

TENSOR NETWORK STATES FOR QUANTUM FIELD THEORY

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(Garching)

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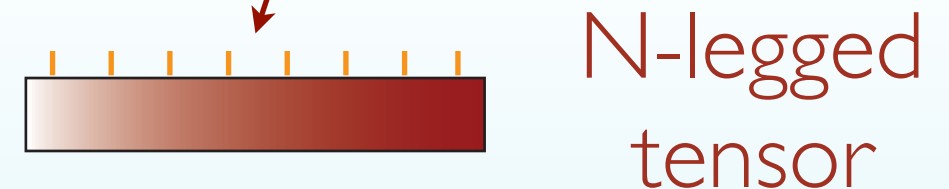
TENSOR NETWORKS

WHAT ARE TNS?

- TNS = Tensor Network States

A general state of the N-body Hilbert space has exponentially many coefficients

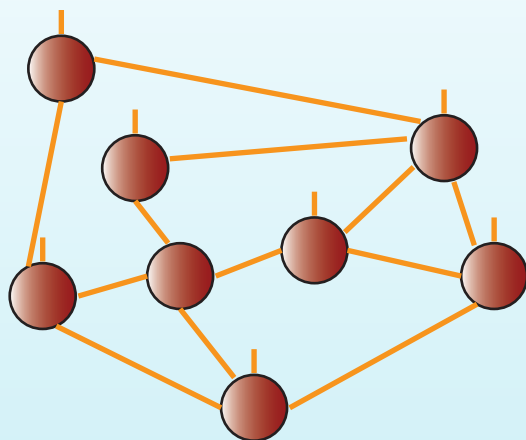
$$|\Psi\rangle = \sum_{i_j} c_{i_1 \dots i_N} |i_1 \dots i_N\rangle$$



$$d^N$$

A TNS has only a polynomial number of parameters

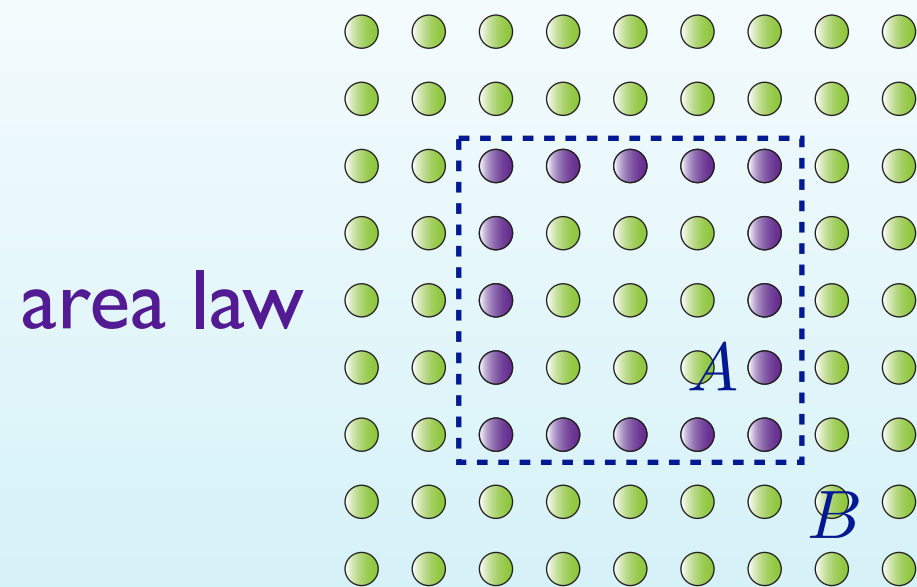
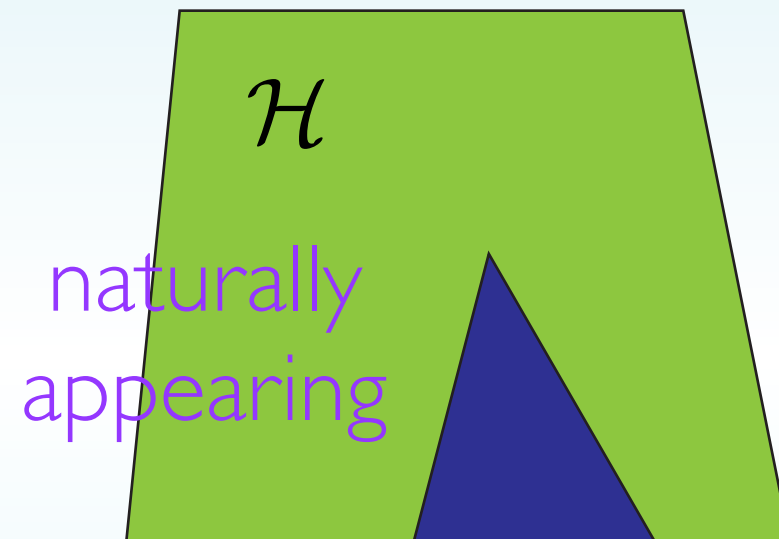
$\text{poly}(N)$



WHY SHOULD TNS BE USEFUL?

States appearing in Nature are peculiar

State at random from Hilbert space is not close to product



We look for states with *little* entanglement

TNS = entanglement based ansatz

WHY SHOULD TNS BE USEFUL?

Which properties characterize ground states of relevant Hamiltonians?

local gapped 1D Hamiltonians
have ground states
with area law of entanglement

$$S_{A_{\max}} \propto |\delta A| \quad \text{Hastings 2007}$$

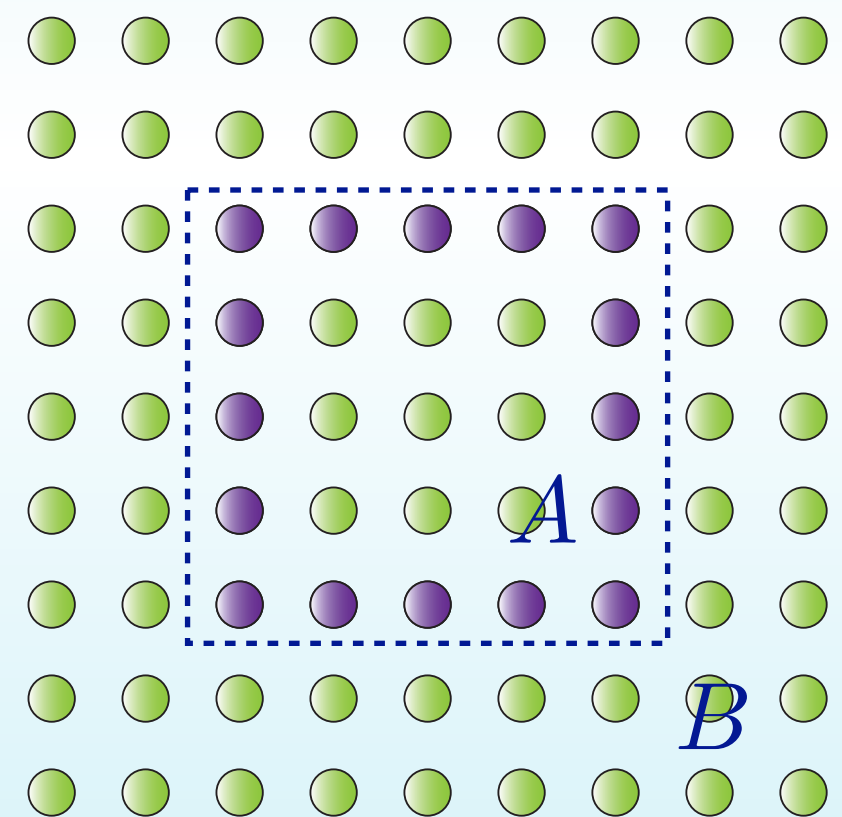
in 1D critical systems,
logarithmic corrections

$$S_{A_{\max}} \propto |\delta A| \log A \quad \text{Calabrese, Cardy 2004}$$

satisfied at finite temperature

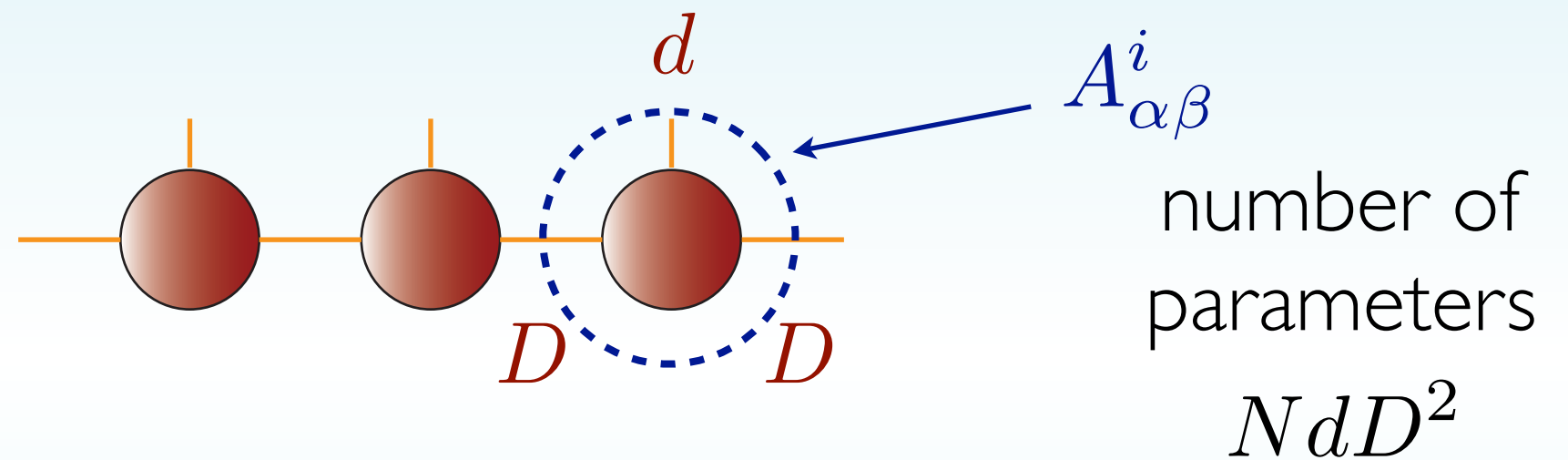
Wolf, Verstraete, Hastings, Cirac, PRL 2008

Area law

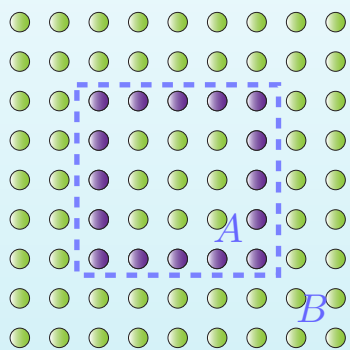


MPS

Matrix Product States



$$|\Psi\rangle = \sum_{i_1 \dots i_N} \text{tr}(A_1^{i_1} A_2^{i_2} \dots A_N^{i_N}) |i_1 \dots i_N\rangle$$



Area law by construction

Bounded entanglement $S(L/2) \leq \log D$

MPS

Matrix Product States

SOME PROPERTIES

good approximation of ground states

area law (ground state)

critical states: $D \sim \text{poly}(N)$

Verstraete, Cirac, PRB 2006

Hastings, J. Stat. Phys 2007

efficient calculation of expectation values

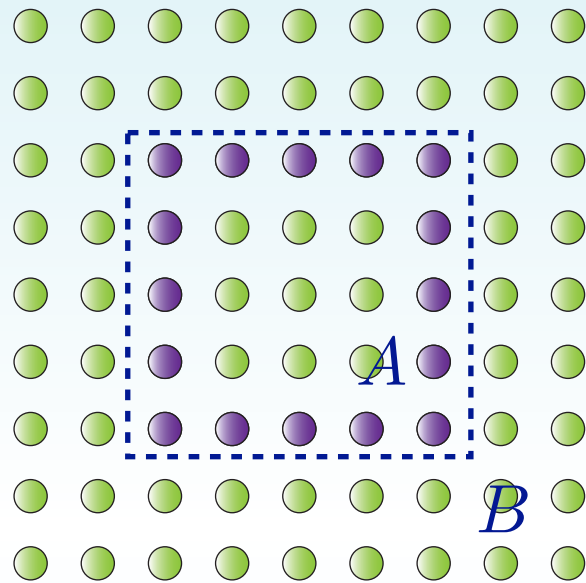
exponentially decaying correlations

can be prepared efficiently

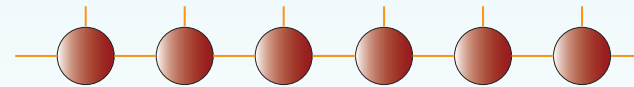
natural generalisation to higher dimensions: PEPS

TNS = entanglement based ansatz

area law

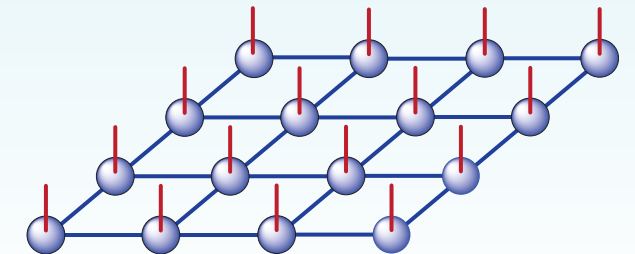


MPS



Schollwöck Ann.Phys.2011

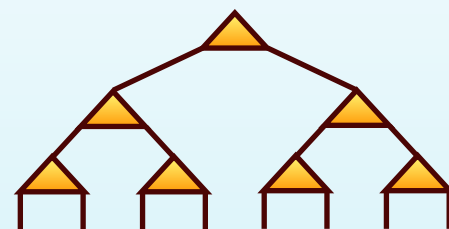
PEPS



Verstraete et al. Adv. Phys. 2008

other TNS

TTN

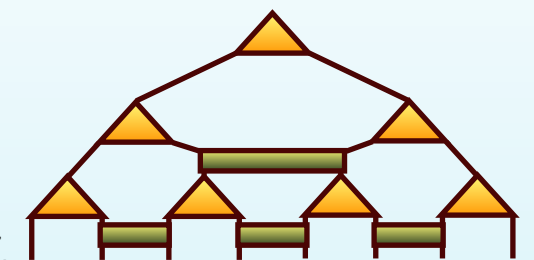


Shi et al PRA 2006

suggested connection
to AdS/CFT

Vidal PRL 2007

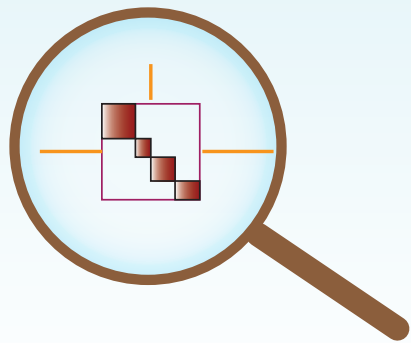
MERA



Swingle PRD 2012
Molina JHEP 2013
Nozaki et al JHEP 2012
Bao et al PRD 2015

USING TNS FOR QMB

a formal approach



classifying tensors

constructing states

great descriptive power: phases,
topological chiral states, anyons...

Chen et al PRB 2011

Schuch et al PRB 2011

Wahl et al PRL 2013; Yang et al PRL 2015

Haegeman et al, Nat. Comm. 2015

no sign problem

numerical algorithms

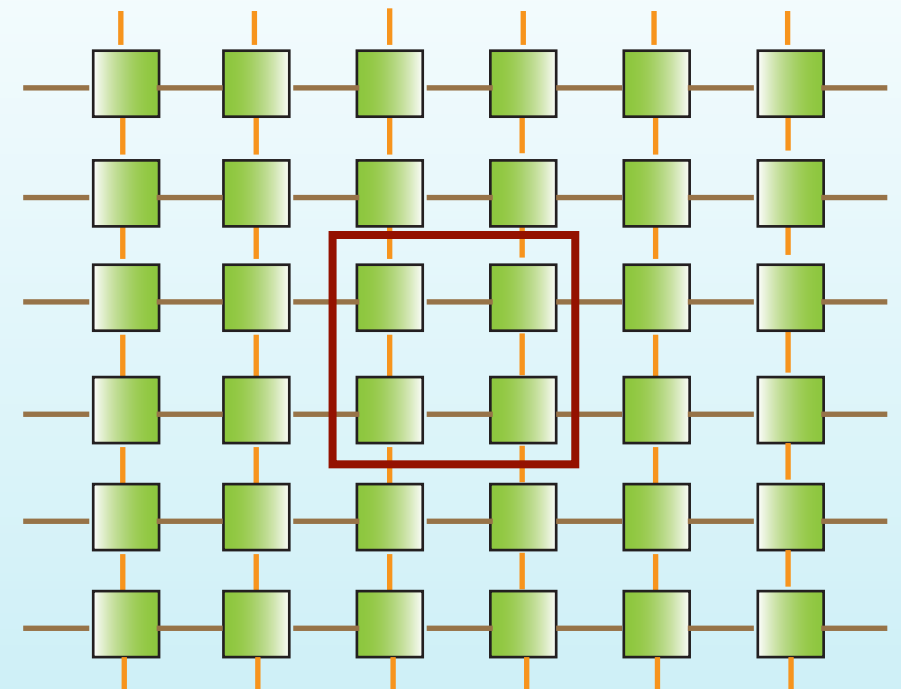
tensor networks describe
partition functions (observables)

need to contract a TN
TRG approaches

Nishino, JPSJ 1995

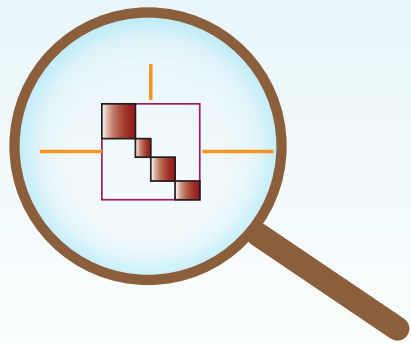
Levin & Wen PRL 2008

Xie et al PRL 2009; Zhao et al PRB 2010



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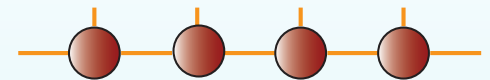
Haegeman et al, Nat. Comm. 2015

no sign problem

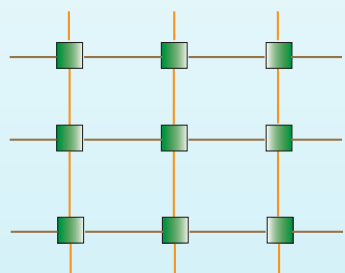
numerical algorithms



tensor networks describe
partition functions (observables)



TNS as ansatz for the state



need to contract a TN
TRG approaches

Nishino, JPSJ 1995

Levin & Wen PRL 2008

Xie et al PRL 2009; Zhao et al PRB 2010

efficient algorithms for GS, low
excited states, thermal, dynamics

White PRL 1992; Schollwöck RMP 2011

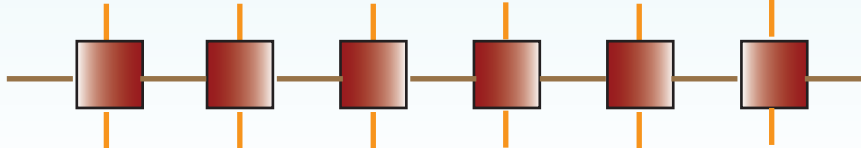
Vidal PRL 2003; Verstraete et al PRL 2004

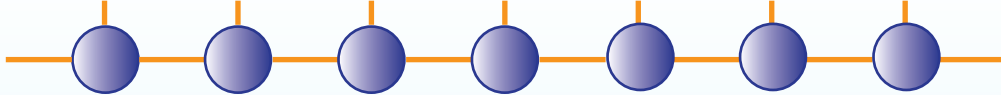
Verstraete et al Adv Phys 2008; Orús Ann Phys 2014

NUMERICAL ALGORITHMS

variational minimization of energy

local
Hamiltonian

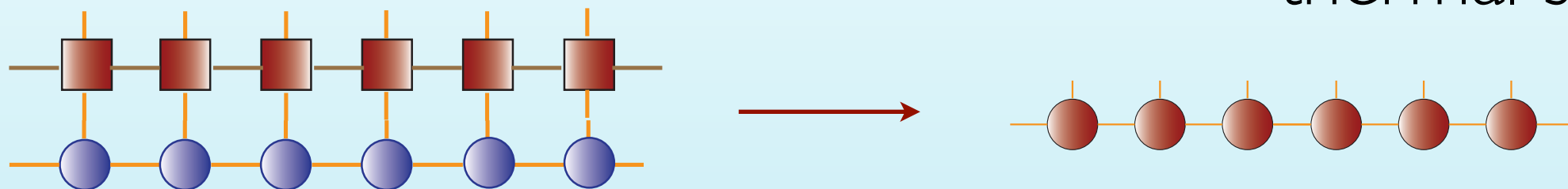
$$H =$$


$$|E_0\rangle \simeq$$


ground state
excitations

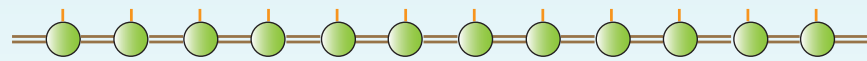
apply local operators $\rightarrow$ simulate time evolution

imaginary time $\rightarrow$ ground state
thermal state



various evolution algorithms

evolving the (pure state) ansatz

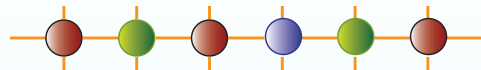


Vidal, PRL 2003, PRL 2007
White, Feiguin, PRL 2004
Daley et al., 2004
Haegeman et al., 2011

entanglement can grow fast!

Osborne, PRL 2006
Schuch et al., NJP 2008

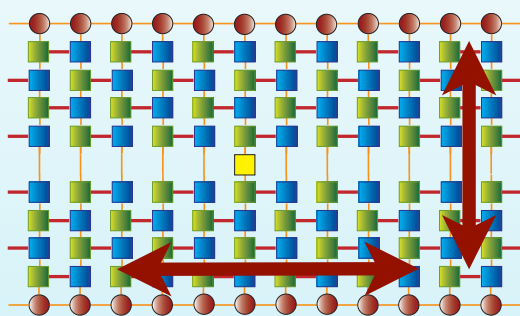
evolving operators: Heisenberg picture Hartmann et al, PRL 2009



also for mixed states
operator space entanglement

Prosen Pizorn, PRL 2008

observables as TN to contract



different *entanglement* quantities

MCB, Hastings, Verstraete, Cirac, PRL 2009
Müller-Hermes et al., NJP 2012
Hastings, Mahajan 2014

others: density of states

Yang, Iblisdir, Cirac, MCB, 1909.01398

USING TNS FOR QFT

<http://heptnseminar.org>

Tensor Networks in HEP Seminar x +

gqfi.aei.mpg.de/node/25

Upcoming Seminars

Frank Verstraete

When: January, 2020
Title: TBA
Abstract: TBA

Past Seminars

1. Ignacio Cirac (Opening seminar of the series)

When: November 8, 2019 @ 15.00 (GMT +1 hour, Berlin time)
Title: *Tensor Networks and Lattice Gauge Theories*
Abstract: Certain Quantum Many-body states can be efficiently described in terms of tensor networks. Those include Matrix Product States (MPS), Projected Entangled-Pair States (PEPS), or the Multi-scale Entanglement Renormalization Ansatz. They play an important role in quantum computing, error correction, or the description of topological order in condensed matter physics, and are widely used in computational physics. In the last years, it has also been realized their suitability to describe Lattice Gauge Theories, at least in the context of MPS in low dimensions. In this talk I will review some of the basic ideas about tensor networks and their applications to lattice gauge theories, and explain current efforts to extend them to higher dimensions using PEPS.
YouTube link: <https://www.youtube.com/watch?v=hdb82b1kazw&feature=youtu.be>

2. Bartlomiej Czech

When: December 6, 2019 @ 15.30 (GMT +1 hour, Berlin time)
Title: *What does the Chern-Simons formulation of AdS3 gravity tell us about complexity?*
Abstract: I will explain how to realize the wavefunction of a CFT2 ground state as a network of Wilson lines in the Chern-Simons

USING TNS FOR LGT

an ongoing LGT-TNS roadmap...

early works with DMRG/TNS

Byrnes PRD2002; Sugihara NPB2004
Tagliacozzo PRB2011; Sugihara JHEP2005
Meurice PRB2013

Schwinger model
 $U(1)$ in 1D
precise equilibrium
simulations,
feasibility of QSim

finite density

S. Kuehn et al, PRL 118 (2017) 071601

Non-Abelian in 1D
string breaking dynamics

S. Kühn et al., JHEP 07 (2015) 130;
Silvi et al., Quantum 2017
S. Kühn et al. PRX 2017

MCB et al JHEP11 (2013) 158;
Rico et al PRL 2014; Buyens et al. PRL 2014;
S. Kühn et al., PRA 90, 042305 (2014);
MCB et al PRD 2015, Buyens et al. PRD 2016;
Pichler et al. PRX 2016;
review Dalmonte, Montangero, Cont. Phys. 2016
MCB, Cichy, Cirac, Jansen, Kühn, arXiv:1810.12838

2+1 dimensions
work in progress

full LQCD in 3+1
dimensions

MCB, K. Cichy 1910.00257
(2019 ROPP)

OTHER MODELS

THIRRING MODEL

MCB, K. Cichy, Y.-J. Kao, C.-J. D. Lin, Y.-P. Lin, D.T.-L. Tan, arXiv:1908.04536

THIRRING MODEL

probing two phases in lattice Thirring model

one is critical, $c=1$, non-vanishing chiral condensate

one is gapped, with exponentially decaying correlations

transition compatible with BKT

outlook

non-perturbative study of scaling (e.g. $\bar{\psi}\psi$)

excitations (cont. limit $g > \pi/2$)

REAL TIME THERMAL FIELD THEORY

with M. Heller, K. Jansen, J. Knaute, V. Svensson, arXiv:1912.08836

THERMAL QUENCHES IN ISING MODELS

Thermal response determined by retarded 2-point function

linear response

source

$$\delta\langle\mathcal{O}(t,p)\rangle = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} d\omega \boxed{G_R(\omega,p)} \delta J(-\omega,-p) e^{-i\omega t}$$

$$G_R(\Delta t, \Delta x) = i \text{tr} \left(Z_{\beta}^{-1} e^{-\beta H} [\mathcal{O}(\Delta t, \Delta x), \mathcal{O}(0,0)] \right)$$

structure of singularities determines time response

→ holographic QFTs: single poles with
imaginary part (transients) BH quasi normal modes

→ CFT in $1+1$ dim, scalar operator

$$\omega = \pm p - i 2\pi T (\Delta + 2n)$$

scaling dimension

THERMAL QUENCHES IN ISING MODELS

Thermal states well described by TNS Hastings 2006; Molnar et al. 2015

real time evolution: thermal response functions Barthel 2013

- use to identify characteristic frequency modes
- connect lattice and field model

Compute 2-point correlators at different times

analyse real-time retarded 2-point function

- zero or finite momentum can be reconstructed
- numerical extraction of frequencies: Prony

THERMAL QUENCHES IN ISING MODELS

Ising spin chain

$$H = -J \left(\sum_{j=1}^{L-1} \sigma_z^j \sigma_z^{j+1} + h \sum_{j=1}^L \sigma_x^j + g \sum_{j=1}^L \sigma_z^j \right) \quad \begin{array}{ll} g = 0 & \text{free fermions} \\ h = 1 & \text{critical} \end{array}$$

near criticality Ising field theory

$$H = \int_{-\infty}^{\infty} dx \left\{ \frac{1}{2\pi} \left[\frac{i}{2} (\psi(x) \partial_x \psi(x) - \bar{\psi}(x) \partial_x \bar{\psi}(x)) - i M_h \bar{\psi}(x) \psi(x) \right] + \bar{g} \sigma(x) \right\}$$

$$M_h = 2J|1 - h|$$

$$M_h = 0 \text{ exactly solvable QFT}$$

$$\bar{g} = \frac{2}{\bar{s}} J^{15/8} g$$

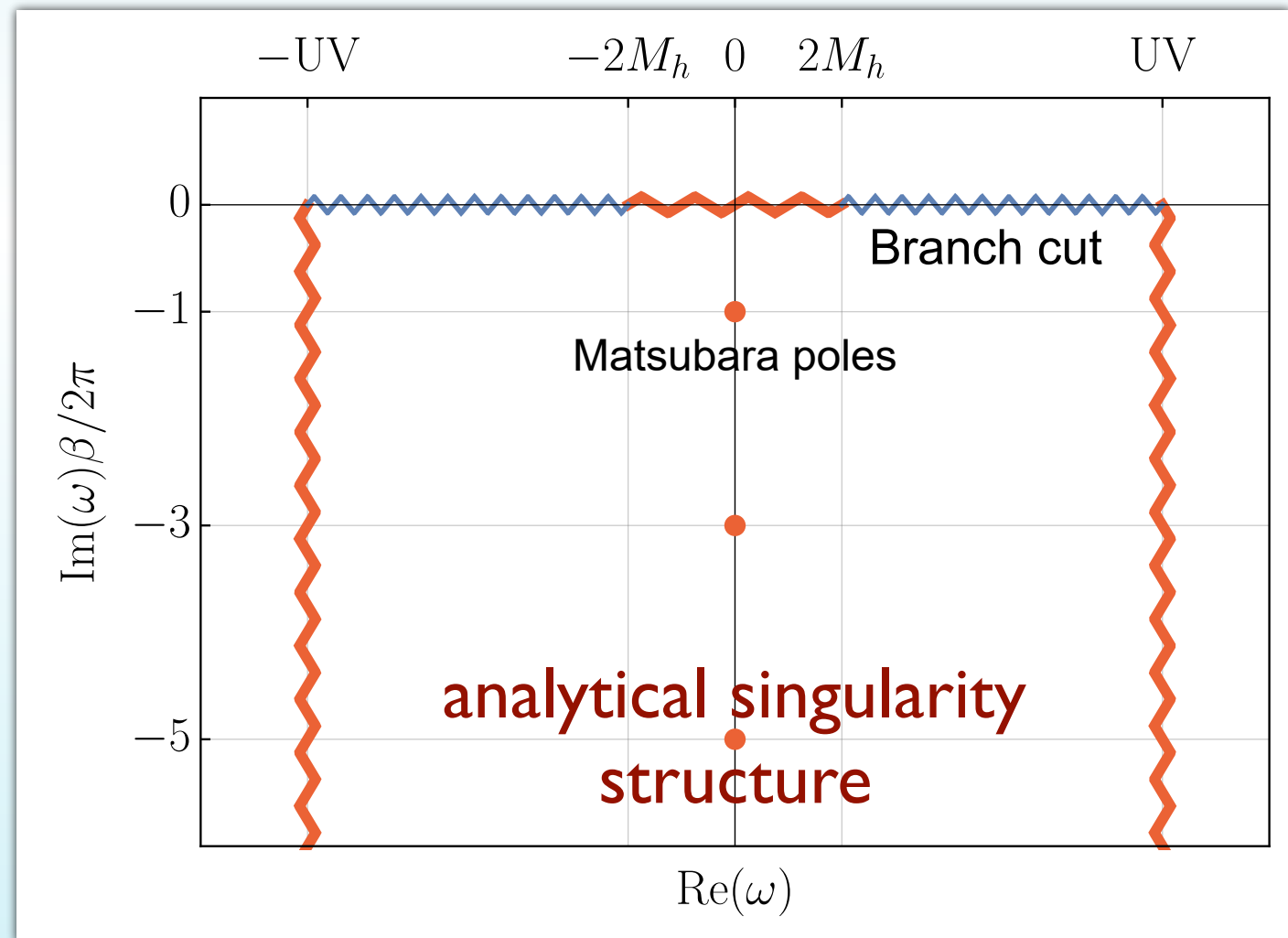
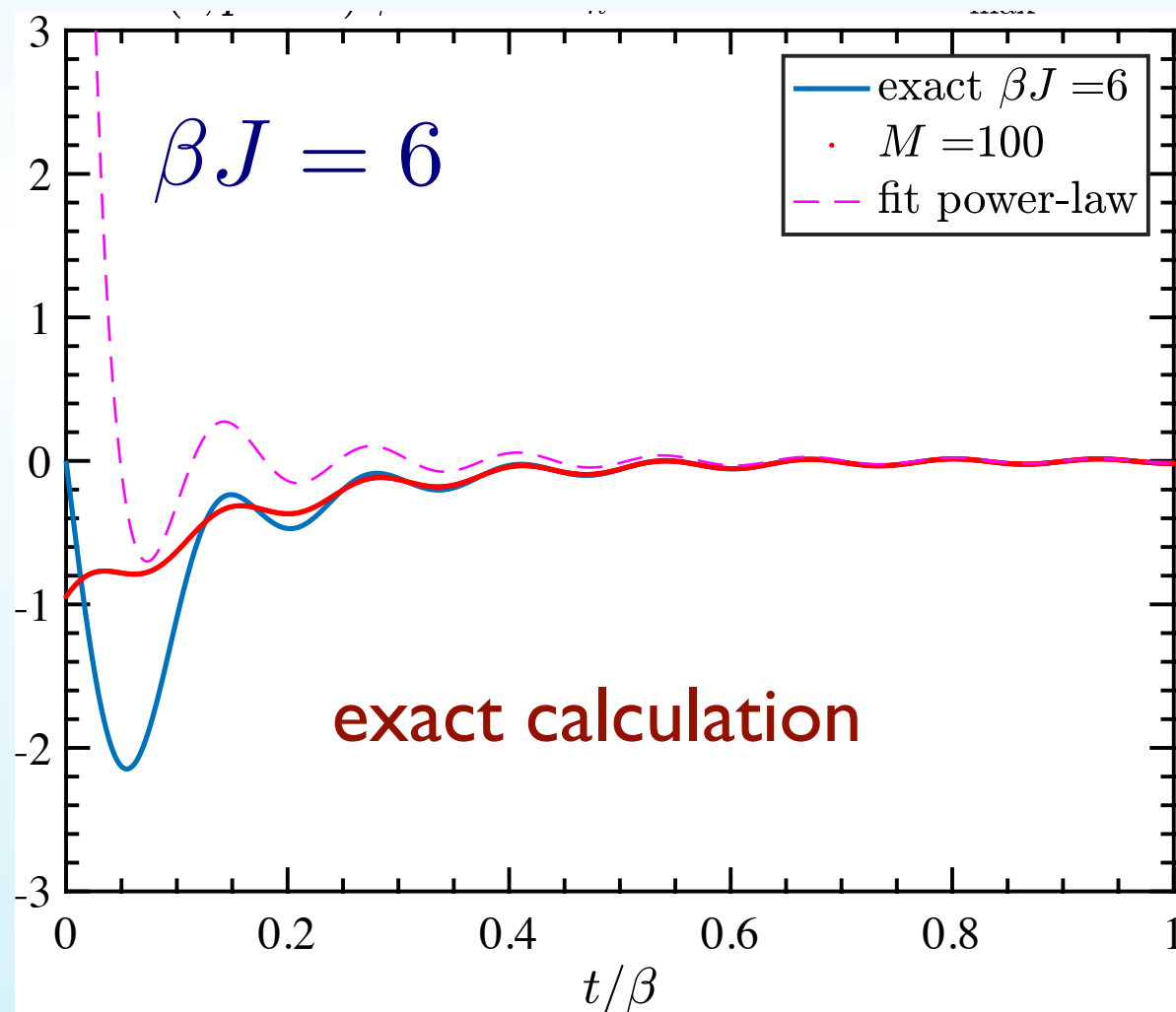
Fonseca, Zamolodchikov 2006

adimensional parameters $(\beta M_h, \beta |\bar{g}|^{8/15}, \beta J)$

THERMAL QUENCHES IN ISING MODELS

VALIDATING THE METHOD: FREE FERMIONIC CASE

transverse magnetization $\sigma_x \longrightarrow \bar{\psi}\psi \quad \Delta_{\bar{\psi}\psi} = 1$



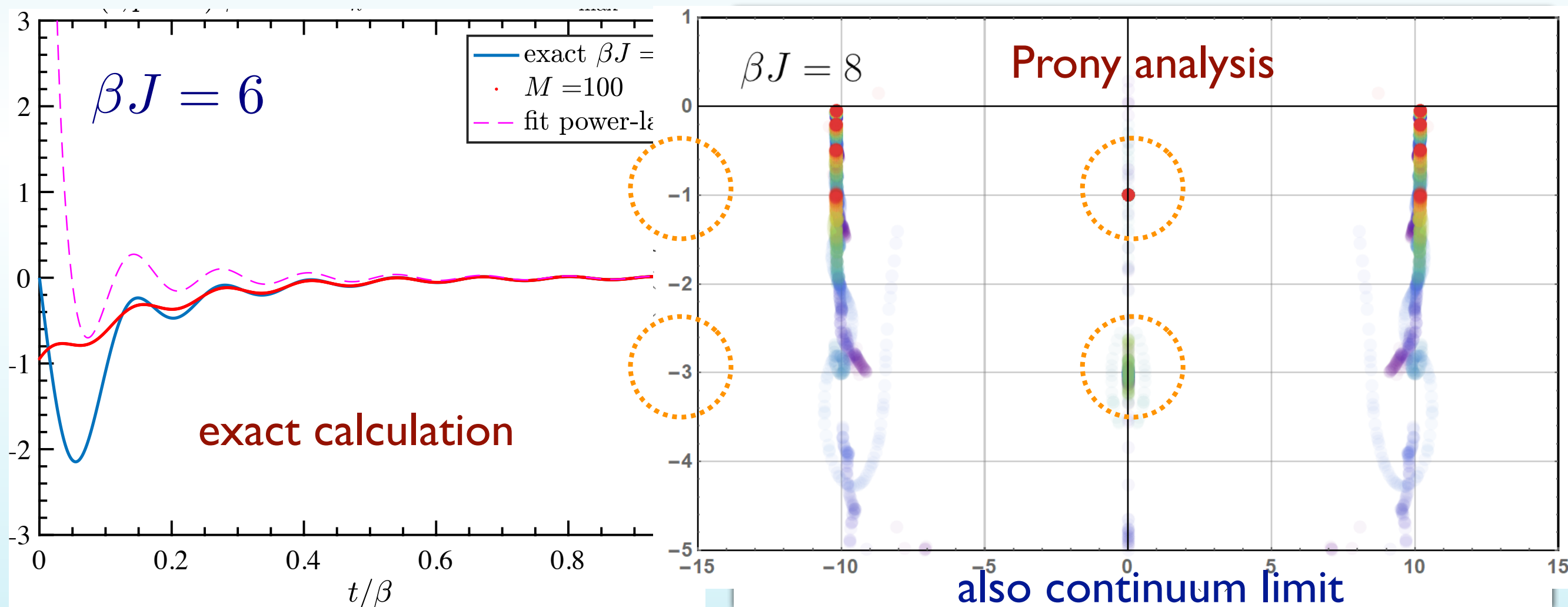
recall CFT $\omega = \pm p - i 2\pi T (\Delta + 2n)$

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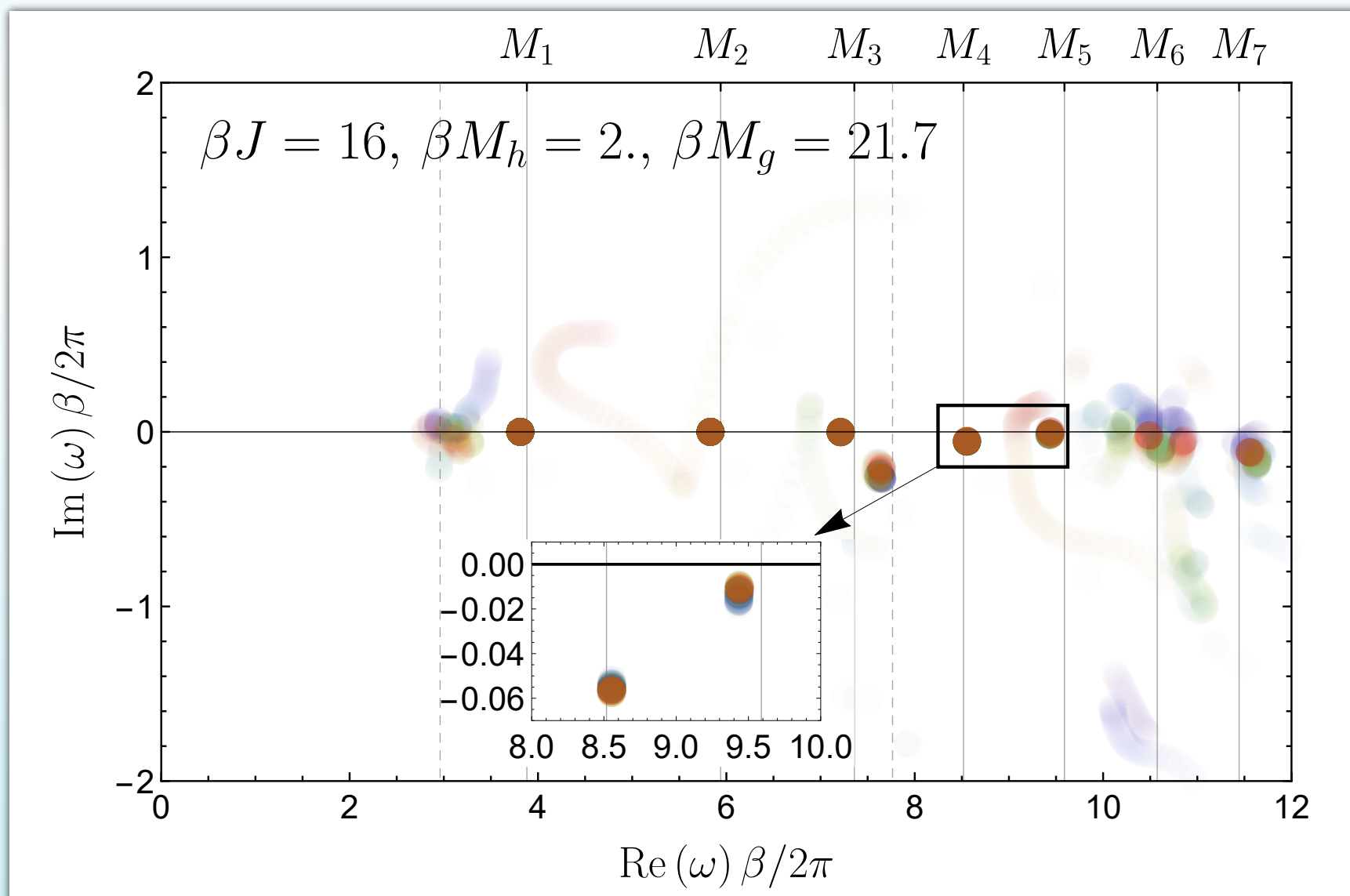
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THERMAL QUENCHES IN ISING MODELS

BEYOND EXACTLY SOLVABLE

transverse magnetization $g \neq 0 \longrightarrow$ bound states (mesons)

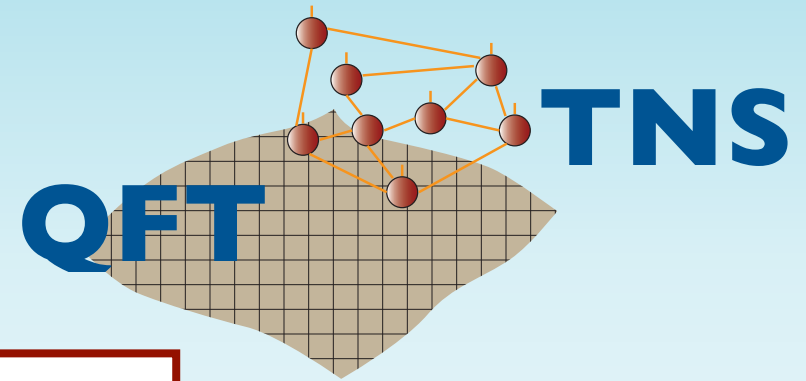


seen as single poles

also imaginary parts!

very good numerical
agreement with
theoretical calculations

THANKS



<http://heptnseminar.org>

TNS = entanglement based ansatz

versatile algorithms: spectrum, thermal equilibrium,
(some) dynamics

applicability for QFT problems

LGT: systematically probed in 1D, progress in 2D

review: MCB, K. Cichy arXiv:1910.00257

suitable for other problems

Thirring model

MCB, K. Cichy, Y.-J. Kao, C.-J. D. Lin, Y.-P. Lin, D.T.-L. Tan, arXiv:1908.04536

thermal response in Ising field theory

MCB, M. Heller, K. Jansen, J. Knaute, V. Svensson, arXiv:1912.08836

