



CSIC



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TOP MASS MEASUREMENT IN RADIATIVE EVENTS

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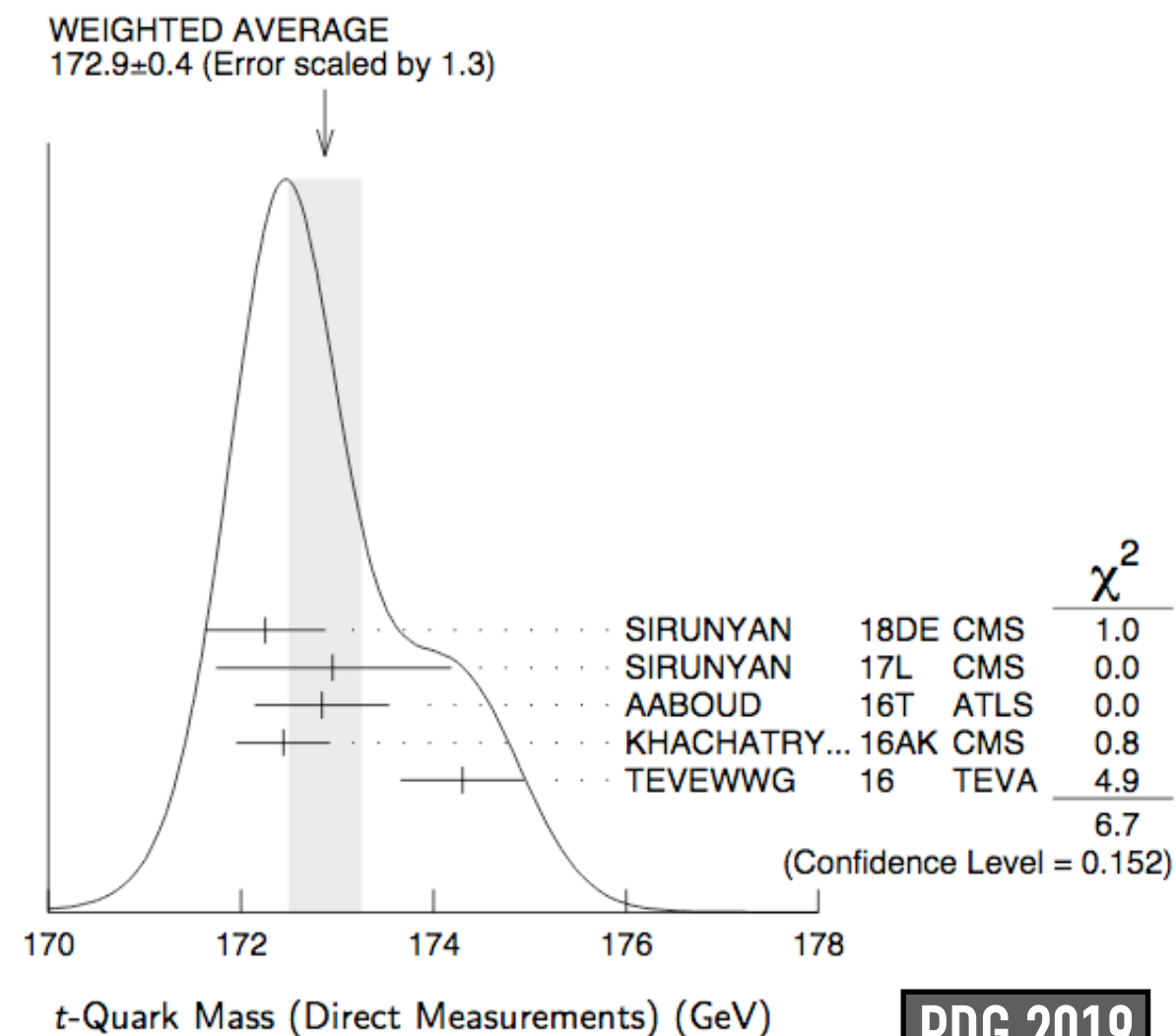
- ▶ Top quark mass
- ▶ Observable definition
- ▶ Theoretical model
- ▶ Experimental strategy
- ▶ Assessment of the uncertainties
- ▶ Summary and the way forward

- ▶ It carries color charge → color confinement → it cannot be measured freely
- ▶ Has a huge mass → tiny lifetime → decays way before hadronization
- ▶ As it does not hadronize, it behaves as a quasi-free quark
- ▶ Its huge mass influence many QCD calculations, and can be related to high precision predictions in perturbative calculations
 - Great for precise tests of the SM
- ▶ As it cannot be spotted freely, its mass has to be measured through its influence on QCD processes
 - Theory dependent measurements → renormalization scheme

► Direct measurements

- Kinematic fit
- MC mass extracted
- Tricky theoretical interpretation

$$m_t = 172.9 \pm 0.4 \text{ GeV}$$



PDG 2019

► Cross section measurements:

- Pole mass
- \overline{MS} mass

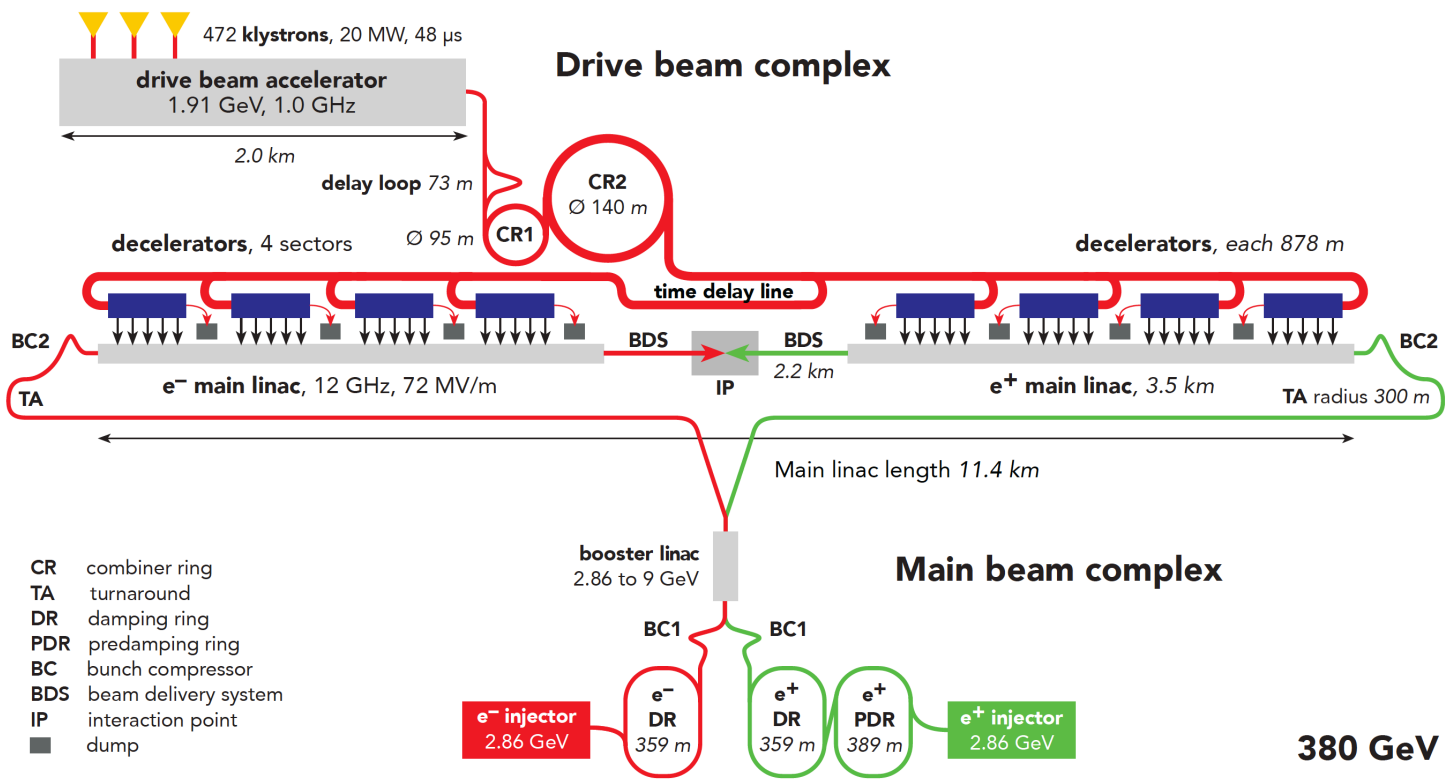
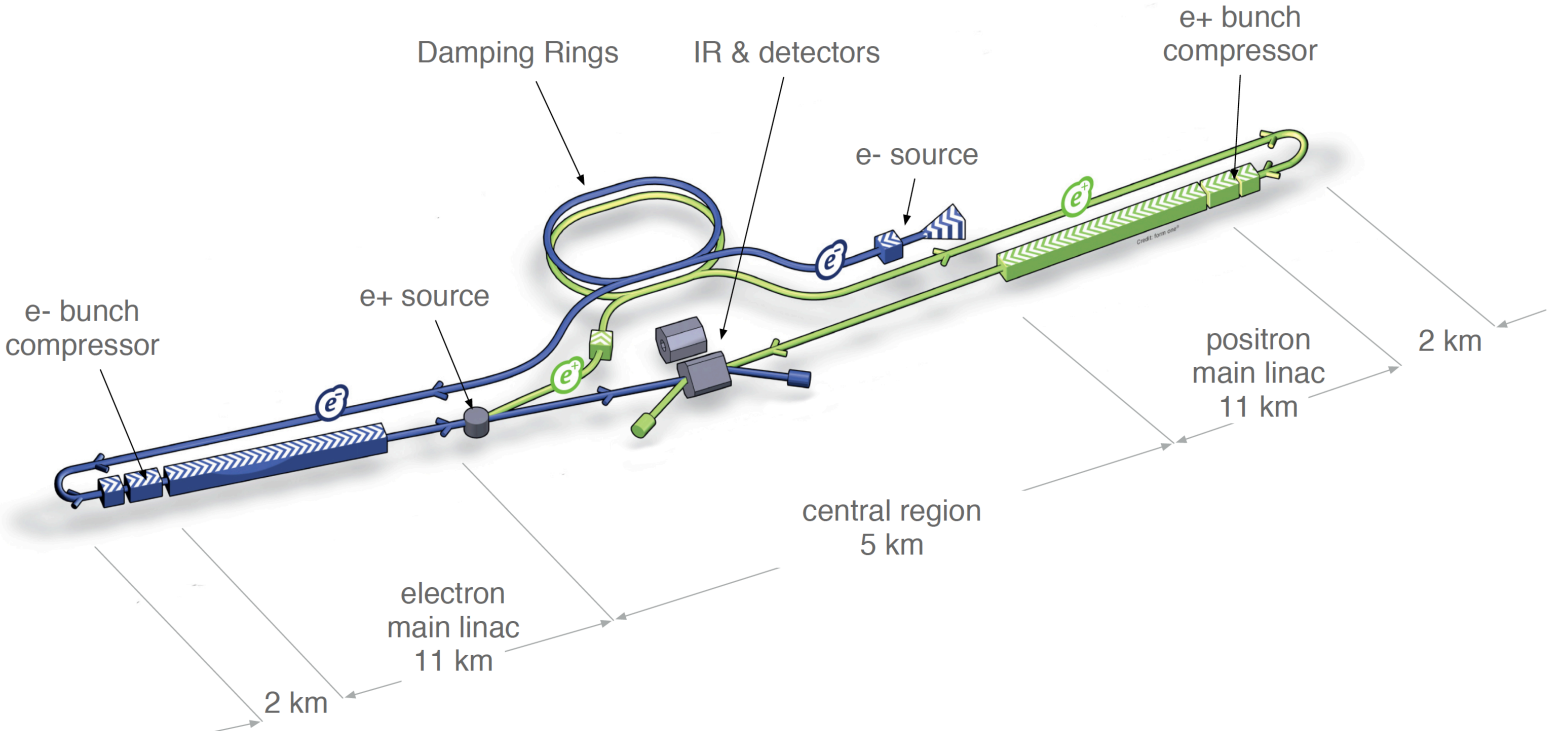
$$m_t^{\text{pole}} = 173.1 \pm 0.9 \text{ GeV}$$

$$\overline{m}_t \equiv \overline{m}_t(\overline{m}_t) = 160.0^{+4.8}_{-4.3} \text{ GeV}$$

Lepton linear colliders offer the following advantages:

- ▶ Knowledge of the initial state e^+e^-
- ▶ Point like interactions \rightarrow known collision energy
- ▶ Lower background environments
 - Reduces the experimental complexity (triggering, clustering...)
 - Reducing the requirements about radiation hardness and occupancy allow for the design of more precise detectors
- ▶ High flexibility to set the energy staging
 - Collision energy tuned by turning on/off sections in the main linac
- ▶ Allows for precision measurements in the TeV scale

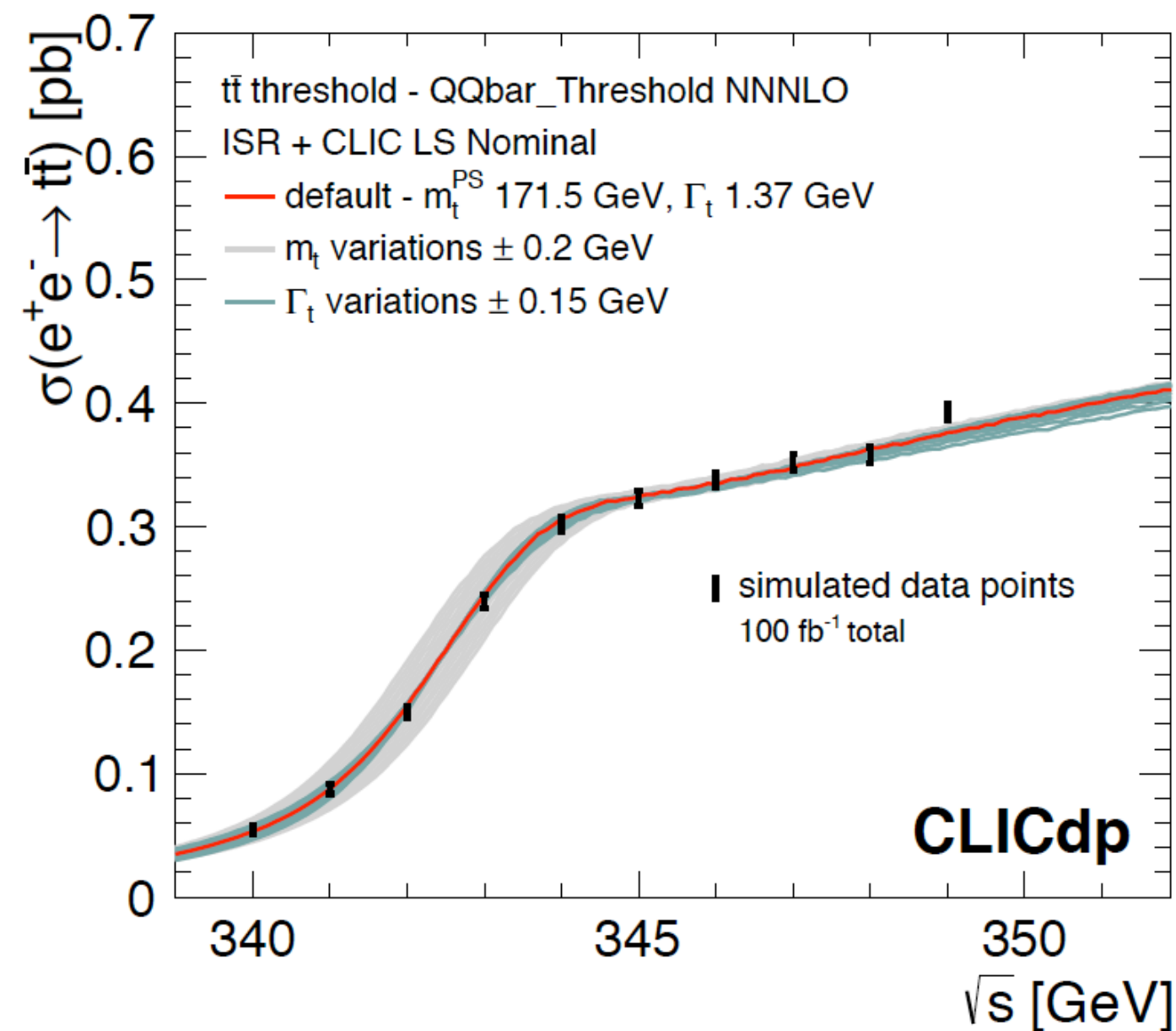
International Linear Collider	
Stage	Integrated luminosity
Initial @ 250 GeV	2000 fb ⁻¹
Upgrade @ 500 GeV	4000 fb ⁻¹



Compact Linear Collider	
Stage	Integrated luminosity
Stage I @ 380 GeV	1000 fb ⁻¹
Stage II @ 1500 GeV	3000 fb ⁻¹
Stage III @ 3000 GeV	5000 fb ⁻¹

380 GeV

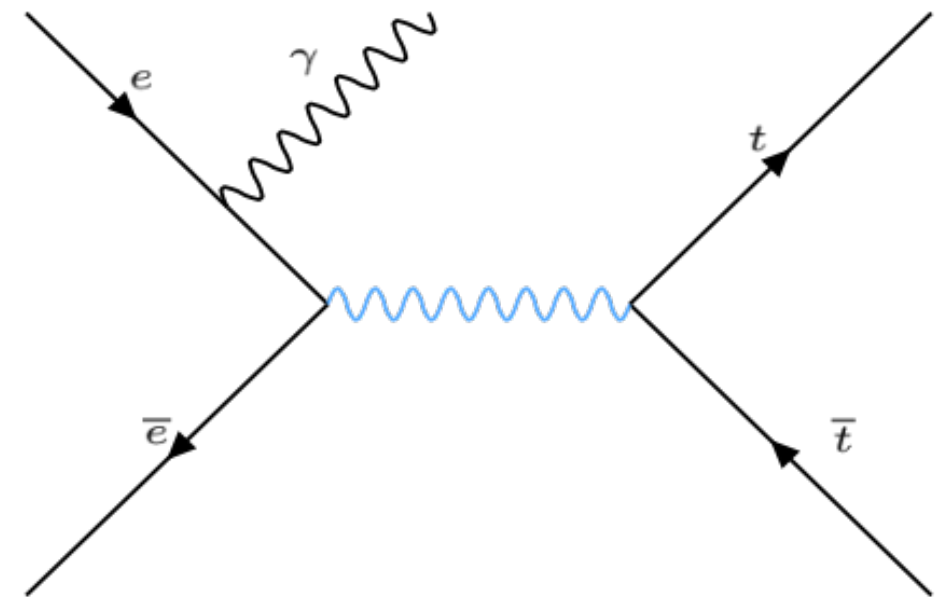
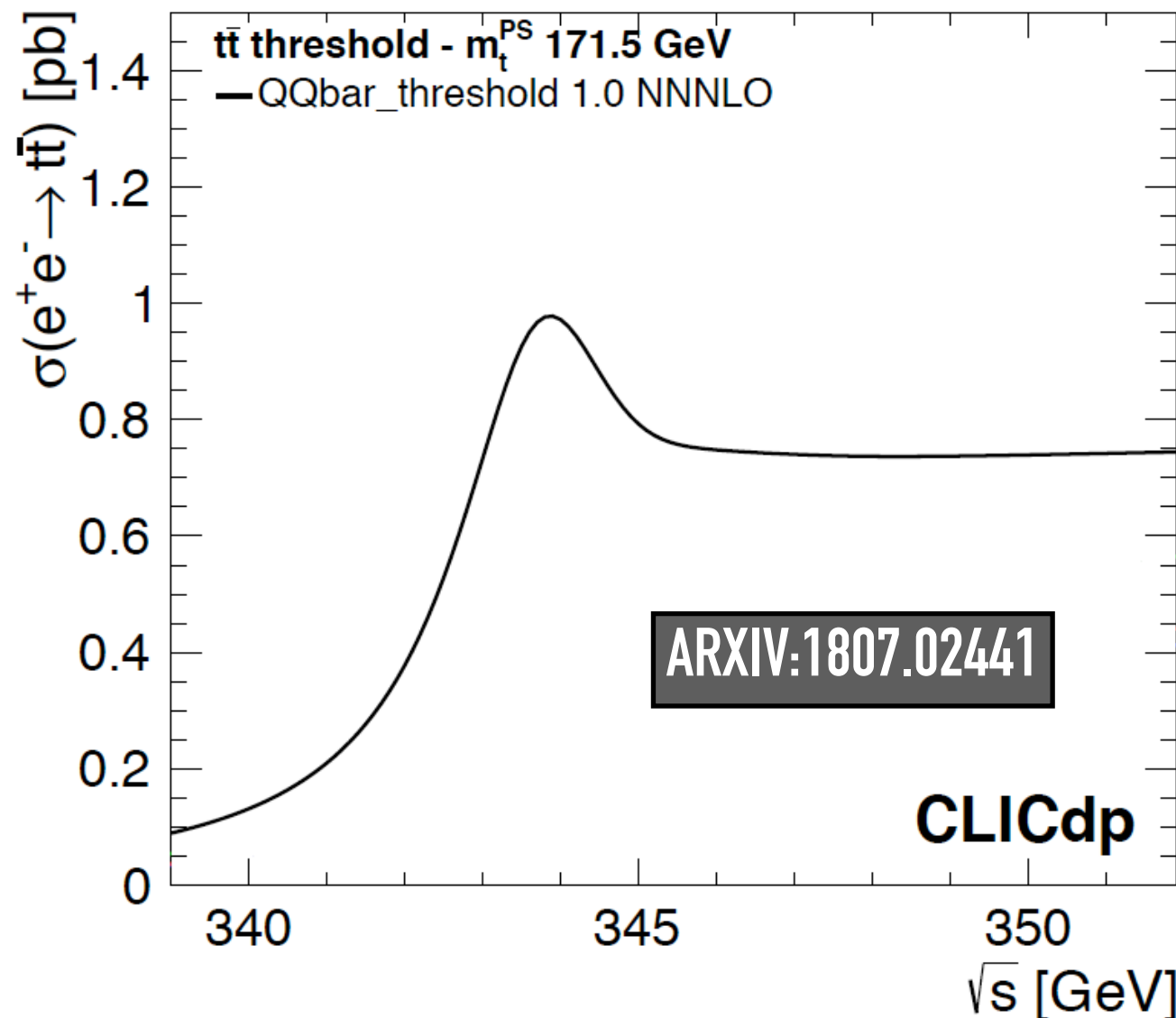
Energy scan of the threshold at steps of 1 GeV, measuring $\sigma_{t\bar{t}}$



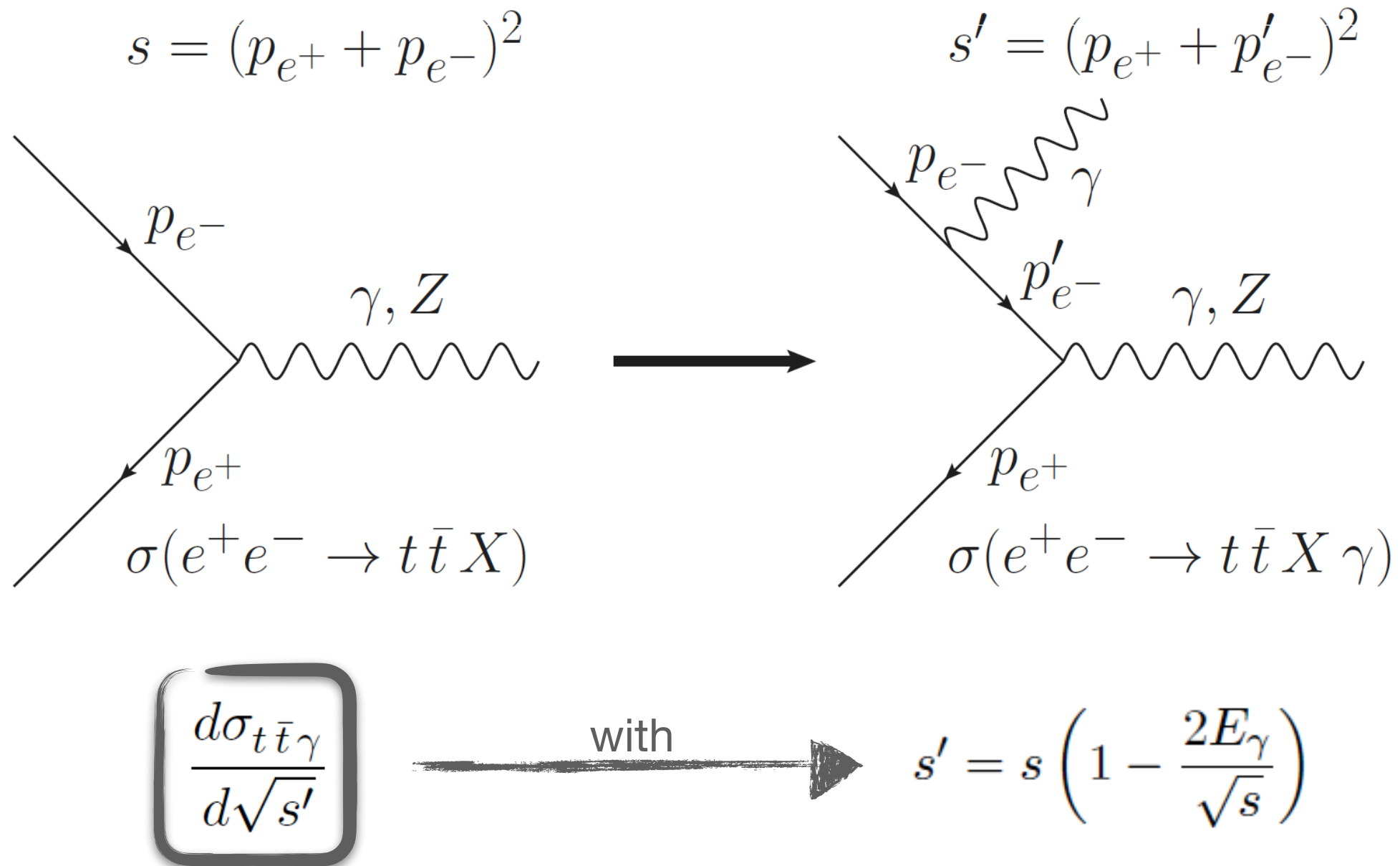
Uncertainties estimation	
Uncertainty	Δm_t^{PS} (MeV)
Statistical @ 100 fb ⁻¹	± 22
Experimental systematics	$\pm 25 \sim 50$
Theoretical systematics	$\pm 30 \sim 50$
Overall	$\pm 45 \sim 75$

ARXIV:1807.02441

- ▶ The idea is to measure the top-quark mass (m_t) in the events $e^+e^- \rightarrow t\bar{t}\gamma_{\text{ISR}}$



- ▶ $t\bar{t}$ pair cross section depends strongly on the production energy
- ▶ The ISR hard photon carries away energy from the $t\bar{t}$ production vertex

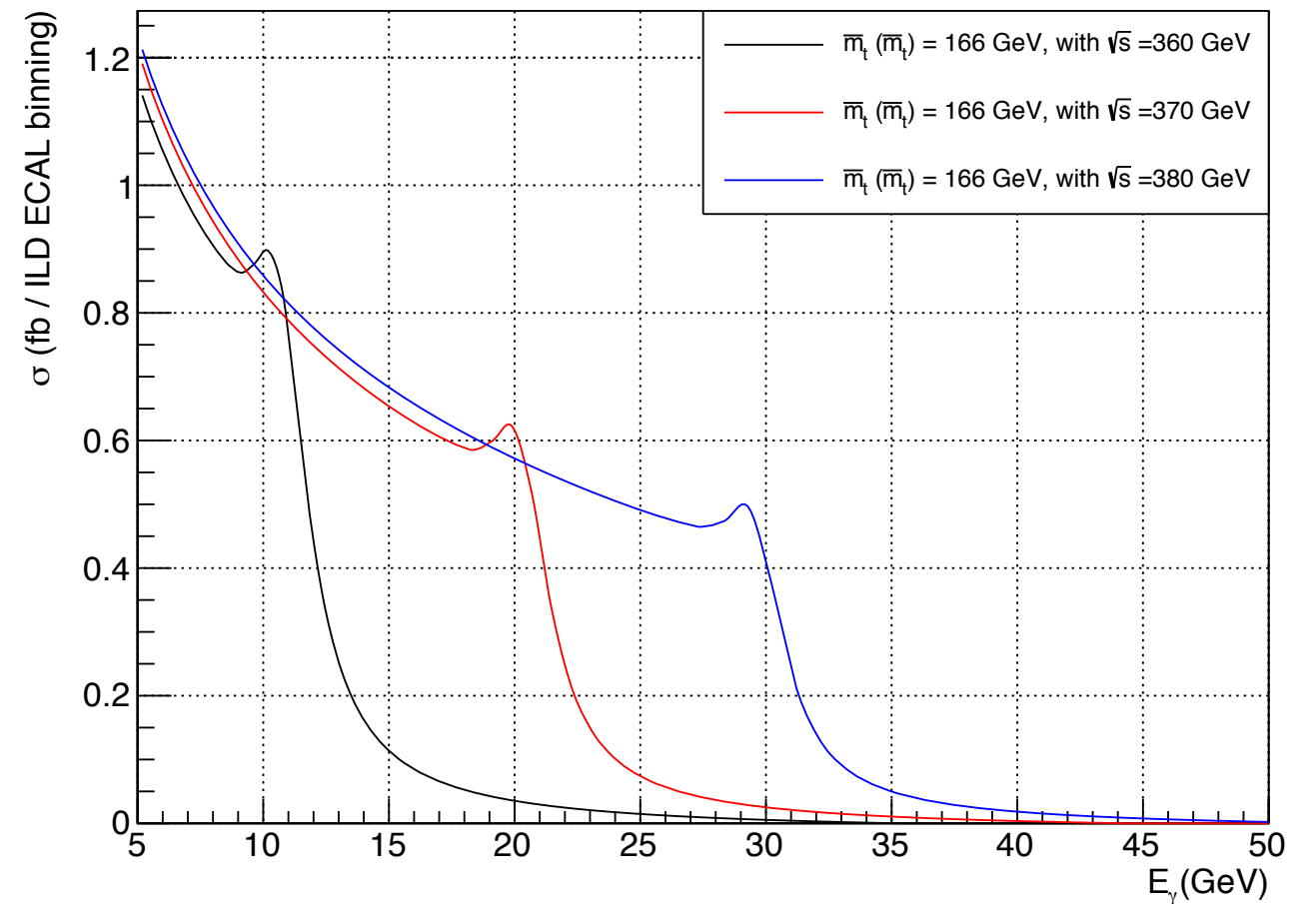


- The observable is defined as the differential cross section of the $t\bar{t}$ production respect to the remaining energy after the ISR emission

- ▶ Factorization theorem by V. Mateu:

$$\sigma_{t\bar{t}\gamma_{ISR}}(\overline{m}_t, s') = \sigma_{ISR}(E_\gamma) * \sigma_{t\bar{t}}(\overline{m}_t, s')$$

- ▶ The model convolutes the ISR emission with the $t\bar{t}$ inclusive cross section.
- ▶ State of the art $t\bar{t}$ calculation, valid in the production threshold and the continuum, by matching:
 - NNLL resummed calculation at the threshold
 - Bound state enhancements included
 - An NNLO calculation in the continuum



- ▶ For more information and details on the $t\bar{t}$ -bar matched calculation:

ANGELIKA WIDL LCWS17 TALK

- ▶ The input mass can be chosen to be any short-distance mass scheme, in this case we chose the \overline{MS} scheme. For the calculation itself the 1S and MSR masses are used.

- ▶ Identify $t\bar{t}$ events with an ISR photon

- ▶ Reconstruct the photon energy

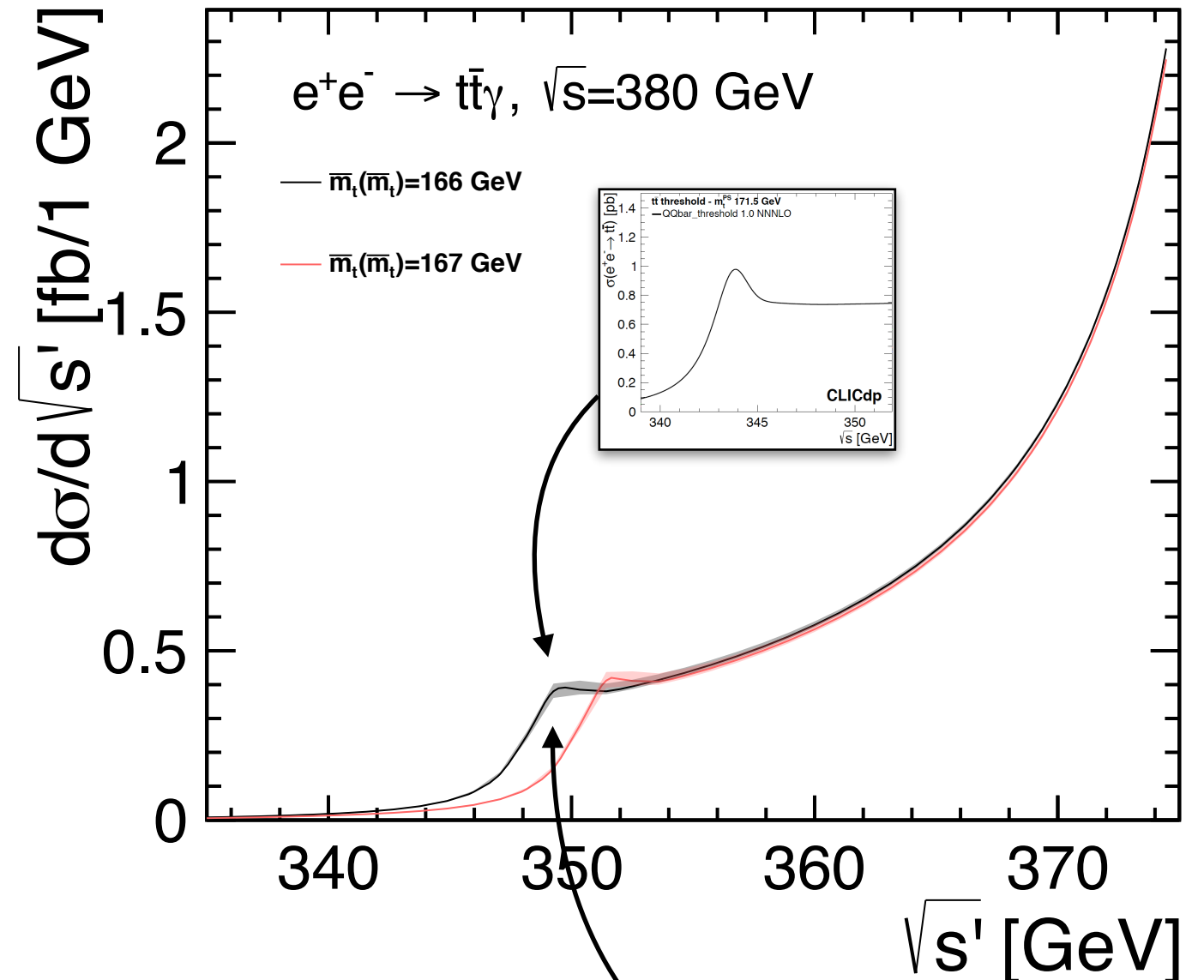
- ▶ Translate the photon energy into remaining center of mass energy

$$s' = s \left(1 - \frac{2E_\gamma}{\sqrt{s}} \right)$$

- ▶ Count the $t\bar{t}$ pairs for a given s'
binning, measuring

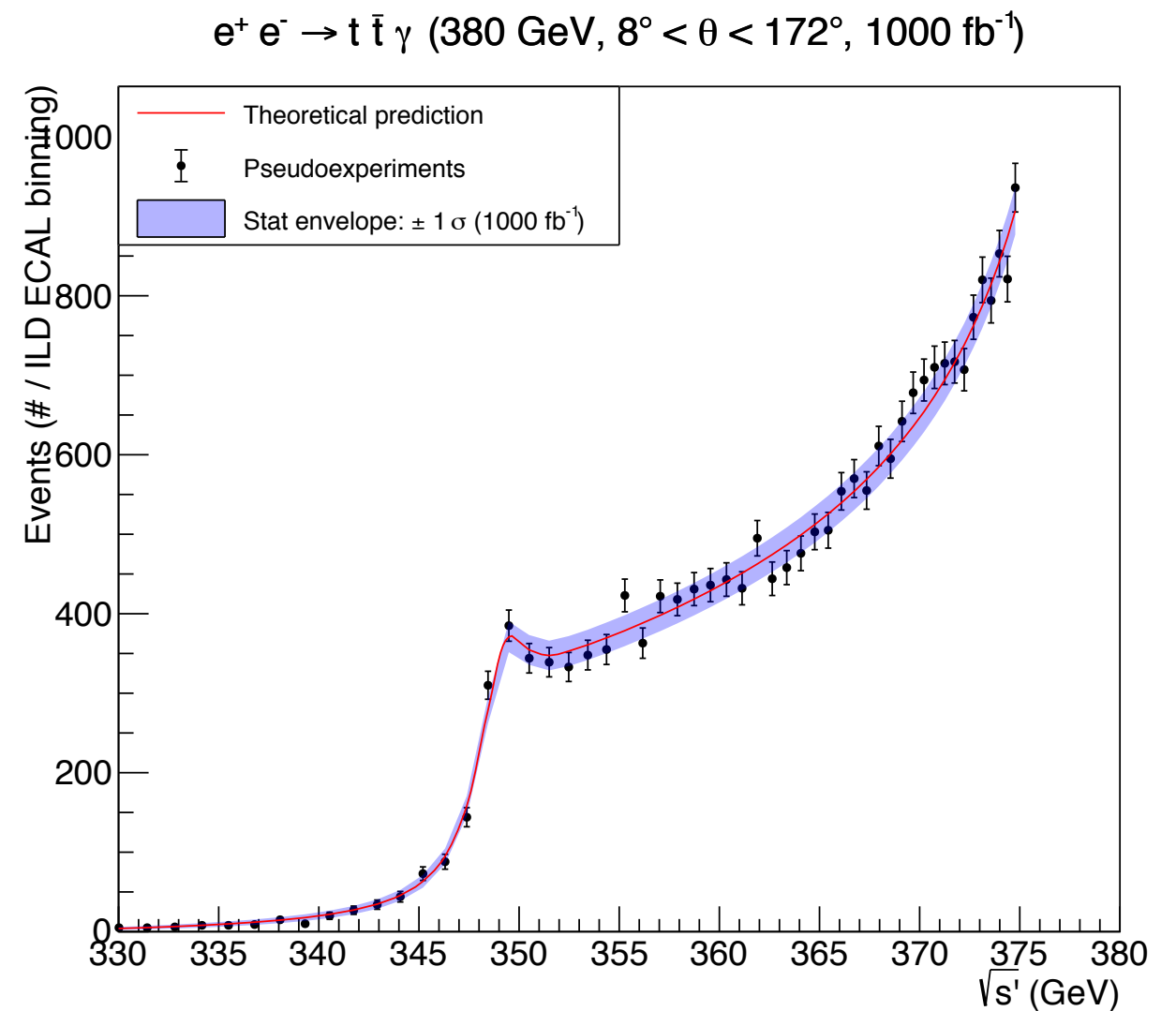
$$\frac{d\sigma_{t\bar{t}\gamma}}{d\sqrt{s'}}$$

- ▶ The maximal sensitivity to the mass is reached at the $t\bar{t}$ production threshold (*radiative return* to the threshold)



$$E_{\gamma,\text{max}} \simeq \frac{s - 4\bar{m}_t^2}{2\sqrt{s}}$$

- ▶ Generation of pseudo-experiments by applying poissonian fluctuations around the expected number of events
- ▶ ~10000 datasets generated
- ▶ Fitting of the pseudo-experiments to the theoretical model with the mass as a free parameter
- ▶ The precision is estimated as the mean of the TMinuit uncertainty estimation
 - Which is in good agreement with the fitted mass distribution

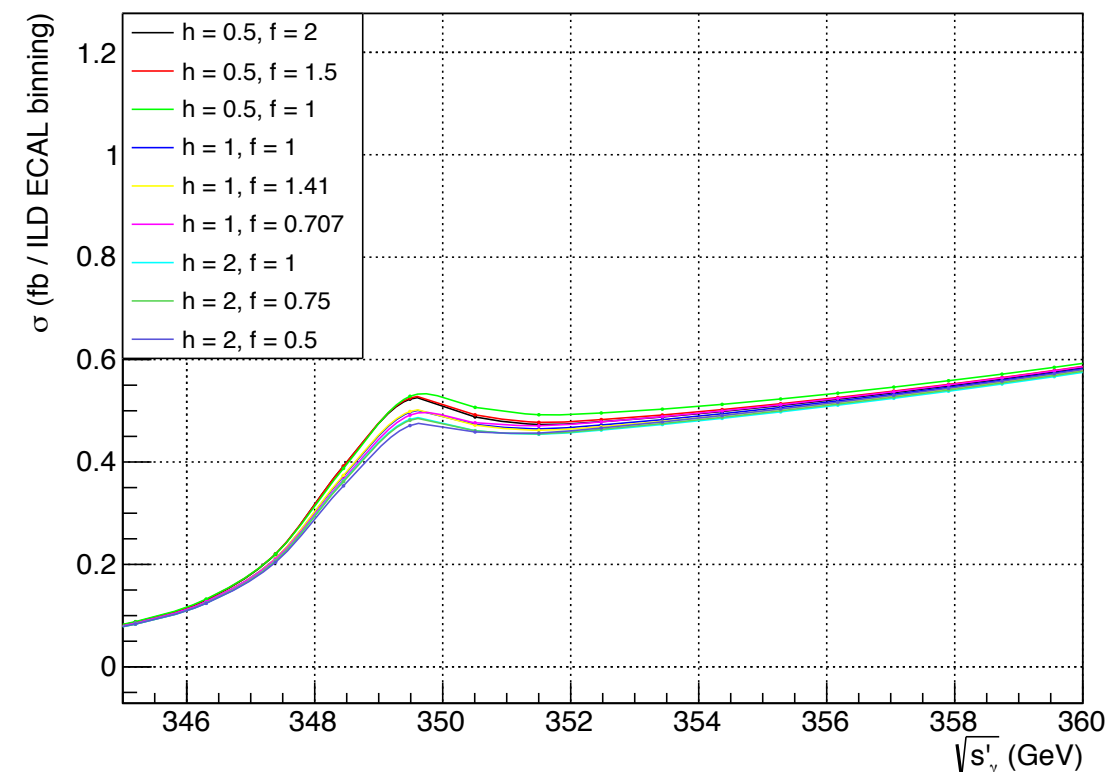


Naive statistical uncertainty estimation		
Nominal energy	Integrated luminosity	Statistical uncertainty
380 GeV	1000 fb^{-1}	41 MeV
500 GeV	4000 fb^{-1}	64 MeV

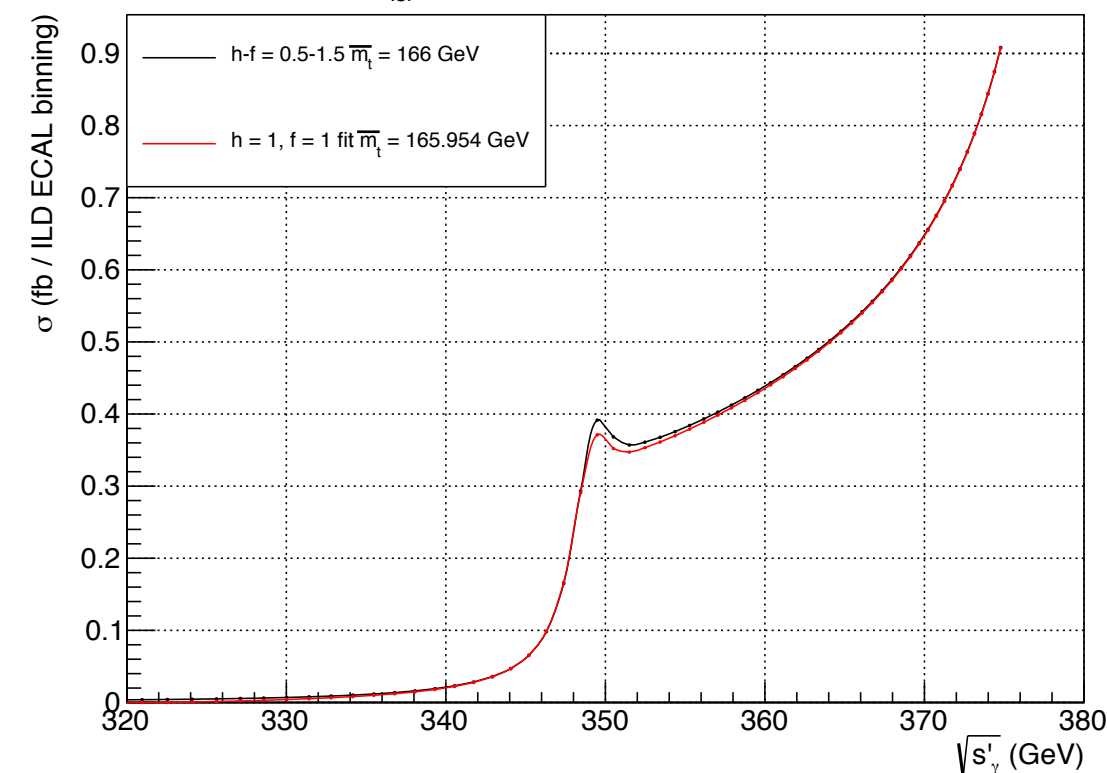
- Assess uncertainty due to missing orders via scale variation in the $t\bar{t}$ matched calculation
- Scales parametrized in h and f parameters:
 - h correspond to the top mass
 - $h f$ to the top momentum
 - $h f^2$ to the kinetic energy of the $t\bar{t}$ pair

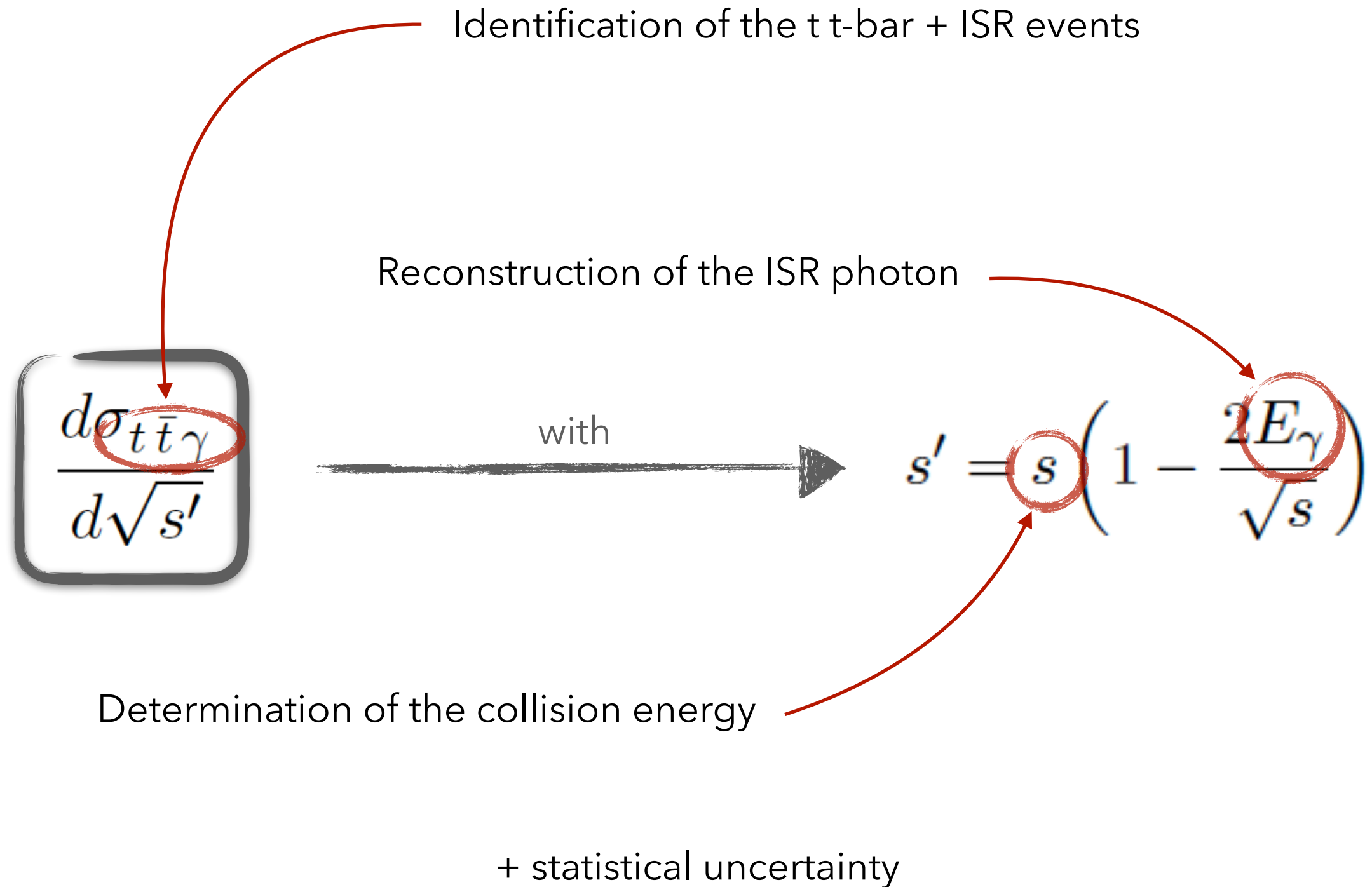
Proposed scale parameters variations (A. Hoang, M. Stahlhofen)									
h	1/2	1/2	1/2	1	1	1	2	2	2
f	2	3/2	1	1	$\sqrt{2}$	$\sqrt{(1/2)}$	1	3/4	1/2
$\Delta\bar{m}_t$ (MeV) @380 GeV	-44	-46	-43	0	-0.3	8	29	30	45
$\Delta\bar{m}_t$ (MeV) @500 GeV	-55	-58	-54	0	-1.5	12	32	34	51

- Scale varied observables are fitted to the model with nominal ($h = f = 1$) scale values, and m_t as a free parameter
- The resulting theory uncertainty is of ± 46 MeV at 380 GeV, and ± 55 MeV at 500 GeV

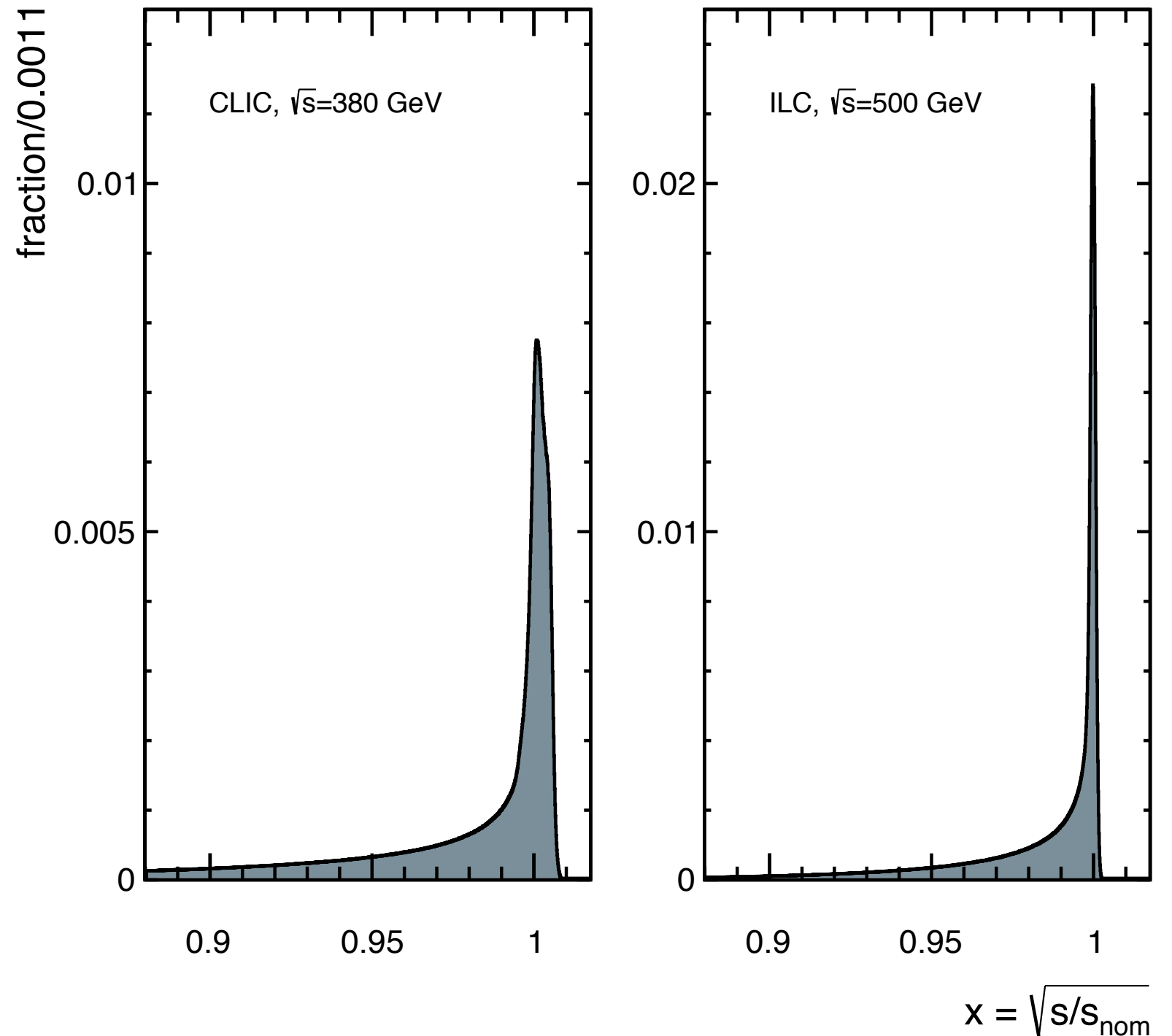


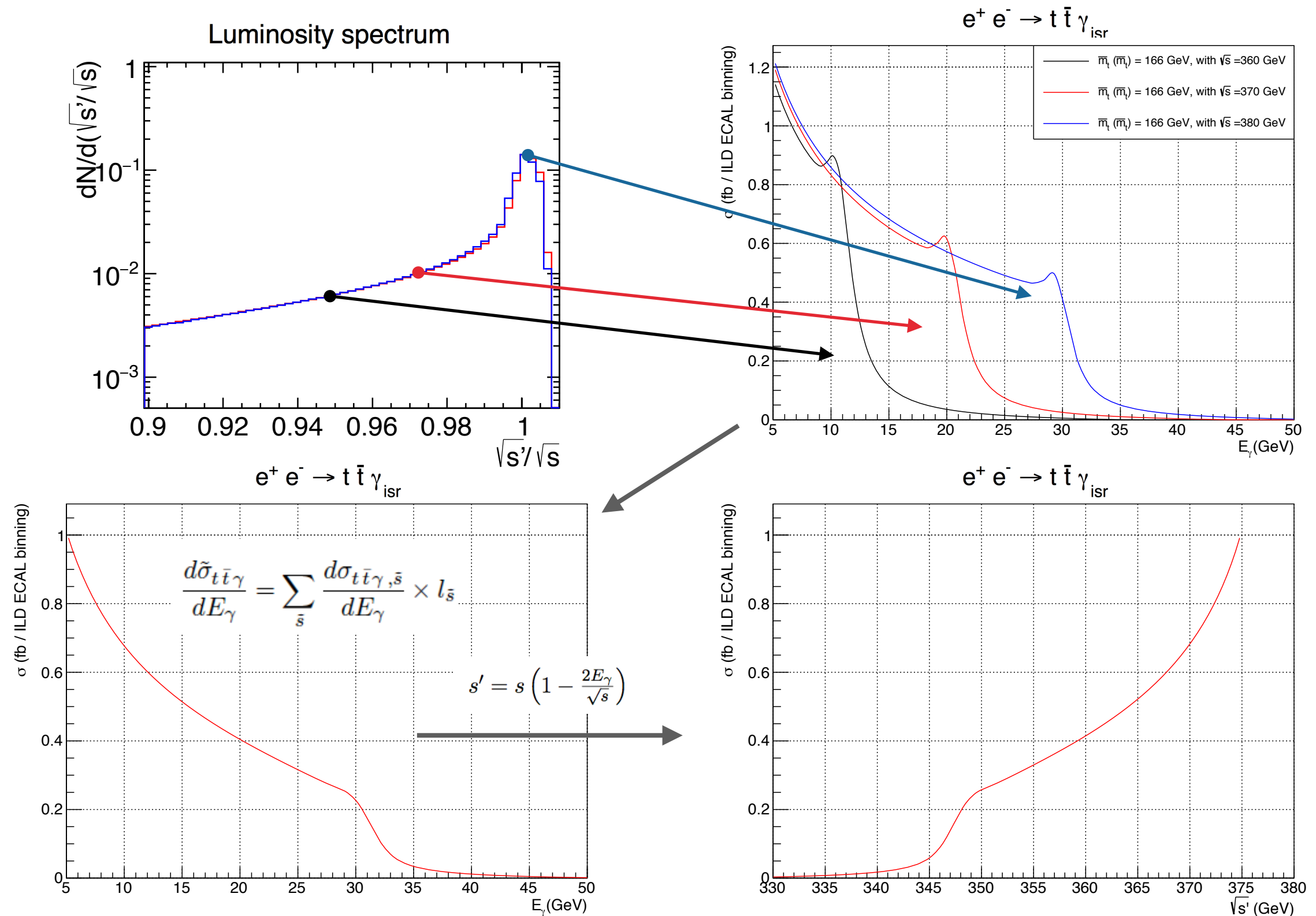
$e^+ e^- \rightarrow t \bar{t} \gamma_{\text{ISR}}$ (matched calc, 380 GeV, $8^\circ < \theta < 172^\circ$)





- ▶ The centre of mass energy does not correspond to a δ centered at \sqrt{s}
- ▶ The collision energy is affected by
 - Beam energy spread
 - Beamstrahlung
- ▶ The actual collision energy is described by the luminosity spectrum

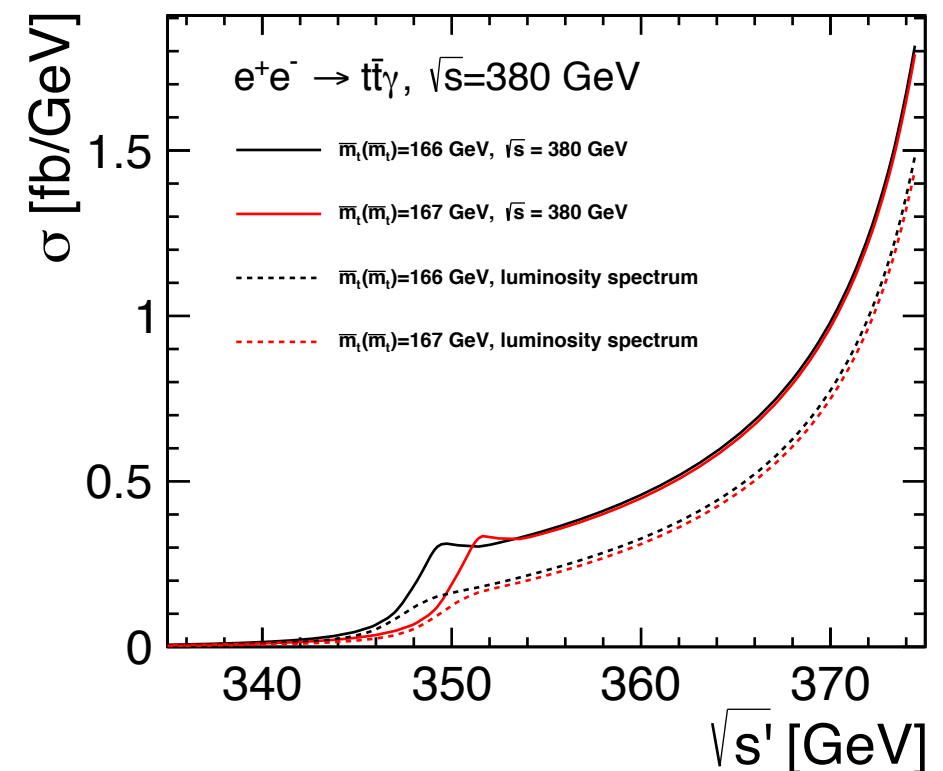




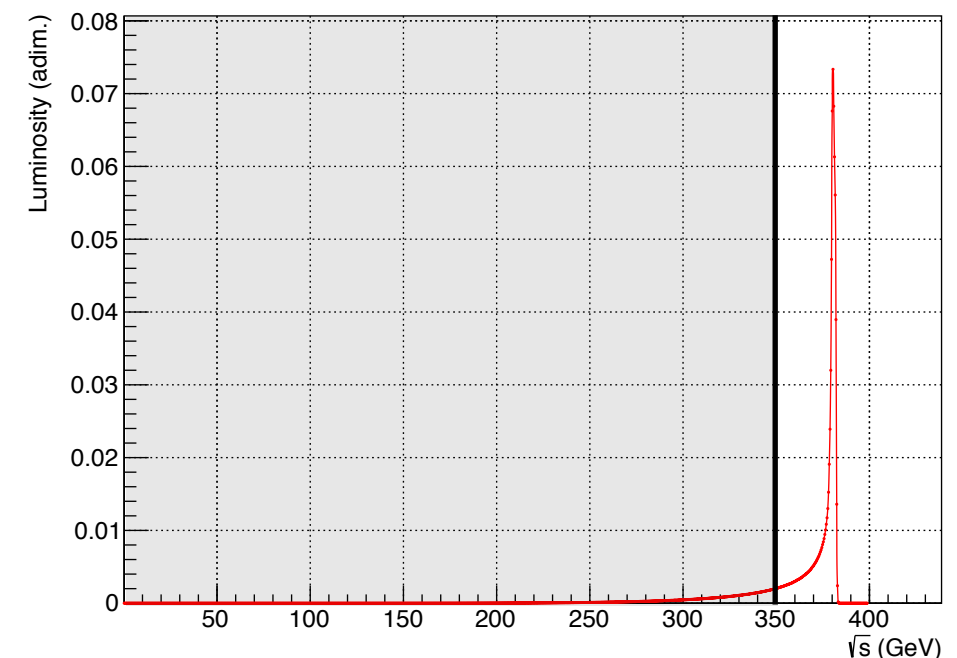
- ▶ When weighting the observable by the luminosity spectrum two things happen:
 - A loss of statistics due to collisions under the $t\bar{t}$ threshold
 - A change in the shape of the observable that reduces its sensitivity to the top mass
- ▶ Repeating the naive statistical estimation, a ~50% deterioration in the sensitivity is found

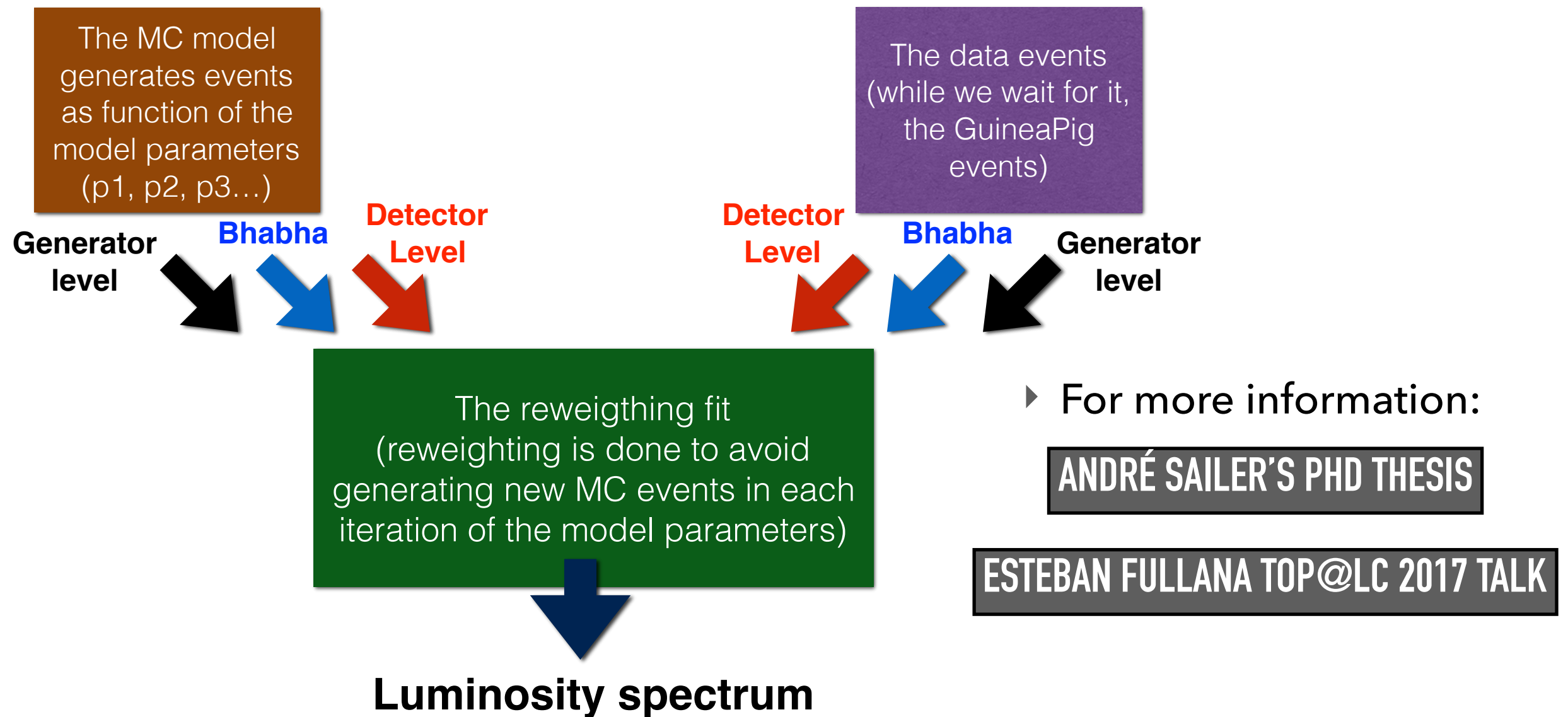
Naive statistical uncertainty estimation	
1000 fb ⁻¹	Statistical uncertainty
380 GeV (δ spectrum)	41 MeV
CLIC @ 380 GeV	65 MeV

- ▶ Work in progress: a method to recover the shape previous to the luminosity spectrum weighting.



Luminosity spectrum





- ▶ The MC model reproduces the 4-momenta of the $e^+ e^-$ pairs.
- ▶ 19 free parameters, optimisation done minimising χ^2 respect to the detector smeared data.
- ▶ 3 stages: GuineaPig beam generator, Bhabha scattering simulator and detector simulator.

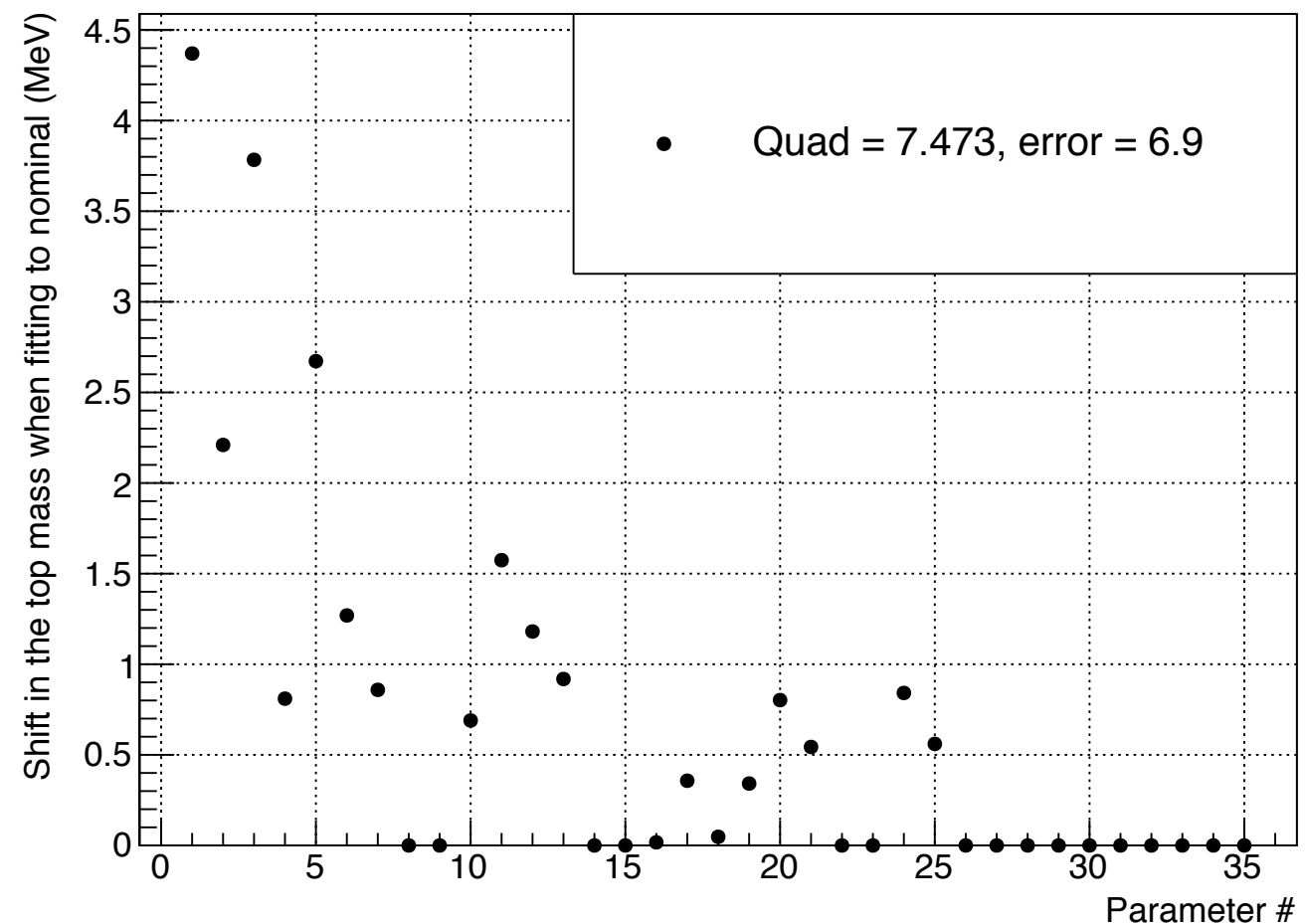
- ▶ Reconstructing the luminosity spectrum, 19 parameters with their error are obtained (+correlations between them)
- ▶ Generating 38 observables within the parameter \pm error and fitting them to the *nominal* weighted observable, the error of each parameter is propagated

- ▶ The total uncertainty is calculated by performing

$$E = E_p \text{Cov } E_p^T.$$

- ▶ A total propagated uncertainty of ± 7 MeV is estimated

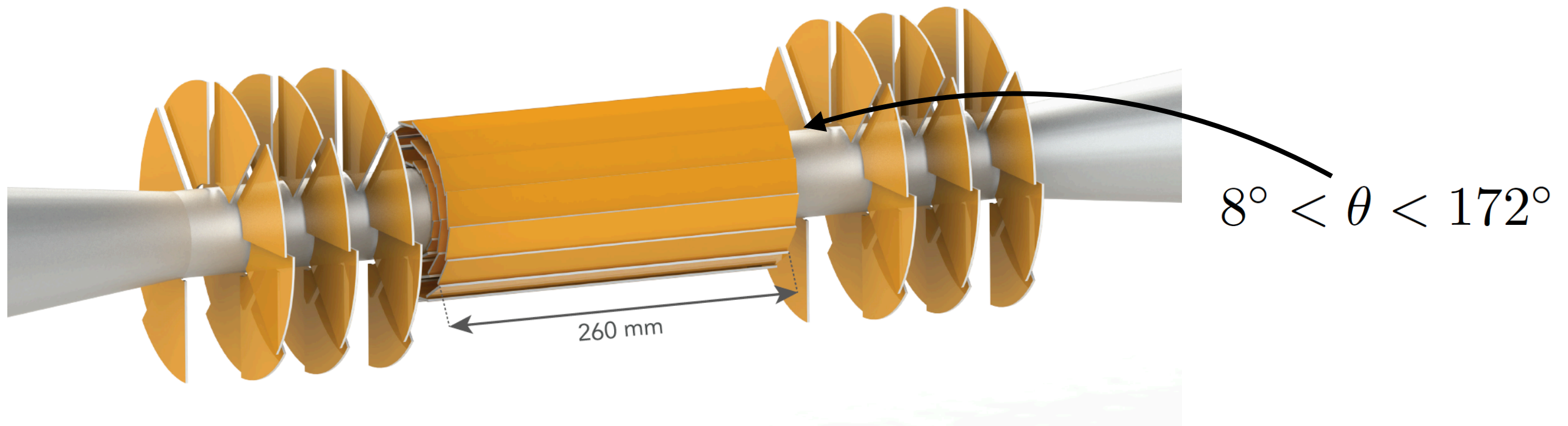
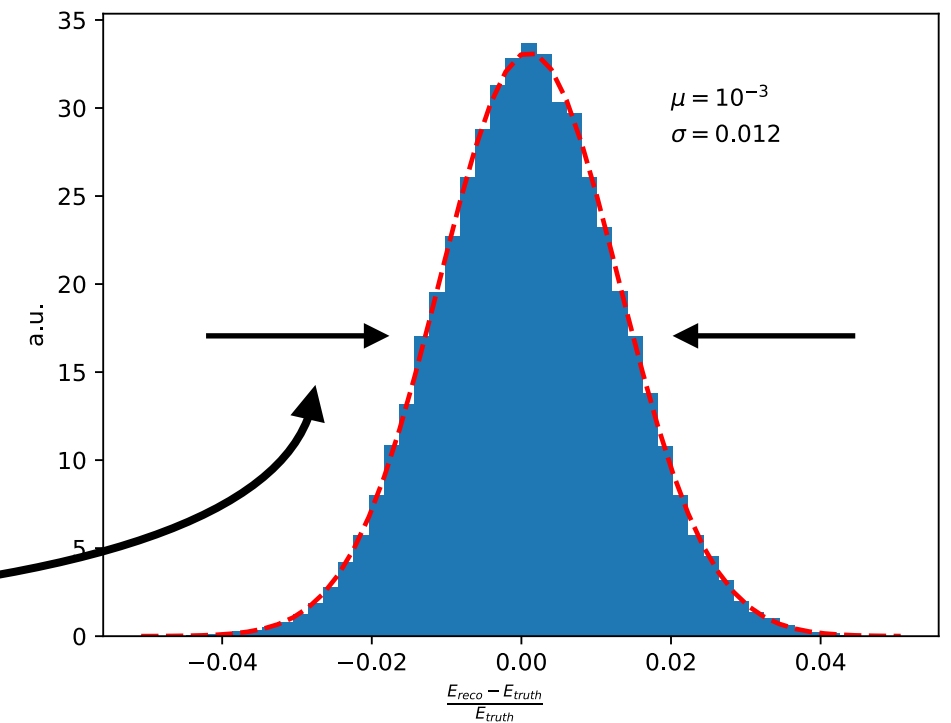
Symmetrized lumi spectra parameter uncertainty



Two major concerns over photon reconstruction

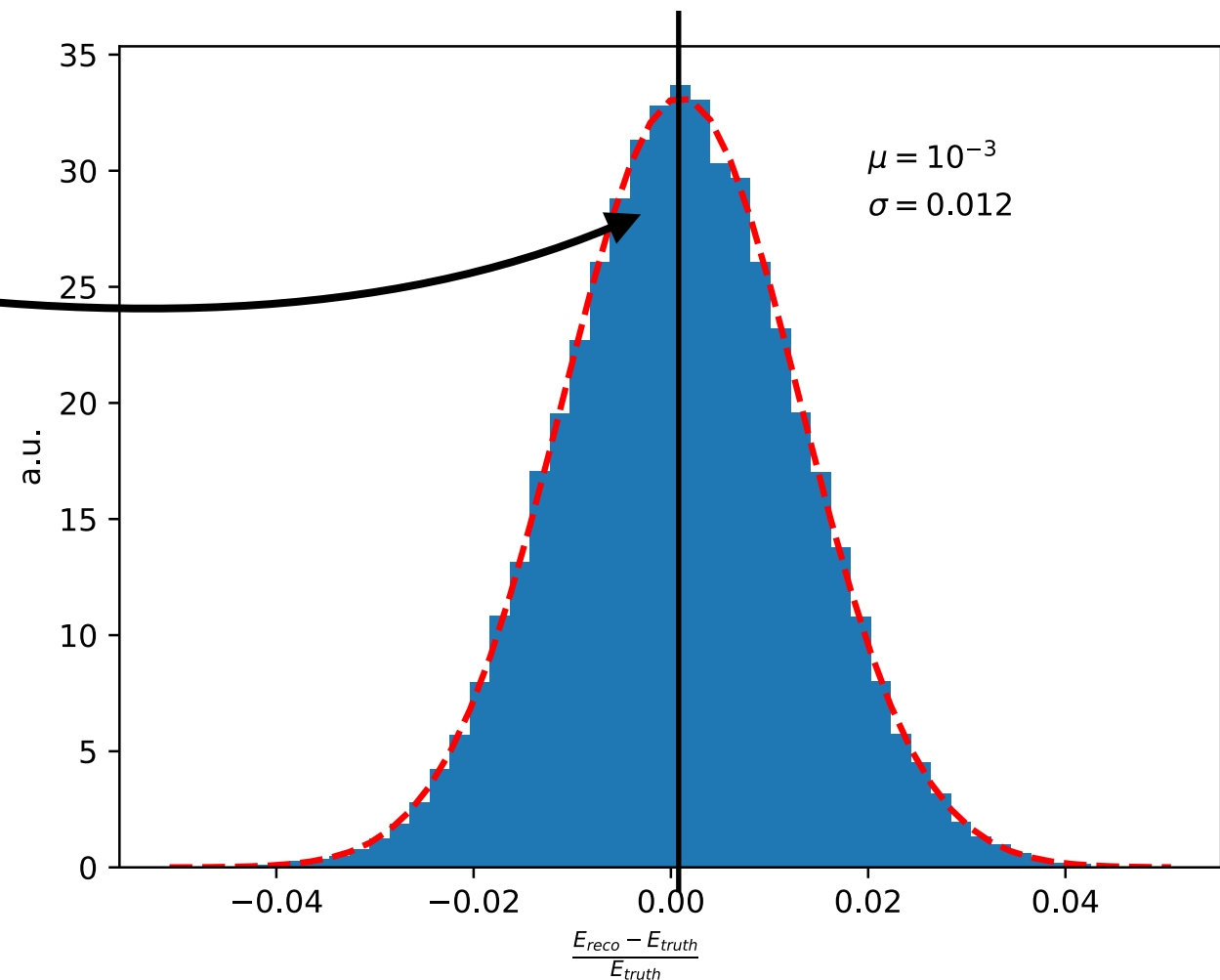
- ▶ Detector coverage → limit acceptance
- ▶ Reconstruction of the photon energy
 - Energy resolution → adopt it in the binning
 - Photon energy scale (calibration offset)

$$\sigma/E = 0.166/\sqrt{E} \oplus 1.1\%$$



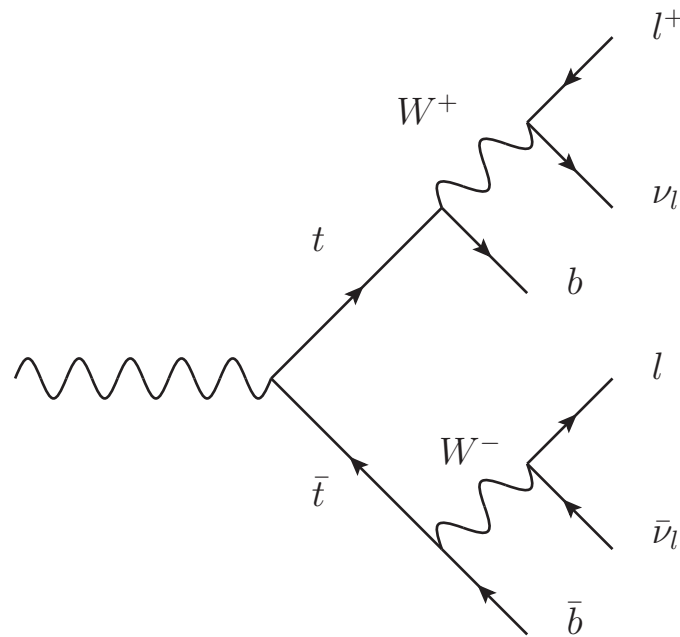
The calibration is not exact, the reconstructed energy is usually over or underestimated by a tiny amount

- Propagate the uncertainty by over(under)-estimating the photon energy in the observable
- Fit to the *nominal* model with m_t as a free parameter

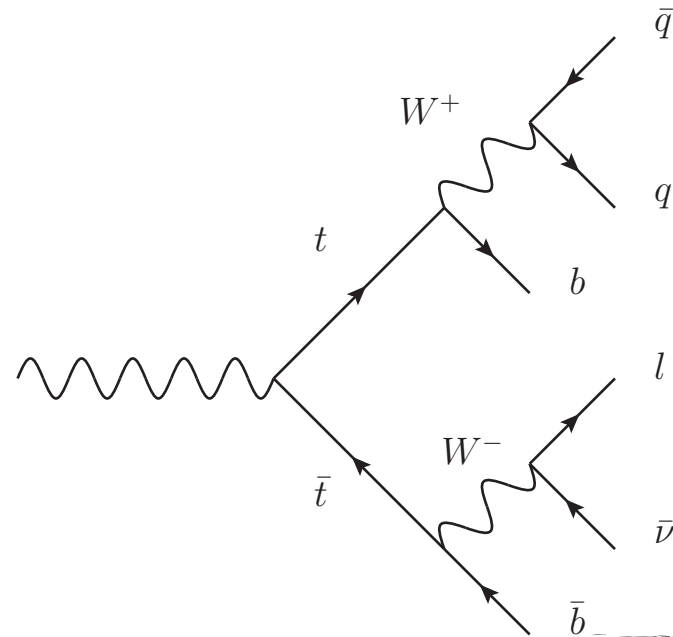


\overline{m}_t shift for different photon energy scale uncertainties			
Nominal energy	10^{-2} (1%)	10^{-3} (0.1%)	10^{-4} (0.01%)
380 GeV	+157 - 160 MeV	± 16 MeV	± 1.6 MeV
500 GeV	+842 - 863 MeV	± 85 MeV	± 9 MeV

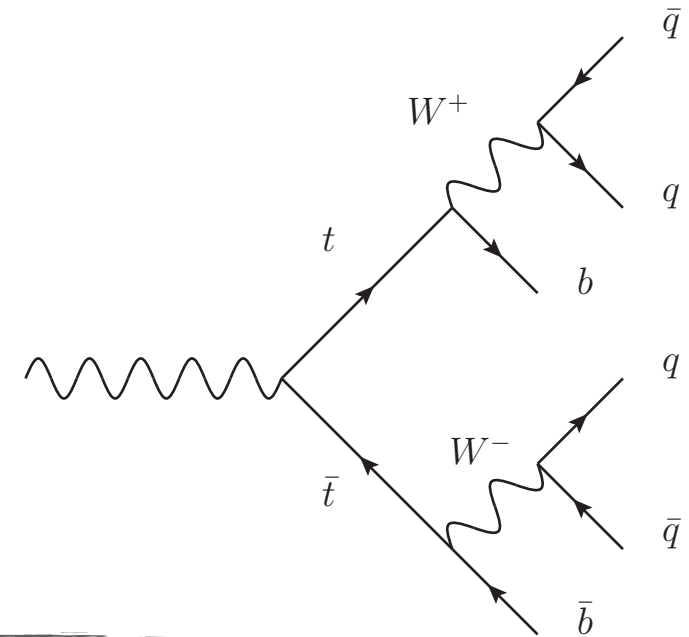
feasible today



Full leptonic (9%)



Semi-leptonic (45%)



Fully hadronic (46%)

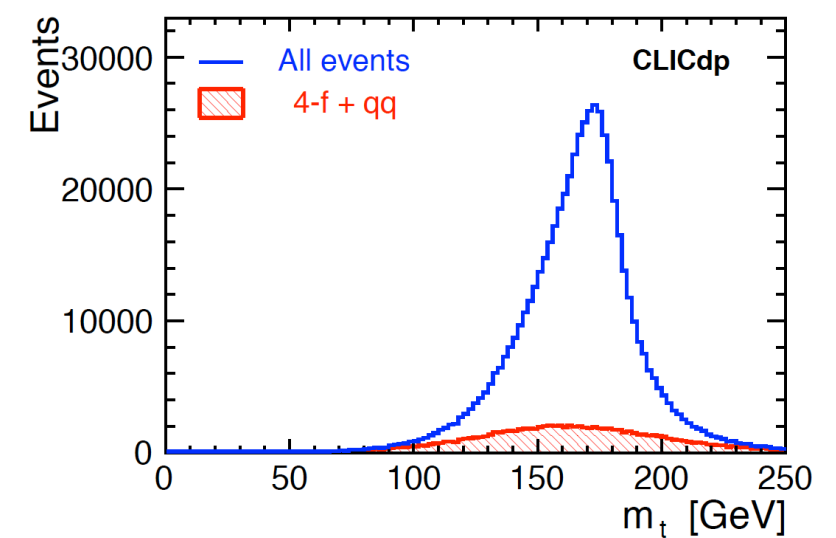
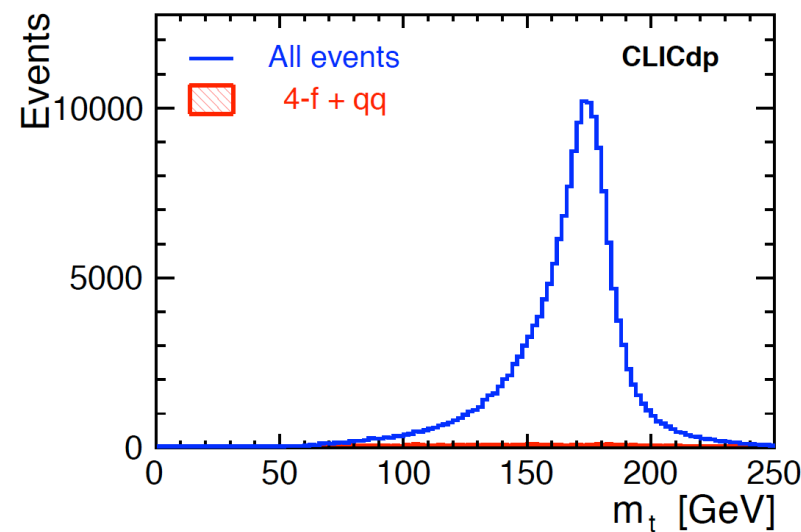
Selection efficiencies
estimated at full simulation

- Threshold scan study

70.2%

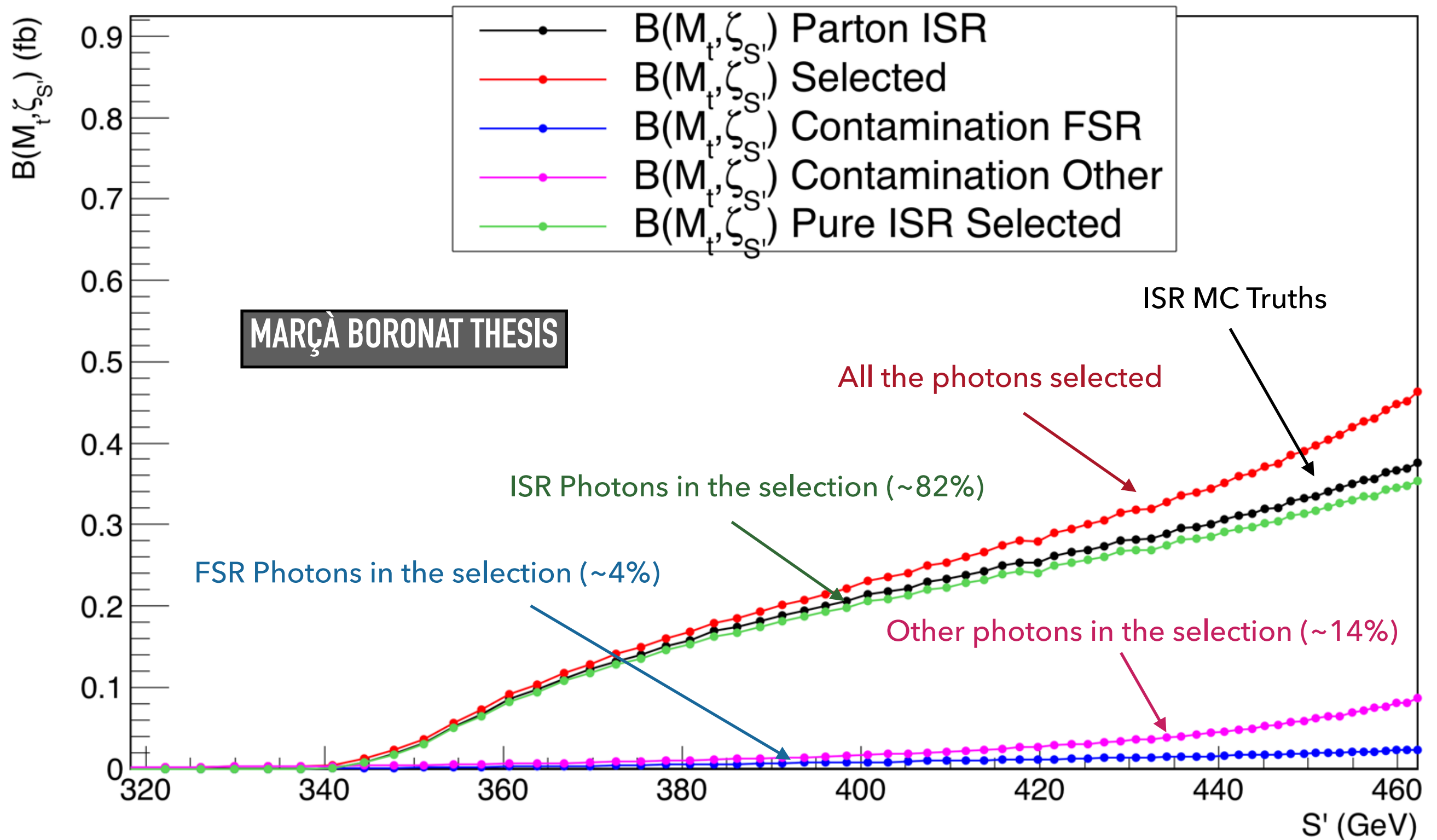
- Boosted decision tree

~90%

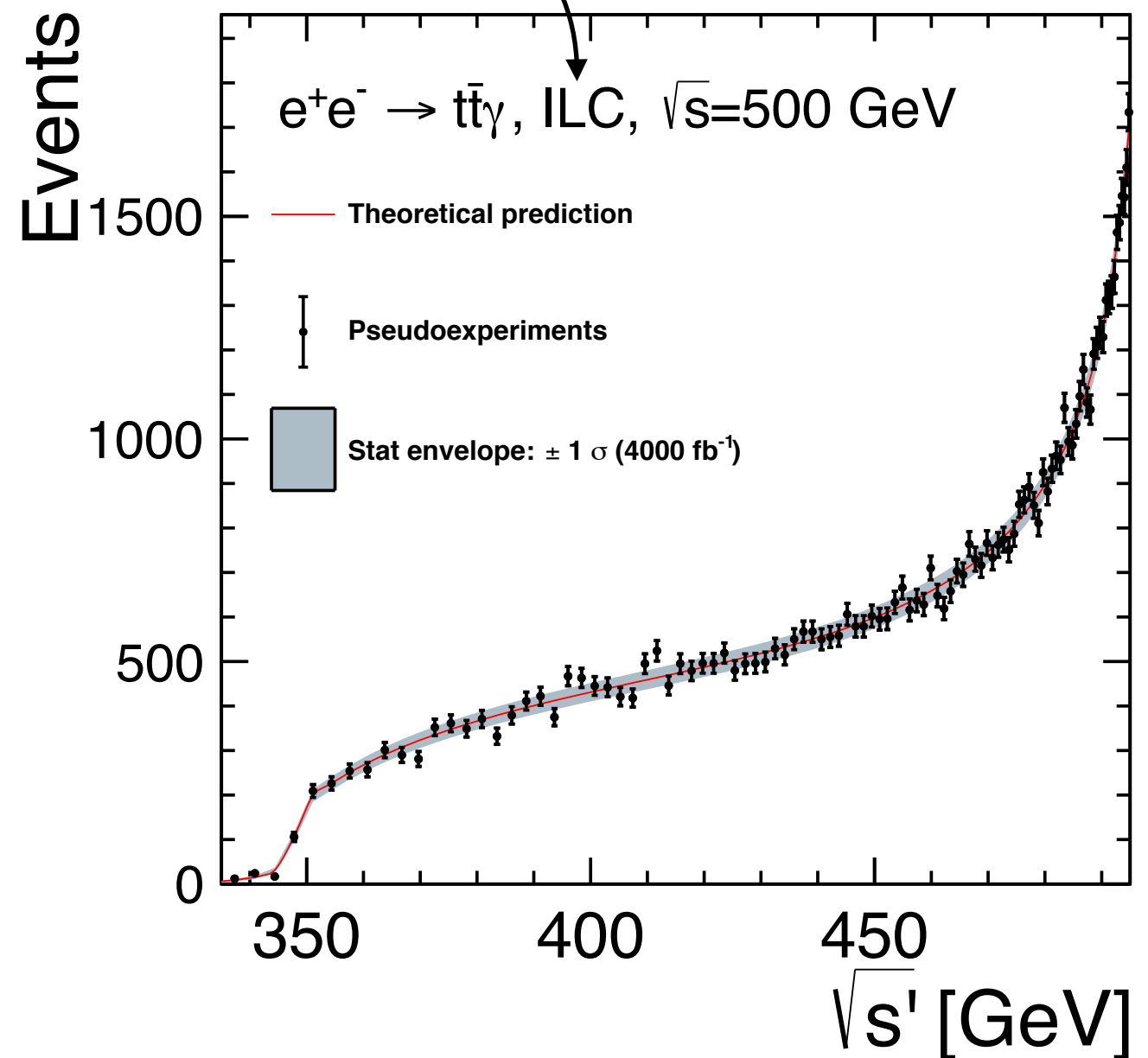
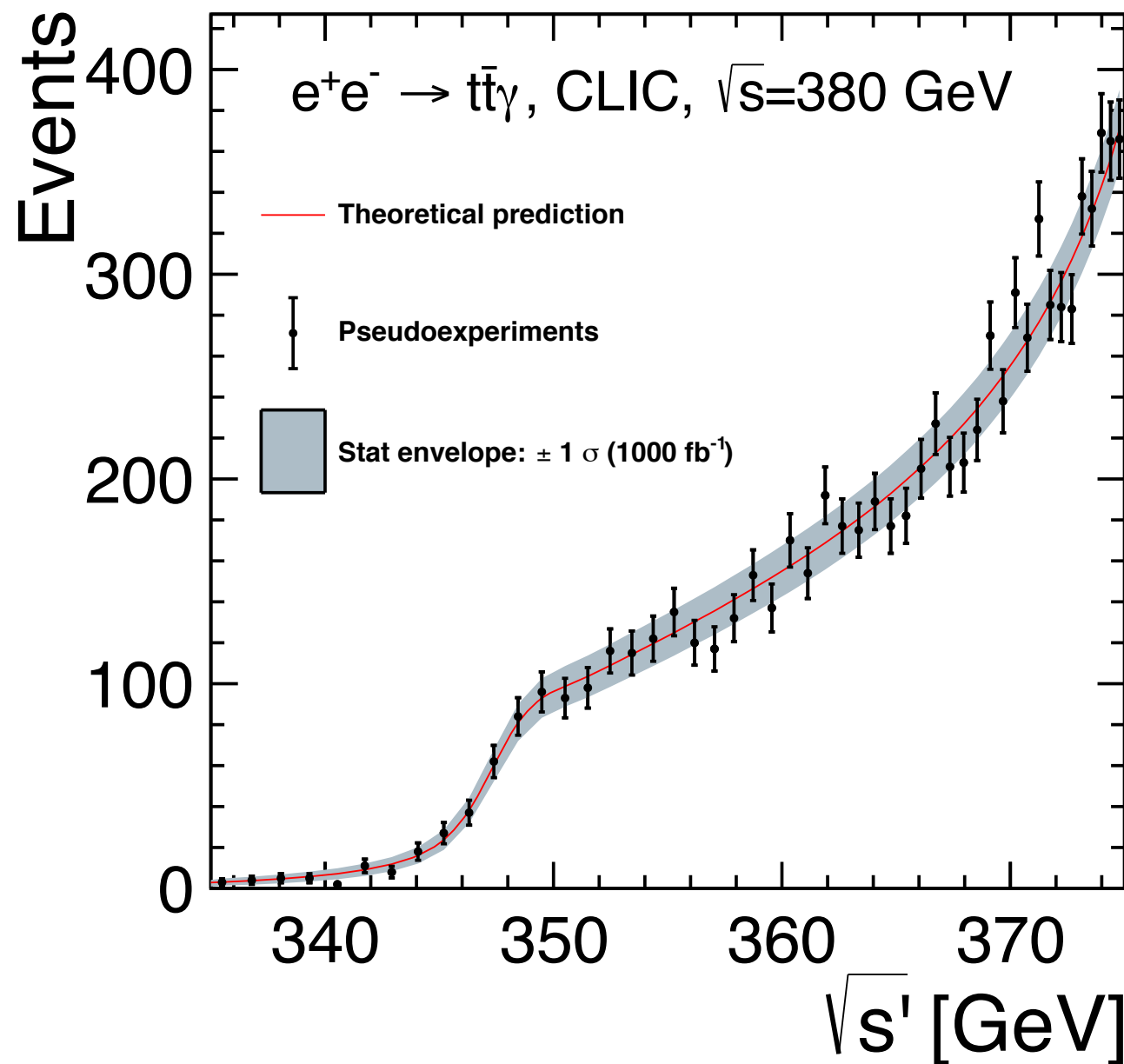


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$B(M_t, \zeta_{S'})$ (fb) Contributions for Selected ISR - EneCri - $m_{\text{top}} = 170$ GeV



24% bump in the statistics
due to beam polarisation



ECAL binning, 50% efficiency and luminosity spectrum accounted for

50% efficiency, ECAL resolution, luminosity spectrum and beam polarisation accounted

Experiment	CLIC, $\sqrt{s} = 380$ GeV		ILC, $\sqrt{s} = 500$ GeV	
L_{int} (fb^{-1})	500	1000	500	4000
Statistical	138 MeV	93 MeV	348 MeV	109 MeV
Theory	46 MeV		55 MeV	
Luminosity spectrum	7 MeV		7 MeV*	
Photon energy scale	16 MeV		85 MeV	
Total	147 MeV	105 MeV	363 MeV	149 MeV

Potential to measure \overline{m}_t with 100-150 MeV precision for the nominal luminosities at CLIC and ILC

- ▶ A new observable to measure the top quark mass with high precision in the continuum of a e^+e^- collider has been introduced
- ▶ A state of the art theoretical calculation was performed to assess the potential of the observable in a well defined theoretical framework
- ▶ The major sources of uncertainties were explored, including
 - Theory uncertainty
 - Photon reconstruction
 - Event selection
 - Luminosity spectrum
- ▶ Measurements in the order of 100 to 150 MeV are achievable for CLIC and ILC respectively for their 380 GeV and 500 GeV runs.

- ▶ The ISR emission is usually collinear to the emitter → extend the polar angle range of the analysis by including forward calorimeters
 - Going down from 8° to 6° increases statistics by a 33%
 - Going down from 8° to 4° would double the statistics
- ▶ Inclusion of the FSR from the top anti-top pair
 - Similar sensitivity to the top quark mass when compared with ISR
- ▶ Progress in the theory calculation to bring the uncertainty down
- ▶ Potential to measure the running of \overline{m}_t , as the measurement can be performed at different scales

THANKS FOR YOUR ATTENTION

¿Dudas? ¿Comentarios?