Heating up the chiral condensate with the Unruh thermal bath

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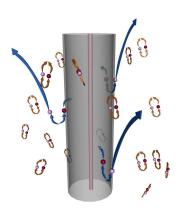
In collaboration with Antonio Dobado



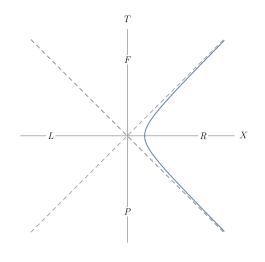
Introduction.

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- In the 1970s, **Quantum Field Theory** was generalized to
 - arbitrary observers in Minkowski,
 - curved spacetimes.
- Several seminal results:
 - Hawking: black holes radiate until they evaporate;
 - horizons have an impact on field quantization.



Accelerated motion in flat spacetime.



Hyperbolic motion:

$$T^2 - X^2 = -\frac{1}{a^2}.$$

Two **branches**: L and R.

Asymptotes:

$$T = \pm X$$
 (null lines \Longrightarrow **horizon**).

The Unruh effect.

Thermalization Theorem (Lee, 1986): the restriction of the Minkowski vacuum state $|\Omega_M\rangle$ of *any* quantum field theory to R is given by

$$\rho_R = \operatorname{tr}_L |\Omega_M\rangle\langle\Omega_M| = \frac{\mathrm{e}^{-2\pi H_R/a}}{\operatorname{tr} \, \mathrm{e}^{-2\pi H_R/a}},$$

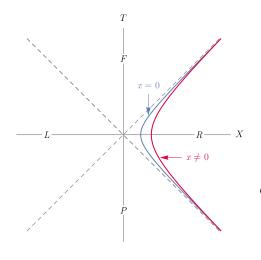
i.e. it is a thermal state at the Unruh temperature

$$T_{\rm U} = \frac{a}{2\pi}$$
 $(\hbar = c = G = k_{\rm B} = 1).$

This is the **generalization of the Unruh effect** to theories of interacting fields of arbitrary spin.

3

Accelerated motion in flat spacetime.



Comoving coordinates:

$$T = a^{-1}e^{ax} \sinh(at),$$

$$X = a^{-1}e^{ax} \cosh(at),$$

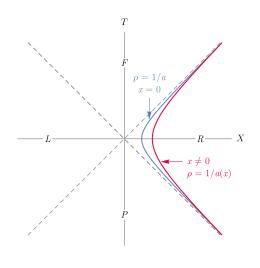
$$X_{\perp} = x_{\perp};$$

$$t, x, y, z \in \mathbb{R}.$$

$$\mathrm{d}s^2 = \mathrm{e}^{2ax}(\mathrm{d}t^2 - \mathrm{d}x^2) - \mathrm{d}x_\perp^2,$$

$$a(0) = a, \quad a(x) = a e^{-ax}.$$

Accelerated motion in flat spacetime.



Rindler coordinate:

$$\rho \equiv \frac{1}{a(x)} = \frac{e^{ax}}{a} \in (0, \infty),$$

$$ds^{2} = a^{2} \rho^{2} dt^{2} - d\rho^{2} - dx_{\perp}^{2}.$$

Euclidization
$$(t_E = it)$$
:

$$ds_{E}^{2} = a^{2} \rho^{2} dt_{E}^{2} + d\rho^{2} + dx_{\perp}^{2}.$$

$$t_{\rm E} \sim t_{\rm E} + \frac{2\pi}{a} \Rightarrow$$
 Thermal!

Spontaneous symmetry breaking and the Unruh effect.

The Unruh temperature is extremely small to be measured directly:

$$T_{\rm U} = \frac{a\hbar}{2\pi c k_{\rm B}} \simeq 3.97 \times 10^{-20} \frac{a}{g_{\oplus}} \text{ K}.$$

Furthermore, two questions arise:

- Is the Unruh effect just a formal result?
- Is acceleration completely analogous to temperature?

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Our goal: trying to shed light on these issues by studying a physically relevant case.

The QCD phase transition.

QCD and its chiral symmetry.

- Quantum Chromodynamics (QCD) is the theory of strong interactions within the Standard Model.
- The Euclidean Lagrangian of two-flavour massless QCD,

$$\mathscr{L}_{\text{QCD}} = -\frac{1}{4} G^a_{\mu\nu} G^a_{\mu\nu} + i \bar{q} \gamma_\mu D_\mu q,$$

is invariant under global $SU(2)_L \times SU(2)_R$ transformations.

- This **chiral symmetry** is actually **spontaneously broken**, since it is *not* a symmetry of the hadronic spectrum.
 - **Nambu-Goldstone Theorem**: 3 NGBs, the **pions**.

QCD and its chiral symmetry.

Chiral symmetry is also **explicitly broken** by operators of the form $\bar{q}Mq = \bar{q}_L Mq_R + \bar{q}_R Mq_L$ (e.g. a **mass term** for the quarks).

The simplest of these operators is $\bar{q}q$, whose expectation value in the Minkowski QCD vacuum state is known as the **quark condensate**, which, at **finite temperature** T in flat spacetime, is given by

$$\frac{\langle \bar{q}q \rangle_T}{\langle \bar{q}q \rangle_0} = \begin{cases} \sqrt{1 - \frac{T^2}{T_c^2}} & \text{if } 0 \le T < T_c, \\ 0 & \text{if } T \ge T_c. \end{cases}$$

Consequently, chiral symmetry is **restored** for temperatures higher than $T_c = 2f_{\pi}$, with a **second order phase transition** taking place at T_c .

We can compute the Euclidean partition function of the **lowest-order effective description of low-energy QCD** in Rindler space,

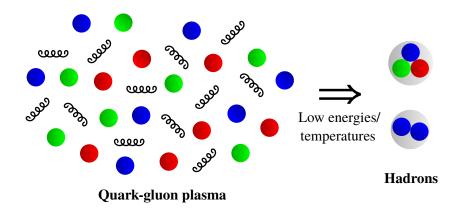
$$Z_{\mathrm{NL}\sigma\mathrm{M}} = \int [\mathrm{d}\pi^a][\mathrm{d}\sigma][\mathrm{d}\lambda] \, \mathrm{e}^{-\Gamma[\pi^a,\sigma,\lambda]},$$

where the effective action in the exponent is

$$\Gamma[\pi^a, \sigma, \lambda] = \int d^4x \sqrt{g} \left(-\frac{1}{2} \pi^a \Box \pi^a - \frac{1}{2} \sigma \Box \sigma + \frac{\lambda}{2} (\pi^a \pi^a + \sigma^2 - f_\pi^2) \right),$$

with $\pi^a \pi^a + \sigma^2 = f_{\pi}^2$ and $\langle \sigma \rangle \propto \langle \bar{q}q \rangle$.

Why an effective description? The two phases of hadronic matter.



The functional integral over the pion fields is Gaussian. \checkmark

What about the remaining integrals in σ , λ ?

- \implies Saddle-point approximation.
 - Large-N limit: $f_{\pi}^2 \equiv NF^2$, $F \neq F(N)$.
 - The fields are expanded around $(\bar{\sigma}, \bar{\lambda})$, with

$$\left. \frac{\delta \Gamma[\sigma, \lambda]}{\delta \sigma} \right|_{\sigma = \bar{\sigma}} = 0, \qquad \left. \frac{\delta \Gamma[\sigma, \lambda]}{\delta \lambda} \right|_{\lambda = \bar{\lambda}} = 0.$$

$$\implies \bar{\sigma}^2(x) = \langle \sigma(x) \rangle_a^2 = \langle \sigma^2(x) \rangle_a.$$

It is easy to check that $\bar{\sigma}$ and $\bar{\lambda}$ are then the solutions of

$$\frac{\delta\Gamma}{\delta\sigma(x)} = -\Box\sigma + \lambda\sigma = 0,$$

$$\frac{\delta\Gamma}{\delta\lambda(x)} = \frac{1}{2}(\sigma^2 - f_\pi^2) + \frac{N}{2}G(x, x; \lambda) = 0,$$

where the Euclidean Green function $G(x, x'; \lambda)$ satisfies

$$(-\Box + \lambda)_x G(x, x'; \lambda) = \frac{\delta^4(x - x')}{\sqrt{g}}.$$

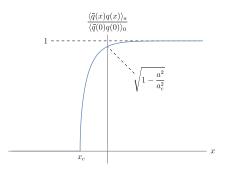
This system of equations can be solved in the limit $\lambda \simeq 0$.

Defining the **critical acceleration** as

$$a_c^2 \equiv \frac{48\pi^2 f_\pi^2}{N} \neq a_c^2(N),$$

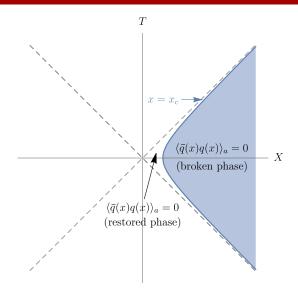
we obtain, for the **quark condensate**,

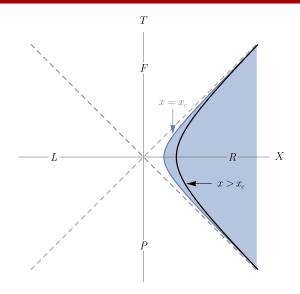
$$\frac{\langle \bar{q}(x)q(x)\rangle_a}{\langle \bar{q}(0)q(0)\rangle_0} = \sqrt{1 - \frac{a^2}{a_c^2}} e^{-2ax}.$$



The condensate **vanishes** if $\begin{cases} x = 0, \ a \ge a_c \text{ (thermal-like!)}, \text{ or} \\ x \le x_c \equiv \frac{1}{a} \ln \left(\frac{a}{a_c} \right) < 0. \end{cases}$

[A. C.-T., A. Dobado, Phys. Rev. **D 99** 125018 (2019), arXiv:1905.11179]





Conclusions.

- The Unruh effect can be extended to any quantum field theory using Lee's **Thermalization Theorem**.
- This powerful functional formalism allows us to determine whether acceleration is able to trigger phase transitions.
- In particular, we have found that **chiral symmetry is restored by acceleration**, with the results being equivalent to the inertial, thermal case: $a_c = 4\pi f_{\pi} = 2\pi T_c$ (for N = 3 pions).
- At least in principle, our results may have **applications** in Cosmology, Astrophysics and (maybe) Heavy-Ion Collisions.

Thank you!

¡Muchas gracias!