#### Physics at a Future Circular Collider

Jornadas (virtuales) de la red española de futuros colisionadores

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### Additional material: recent reports on Future Circular Colliders

#### • FCC CDR:

- Vol.1: Physics Opportunities (CERN-ACC-2018-0056) <a href="http://cern.ch/go/Nqx7">http://cern.ch/go/Nqx7</a>
- Vol.2: The Lepton Machine (CERN-ACC-2018-0057) <a href="http://cern.ch/go/7DH9">http://cern.ch/go/7DH9</a>
- Vol.3: The Hadron Machine (CERN-ACC-2018-0058), http://cern.ch/go/Xrg6
- Vol.4: High-Energy LHC (CERN-ACC-2018-0059) <a href="http://cern.ch/go/S9Gq">http://cern.ch/go/S9Gq</a>
- "Physics at 100 TeV", CERN Yellow Report: <a href="https://arxiv.org/abs/1710.06353">https://arxiv.org/abs/1710.06353</a>
- CEPC CDR: Physics and Detectors

#### The important questions

#### Data driven:

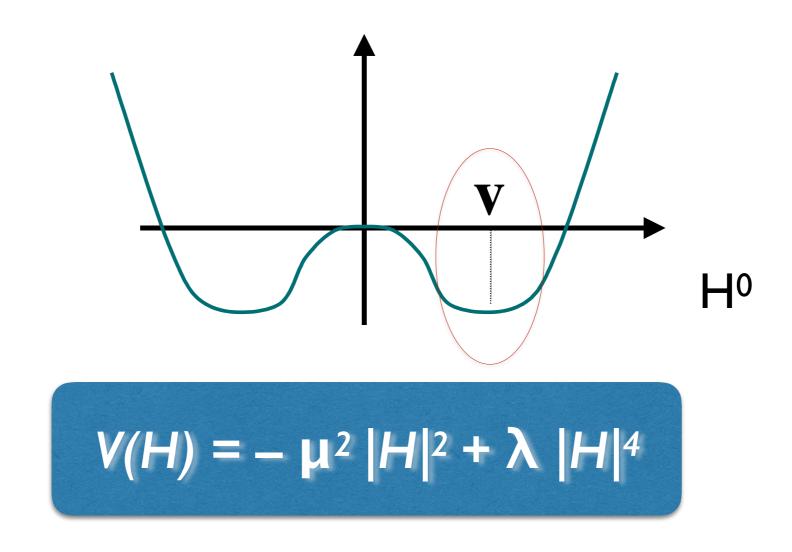
- DM
- Neutrino masses
- Matter vs antimatter asymmetry
- Dark energy
- ...

#### Theory driven:

- The hierarchy problem and naturalness
- The flavour problem (origin of fermion families, mass/mixing pattern)
- Quantum gravity
- Origin of inflation
- ...

We have no guarantees as to where answers to these questions will come from, and what are the experiments that will eventually answer them.

But there is one question that can only be addressed by colliders, and future collider efforts must focus on its thorough exploration



#### Where does this come from?

## Other important open issues on the Higgs sector

- Is the Higgs the only (fundamental?) scalar field, or are there other Higgs-like states (e.g. H<sup>±</sup>, A<sup>0</sup>, H<sup>±±</sup>, ..., EW-singlets, ....) ?
  - Do all SM families get their mass from the **same** Higgs field?
  - Do  $I_3=1/2$  fermions (up-type quarks) get their mass from the <u>same</u> Higgs field as  $I_3=-1/2$  fermions (down-type quarks and charged leptons)?
  - Do Higgs couplings conserve flavour? H→μτ? H→eτ? t→Hc?
- Is there a deep reason for the apparent metastability of the Higgs vacuum?
- Is there a relation among Higgs/EWSB, baryogenesis, Dark Matter, inflation?
- What happens at the EW phase transition (PT) during the Big Bang?
  - what's the order of the phase transition?
  - are the conditions realized to allow EW baryogenesis?

The importance of the in-depth exploration of the Higgs properties was acknowledged by the 2020 update of the European Strategy for Particle Physics:

"An electron-positron Higgs factory is the highest-priority next collider"

# Key question for the future developments of HEP: Why don't we see the new physics we expected to be present around the TeV scale?

- Is the mass scale beyond the LHC reach?
- Is the mass scale within LHC's reach, but final states are elusive to the direct search?

These two scenarios are a priori equally likely, but impact in different ways the assessment of the physics potential of possible future facilities

Readiness to address both scenarios is the best hedge for the field:

- precision
- sensitivity (to elusive signatures)
- extended energy/mass reach

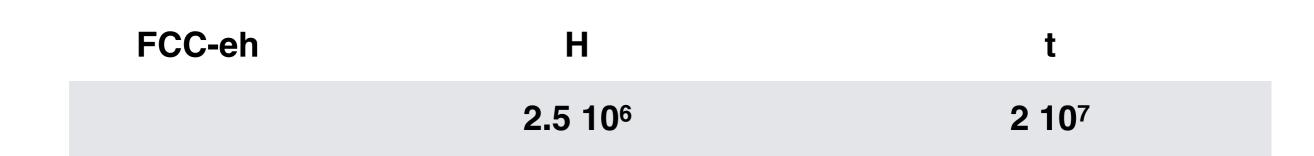
#### What the future circular collider can offer

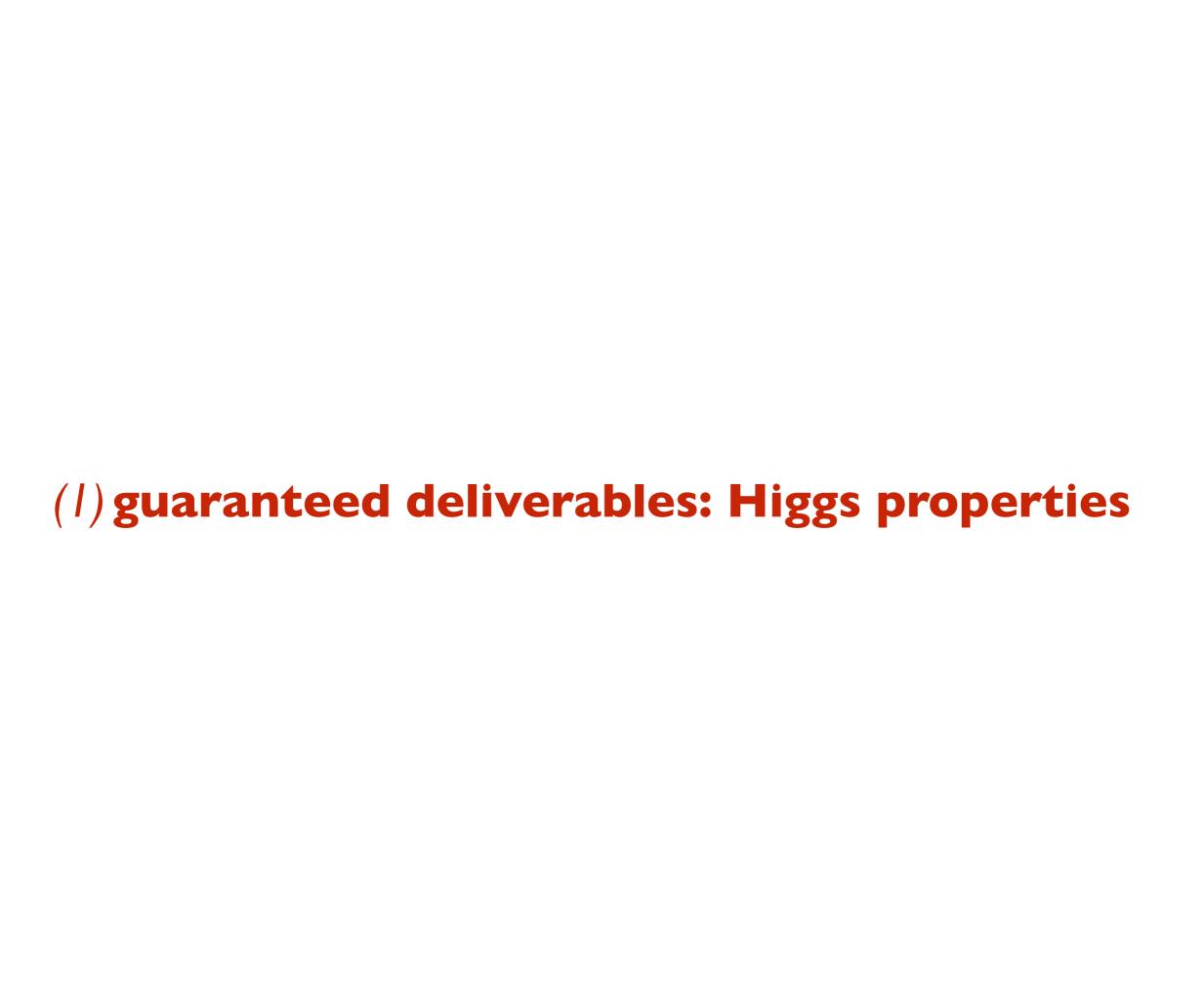
- Guaranteed deliverables:
  - study of Higgs and top quark properties, and exploration of EWSB phenomena, with the best possible precision and sensitivity
- Exploration potential:
  - exploit both direct (large Q2) and indirect (precision) probes
  - enhanced mass reach for direct exploration at 100 TeV
    - E.g. match the mass scales for new physics that could be exposed via indirect precision measurements in the EW and Higgs sector
- Provide firm Yes/No answers to questions like:
  - is there a TeV-scale solution to the hierarchy problem?
  - is DM a thermal WIMP?
  - could the cosmological EW phase transition have been 1st order?
  - could baryogenesis have taken place during the EW phase transition?
  - could neutrino masses have their origin at the TeV scale?
  - ...

#### Event rates: examples

FCC-ee	Н	Z	W	t	<b>τ(←Z)</b>	b( <b>←</b> Z)	c(←Z)
	10 <sup>6</sup>	5 10 <sup>12</sup>	108	<b>10</b> <sup>6</sup>	3 1011	1.5 10 <sup>12</sup>	<b>10</b> <sup>12</sup>

FCC-hh	Н	b	t	W(←t)	<b>τ(←W←t)</b>
	2.5 10 <sup>10</sup>	<b>10</b> <sup>17</sup>	<b>10</b> <sup>12</sup>	<b>10</b> <sup>12</sup>	1011





# Sensitivity of various Higgs couplings to <u>examples</u> of beyond-the-SM phenomena

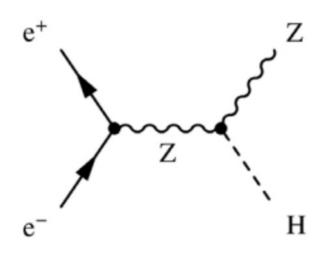
arXiv:1310.8361

Model	$\kappa_V$	$\kappa_b$	$\kappa_{\gamma}$
Singlet Mixing	$\sim 6\%$	$\sim 6\%$	$\sim 6\%$
2HDM	$\sim 1\%$	$\sim 10\%$	$\sim 1\%$
Decoupling MSSM	$\sim -0.0013\%$	$\sim 1.6\%$	$\sim4\%$
Composite	$\sim -3\%$	$\sim -(3-9)\%$	$\sim -9\%$
Top Partner	$\sim -2\%$	$\sim -2\%$	$\sim +1\%$

=> for evidence of 3σ deviations from SM, the precision goal should be (sub)percent!

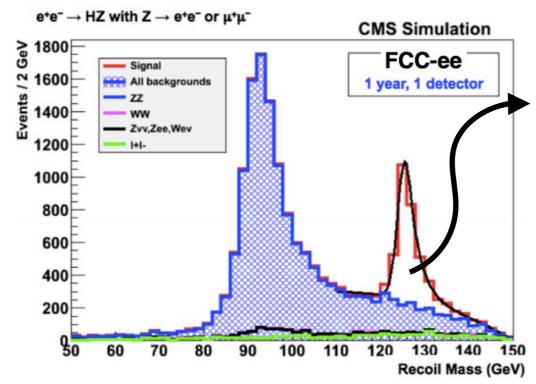
#### The absolutely unique power of $e^+e^- \rightarrow ZH$ (circular or linear):

- the model independent % measurement of  $\Gamma(H)$ , which allows the subsequent:
  - sub-% measurement of couplings to W, Z, b, T
  - % measurement of couplings to gluon and charm



$$p(H) = p(e^-e^+) - p(Z)$$
  
=> [  $p(e^-e^+) - p(Z)$  ]<sup>2</sup> peaks at m<sup>2</sup>(H)

reconstruct Higgs events independently of the Higgs decay mode!



$$N(ZH) \propto \sigma(ZH) \propto g_{HZZ}^2$$

$$N(ZH[\rightarrow ZZ]) \propto$$
 $\sigma(ZH) \times BR(H\rightarrow ZZ) \propto$ 
 $g_{HZZ^2} \times g_{HZZ^2} / \Gamma(H)$ 

=> absolute measurement of width and couplings

$$m_{recoil} = \sqrt{[p(e^-e^+) - p(Z)]^2}$$

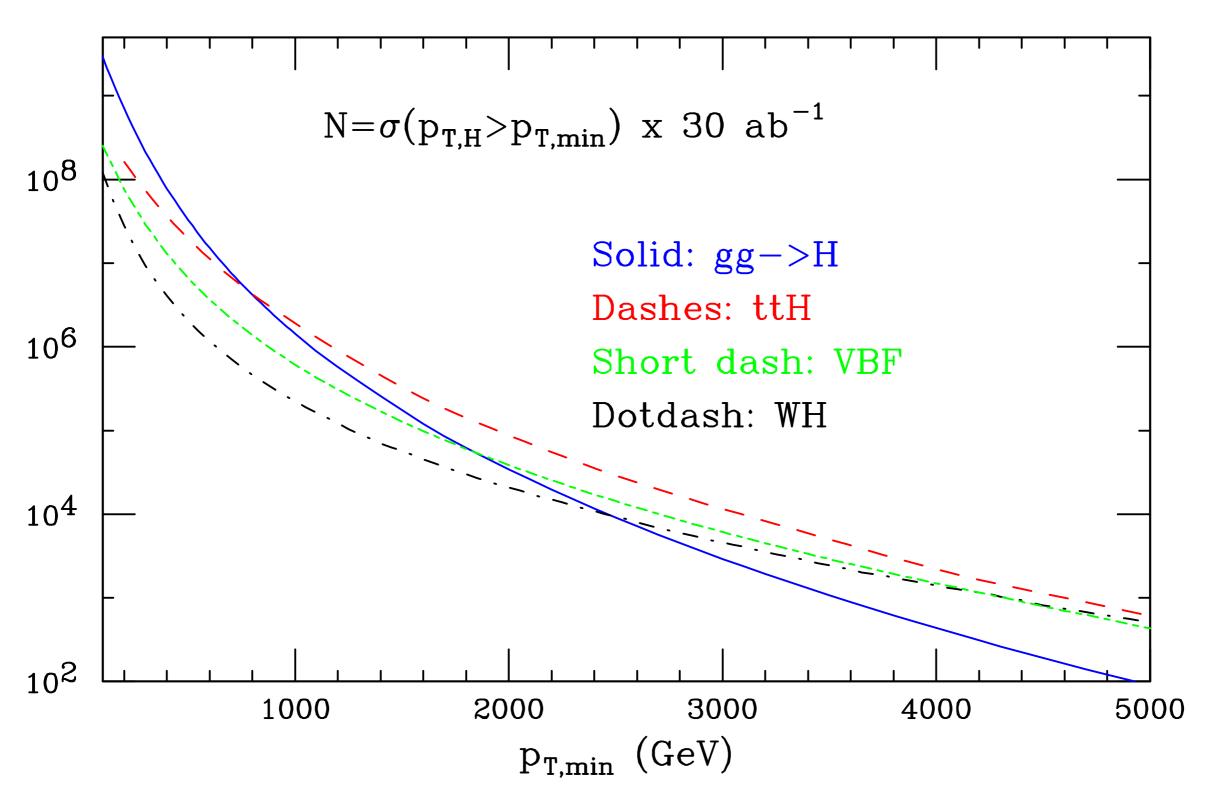
#### The absolutely unique power of pp $\rightarrow$ H+X:

- the extraordinary statistics that, complemented by the per-mille e<sup>+</sup>e<sup>-</sup> measurement of eg BR( $H \rightarrow ZZ^*$ ), allows
  - the sub-% measurement of rarer decay modes
  - the ~5% measurement of the Higgs trilinear selfcoupling
- the huge dynamic range (eg pt(H) up to several TeV), which allows to
  - probe d>4 EFT operators up to scales of several TeV
  - search for multi-TeV resonances decaying to H, or extensions of the Higgs sector

	gg→H	VBF	WH	ZH	ttH	нн
N <sub>100</sub>	24 x 10 <sup>9</sup>	2.1 x 10 <sup>9</sup>	4.6 x 10 <sup>8</sup>	3.3 x 10 <sup>8</sup>	9.6 x 10 <sup>8</sup>	$3.6 \times 10^7$
N <sub>100</sub> /N <sub>14</sub>	180	170	100	110	530	390

$$N_{100} = \sigma_{100 \text{ TeV}} \times 30 \text{ ab}^{-1}$$
  
 $N_{14} = \sigma_{14 \text{ TeV}} \times 3 \text{ ab}^{-1}$ 

#### H at large pt



- Hierarchy of production channels changes at large  $p_T(H)$ :
  - $\sigma(ttH) > \sigma(gg \rightarrow H)$  above 800 GeV
  - $\sigma(VBF) > \sigma(gg \rightarrow H)$  above 1800 GeV

#### Three kinematic regimes

- Inclusive production,  $p_T > 0$ :
  - largest overall rates
  - most challenging experimentally:
    - $\bullet$  triggers, backgrounds, pile-up  $\Rightarrow$  low efficiency, large systematics
  - det simulations challenging, likely unreliable \Rightarrow regime not studied so far

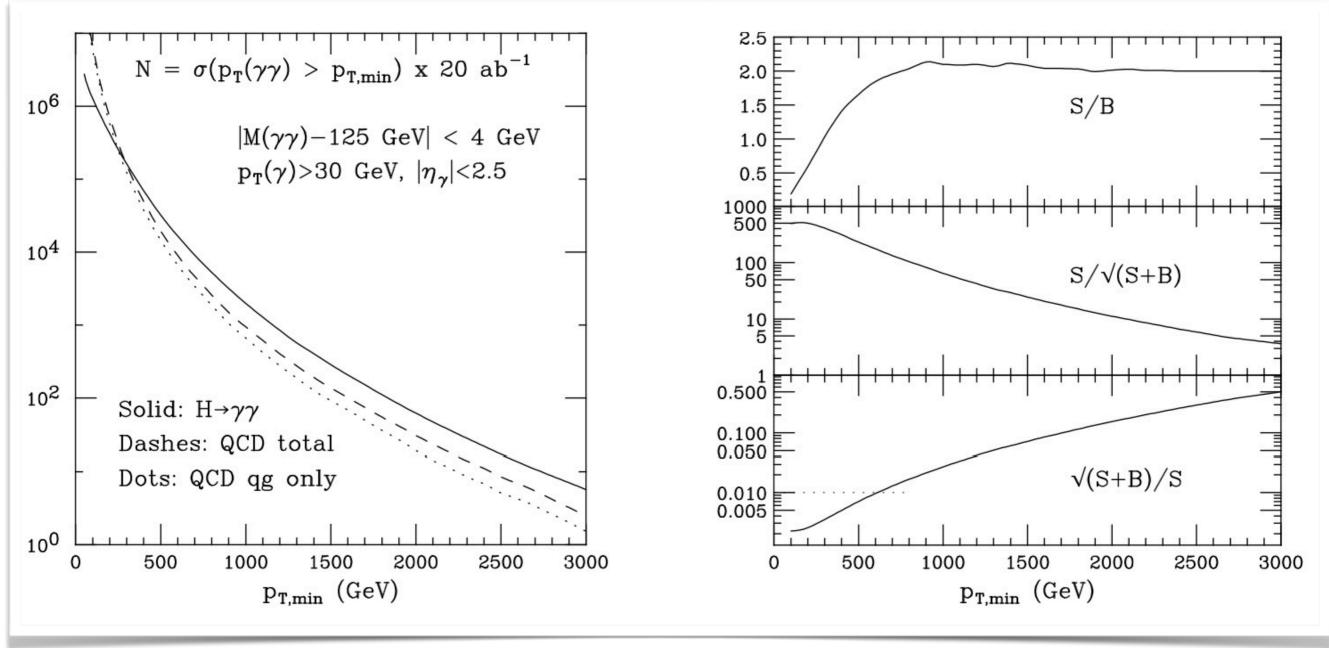
#### • p<sub>T</sub> ≥ 100 GeV :

- stat uncertainty ~few ×  $10^{-3}$  for H  $\rightarrow$  4I,  $\gamma\gamma$ , ...
- improved S/B, realistic trigger thresholds, reduced pile-up effects?
- current det sim and HL-LHC extrapolations more robust
- focus of FCC CDR Higgs studies so far
- sweet-spot for precision measurements at the sub-% level

#### • <u>p</u><sub>T</sub> ≳ TeV :

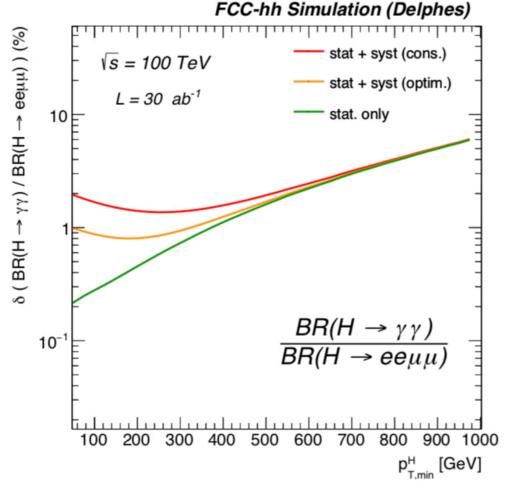
- stat uncertainty O(10%) up to 1.5 TeV (3 TeV) for  $H \rightarrow 4I$ ,  $\gamma\gamma$  ( $H \rightarrow bb$ )
- new opportunities for reduction of syst uncertainties (TH and EXP)
- different hierarchy of production processes
- indirect sensitivity to BSM effects at large Q<sup>2</sup>, complementary to that emerging from precision studies (eg decay BRs) at Q~m<sub>H</sub>

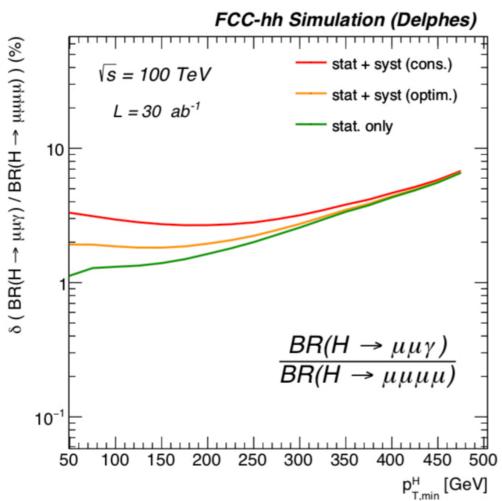
#### gg→H→γγ at large p<sub>T</sub>

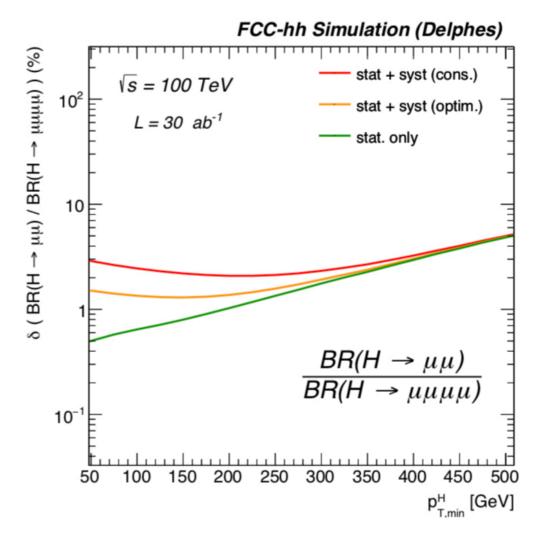


- At LHC, S/B in the  $H \rightarrow \gamma \gamma$  channel is O( few %)
- At FCC, for  $p_T(H)>300$  GeV,  $S/B\sim I$
- Potentially accurate probe of the H pt spectrum up to large pt

р <sub>т,min</sub> (GeV)	$\delta_{stat}$
100	0.2%
400	0.5%
600	1%
1600	10%







Normalize to BR(4I) from ee => sub-% precision for absolute couplings

Future work: explore in more depth data-based techniques, to <u>validate and</u> then reduce the systematics in these ratio measurements, possibly moving to lower pt's and higher stat

#### Higgs couplings after FCC-ee / hh

	HL-LHC	FCC-ee	FCC-hh
δΓ <sub>H</sub> / Γ <sub>H</sub> (%)	SM	1.3	tbd
δg <sub>HZZ</sub> / g <sub>HZZ</sub> (%)	1.5	0.17	tbd
δg <sub>HWW</sub> / g <sub>HWW</sub> (%)	1.7	0.43	tbd
δg <sub>Hbb</sub> / g <sub>Hbb</sub> (%)	3.7	0.61	tbd
δg <sub>Hcc</sub> / g <sub>Hcc</sub> (%)	~70	1.21	tbd
δg <sub>Hgg</sub> / g <sub>Hgg</sub> (%)	2.5 (gg->H)	1.01	tbd
δднττ / днττ (%)	1.9	0.74	tbd
δg <sub>Ημμ</sub> / g <sub>Ημμ</sub> (%)	4.3	9.0	0.65 (*)
δg <sub>Hγγ</sub> / g <sub>Hγγ</sub> (%)	1.8	3.9	0.4 (*)
δg <sub>Htt</sub> / g <sub>Htt</sub> (%)	3.4	~10 (indirect)	0.95 (**)
δg <sub>HZγ</sub> / g <sub>HZγ</sub> (%)	9.8	_	0.9 (*)
δдннн / дннн (%)	50	~44 (indirect)	5
BR <sub>exo</sub> (95%CL)	BR <sub>inv</sub> < 2.5%	< 1%	BR <sub>inv</sub> < 0.025%

#### NB

BR(H $\rightarrow$ Z $\gamma$ , $\gamma\gamma$ ) ~O(10<sup>-3</sup>)  $\Rightarrow$  O(10<sup>7</sup>) evts for  $\Delta_{stat}$ ~%
BR(H $\rightarrow\mu\mu$ ) ~O(10<sup>-4</sup>)  $\Rightarrow$  O(10<sup>8</sup>) evts for  $\Delta_{stat}$ ~%

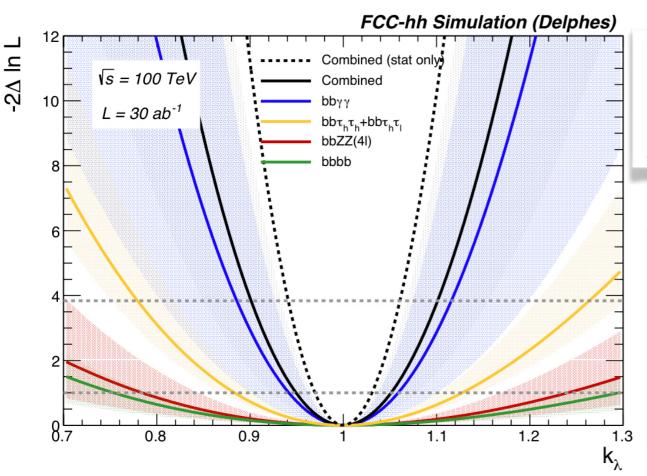


pp collider is essential to beat the % target, since no proposed ee collider can produce more than O(106) H's

<sup>\*</sup> From BR ratios wrt B(H→ZZ\*) @ FCC-ee

<sup>\*\*</sup> From pp→ttH / pp→ttZ, using B(H→bb) and ttZ EW coupling @ FCC-ee

#### The Higgs self-coupling at FCC-hh



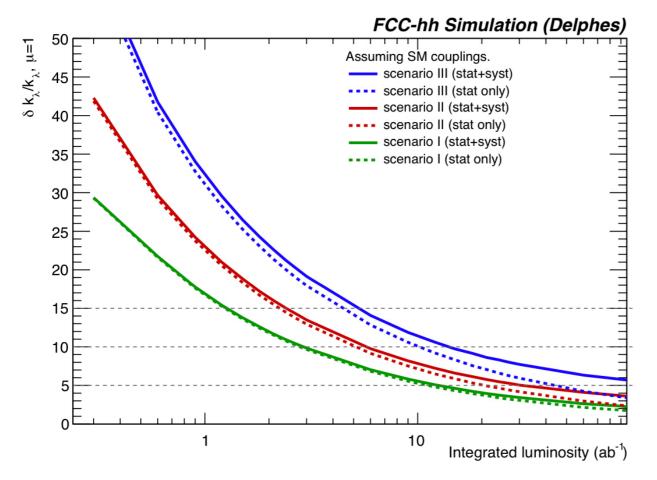


Figure 13. Expected negative log-Likelihood scan as a function of the trilinear self-coupling modifier  $\kappa_{\lambda} = \lambda_3/\lambda_3^{\rm SM}$  in all channels, and their combination. The solid line corresponds to the scenario II for systematic uncertainties. The band boundaries represent respectively scenario I and III. The dashed line represents the sensitivity obtained including statistical uncertainties only, under the assumptions of scenario I.

C.1	10t	SC	SINC	VI/	20
.71	/51	51 1	-11/2	4 I I (	15
$\sim$ 1				411	$\mathcal{I} \cup$
	,				

	@68% CL	scenario I	scenario II	scenario III
	stat only	2.2	2.8	3.7
$\delta_{\mu}$	stat + syst	2.4	3.5	5.1
-2	stat only	3.0	4.1	5.6
$\delta_{\kappa_{\lambda}}$	stat + syst	3.4	5.1	7.8

**Table 7**. Combined expected precision at 68% CL on the di-Higgs production cross- and Higgs self coupling using all channels at the FCC-hh with  $\mathcal{L}_{int} = 30 \text{ ab}^{-1}$ . The symmetrized value  $\delta = (\delta^+ + \delta^-)/2$  is given in %.

- I. Target det performance: LHC Run 2 conditions
- II. Intermediate performance
- III. Conservative: extrapolated HL-LHC performance, with today's algo's (eg no timing, etc)

Expected precision on the Higgs self-coupling as a function of the integrated luminosity.

3-5 ab<sup>-1</sup> are sufficient to get below the 10% level

- => within the reach of the first 5yrs of FCC-hh running, in the "low" luminosity / low pileup phase
- => compatible with the timescale for a similar precision measurement by CLIC @ 3 TeV

#### Higgs as a BSM probe: precision vs dynamic reach

$$L = L_{SM} + \frac{1}{\Lambda^2} \sum_{k} \mathcal{O}_k + \cdots$$

$$O = |\langle f|L|i\rangle|^2 = O_{SM} \left[1 + O(\mu^2/\Lambda^2) + \cdots\right]$$

For H decays, or inclusive production,  $\mu \sim O(v, m_H)$ 

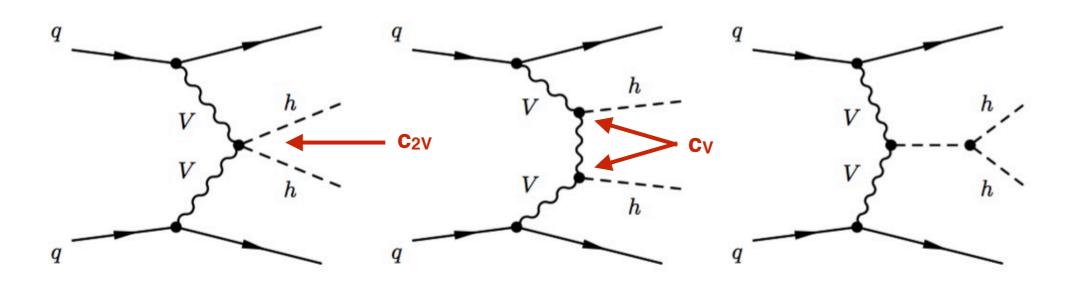
$$\delta O \sim \left(\frac{v}{\Lambda}\right)^2 \sim 6\% \left(\frac{\text{TeV}}{\Lambda}\right)^2$$
  $\Rightarrow$  precision probes large  $\Lambda$  e.g.  $\delta O = 1\% \Rightarrow \Lambda \sim 2.5 \,\text{TeV}$ 

For H production off-shell or with large momentum transfer Q,  $\mu\sim O(Q)$ 

$$\delta O \sim \left(\frac{Q}{\Lambda}\right)^2$$
  $\Rightarrow$  kinematic reach probes large  $\Lambda$  even if precision is low e.g.  $\delta O=15\%$  at Q=1 TeV  $\Rightarrow \Lambda\sim2.5$  TeV

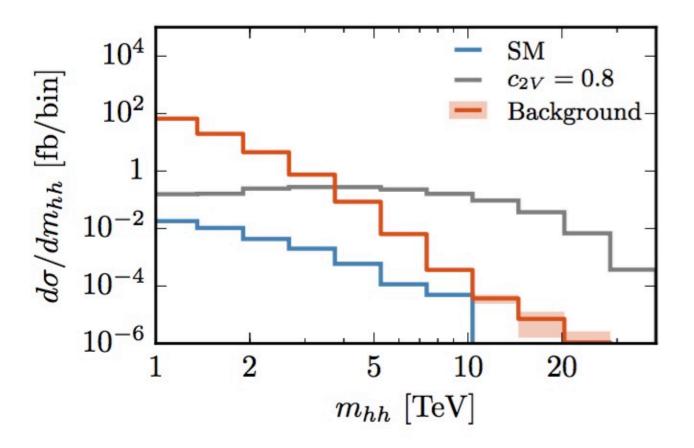
Precision and extensive kinematic reach provide unique complementarity and redundancy, crucial to interpret possible SM deviations manifest in either of these observabes

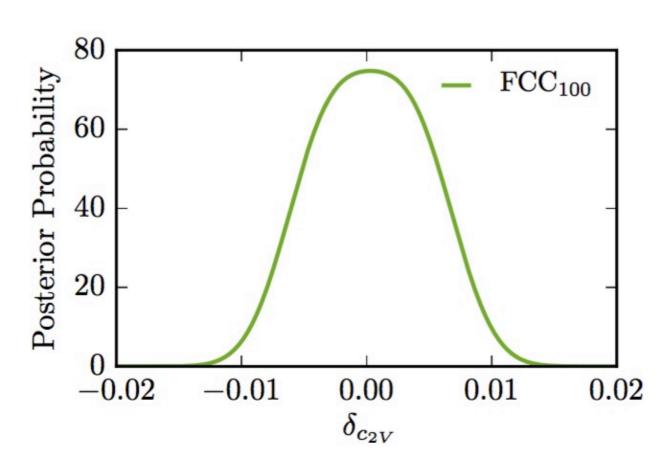
#### **Example:** high mass VV → HH



$$A({
m V_LV_L}
ightarrow {
m HH}) \sim rac{\hat s}{v^2}(c_{2V}-c_V^2)$$
  $\cdot$  where

$$A(\mathbf{V_LV_L} \to \mathbf{HH}) \sim \frac{\hat{s}}{v^2}(c_{2V} - c_V^2) \cdot \text{where} \qquad \begin{cases} c_V = g_{HVV}/g_{HVV}^{SM} \\ c_{2V} = g_{HHVV}/g_{HHVV}^{SM} \end{cases} \implies \left(c_{2V} - c_V^2\right)_{SM} = 0$$





#### (1) guaranteed deliverables: EW observables

#### The absolutely unique power of **Circular** e<sup>+</sup>e<sup>-</sup>:

$e^+e^- \rightarrow Z$	$e^+e^- \rightarrow WW$	<b>τ(←Z)</b>	b( <b>←Z</b> )	c(←Z)
5 10 <sup>12</sup>	<b>10</b> <sup>8</sup>	3 1011	1.5 10 <sup>12</sup>	<b>10</b> <sup>12</sup>

=> O(105) larger statistics than LEP at the Z peak and WW threshold

## EW parameters @ FCC-ee

Observable	present value ± error	FCC-ee stat.	FCC-ee syst.
$m_Z (keV)$	91186700±2200	5	100
$\Gamma_{\rm Z}$ (keV)	2495200±2300	8	100
$R_l^Z \ (\times 10^3)$	20767±25	0.06	0.2-1.0
$\alpha_s$ (mz) (×104)	1196±30	0.1	0.4-1.6
R <sub>b</sub> (×10 <sup>6</sup> )	216290±660	0.3	<60
$\sigma_{had}^{0} \; (\times 10^{3}) \; (nb)$	41541±37	0.1	4
$N_{\nu} \ (\times 10^{3})$	2991±7	0.005	1
$\sin^2 \theta_W^{eff} (\times 10^6)$	231480±160	3	2-5
$1/\alpha_{\text{QED}}(m_{\text{Z}}) (\times 10^3)$	128952±14	4	Small
$A_{\rm FB}^{b,0}~(\times 10^4)$	992±16	0.02	1-3
$A_{\rm FB}^{{\rm pol},\tau}~(\times 10^4)$	1498±49	0.15	<2
m <sub>W</sub> (MeV)	80350±15	0.6	0.3
Γ <sub>W</sub> (MeV)	2085±42	1.5	0.3
$\alpha_s$ (m <sub>W</sub> ) (×10 <sup>4</sup> )	1170±420	3	Small
$N_{\nu}(\times 10^3)$	2920±50	0.8	Small
m <sub>top</sub> (MeV)	172740±500	20	Small
Γ <sub>top</sub> (MeV)	1410±190	40	Small
$\lambda_{\mathrm{top}}/\lambda_{\mathrm{top}}^{\mathrm{SM}}$	1.2±0.3	0.08	Small
ttZ couplings	±30%	0.5 - 1.5%	Small

#### Precision W physics with pp→tt[→Wb]

#### MLM @ SEARCH2016

#### A concrete application: testing lepton universality in W decays

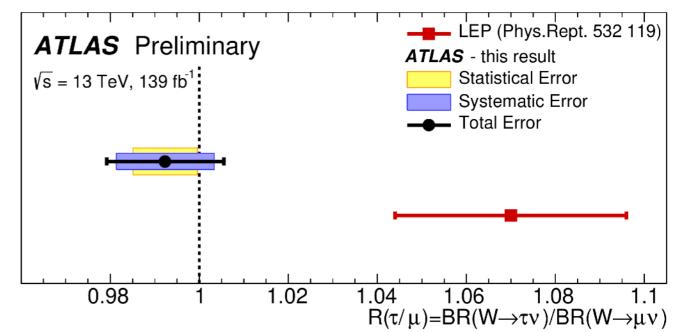
PDG entries dominated by LEP2 data

W+ DECAY MODES	F	Fraction $(\Gamma_i/\Gamma)$		Confidence level	(MeV/c)
$\ell^+ \nu$	[b]	(10.86±	0.09) %		_
$e^+ u$		$(10.71 \pm$	0.16) %		40192
$\mu^+ \nu$		$(10.63 \pm$	0.15) %		40192
$ au^+ u$		$(11.38 \pm$	0.21) %		40173

BR(
$$\tau$$
) / BR( $e/\mu$ ) ~ 1.066 ± 0.025 => ~ 2.5  $\sigma$ 

can the LHC clarify this issue with its eventual  $10^7$  leptonic W decays from the top?

#### **ATLAS 2020:**



LEP:

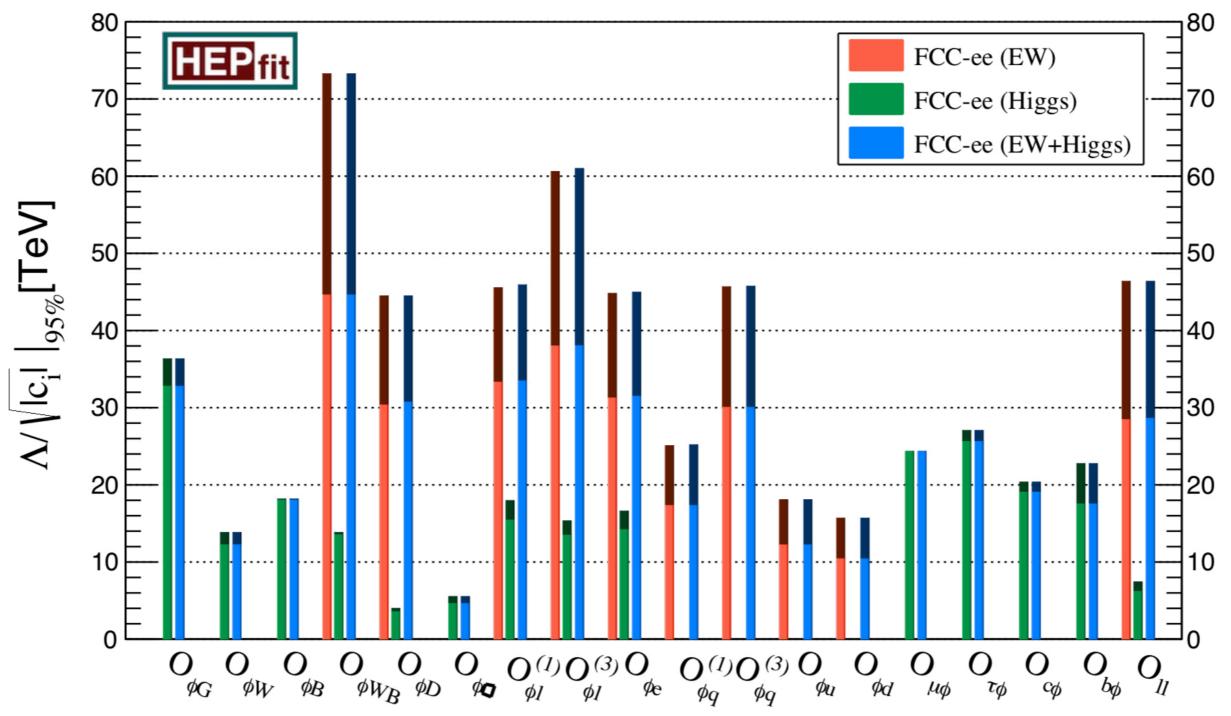
$$BR(W \to \tau v)/BR(W \to \mu v) = 1.066 \pm 0.025$$

**ATLAS:** 

$$BR(W \to \tau v)/BR(W \to \mu v) = 0.992(\pm 0.013)$$

(2) Direct discovery reach at high mass: the power of 100 TeV

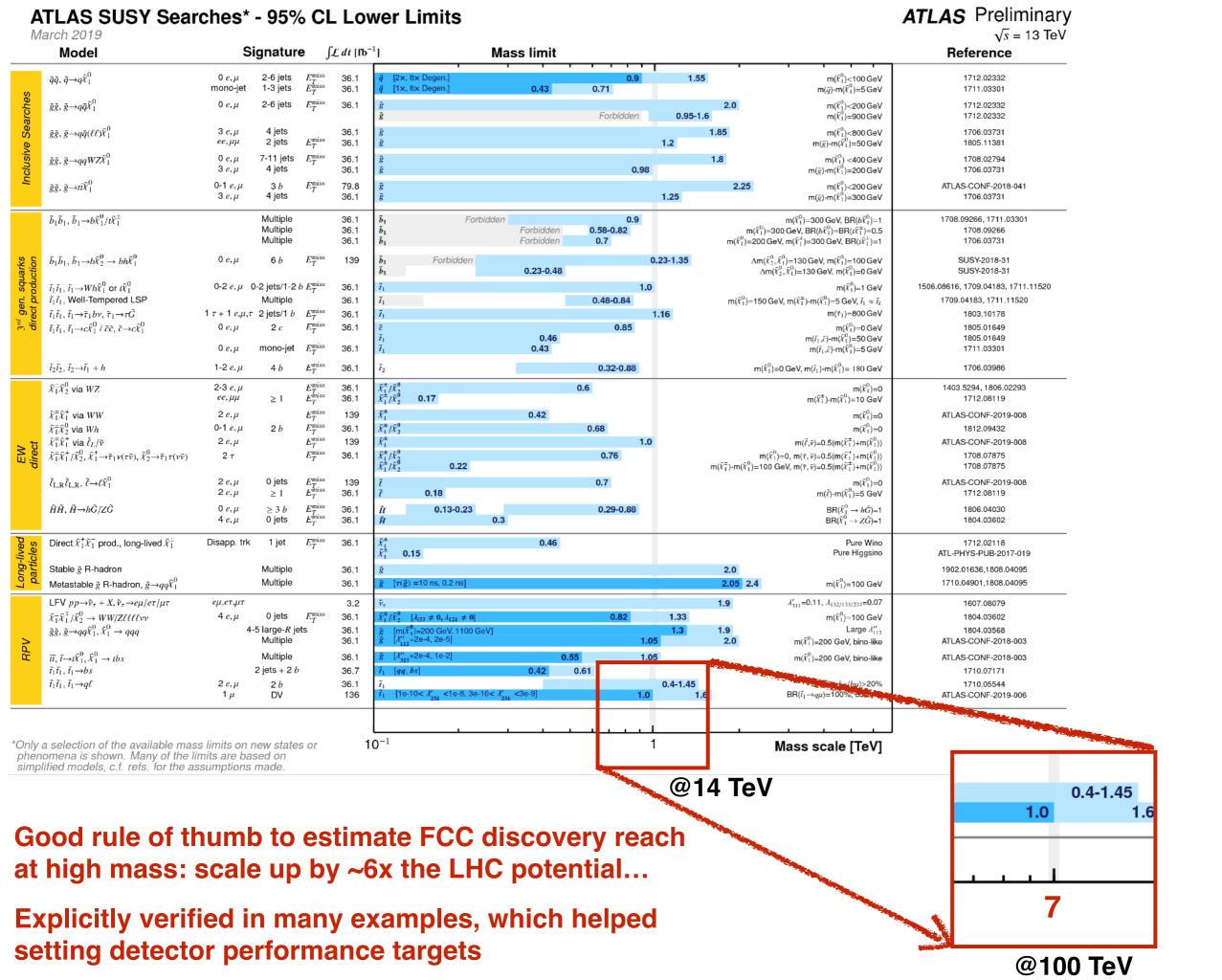
#### Global EFT fits to EW and H observables at FCC-ee



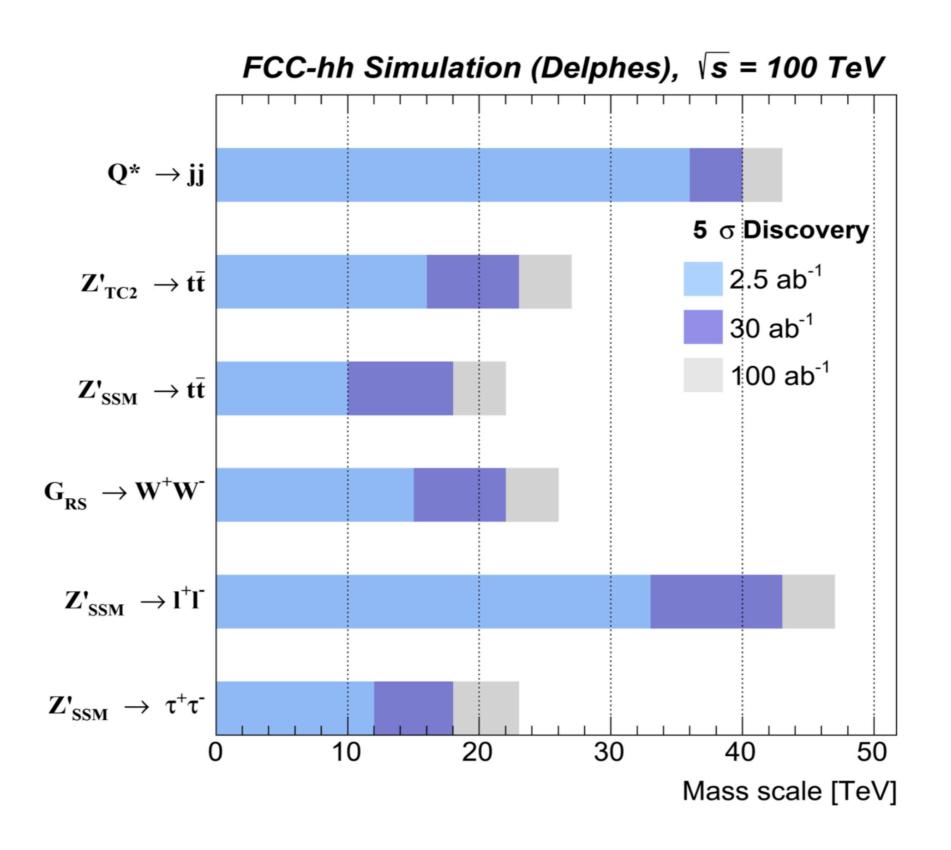
Constraints on the coefficients of various EFT op's from a global fit of (i) EW observables, (ii) Higgs couplings and (iii) EW+Higgs combined. Darker shades of each color indicate the results neglecting all SM theory uncertainties.



100 TeV is the appropriate CoM energy to directly search for new physics appearing indirectly through precision EW and H measurements at the future ee collider

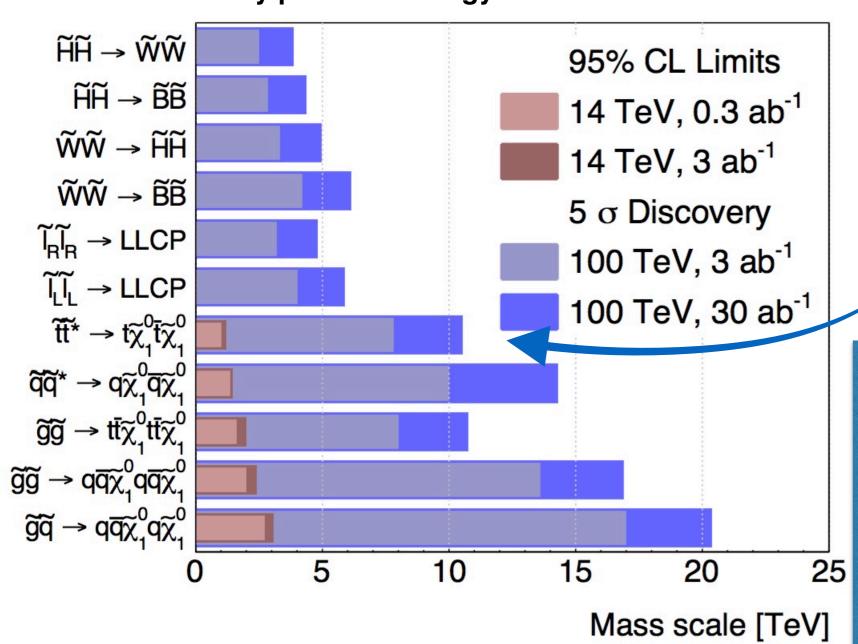


#### s-channel resonances

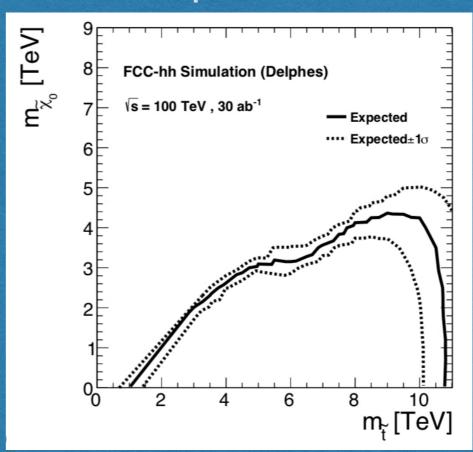


#### SUSY reach at 100 TeV

#### Early phenomenology studies



#### **New detector performance studies**



(3) The potential for yes/no answers to important questions

#### WIMP DM theoretical constraints

For particles held in equilibrium by pair creation and annihilation processes, ( $\chi \chi \leftrightarrow SM$ )

$$\Omega_{\mathrm{DM}} h^2 \sim \frac{10^9 \mathrm{GeV}^{-1}}{M_{\mathrm{pl}}} \frac{1}{\langle \sigma v \rangle}$$

For a particle annihilating through processes which do not involve any larger mass scales:

$$\langle \sigma v \rangle \sim g_{\rm eff}^4 / M_{\rm DM}^2$$

$$\Omega_{\rm DM} h^2 \sim 0.12 \times \left(\frac{M_{\rm DM}}{2 \, {\rm TeV}}\right)^2 \left(\frac{0.3}{g_{\rm eff}}\right)^4$$

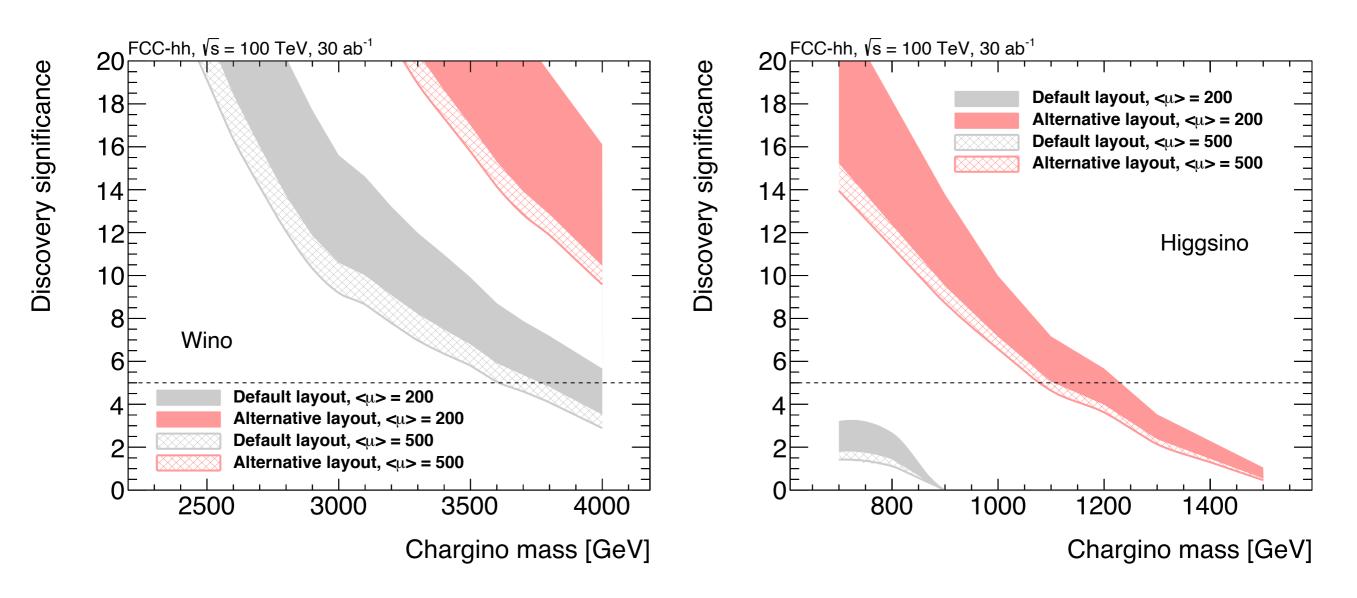
$$\Omega_{wimp} h^2 \lesssim 0.12$$



$$M_{wimp} \lesssim 2 \text{ TeV} \left(\frac{g}{0.3}\right)^2$$

#### DM WIMP searches in the most elusive, compressed scenarios:

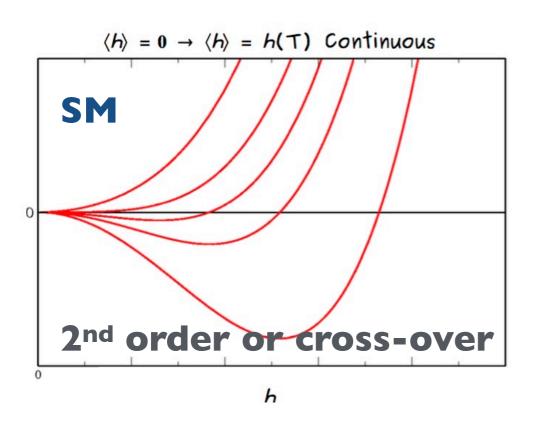
#### Disappearing charged track analyses (at ~full pileup)

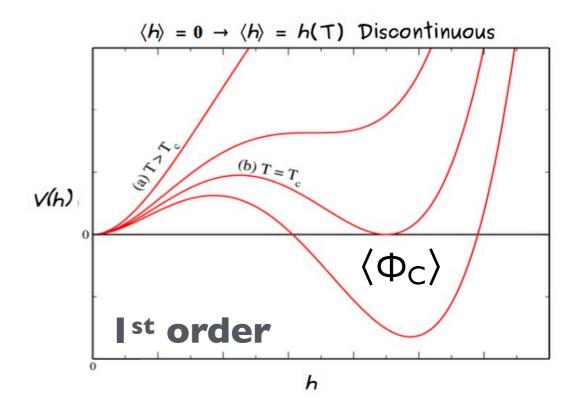


=> coverage beyond the upper limit of the thermal WIMP mass range for both higgsinos and winos !!

$$M_{wimp} \lesssim 2 \text{ TeV} \left(\frac{g}{0.3}\right)^2$$

#### The nature of the EW phase transition





Strong Ist order phase transition is required to induce and sustain the out of equilibrium generation of a baryon asymmetry during EW symmetry breaking

**Strong** Ist order phase transition  $\Rightarrow \langle \Phi_C \rangle > T_C$ 

In the SM this requires  $m_H \lesssim 80$  GeV, else transition is a smooth crossover.

Since  $m_H = 125$  GeV, **new physics**, coupling to the Higgs and effective at **scales** O(TeV), must modify the Higgs potential to make this possible



- Probe higher-order terms of the Higgs potential (selfcouplings)
- Probe the existence of other particles coupled to the Higgs

#### Constraints on models with 1st order phase transition at the FCC

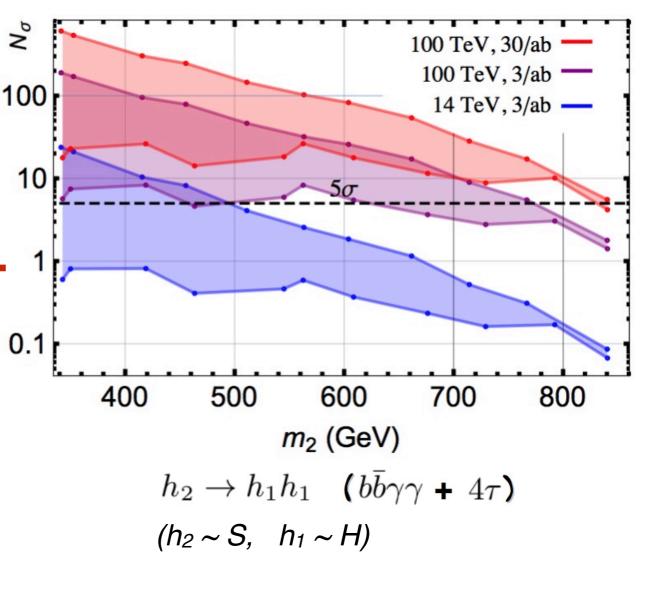
$$V(H,S) = -\mu^{2} (H^{\dagger}H) + \lambda (H^{\dagger}H)^{2} + \frac{a_{1}}{2} (H^{\dagger}H) S$$
$$+ \frac{a_{2}}{2} (H^{\dagger}H) S^{2} + \frac{b_{2}}{2} S^{2} + \frac{b_{3}}{3} S^{3} + \frac{b_{4}}{4} S^{4}.$$

#### Combined constraints from precision Higgs measurements at FCC-ee and FCC-hh

# Real Scalar Singlet Model $1 - \frac{1}{1000}$ 0.100 100

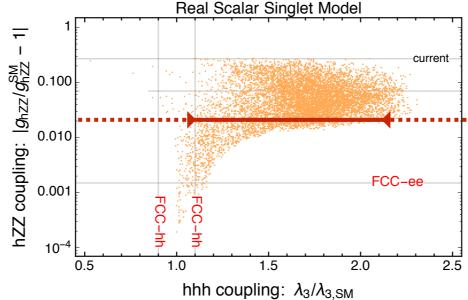
Parameter space scan for a singlet model extension of the Standard Model. The points indicate a first order phase transition.

#### Direct detection of extra Higgs states at FCC-hh



#### Remarks

- Apparently, adding the self-coupling constraint does not add much in terms of exclusion power, wrt the HZZ coupling measurement ...
- ... BUT, should HZZ deviate from the SM,  $\lambda_{HHH}$  is necessary to break the degeneracy among all parameter sets leading to the same HZZ prediction



- The concept of "which experiment sets a better constraint on a given parameter" is a very limited comparison criterion, which looses value as we move from "setting limits" to "diagnosing observed discrepancies"
- Likewise, it's often said that some observable sets better limits than others: "all known model predict deviations in X larger than deviations in Y, so we better focus on X". But once X is observed to deviate, knowing the value of Y could be absolutely crucial ....
- Redundancy and complementarity of observables is of paramount importance
- The full, integrated, FCC programme, is the only proposed facility capable of providing such a complementarity

#### **Not covered**

- Countless studies of discovery potential for multiple BSM scenarios, from SUSY to heavy neutrinos, from very low masses to very high masses, LLPs, DM, etcetcetc, at FCC-ee, FCC-hh and FCC-eh
- Sensitivity studies to SM deviations in the properties of top quarks, flavour physics in Z decays: huge event rates offer unique opportunities, that cannot be matched elsewhere
- •
- Operations with heavy ions: new domains open up at 100 TeV in the study of high-T/high-density QCD. Broaden the targets, the deliverables, extend the base of potential users, and increase the support beyond the energy frontier community

#### Final remarks

- The study of the SM will not be complete until we clarify the nature of the Higgs mechanism and exhaust the exploration of phenomena at the TeV scale: many aspects are still obscure, many questions are still open.
- The combination of a versatile high-luminosity e<sup>+</sup>e<sup>-</sup> circular collider, with a follow-up pp collider in the 100 TeV range, appears like the ideal facility for the post-LHC era
  - complementary and synergetic precision studies of EW, Higgs and top properties
  - energy reach to allow direct discoveries at the mass scales possibly revealed by the precision measurements