

BSM physics - Effective

Lagrangians - Naturalness -

Relaxion

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Monday: precision, \mathcal{L}_{eff} .

Tuesday: naturalness, partners

Wednesday: relaxion

BSM at colliders

1. precision measurements
(effective Lagrangians)
2. resonances (particles)

Effective Lagrangians for BSM and precision measurements

Example:

impressive!

ATLAS: $m_W = 80.370 \pm 0.019$ GeV
(Run 1)

Tevatron + LEP 80.385 ± 0.015 GeV

totally uninteresting (by itself)

Why? We don't care about the values of masses, parameters. $m_W = 75 \text{ GeV}$ is equally uninteresting

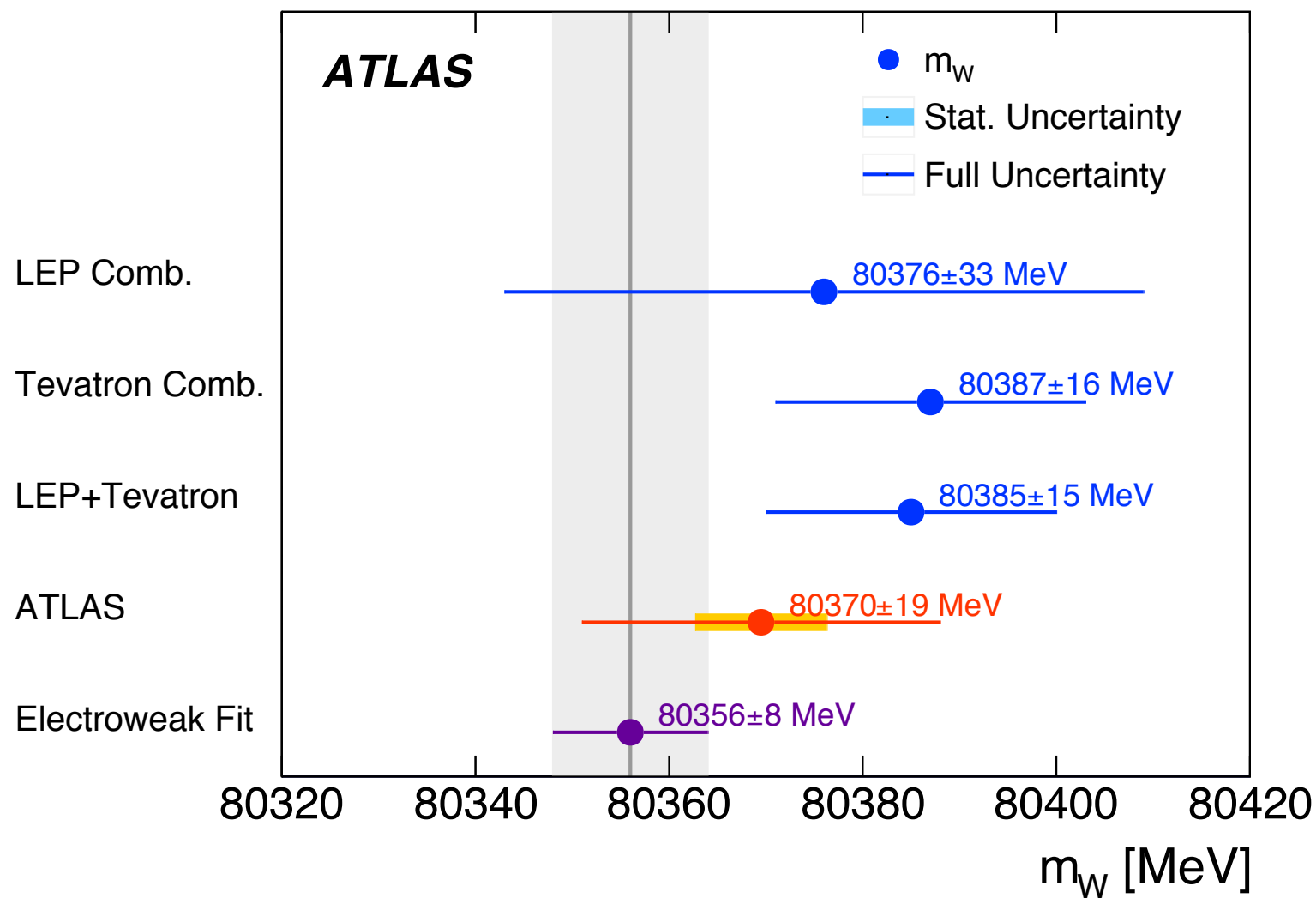
We want to understand physics \Rightarrow need to test predictions!

$$\text{in SM: } m_W = m_Z \cos \theta_W \quad (\text{tree level})$$

LEP ↑ ↑ SLC

Electroweak
fit:

$$m_Z \cos \theta_w = 80.356 \pm 0.008 \text{ GeV}$$



1701.07240v2 [hep-ex]

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a 0.2 per mille test of the Standard Model!

Sub-permille accuracy requires

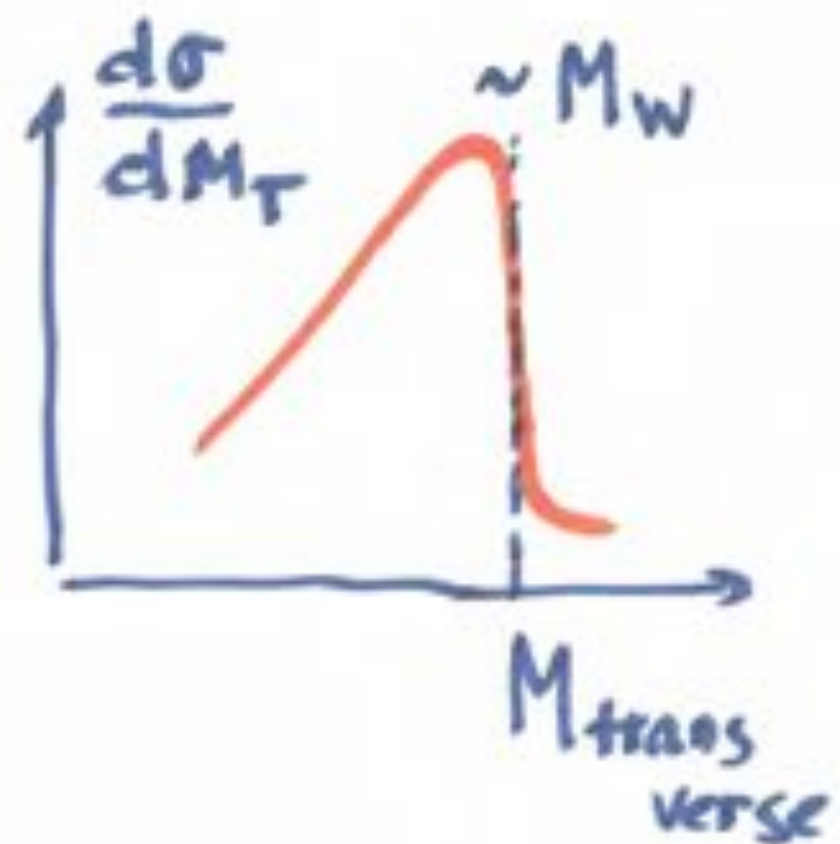
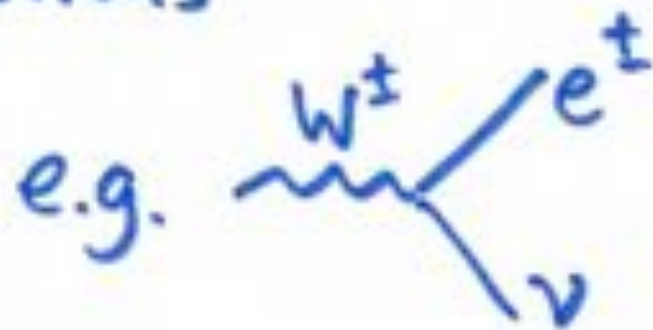
- loops



$$\frac{\delta m^2}{m^2} \sim \frac{g^2}{16\pi^2} \sim 1\%$$

\Rightarrow need 2 loop accuracy!

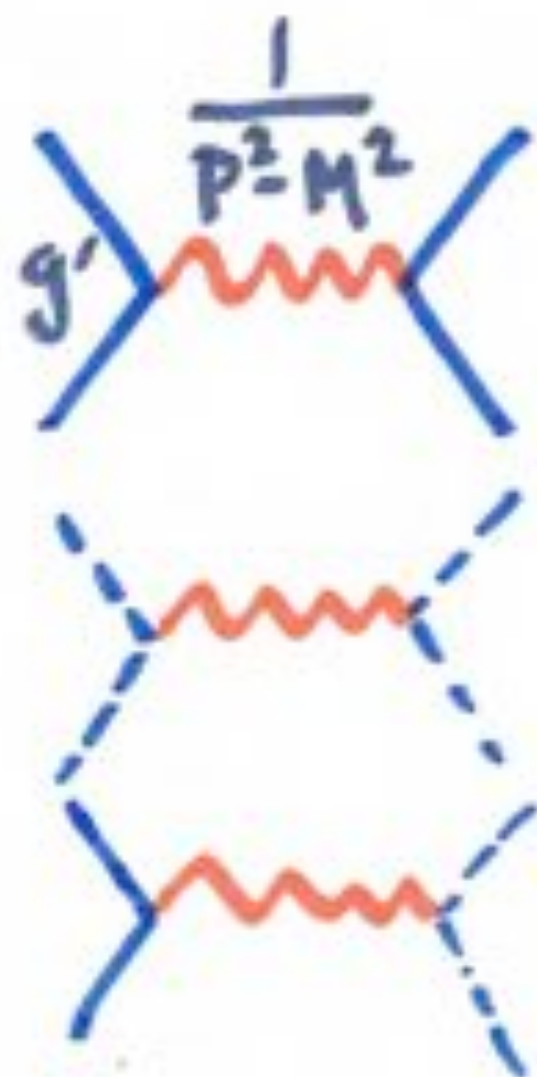
- Careful definitions



- careful experiments

Connection to BSM?

Z' example:



$M_{Z'} \gg p^2$



$\frac{g'^2}{M_{Z'}^2}$

$$\mathcal{L}_{\text{eff}} \sim \frac{(\bar{f} \gamma_{\mu} f + H^{\dagger} D_{\mu} H)^2}{\Lambda^2} \quad \Lambda \equiv \frac{M'}{g'}$$

Lesson: Heavy new physics ($M^2 \gg p^2$) parameterized by effective couplings suppressed by heavy scale

Strategy: parameterize NP by writing all possible effective couplings of SM fields and bound coefficients from experiment.

Example:

$$\frac{H^\dagger D_\mu H H^\dagger D^\mu H}{\Lambda^2}$$

$\partial_\mu + ig Z_\mu$ $v+h$

four Higgs scattering! ✗

but also ... $\underbrace{g^2 v^2}_{m_z^2} \frac{v^2}{\Lambda^2} Z_\mu Z^\mu$... a Z mass correction

$$\frac{\delta m_z^2}{m_z^2} \approx \frac{v^2}{\Lambda^2} \quad v = 246 \text{ GeV}$$

no shift in W-mass $\Rightarrow m_W = m_Z \cos\theta$ violated

$$\frac{\delta m^2}{m^2} = \frac{v^2}{\Lambda^2}$$

Experimental precision: $\frac{\delta m_w}{m_w} \sim 2.4 \cdot 10^{-4}$

$$\Rightarrow \frac{v^2}{\Lambda^2} \lesssim 2.4 \cdot 10^{-4}$$

$$\Rightarrow \Lambda \gtrsim v \cdot 46 \approx 11 \text{ TeV}$$

indirect precision test probes 11 TeV!

$$\left\{ \Lambda = \frac{M'}{g'}, \text{ "typical" } Z' \quad g' \sim \frac{1}{2}, \quad \frac{1}{4} \text{ dropped in calculation} \right.$$

$$\Rightarrow M_{Z'} \gtrsim 3 \text{ TeV} \quad \text{similar to direct limits} \quad \left. \right\}$$

systematically... order terms by mass dimension

$$\mathcal{L}_{\text{eff}} = \Lambda^4 + \Lambda^2 H^\dagger H + \lambda (H^\dagger H)^2 + \underbrace{\mathcal{L}_{\text{kinetic}}^{\text{SM}} + \mathcal{L}_{\text{Yukawa}}^{\text{SM}}}_{\text{dim 4, coefficients dimensionless}} \leftarrow \text{the SM}$$

dim 5 $+$ $\frac{(L_2 H)^2}{\Lambda}$

dim 6 $\left\{ \begin{array}{l} + \frac{H^\dagger \not{D}_\mu H H^\dagger \not{D}^\mu H}{\Lambda^2} + \frac{H^\dagger \not{D}_\mu H \bar{e}_R \gamma^\mu e_R}{\Lambda^2} + \frac{H^\dagger \sigma^a H W_{\mu\nu}^a B^{\mu\nu}}{\Lambda^2} \\ + \frac{H^\dagger H B_{\mu\nu} B^{\mu\nu}}{\Lambda^2} + \frac{H^\dagger H W_{\mu\nu}^a W_a^{\mu\nu}}{\Lambda^2} + \dots \end{array} \right.$ > 80 more at dimension 6

$+ \text{dim} > 6$


$$\begin{aligned} \mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i\bar{\psi} \gamma^\mu \partial_\mu \psi + h.c. \\ & + \bar{\psi} \gamma^\mu \gamma_5 \partial_\mu \psi + h.c. \\ & - R + \mathcal{L}_m - V(\phi) \end{aligned}$$

\mathcal{L}_{eff} has ∞ number of undetermined coefficients

• useful as expansion $\frac{\partial_\mu \partial^\mu}{\Lambda^2} \rightarrow \frac{p^2}{\Lambda^2} \ll 1$

• also $\frac{v^2}{\Lambda^2}$

• δm_Z

• $H^\dagger \frac{D_\mu H \bar{e}_R \gamma^\mu e_R}{\Lambda^2} \rightarrow g \frac{v^2}{\Lambda^2} \bar{e}_R \tilde{Z}_\mu \gamma^\mu e_R$ 

$$\frac{\delta g}{g} \sim \frac{v^2}{\Lambda^2} \Rightarrow \Lambda \gtrsim \text{few TeV} \quad (\text{LEP})$$

non-trivial example:

$$\mathcal{L} \sim S_1 g'^2 \frac{H^\dagger H B_{\mu\nu} B^{\mu\nu}}{\Lambda^2} + S_2 g^2 \frac{H^\dagger H W_{\mu\nu}^a W^{a\mu\nu}}{\Lambda^2} + S_{12} g g' \frac{H^\dagger \sigma^a H W_{\mu\nu}^a B^{\mu\nu}}{\Lambda^2}$$

↑ Hypercharge
 ↑ SU(2)

$$\langle H \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v+h \end{pmatrix}$$

→

$$\sim \frac{v^2}{\Lambda^2} B_{\mu\nu} B^{\mu\nu} + \frac{v^2}{\Lambda^2} W_{\mu\nu}^a W^{a\mu\nu} + S_{12} g g' \frac{v^2}{2\Lambda^2} W_{\mu\nu}^3 B^{\mu\nu}$$

absorb into B_μ, W_μ^a field-redefinitions

↑ SU(2) violating Z-γ mixing

$$"S" \equiv 16\pi S_{12} \frac{v^2}{\Lambda^2}$$

+ linear h-terms

$$+ S_{12} g g' \frac{v^2}{2\Lambda^2} W_{\mu\nu}^3 B^{\mu\nu}$$

← $SU(2)$ violating Z - γ mixing

$$"S" \equiv 16\pi S_{12} \frac{v^2}{\Lambda^2}$$

linear h -terms

$$+ S_1 g'^2 \frac{v^2}{\Lambda^2} \frac{h}{v} B_{\mu\nu} B^{\mu\nu} + S_2 g^2 \frac{v^2}{\Lambda^2} \frac{h}{v} W_{\mu\nu}^a W_{\mu\nu}^a + S_{12} g g' \frac{v^2}{\Lambda^2} \frac{h}{v} W_{\mu\nu}^3 B^{\mu\nu}$$

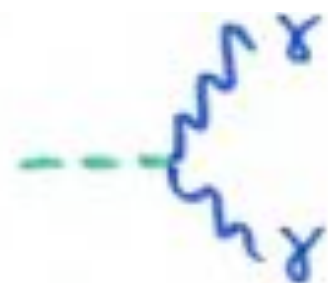
$$= \underbrace{4e^2 \frac{v^2}{\Lambda^2} (S_1 + S_2 + S_{12}) \frac{h}{v}}_{\equiv \frac{\alpha}{2\pi} k_\gamma} \frac{F_{\mu\nu} F^{\mu\nu}}{4} + \dots \frac{h}{v} F_{\mu\nu} Z^{\mu\nu} + \dots \frac{h}{v} Z_{\mu\nu} Z^{\mu\nu}$$

$$\equiv \frac{\alpha}{2\pi} k_\gamma$$

↑
 k_Z

↑
 $k_{Z\gamma}$

Higgs decays:



$\gamma\gamma$

ZZ^*

the data (95% confidence):

PDG 2018 precision EW fit : $S = 0.02 \pm 0.14$

PDG 2018 Higgs physics : $k_Y = 0.87 \pm 0.24$

$$k_Y^{SM} = 1$$

$$\Rightarrow \delta k_Y^{NP} = -0.13 \pm 0.24$$

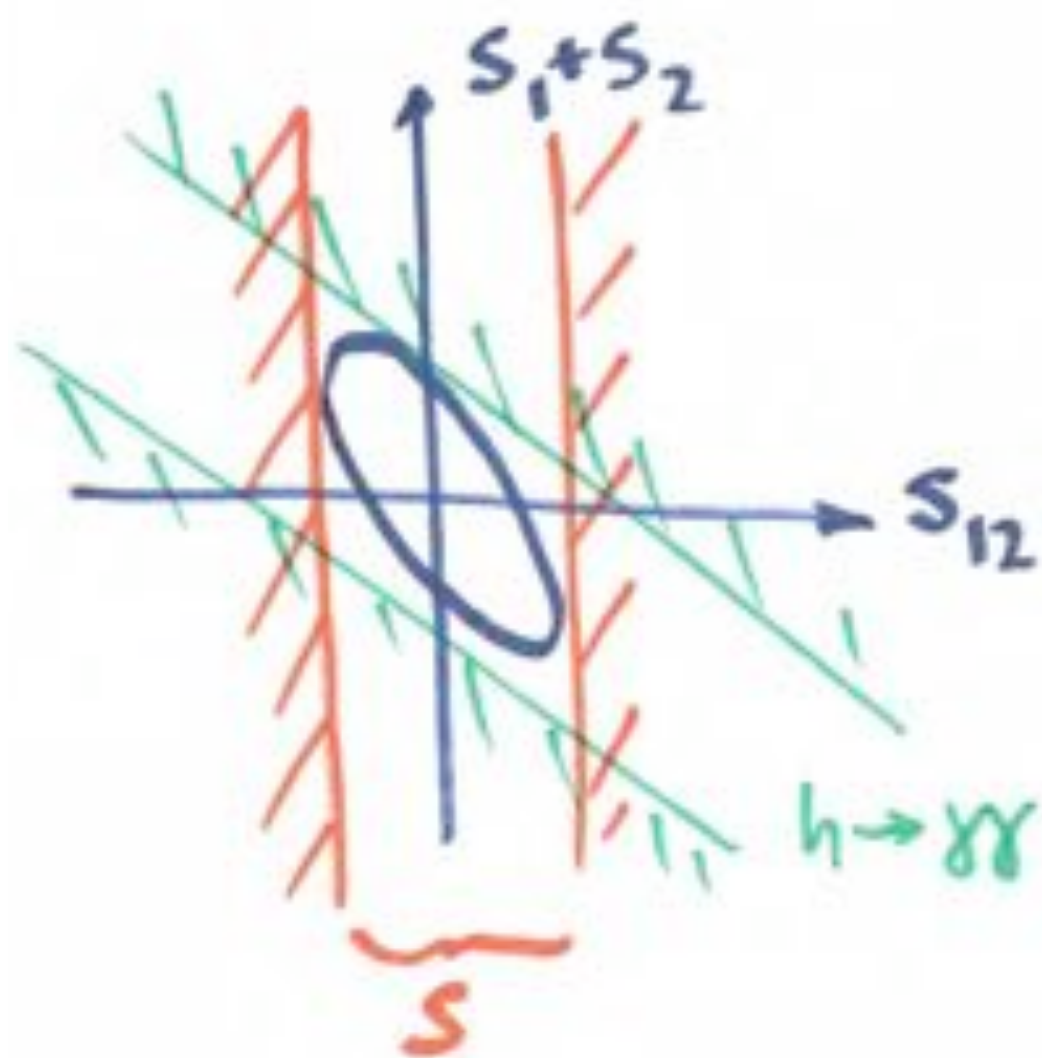
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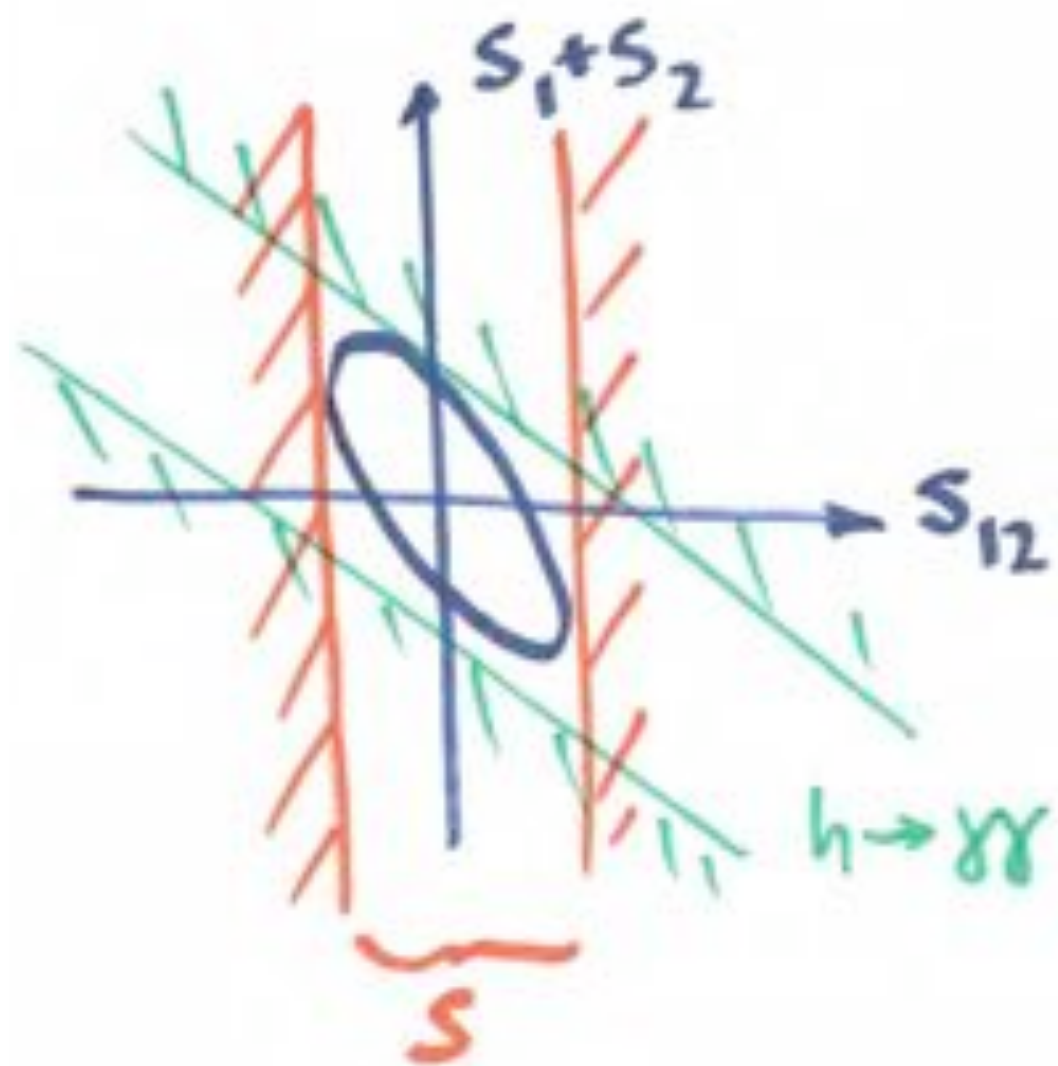
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New physics bounds ?

trees $S_i \approx 1 \Rightarrow \Lambda \approx 10 \text{ TeV}$

loops $S_i \approx \frac{1}{16\pi^2} \Rightarrow \Lambda \approx 1 \text{ TeV}$

back to Effective SM

$$\mathcal{L}_{\text{eff}} = \Lambda^4 + \Lambda^2 H^\dagger H + \mathcal{L}_4 + \frac{\mathcal{L}_5}{\Lambda} + \frac{\mathcal{L}_6}{\Lambda^2} + \dots$$

- expansion in $1/\Lambda$, valid when $p, m_{\text{SM}}, v \ll \Lambda$

- $\Lambda \sim$ scale of NP

- coefficients free parameters, determined by experiment.

If UV physics known, calculate coeffs in terms of UV parameters



Assume a non-zero coeff in \mathcal{L}_6 measured

$\Rightarrow \Lambda$ known

\Rightarrow guarantee of new physics

$$\Lambda \sim \frac{M}{g} \Rightarrow M \sim g\Lambda \leq 4\pi\Lambda$$

upper bound!

- Historical example: muon decay



$$\frac{1}{\Lambda^2} \sim \frac{1}{(200 \text{ GeV})^2}$$

Nature was nice: $M_W = 80 \text{ GeV}$, $g < 1$

current situation:

- neutrino mass $(LH)^2/\Lambda$ $m_\nu \sim \frac{v^2}{\Lambda} \sim 0.1 \text{ eV}$

$$\Rightarrow \Lambda \sim 10^{14} \text{ GeV} \quad \Rightarrow M \lesssim 10^{15} \text{ GeV}$$

(or neutrino is Dirac $LH\nu_R$)

- gravity $g_{\mu\nu} \frac{\partial^\mu H^\dagger \partial^\nu H}{M_{\text{pl}}}$ $\Lambda \sim M_{\text{pl}}$ new physics @ M_{pl}
 ↑
 graviton

- B physics $R_{K^{(*)}}$?? $\Lambda \approx 30 \text{ TeV}$

⇒ currently no direct evidence for a low NP scale ☹

indirect? Naturalness