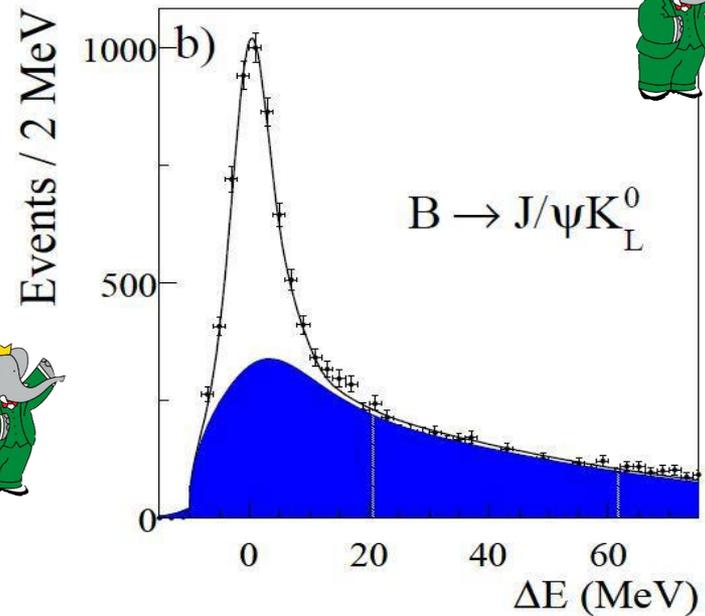
signal sample

Category	Decay(s)
$c\bar{c}K_S^0$	$B^0 \rightarrow J/\psi K_S^0$
	$B^0 \rightarrow \psi(2S)K_S^0$
	$B^0 \rightarrow \chi_{c1}K_S^0$
$c\bar{c}K_L^0$	$B^0 \rightarrow J/\psi K_L^0$
B_{flav} (high statistics)	$B^0 \rightarrow D^* \pi(\rho, a_1)$
	$B^0 \rightarrow J/\psi K^{*0}$
Control sample $c\bar{c}K^\pm, J/\psi K^{*\pm}$	$B^+ \rightarrow J/\psi K^+$
	$B^+ \rightarrow \psi(2S)K^+$
	$B^+ \rightarrow J/\psi K^{*+}$

$c\bar{c}K_S$ sample



7796 events, purity 87–96%



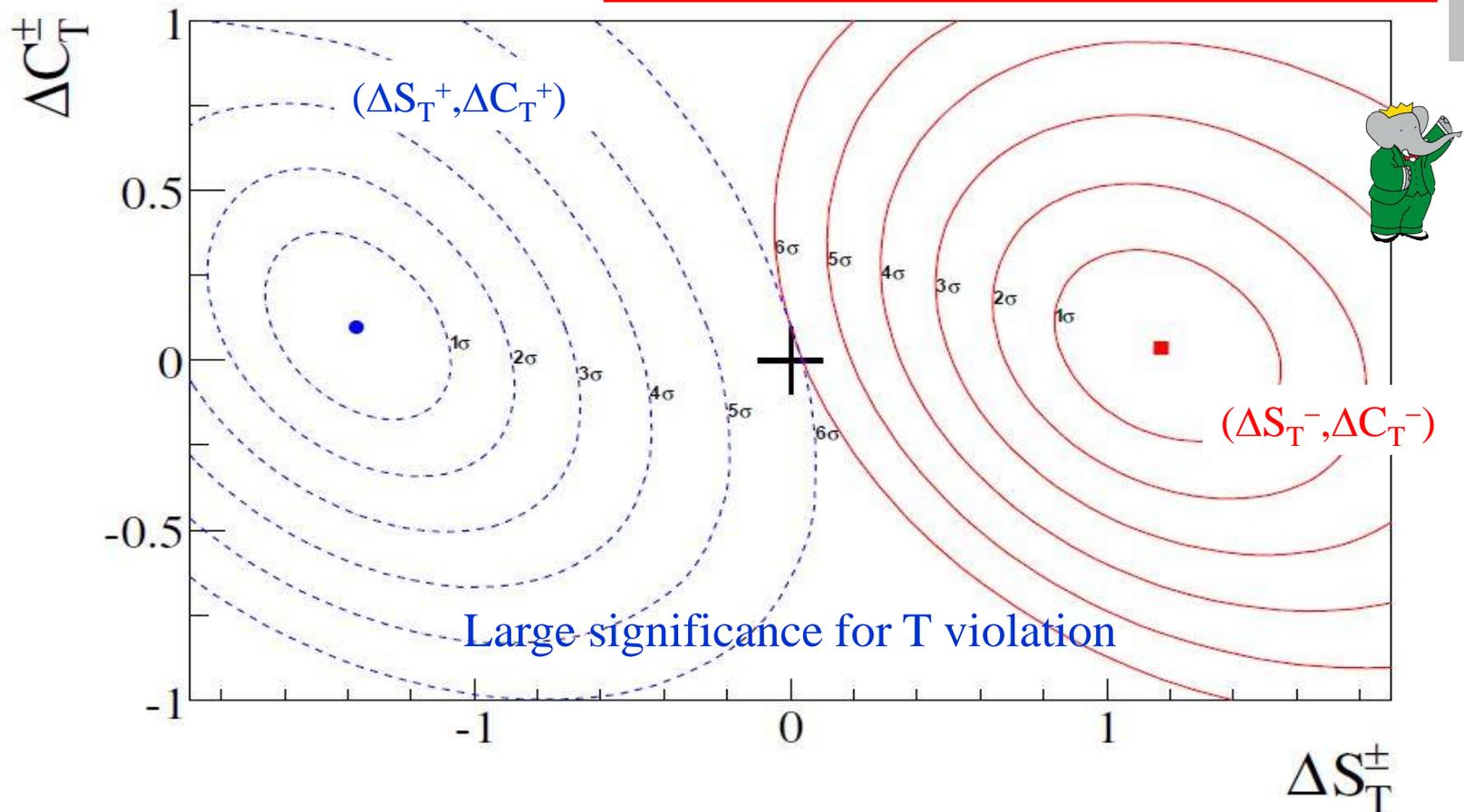
5813 events, purity $\approx 56\%$

$J/\psi K_S$ sample

Results

$\Delta S_T^+ = S_{\ell^-, K_L^0}^- - S_{\ell^+, K_S^0}^+$	$-1.37 \pm 0.14 \pm 0.06$	$-2\sin 2\beta$
$\Delta S_T^- = S_{\ell^-, K_L^0}^+ - S_{\ell^+, K_S^0}^-$	$1.17 \pm 0.18 \pm 0.11$	$+2\sin 2\beta$
$\Delta C_T^+ = C_{\ell^-, K_L^0}^- - C_{\ell^+, K_S^0}^+$	$0.10 \pm 0.14 \pm 0.08$	0
$\Delta C_T^- = C_{\ell^-, K_L^0}^+ - C_{\ell^+, K_S^0}^-$	$0.04 \pm 0.14 \pm 0.08$	0

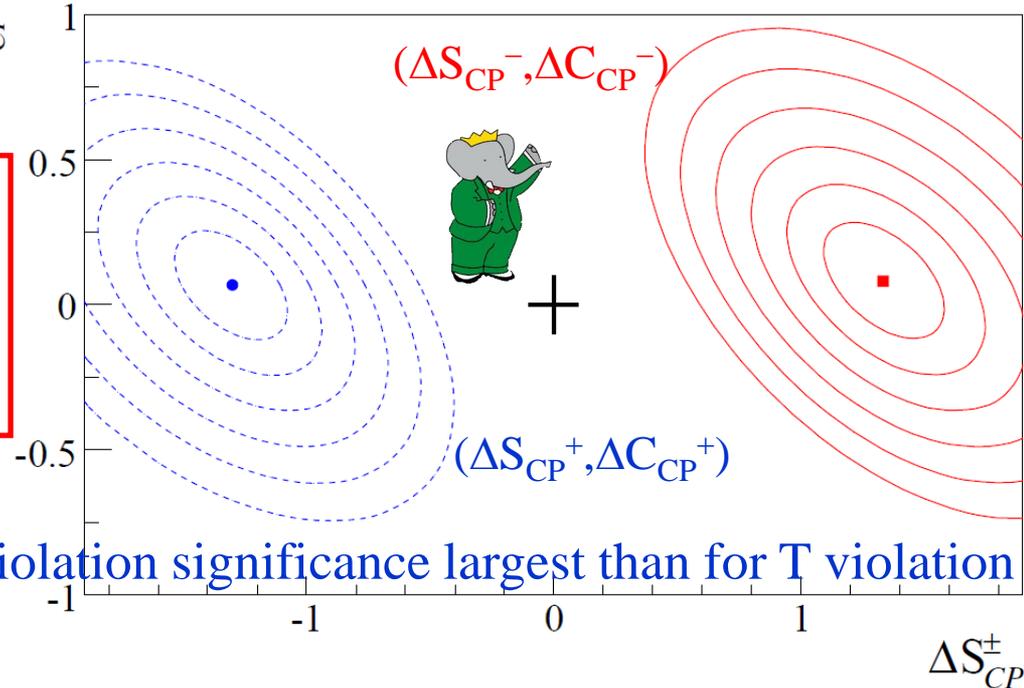
expectation from
canonical CP



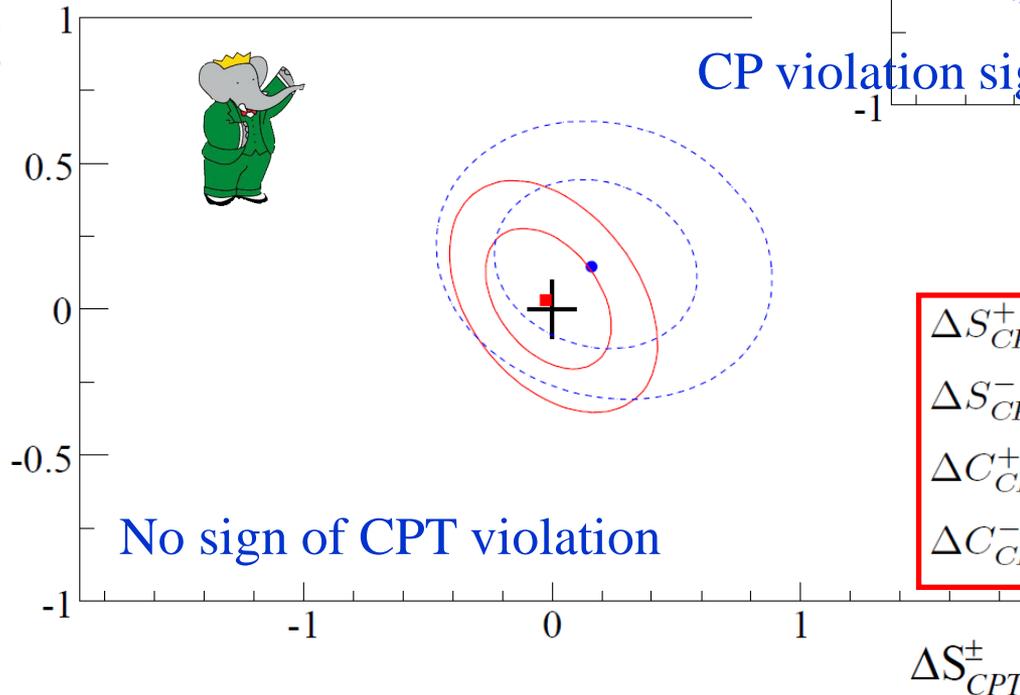
Results

CP-violating parameters

$\Delta S_{CP}^+ = S_{\ell^-, K_S^0}^+ - S_{\ell^+, K_S^0}^+$	$-1.30 \pm 0.11 \pm 0.07$
$\Delta S_{CP}^- = S_{\ell^-, K_S^0}^- - S_{\ell^+, K_S^0}^-$	$1.33 \pm 0.12 \pm 0.06$
$\Delta C_{CP}^+ = C_{\ell^-, K_S^0}^+ - C_{\ell^+, K_S^0}^+$	$0.07 \pm 0.09 \pm 0.03$
$\Delta C_{CP}^- = C_{\ell^-, K_S^0}^- - C_{\ell^+, K_S^0}^-$	$0.08 \pm 0.10 \pm 0.04$

 ΔC_{CP}^\pm


CP violation significance largest than for T violation

 ΔC_{CPT}^\pm


No sign of CPT violation

CPT-violating parameters

$\Delta S_{CPT}^+ = S_{\ell^+, K_L^0}^- - S_{\ell^+, K_S^0}^+$	$0.16 \pm 0.21 \pm 0.09$
$\Delta S_{CPT}^- = S_{\ell^+, K_L^0}^+ - S_{\ell^+, K_S^0}^-$	$-0.03 \pm 0.13 \pm 0.06$
$\Delta C_{CPT}^+ = C_{\ell^+, K_L^0}^- - C_{\ell^+, K_S^0}^+$	$0.14 \pm 0.15 \pm 0.07$
$\Delta C_{CPT}^- = C_{\ell^+, K_L^0}^+ - C_{\ell^+, K_S^0}^-$	$0.03 \pm 0.12 \pm 0.08$

Observed T violation as due to compensate CP violation

Significance of T violation

1. We obtain the likelihood value of the fit to S, C for the 8 independent samples (Standard Fit).
2. We repeat the fit, reassembling the parameters for T-conjugated processes, to forbid T violation.
3. Significance of T violation evaluated from the difference of the likelihood values.
4. Raw asymmetries and fit projections can be now plotted in the standard way.

T invariance
$\Delta S_T^+ = 0$
$\Delta S_T^- = 0$
$\Delta S_{CP}^+ = \Delta S_{CPT}^+$
$\Delta S_{CP}^- = \Delta S_{CPT}^-$
$\Delta C_T^+ = 0$
$\Delta C_T^- = 0$
$\Delta C_{CP}^+ = \Delta C_{CPT}^+$
$\Delta C_{CP}^- = \Delta C_{CPT}^-$

$$\Delta\chi^2 = -2(\ln L_{No_T_Violation} - \ln L)$$

$$\Delta\nu = 8$$

5. CP, and CPT significance is evaluated similarly.
6. Using Gaussian approximation, we evaluate the change of likelihood in 1σ systematic variation.

$$m_j^2 = -2[\ln L(q_j, o_j) - \ln L(p_0)] / s_{stat,j}^2$$

7. We take the $\max\{m_j^2\}$ and we divide our significance (s^2) by $(1 + \max\{m_j^2\})$

Asymmetries

- We build an asymmetry for these four reference transitions:

$$\bar{B}^0 \rightarrow B_- \quad \bar{B}^0 \rightarrow B_+ \quad B_+ \rightarrow B^0 \quad B_- \rightarrow B^0$$

- For the first reference:

$$A_T(\Delta t) \equiv \frac{H_{l^-,K_L}^-(\Delta t) - H_{K_S,l^+}^+(\Delta t)}{H_{l^-,K_L}^-(\Delta t) + H_{K_S,l^+}^+(\Delta t)} \approx \frac{\Delta C_T^+}{2} \cos \Delta m \Delta t + \frac{\Delta S_T^+}{2} \sin \Delta m \Delta t$$

Assuming no experimental effects

$H_{\alpha,\beta}^\pm(\Delta t)$ is the intensity for each sample with $\Delta t > 0$ and all experimental effects.

T raw asymmetries + significance

Significance test

$$s_{NoT}^2 = 226$$

$$14\sigma$$

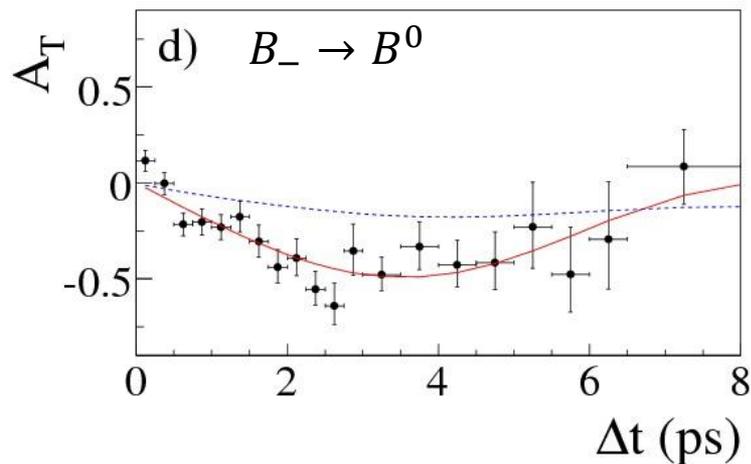
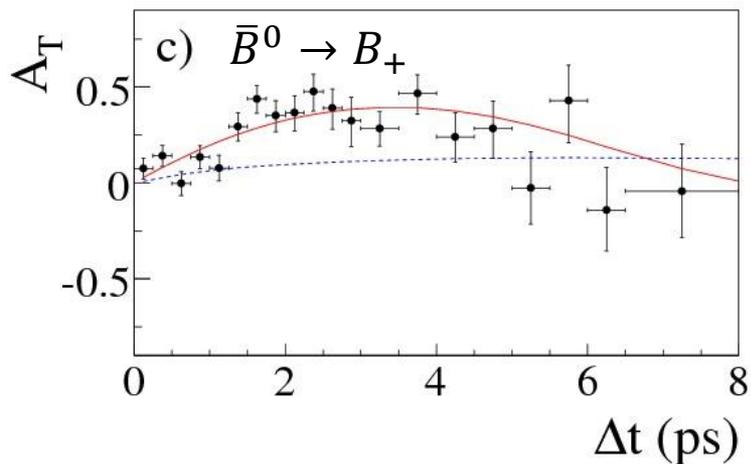
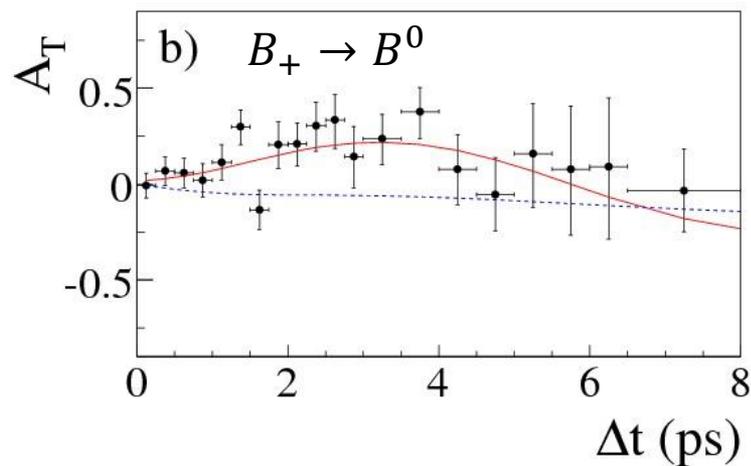
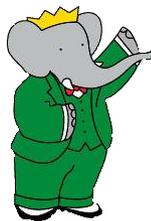
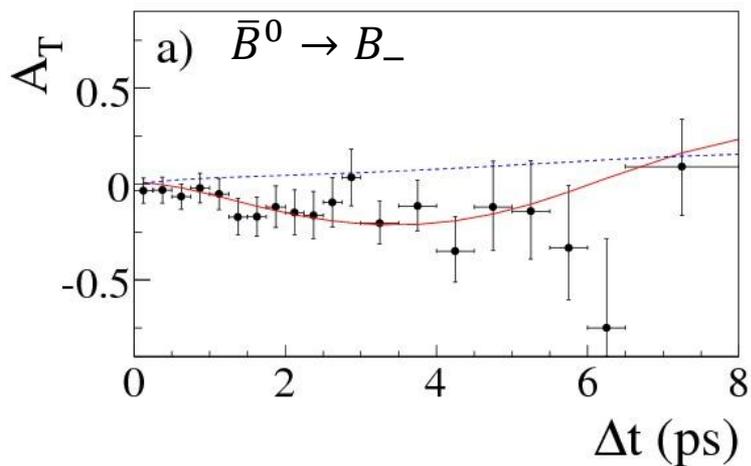
$$s_{NoCP}^2 = 307$$

$$16.6\sigma$$

$$s_{NoCPT}^2 = 5$$

$$0.33\sigma$$

Stat. and
Syst.
 $\Delta v=8$



Conclusions

- We have measured T-violating parameters in the time evolution of neutral-B mesons.
- These parameters have been measured:
 - Directly: without exp. connection to CP and CPT.
 - Genuinely: exchanging *in-states* and *out-states*.
- We observe a large deviation of T invariance at 14σ level.
- Our result is consistent with CP-violating measurements assuming CPT invariance.
- **This constitutes the first direct observation of Time Reversal Violation in the evolution of the B mesons.**

$$\underline{B \rightarrow D^{(*)} \tau \nu}$$

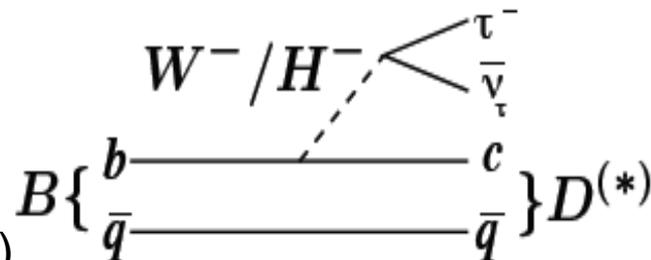
J. P. Lees et al. Phys. Rev. Lett. 109, 101802 (2012)

Motivation

- Observables: $R(D)$ and $R(D^*)$ ratios

$$R(D^{(*)}) = \frac{\mathcal{B}(B \rightarrow D^{(*)}\tau\nu_\tau)}{\mathcal{B}(B \rightarrow D^{(*)}l\nu_\tau)}$$

- can be enhanced by the charged-Higgs ($\tan\beta/m_H$)
- several syst. and theo. uncertainties cancel out

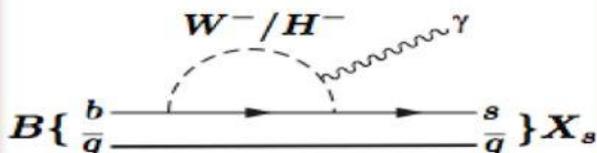


- Involve form factors which can be measured in $B \rightarrow D^{(*)}e/\mu\nu$ decays

- Sensitive to charged-Higgs effects. τ hadronic decays can give interesting information on charged Higgs couplings (Phys. Rev. D78:015006, 2008)

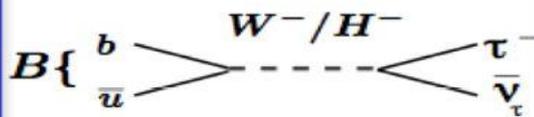
$B \rightarrow X_s \gamma$

- Small $\sigma_{SM} \sim 7\%$
- H^- enters in a loop
- BF $\sim 0.03\%$
- Inclusive measurement difficult



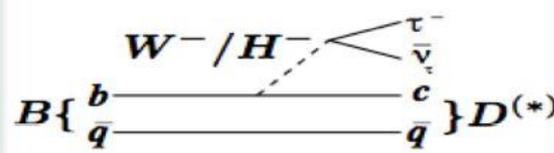
$B^- \rightarrow \tau^- \bar{\nu}_\tau$

- Large $\sigma_{SM} \sim 25\%$
- H^- enters at tree level
- BF $\sim 0.01\%$
- Helicity suppressed

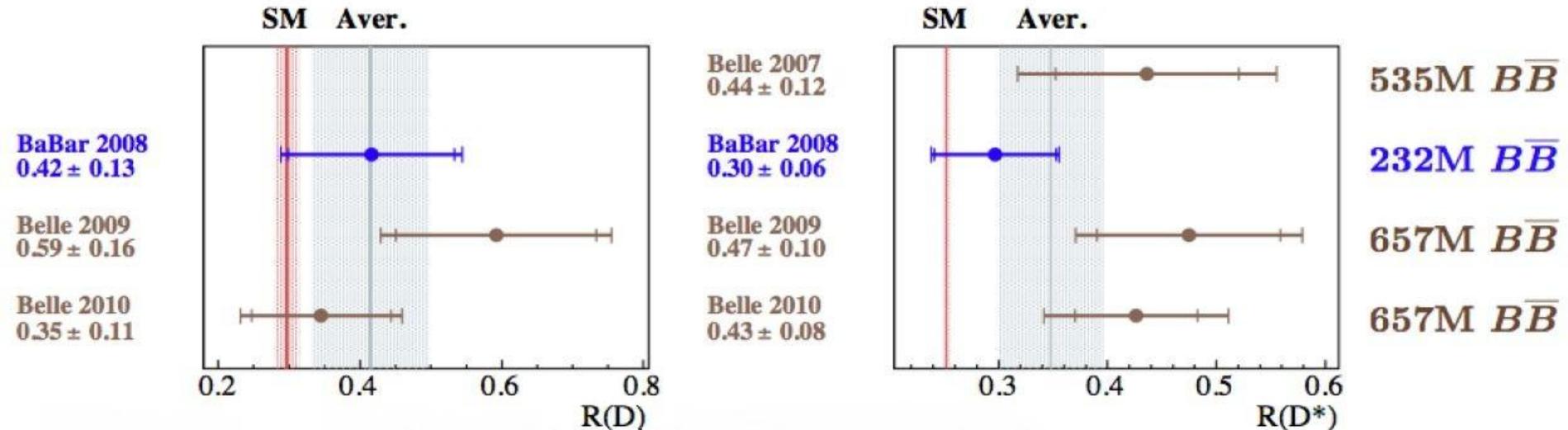


$\bar{B} \rightarrow D^{(*)}\tau^- \bar{\nu}_\tau$

- Small $\sigma_{SM} \sim 2-5\%$
- H^- enters at tree level
- BF $\sim 1-2\%$
- $D^{(*)}$ provides constraint



Previous situation



- Consistent slightly excess over the SM prediction.
- Large uncertainties.
- Update the BaBar 2008 measurement to the full dataset

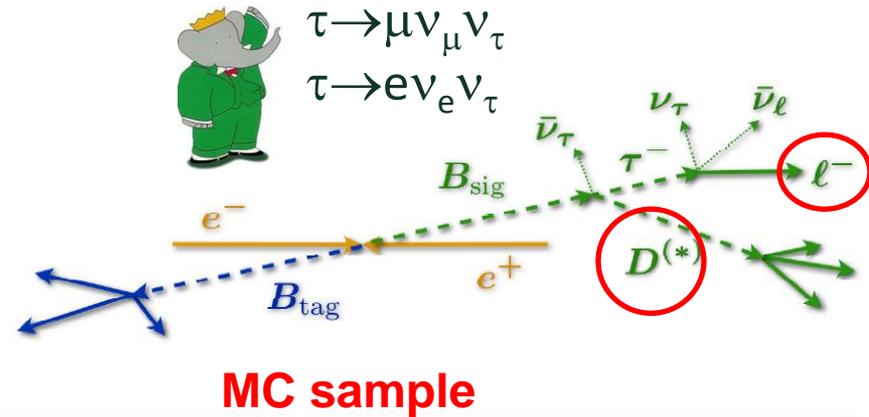
SM predictions:

$$R(D)_{\text{SM}} = 0.297 \pm 0.017$$
$$R(D^*)_{\text{SM}} = 0.252 \pm 0.003$$

[S. Fajfer, J. F. Kamenik, and I. Nisandzic, [arXiv:1203.2654](https://arxiv.org/abs/1203.2654);
J. F. Kamenik and F. Mescia, *Phys. Rev. D* **78**, 014003, (2008)]

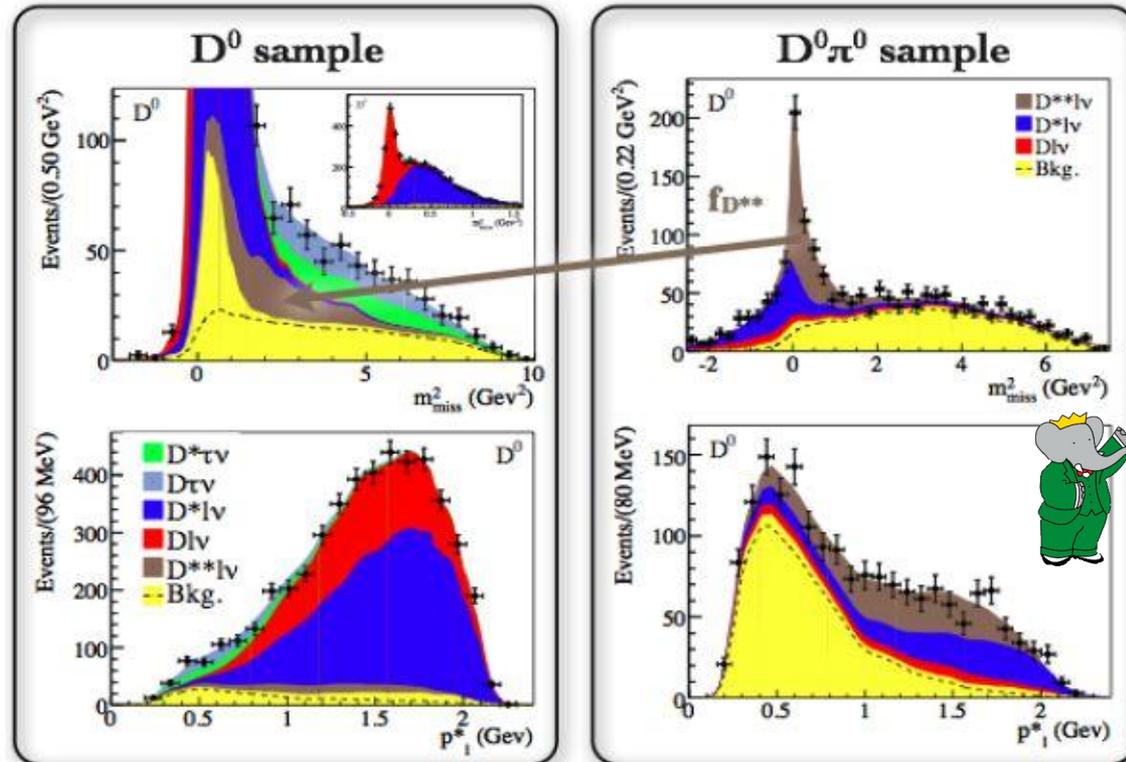
Experimental method

- BaBar: Btag fully reconstructed into hadrons
(Improved efficiencies (lepton and Btag))
Efficiency: 0.21% → 0.40% **2x higher**
- Bsig: $D^{(*)}$ and lepton (μ, e)
 - 4 signal samples: $(D^0, D^+, D^{*0}, D^{*+})\ell\nu$
(to extract $B \rightarrow D^{(*)}\tau\nu$)
 - 4 control samples: $(D^0, D^+, D^{*0}, D^{*+})\pi^0\ell\nu$
(to derive $D^{**}\ell\nu$ bkg)



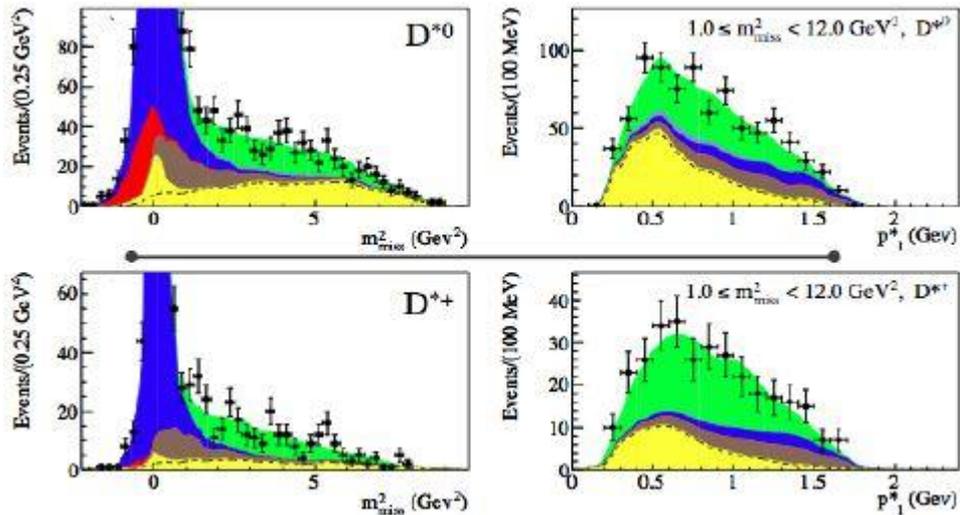
- 2D unbinned ML fit $m_{\text{miss}}^2 - p_\ell^*$
- Yields for:
- $B \rightarrow (D^0, D^+, D^{*0}, D^{*+})\tau\nu$
 - $B \rightarrow (D^0, D^+, D^{*0}, D^{*+})\ell\nu$
 - $B \rightarrow (D^0, D^+, D^{*0}, D^{*+})\pi^0\ell\nu$

$$m_{\text{miss}}^2 = (p_{e+e-} - p_{\text{Btag}} - p_{D^{(*)}} - p_\ell)^2$$

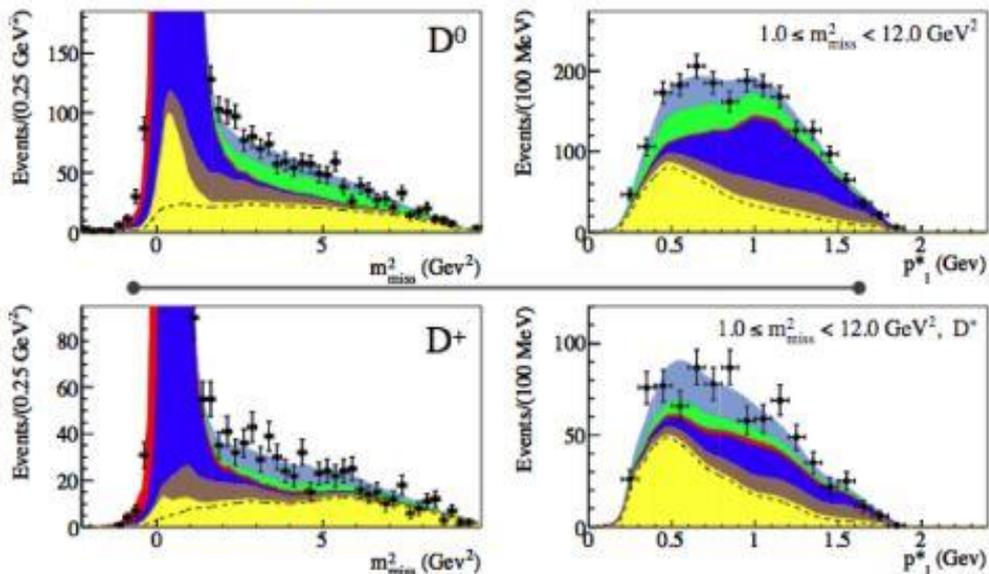


Fit results

D^*

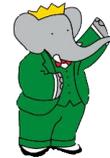


D



Statistical only

	$D^{*0}\tau\nu$	$D^{*+}\tau\nu$	$D^*\tau\nu$
N_{sig}	639±62	245±27	888±63
Sig.	11.3 σ	11.6 σ	16.4 σ
R(D^*)	0.322±0.032	0.355±0.039	0.332±0.024



First observation

	$D^0\tau\nu$	$D^+\tau\nu$	$D\tau\nu$
N_{sig}	314±60	177±31	489±63
Sig.	5.5 σ	6.1 σ	8.4 σ
R(D)	0.429±0.082	0.469±0.084	0.440±0.058

Fit results

Decay	N_{sig}	N_{norm}	$\mathcal{R}(D^{(*)})$	$\mathcal{B}(B \rightarrow D^{(*)}\tau\nu)$ (%)	Σ_{tot}
$B^- \rightarrow D^0\tau^-\bar{\nu}_\tau$	314 ± 60	1995 ± 55	$0.429 \pm 0.082 \pm 0.052$	$0.99 \pm 0.19 \pm 0.13$	4.7
$B^- \rightarrow D^{*0}\tau^-\bar{\nu}_\tau$	639 ± 62	8766 ± 104	$0.322 \pm 0.032 \pm 0.021$	$1.71 \pm 0.17 \pm 0.13$	9.6
$\bar{B}^0 \rightarrow D^+\tau^-\bar{\nu}_\tau$	177 ± 31	986 ± 35	$0.469 \pm 0.084 \pm 0.052$	$1.01 \pm 0.18 \pm 0.12$	5.3
$\bar{B}^0 \rightarrow D^{*+}\tau^-\bar{\nu}_\tau$	245 ± 27	3186 ± 61	$0.355 \pm 0.039 \pm 0.020$	$1.74 \pm 0.19 \pm 0.12$	10.5
$\bar{B} \rightarrow D\tau^-\bar{\nu}_\tau$	489 ± 63	2981 ± 65	$0.440 \pm 0.058 \pm 0.042$	$1.02 \pm 0.13 \pm 0.10$	6.8
$\bar{B} \rightarrow D^*\tau^-\bar{\nu}_\tau$	888 ± 63	11953 ± 122	$0.332 \pm 0.024 \pm 0.017$	$1.76 \pm 0.13 \pm 0.11$	13.4

• Comparison with the SM:

$$R(D) = 0.440 \pm 0.071 \quad \boxed{2.0\sigma}$$

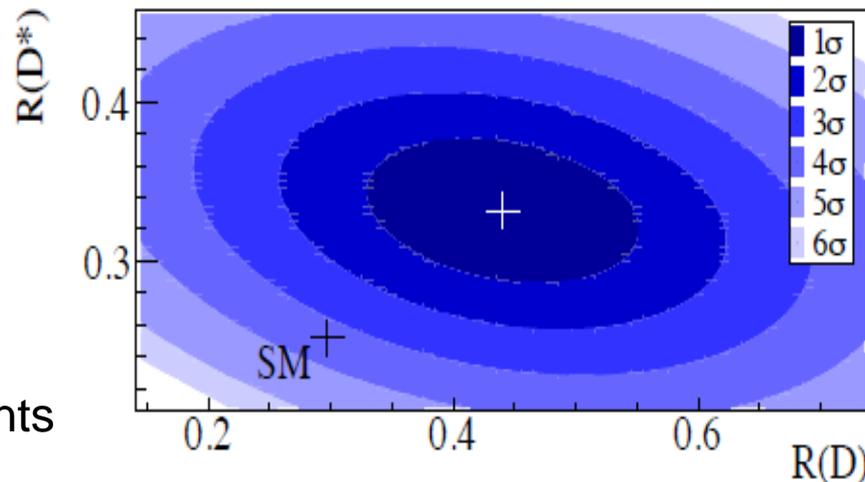
$$\text{SM} = 0.297(17)$$

$$R(D^*) = 0.332 \pm 0.029 \quad \boxed{2.7\sigma}$$

$$\text{SM} = 0.252(3)$$

The combination of the two measurements (-0.27 correlation) yields $\chi^2/\text{NDF}=14.6/2$ (i.e. Prob. = 6.9×10^{-4}) \rightarrow

The SM prediction is excluded at 3.4σ



2HDM Scan

- Interpretation Beyond the Standard Model:

A charged Higgs (2HDM type II) could enhance or decrease the $R(D)$ and $R(D^*)$ ratios depending on $\tan\beta/m_H$

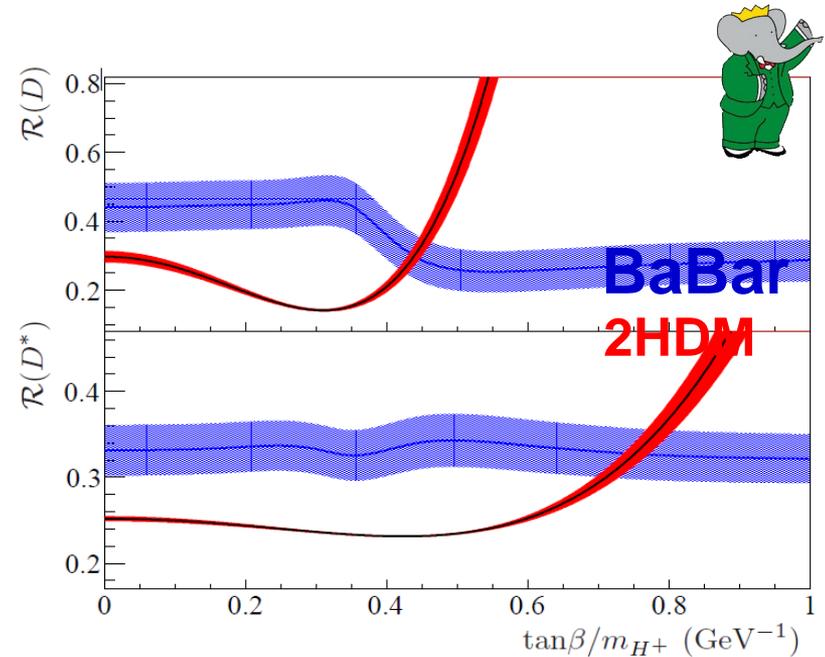
$$R(D^{(*)}) = R(D^{(*)})_{\text{SM}} + A^{(*)} \frac{\tan^2 \beta}{m_{H^+}^2} + B^{(*)} \frac{\tan^4 \beta}{m_{H^+}^4}$$

Effect of 2DHM (accounting for difference in efficiency):

$\tan\beta/m_H = 0.44 \pm 0.02$ for $R(D)$
 $\tan\beta/m_H = 0.75 \pm 0.04$ for $R(D^*)$



The combination of $R(D)$ and $R(D^*)$ excludes the Type II 2HDM in the full $\tan\beta$ - m_H parameter space with a probability of $>99.8\%$ ($M_H > 10\text{GeV}$)



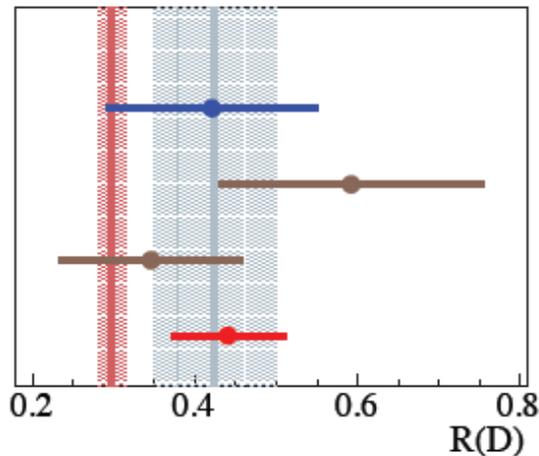
Conclusions

This is the **first** measurement of $D^{(*)} \rightarrow \tau\nu$ with the **full BaBar data set**

➤ **Larger improvement than \sqrt{L}**

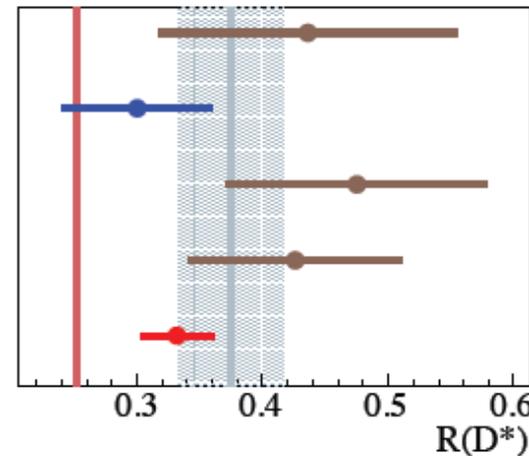
- Comparison with previous measurements:

SM Average
(w/o BaBar 2012)



SM Average
(w/o BaBar 2012)

Belle 2007
 0.44 ± 0.12
BaBar 2008
 0.30 ± 0.06
Belle 2009
 0.47 ± 0.10
Belle 2010
 0.43 ± 0.09
BaBar 2012
 0.332 ± 0.029



535M $B\bar{B}$
232M $B\bar{B}$
657M $B\bar{B}$
657M $B\bar{B}$
471M $B\bar{B}$

Scorecard:

- **First observation** of $D \rightarrow \tau\nu$ with 6.8σ
- 3.4σ **excess** over the **SM** prediction
- **Not compatible** with Type II **2HDM** model

General conclusions

- Despite BaBar stopped taking data at 2008, it is producing publications.
- Not only because of repeating the old analysis with the whole data sample.
- Great improvement in the analysis tools and the machine performance.
- New ideas are coming to analyze the data.
- Collaboration with theorist
 - M. C. Bañuls and J. Bernabeu, Phys. Lett. B 464, 117 (1999); Nucl. Phys. B 590, 19 (2000).