

An introduction to HEPfit and how to constraint EW couplings using it

GenT workshop on Precision
probes of New Physics

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Ana Peñuelas (Instituto de Física Corpuscular, València)

Based on [arXiv:1910.14012](#) and [arXiv:1907.10619](#)



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Outline of the talk

- HEPfit
 - What is it?
 - Why is it a nice tool?
 - Projects using HEPfit
- An application: Global fit and future prospects on electroweak couplings of top and bottom quarks

What is HEPfit?

Usual “definition”

Flexible [open-source](#) C++ code to perform fits in the SM and as many models as possible featuring [flavour](#), [electroweak](#) and [Higgs observables](#)

Open source, **NO** HEPfit collaboration



What is HEPfit?

How does a phenomenologist work? (simplification!)

1. Take your favourite mode

2HDM,
Manohar-Wise

2. Choose one/some observables

Higgs signal
strengths, $\mathcal{R}_{D(*)}$

3. Compare theory and experimental
data

LHC data

4. Write a paper

arXiv:19xx.xxxx

[From O. Eberhardt]

What is HEPfit?

How does a phenomenologist work? (closer to reality)

1. Take your favourite mode
2. **Choose one/some observables**
3. Compare theory and experimental data
4. Write a paper

Observables may depend on **multiple parameters**

2HDM: Yukawa couplings $y_f^{\varphi_i^0}$, mass of the scalars $M_{H,A,H^\pm}$, $\tilde{\alpha}$, $\varsigma_f \dots$

What is HEPfit?

How does a phenomenologist work? (closer to reality)

1. Take your favourite mode
2. Choose one/some observables
3. **Compare theory and experimental data**
4. Write a paper

Experimental observables can be correlated

Correlation between several signal strengths or between $\mathcal{R}_D$ and $\mathcal{R}_{D^*}$

What is HEPfit?

How does a phenomenologist work? (closer to reality)

1. Take your favourite mode
2. Choose one/some observables
3. **Compare theory and experimental data**
4. Write a paper

Theoretical
uncertainties

PDFs,
form factors..

What is HEPfit?

How does a phenomenologist work? (closer to reality)

1. Take your favourite mode
2. Choose one/some observables
3. Compare theory and experimental data
4. **Write a paper**

Problems with
referees

Why is HEPfit a nice tool?

To solve this issues...

1. Parameter scan: Generate random points, evaluate observables and take best points
 - Inefficient
 - You can miss features/solutions (global mode)
 - No statistical meaning
2. Random scans with χ^2 determination
 - Still one can miss the global mode
3. Fitter
 - Statistical framework which fits the model parameters to the experimentally measured observables by maximising the likelihood

Fitters

Frequentist

vs.

Bayesian

Determine the probability
of data for given
parameters $\rightarrow p$ -value

Bayes' theorem

$$P(\mathbf{x}|D) = \frac{P(D|\mathbf{x})P_0(\mathbf{x})}{\int P(D|\mathbf{x})P_0(\mathbf{x})d\mathbf{x}} \rightarrow$$

posterior probability

(Frequentist vs. Bayesian)

Example-unfair coin

We flip an unfair coin that shows tails more than heads..

Frequentist

10 tries $\rightarrow$ probability heads

5/10. Observed 3/10

100 tries $\rightarrow$ probability heads

50/100. Observed 27/100

...

As tries increase, the *observed value* is closer to the *true value*.

We do not incorporate our belief to the fairness of the coin.

Bayesian

10 tries $\rightarrow$ probability heads

5/10. Observed 3/10

100 tries $\rightarrow$ probability heads

30/100. Observed 27/100

...

As tries increase, we modify the *prior belief* that the coin is fair. The *posterior belief* is modified.

Fitters

Frequentist

vs.

Bayesian

Determine the probability
of data for given
parameters $\rightarrow p$ -value

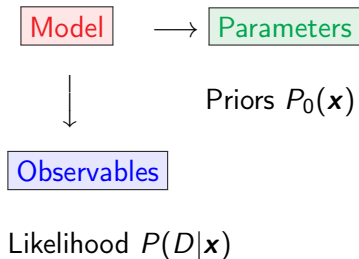
Bayes' theorem

$$P(\mathbf{x}|D) = \frac{P(D|\mathbf{x})P_0(\mathbf{x})}{\int P(D|\mathbf{x})P_0(\mathbf{x})d\mathbf{x}} \rightarrow$$
 posterior probability

CKMfitter, GAMBIT, GAPP, GFitter, HiggsSignals, MasterCode,
UTfit, HEPfit...

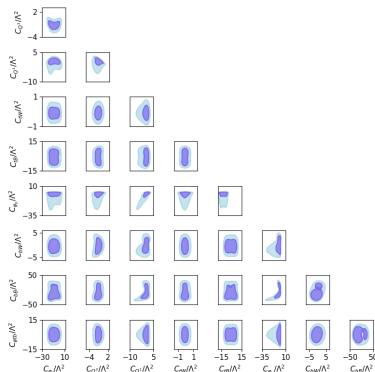


HEPfit

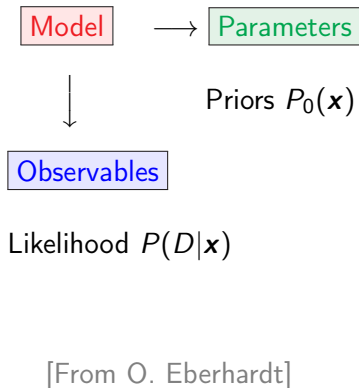


[From O. Eberhardt]

- Posterior distribution of all the parameters and observables



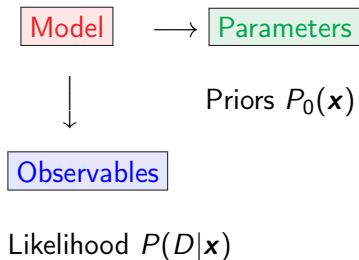
HEPfit



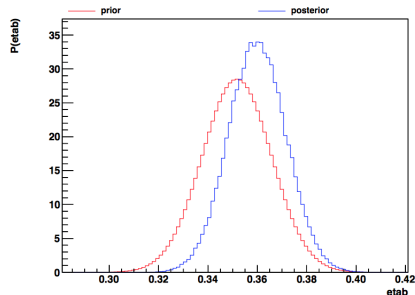
Correlation between parameters



HEPfit



- Comparison of prior and posterior probability



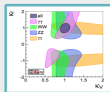
[From O. Eberhardt]

HEPfit

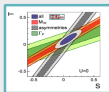


home developers samples documentation

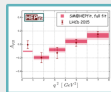
HEPfit: a Code for the Combination of Indirect and Direct Constraints on High Energy Physics Models.



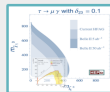
Higgs Physics
HEPfit can be used to study Higgs couplings and analyze data on signal strengths.



Precision Electroweak
Electroweak precision observables are included in HEPfit



Flavour Physics
The Flavour Physics menu in HEPfit includes both quark and lepton flavour dynamics.



BSM Physics
Dynamics beyond the Standard Model can be studied by adding models in HEPfit.

<http://hepfit.roma1.infn.it>



CERN-TH-2019-178 CPHT-RR060.102019
 DESY 19-184 FTUV/19-1031
 IFIC/19-44 KEK-TH-2163
 LPT-Orsay-19-36 PSI-PR-19-22
 UCI-TR-2019-26

HEPfit: a Code for the Combination of Indirect and Direct Constraints on High Energy Physics Models

J. de Blas,^{a,b} D. Chowdhury,^{c,d} M. Ciuchini,^e A. M. Coutinho,^f O. Eberhardt,^g M. Fedele,^h
 E. Franco,ⁱ G. Grilli di Cortona,^j V. Miralles,^g S. Mishima,^k A. Paul,^{l,m} A. Peñuelas,^g
 M. Pierini,ⁿ L. Reina,^o L. Silvestrini,^{i,p} M. Valli,^q R. Watanabe^r and N. Yokozaki^r

^aDipartimento di Fisica e Astronomia “Galileo Galilei”, Università di Padova, Via Marzolo 8, I-35131 Padova, Italy

^bINFN, Sezione di Padova, Via Marzolo 8, I-35131 Padova, Italy

^cCentre de Physique Théorique, CNRS, École Polytechnique, IP Paris, 91128 Palaiseau, France

^dLaboratoire de Physique Théorique (UMR8627), CNRS, Univ. Paris-Sud, Université Paris-Saclay, 91405 Orsay, France

^eINFN, Sezione di Roma Tre, Via della Vasca Navale 84, I-00146 Roma, Italy

^fPaul Scherrer Institut, CH-5232 Villigen PSI, Switzerland

^gIFIC, Universitat de València - CSIC, Apt. Correus 22085, E-46071, València, Spain

^hDept. de Física Quàntica i Astrofísica, Institut de Ciències del Cosmos (ICCUB), Universitat de Barcelona, Martí Franquès 1, E-08028 Barcelona, Spain

ⁱINFN, Sezione di Roma, Piazzale A. Moro 2, I-00185 Roma, Italy

^jInstitute of Theoretical Physics, Faculty of Physics, University of Warsaw, ul. Pasteura 5, PL-02-093 Warsaw, Poland

^kTheory Center, IPNS, KEK, Tsukuba 305-0801, Japan

^lDESY, Notkestraße 85, D-22607 Hamburg, Germany

^mInstitut für Physik, Humboldt-Universität zu Berlin, D-12489 Berlin, Germany

ⁿCERN, Geneva, Switzerland

^oPhysics Department, Florida State University, Tallahassee, FL 32306-4350, USA

^pTheoretical Physics Department, CERN, Geneva, Switzerland

^qDepartment of Physics and Astronomy, University of California, Irvine, CA 92697 USA

^rDepartment of Physics, Tohoku University, Sendai, Miyagi 980-8578, Japan

E-mail: hepfit-support@roma1.infn.it

arXiv:1910.14012

Models and observables

Models implemented in HEPfit:

- Standard Model 18 parameters
- Two-Higgs doublets 7/66 parameters
- Georgi-Machacek model 8 parameters
- Standard Model Effective Field Theory 59 parameters
- Manohar-Wise models 14 parameters
- MSSM with complex couplings 108 parameters
- Left-Right symmetric model 13 parameters
- Other generic SM extensions

Models and observables

Observables implemented in HEPfit:

- Electroweak precision observables
 - STU
 - W mass and width
 - ...
- Higgs physics
 - Strengths
 - Direct searches
- Flavour observables
 - $\Delta F = 2, \Delta F = 1$
 - $B \rightarrow \tau \nu$
 - $B \rightarrow D^{(*)} \ell \nu$
 - ...
- Theoretical constraints
- Electroweak observables
- Model-dependent observables

HEPfit at IFIC

Ongoing projects

"Global fits in the coloured scalar model" O. Eberhardt, V. Miralles, A. Pich

"Lepton Flavour Violation in hadron tau decays and $\ell - \tau$ conversion in nuclei" T. Husek, K. Monsálvez-Pozo, J. Portolés

"Global fits in the A2HDM" O. Eberhardt, A. Peñuelas, A. Pich

Published

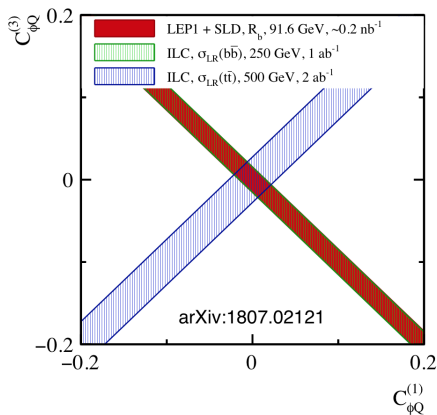
"The electro-weak couplings of the top and bottom quarks - global fit and future prospects" G. Durieux, A. Irles, V. Miralles, A. Peñuelas, R. Pöschl, M. Perelló, M. Vos

The electro-weak couplings of the top and bottom quarks - global fit and future prospects

- NP model predict deviations for the electroweak (EW) couplings of top and bottom quarks [G. Durieux, O. Matsedonskyi, 2019]
- Top quark EW sector poorly constrained
- LHC measurements provide the first constraints on top EW couplings
- Data from LHC, LEP and SLC + EFT $\rightarrow$ current constraints and prospects in future colliders

Theoretical framework

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda^2} \sum_i C_i O_i + \mathcal{O}(\Lambda^{-4})$$



Complementarity between $b\bar{b}$ (LEP) and $t\bar{t}$ (future e^+e^- collider)

$$\delta g_L^t = - (C_{\varphi Q}^1 - C_{\varphi Q}^3) m_t^2 / \Lambda^2$$

$$\delta g_L^b = - (C_{\varphi Q}^1 + C_{\varphi Q}^3) m_t^2 / \Lambda^2$$

Theoretical framework

$$\begin{aligned}
 O_{\varphi Q}^1 &\equiv \frac{y_t^2}{2} \quad \bar{q}\gamma^\mu q \quad \varphi^\dagger i \overleftrightarrow{D}_\mu \varphi, \\
 O_{\varphi Q}^3 &\equiv \frac{y_t^2}{2} \quad \bar{q}\tau^I \gamma^\mu q \quad \varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi, \\
 O_{\varphi u} &\equiv \frac{y_t^2}{2} \quad \bar{u}\gamma^\mu u \quad \varphi^\dagger i \overleftrightarrow{D}_\mu \varphi, \\
 O_{\varphi d} &\equiv \frac{y_t^2}{2} \quad \bar{d}\gamma^\mu d \quad \varphi^\dagger i \overleftrightarrow{D}_\mu \varphi,
 \end{aligned}$$

→ left and right-handed couplings of the top and bottom quark to the Z

$$O_{\varphi ud} \equiv \frac{y_t^2}{2} \quad \bar{u}\gamma^\mu d \quad \varphi^T \epsilon i D_\mu \varphi,$$

→ charged current interaction

$$\begin{aligned}
 O_{uW} &\equiv y_t g_W \quad \bar{q}\tau^I \sigma^{\mu\nu} u \quad \epsilon \varphi^* W_{\mu\nu}^I, \\
 O_{dW} &\equiv y_t g_W \quad \bar{q}\tau^I \sigma^{\mu\nu} d \quad \epsilon \varphi^* W_{\mu\nu}^I, \\
 O_{uB} &\equiv y_t g_Y \quad \bar{q}\sigma^{\mu\nu} u \quad \epsilon \varphi^* B_{\mu\nu}, \\
 O_{dB} &\equiv y_t g_Y \quad \bar{q}\sigma^{\mu\nu} d \quad \epsilon \varphi^* B_{\mu\nu},
 \end{aligned}$$

→ EW dipole operators

$$O_{u\varphi} \equiv \bar{q}u \quad \epsilon \varphi^* \varphi^\dagger \varphi,$$

→ top yukawa

LHC and LEP data

Process	Observable	$\sqrt{s}$	Experiment
$pp \rightarrow t\bar{t}H$	cross section	13 TeV	ATLAS 2016
$pp \rightarrow t\bar{t}Z/W$	cross section	13 TeV	ATLAS 2019
$pp \rightarrow t\bar{t}\gamma$	fid. x-sec.	13 TeV	ATLAS 2018
single-top (t-ch)	cross section	13 TeV	CMS 2018
single-top (Wt)	cross section	13 TeV	CMS 2018
single-top (tZq)	cross section	13 TeV	CMS 2018
$t \rightarrow W^+ b$	F_0, F_L	8 TeV	ATLAS 2016
$e^-e^+ \rightarrow b\bar{b}$	R_b, A_{FBLR}^{bb}	~ 91 GeV	LEP, SLC

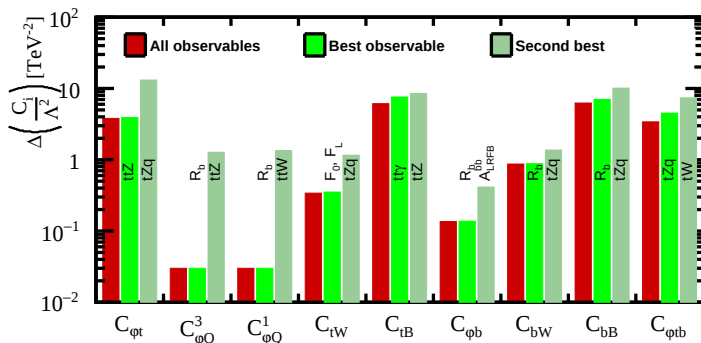
Observables
with
tree-level
dependence
on the dim-6
operators

EWPO and flavour observables not included



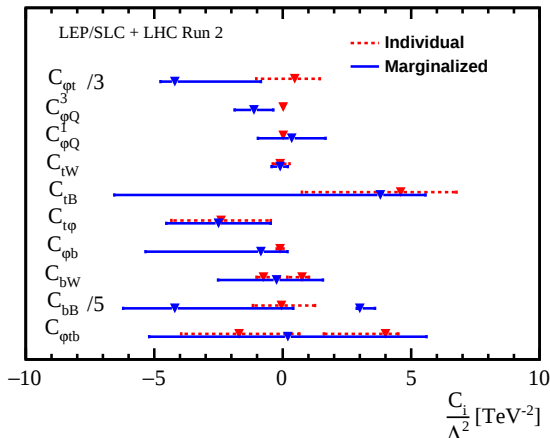
Results - LHC and LEP/SLC data

68% probability for each Wilson coefficient



- Good interplay between the parameters and chosen observables
- LHC: good observables for a LHC EW top fit

Results - LHC and LEP/SLC data



- Robust fit with all operators
- Global limits weaker for some operators (large observable hierarchy)

Results - Impact of Λ^{-4} terms

	Λ^{-2} and Λ^{-4} terms	Λ^{-2} term only
$C_{\varphi t}/\Lambda^2$	$(-16, -2.4)$	$(-2.1, +4.5)$
$C_{\varphi Q}^3/\Lambda^2$	$(-1.9, -0.4)$	$(-0.7, +0.5)$
$C_{\varphi Q}^1/\Lambda^2$	$(-1, +1.7)$	$(-0.6, +0.7)$
C_{tW}/Λ^2	$(-0.4, +0.2)$	$(-0.42, +0.24)$
C_{tB}/Λ^2	$(-6.8, +5.6)$	$(-9.6, +38.4)$
$C_{t\varphi}/\Lambda^2$	$(-4.6, -0.4)$	$(-4.42, 0)$
$C_{\varphi b}/\Lambda^2$	$(-5.4, +0.2)$	$(-0.6, +0.2)$
C_{bW}/Λ^2	$(-2.6, +2.1)$	—
C_{bB}/Λ^2	$(-31.2, +2.4), (+14.4, +18)$	—
$C_{\varphi tb}/\Lambda^2$	$(-5.2, 5.6)$	—

- Λ^{-4} terms $\rightarrow$ two minima for C_{bB}
- C_{tB} : Λ^{-2} term suppressed in $t\bar{t}X$

Results - Impact of Λ^{-4} terms

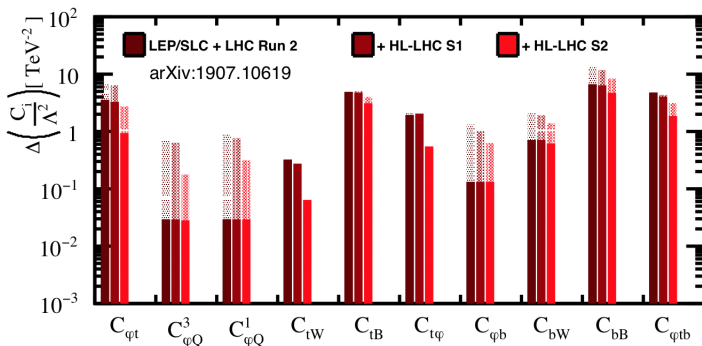
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C_{bW}/Λ^2	(-2.6, +2.1)	—
C_{bB}/Λ^2	(-31.2, +2.4), (+14.4, +18)	—
$C_{\varphi tb}/\Lambda^2$	(-5.2, 5.6)	—

- C_{bW} , C_{bb} , $C_{\varphi tb}$: interference with the SM (Λ^{-2}) vanish if $m_b = 0$ (no limits)
- Λ^{-4} terms $\rightarrow$ validity of EFT expansion carefully!
- Concret BSM scenario: verify dim-8 operators are subdominant comparing to dim-6

Future prospects: HL-LHC

HL-LHC: 300 fb^{-1} , 14 TeV

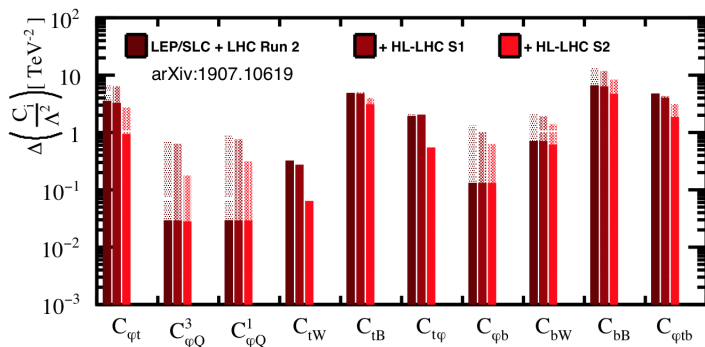
- **S1**: Stat. uncertainty scaled with luminosity
- **S2**: Stat. and syst. uncertainty scaled with (integrated luminosity) $^{1/2}$. Theory uncertainty reduced factor of 2



Solid: individual fit. Dotted: 10-parameter fit

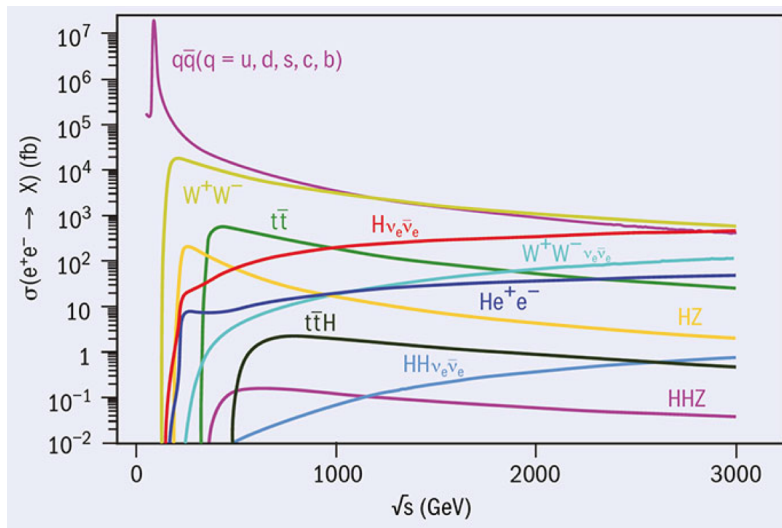
Future prospects: HL-LHC

- **S1** Wilson coefficients only improve marginally $\rightarrow$ conservative assumptions on systematic uncertainties
- **S2** almost all bounds improve
- C_{tb} : limited sensibility to LHC observables



Solid: individual fit. Dotted: 10-parameter fit

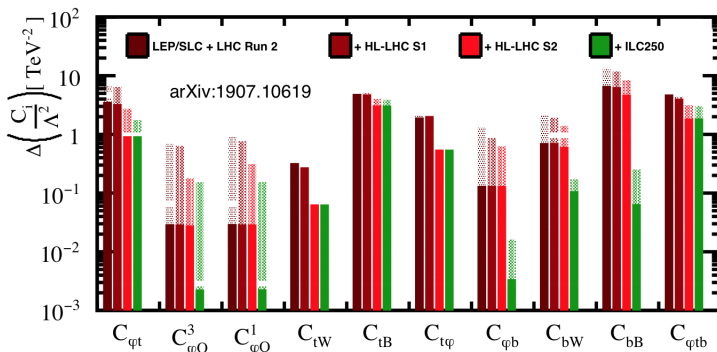
Future prospects: ILC



Future prospects: ILC 250

ILC250: $\sqrt{s} = 250$ GeV

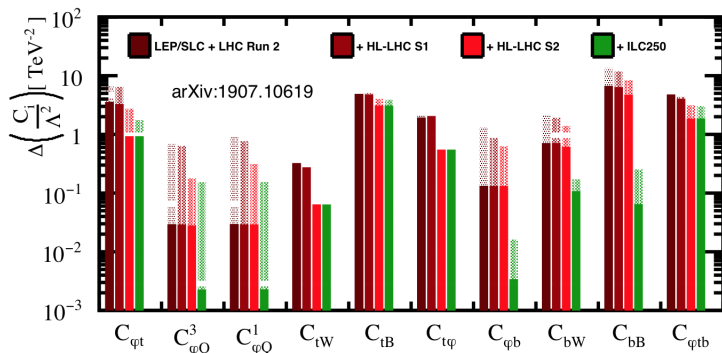
- Including **cross-section** and A_{fb} from $b\bar{b}$ production



Solid: individual fit. Dotted: 10-parameter fit

Future prospects: ILC 250

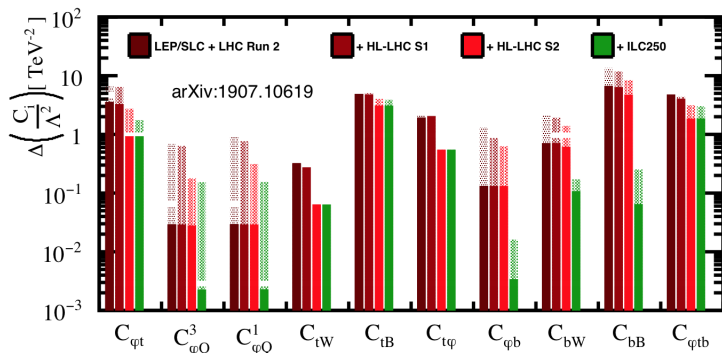
- b -quark operator bounds improve one order of magnitude
- Left-handed couplings barely improve compared to LEP



Solid: individual fit. Dotted: 10-parameter fit

Future prospects: ILC 250

- Right-handed and dipole operators improve significantly
- t -quark operators improve somewhat due to the correlation

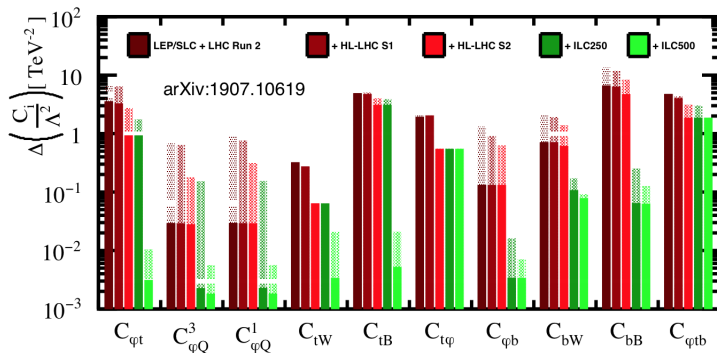


Solid: individual fit. Dotted: 10-parameter fit

Future prospects: ILC 500

ILC500: $\sqrt{s} = 500$ GeV

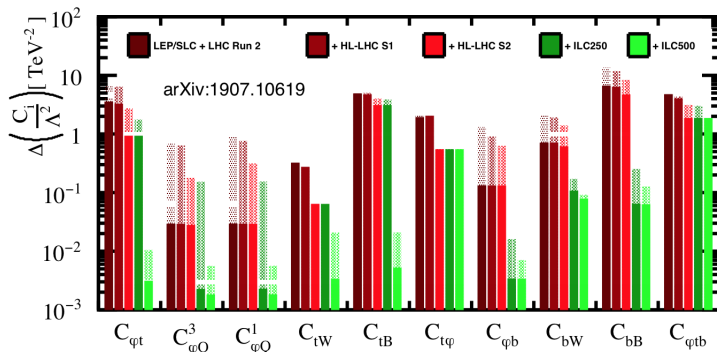
- Including **cross-section** and **A_{fb}** from $b\bar{b}$ production and **optimal $t\bar{t}$ observables**



Solid: individual fit. Dotted: 10-parameter fit

Future prospects: ILC 500

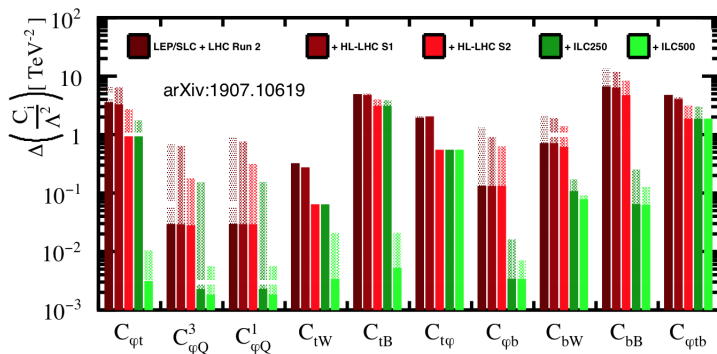
- b-sensitivity similar to ILC250: $b\bar{b}$ production decreases with center of mass energy
- Additional $e^+e^- \rightarrow t\bar{t}$: access to $Z/\gamma t\bar{t}$ vertices



Solid: individual fit. Dotted: 10-parameter fit

Future prospects: ILC 500

- Bounds on top-quark operator reduced 1-2 orders of magnitude
- Bounds for the left-coupling reduced in the global fit



Solid: individual fit. Dotted: 10-parameter fit

Top Yukawa coupling

Why top-quark Yukawa is important?

- $\delta y_t \equiv -\frac{C_{t\varphi} v^2}{\Lambda^2}$
- Larger coupling in the SM
- NP model can shift it (2HDM, supersymmetric models, composite Higgs...)
- $t\bar{t}H$ production provides a direct constraint

Top Yukawa coupling

scenario	LHC Run 2 +LEP/SLC	HL-LHC S2 +LEP/SLC	ILC500
$\sqrt{s}$, $\int \mathcal{L}$	13 TeV, 36 fb ⁻¹	14 TeV, 3 ab ⁻¹	500 GeV, 4 ab ⁻¹
68% probability interval for effective operator coefficient $C_{t\varphi}/\Lambda^2$ [TeV ⁻²]			
individual	[-4.4, +0.0]	[-0.55, +0.55]	[-1.06, +1.06]
marginalized	[-4.6, -0.2]	[-0.55, +0.55]	[-1.07, +1.07]
corresponding relative uncertainty on top-quark Yukawa coupling $\Delta y_t/y_t$ [%]			
individual	13.2	3.3	6.4
marginalized	13.2	3.3	6.4

- Top Yukawa in the global fit is robust but improvable
- Operators affecting $t\bar{t}H$ need to be controlled with other operators

Final remarks

- HEPfit is an useful to calculate different observables in different models
- An application to electro-weak couplings of the top and bottom quark has been shown
 - We obtain constraints from LHC + LEP+SLC data
 - We present prospects for HL-LHC and ILC

Thank you!

Bayesian framework

Bayes theorem (model parameters $\mathbf{x}$, data D):

$$P(\mathbf{x}|D) = \frac{P(D|\mathbf{x})P_0(\mathbf{x})}{\int P(D|\mathbf{x})P_0(\mathbf{x})d\mathbf{x}}$$

$P_0(\mathbf{x}) \rightarrow$ prior distribution of the parameters. Prior knowledge of the parameters (can come from experiments or theory)

$\int P(D|\mathbf{x})P_0(\mathbf{x})d\mathbf{x} \rightarrow$ normalization or evidence

$P(D|\mathbf{x}) \rightarrow$ likelihood

Marginalized distributions (1D)

$$P(x_i|D) = \int P(\mathbf{x}|D) \prod_{j \neq i} dx_j$$



Markov Chain Monte Carlo

- Posterior distribution difficult to calculate when there is a proliferation in model parameters
- Using naive Monte Carlo sampling algorithm can lead to unacceptable execution
- MCMC in BAT
 1. Start at a random point in the parameter space $\mathbf{x}$
 2. Generate a proposal point $\mathbf{y}$ according to a symmetric probability distribution $g(\mathbf{x}, \mathbf{y})$
 3. Compare the value of the function f at proposal point $\mathbf{y}$ with the value at the current point $\mathbf{x}$. The proposal point is accepted if:
 - $f(\mathbf{y}) \geq f(\mathbf{x})$
 - otherwise, generate a random number r from a uniform distribution in the range $[0, 1]$ and accept the proposal if $f(\mathbf{y})/f(\mathbf{x}) > r$. If neither conditions are satisfied the proposal is rejected.
 4. Continue from step 1

MPI implementation

Physics Problem	Hardware	Run Configuration	Time
Unitarity Triangle Fit	3 nodes, 120 CPUs	120 chains, 1.4M iterations	00:02:10
	1 nodes, 40 CPUs	40 chains, 600K iterations	00:00:21
$b \rightarrow s$ decays in SMEFT [†] [?]	6 nodes, 240 CPUs	240 chains, 12.5K iterations	02:05:00
	6 nodes, 240 CPUs	240 chains, 39K iterations	05:20:00
combination of Higgs signal strengths and EWPO[?]	1 node, 16 CPUs	16 chains, 5M iterations	00:14:15
	1 node, 16 CPUs	16 chains, 24M iterations	02:08:00
$D \rightarrow PP$ decays and CP asymmetry[?]	3 nodes, 240 CPUs	240 chains, 4M iterations	00:18:30
	1 node, 8 CPUs	8 chains, 200K iterations	00:00:10

Table: Some representative runs with HEPfit to show the advantages of the MPI implementation. Times are given in DD:HH:MM. The number of iterations refer to the sum total of pre-run and main-run iterations. The number of chains are equal to the number of CPUs by choice. [†]The $b \rightarrow s$ analysis is done with factorized priors, hence the number of iterations should be multiplied by the number of parameters (~ 50) to get a comparative estimate with the other cases. All runs performed in the BIRD or Maxwell clusters at DESY, Hamburg.



Including two-lepton and two-fermion operators

$$\begin{aligned}
O_{lq}^1 &\equiv \frac{1}{2} \bar{q} \gamma_\mu q \quad \bar{l} \gamma^\mu l, \\
O_{lq}^3 &\equiv \frac{1}{2} \bar{q} \tau^I \gamma_\mu q \quad \bar{l} \tau^I \gamma^\mu l, \\
O_{lu} &\equiv \frac{1}{2} \bar{u} \gamma_\mu u \quad \bar{l} \gamma^\mu l, \\
O_{ld} &\equiv \frac{1}{2} \bar{d} \gamma_\mu d \quad \bar{l} \gamma^\mu l, \\
O_{eq} &\equiv \frac{1}{2} \bar{q} \gamma_\mu q \quad \bar{e} \gamma^\mu e, \\
O_{eu} &\equiv \frac{1}{2} \bar{u} \gamma_\mu u \quad \bar{e} \gamma^\mu e, \\
O_{ed} &\equiv \frac{1}{2} \bar{d} \gamma_\mu d \quad \bar{e} \gamma^\mu e
\end{aligned}$$

Including two-lepton and two-fermion operators

	10-parameter fit ILC250 + ILC500	17-parameter fit + ILC1000
$C_{\varphi t}/\Lambda^2$	0.01	0.09
$C_{\varphi Q}^3/\Lambda^2$	0.005	0.04
$C_{\varphi Q}^1/\Lambda^2$	0.005	0.04
C_{tW}/Λ^2	0.02	0.014
C_{tB}/Λ^2	0.02	0.015
$C_{t\varphi}/\Lambda^2$	0.54	0.54
$C_{\varphi b}/\Lambda^2$	0.007	0.008
C_{bW}/Λ^2	0.09	0.17
C_{bB}/Λ^2	0.13	0.17
$C_{\varphi tb}/\Lambda^2$	1.9	1.9
C_{eu}/Λ^2	—	0.0006
C_{ed}/Λ^2	—	0.0005
C_{eq}/Λ^2	—	0.0004
C_{lu}/Λ^2	—	0.0006
C_{ld}/Λ^2	—	0.0009
C_{lq}^-/Λ^2	—	0.0006
C_{lq}^+/Λ^2	—	0.0005