

Light-Quark Dipole Operators at LHC

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Based on arXiv:1905.05187

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arXiv.org > hep-ex

High Energy Physics - Experiment

New submissions

Submissions received from Fri 7 Jun 19 to Mon 10 Jun 19, announced Tue, 11 Jun 19

- New submissions
- Cross-lists
- Replacements

[total of 31 entries: 1-31] [showing up to 2000 entries per page: fewer | more]

New submissions for Tue, 11 Jun 19

[1] arXiv:19 [pdf, other]

Another day, another "no particle" @ LHC

ATLAS, CMS collaboration

We continue to not find new particles at the LHC. Should we start to hate the Standard Model?

- > The search for direct NP did not show any new state.
- Probably they are heavy and there might exist a gap between them and the SM states.

The remarkable success of the SM and the lack of unexpected particles motivates to use a different approach to study NP.

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Effective Lagrangian: Linear realization

We parametrize new physics in terms of a linear effective Lagrangian, with a light Higgs:



Particle content: Same as the SM. No undiscovered particle at low energy. **Symmetries:** The SM gauge symmetry $SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$ is linearly realized. The lepton and baryon numbers are conserved.

There exist 59 Dimension-6 operators.

[Grzadkowski et al. arXiv: 1008.4884]



Light-Quark Dipole Operator @ LHC

Study the operators which include dipole-like couplings for the quarks.

$$\begin{split} \mathcal{O}_{uW,ij} &= i\bar{\mathbf{Q}}_i \sigma^{\mu\nu} u_{R,j} \hat{W}_{\mu\nu} \tilde{\Phi} , \qquad \mathcal{O}_{uB,ij} &= i\bar{\mathbf{Q}}_i \sigma^{\mu\nu} u_{R,j} \hat{B}_{\mu\nu} \tilde{\Phi} , \\ \mathcal{O}_{dW,ij} &= i\bar{\mathbf{Q}}_i \sigma^{\mu\nu} d_{R,j} \hat{W}_{\mu\nu} \Phi , \qquad \mathcal{O}_{dB,ij} &= i\bar{\mathbf{Q}}_i \sigma^{\mu\nu} d_{R,j} \hat{B}_{\mu\nu} \Phi \end{split}$$

 $\mathcal{L} = -\frac{ev}{\sqrt{2}\hbar^2} \left[F_{f\gamma} \bar{f} \gamma^{\mu\nu} f \partial_{\mu} \mathbf{A}_{\nu} + F_{fZ} \bar{f} \gamma^{\mu\nu} f \partial_{\mu} \mathbf{Z}_{\nu} + \left(\bar{f} \sigma^{\mu\nu} \left(F_{ff'W}^L L + F_{ff'W}^R R \right) f' \partial_{\mu} W_{\nu}^+ + h.c. \right) \right]$

$$F_{u\gamma} = f_{uW} + f_{uB} , F_{d\gamma} = f_{dW} - f_{dB} \Rightarrow \text{Anomalous magnetic moment}$$

$$F_{udW}^{R} = \frac{1}{s_{W}} f_{uW} , F_{udW}^{L} = \frac{1}{s_{W}} f_{dW} \Rightarrow W \text{ boson width decay}$$

$$F_{uZ} = \frac{c_{W}}{s_{W}} f_{uW} - \frac{s_{W}}{c_{W}} f_{uB} , F_{dZ} = \frac{c_{W}}{s_{W}} f_{dW} - \frac{s_{W}}{c_{W}} f_{dB} \Rightarrow Z \text{ boson width decay}$$

- Can be constrainted by the fit to the LEP observables.
- ► They also take part in electroweak diboson production $pp \rightarrow W^+W^-$ and $pp \rightarrow ZW^{\pm}$.

Have these operators any impact on the TGV analysis?

TGV analysis

$$\begin{split} \mathcal{O}_{WWW} &= \mathrm{Tr}[\widehat{W}^{\nu}_{\mu}\widehat{W}^{\rho}_{\nu}\widehat{W}^{\mu}_{\rho}]\\ \mathcal{O}_{W} &= (D_{\mu}\Phi)^{\dagger}\widehat{W}^{\mu\nu}(D_{\nu}\Phi)\\ \mathcal{O}_{B} &= (D_{\mu}\Phi)^{\dagger}\widehat{B}^{\mu\nu}(D_{\nu}\Phi) \end{split}$$

The tradicional operators of TGC couplings

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The tradicional operators of TGC couplings

 $\mathcal{O}_{\Phi,1} = (D_{\mu}\Phi)^{\dagger}\Phi\Phi^{\dagger}(D^{\mu}\Phi)$ The "omnipresent" operators. $\mathcal{O}_{BW} = \Phi^{\dagger}\widehat{B}_{\mu\nu}\widehat{W}^{\mu\nu}\Phi$

 $\mathcal{O}_{LLLL} = (\overline{L}\gamma^{\mu}L)(\overline{L}\gamma^{\mu}L)$ Gives a finite contribution to the Fermi constant.

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$$\begin{split} \mathcal{O}_{\Phi Q,ij}^{(1)} &= \Phi^{\dagger}(i\overset{\leftrightarrow}{D_{\mu}}\Phi)(\overline{Q}_{i}\gamma^{\mu}Q_{j}) \ , \quad \mathcal{O}_{\Phi Q,ij}^{(3)} &= \Phi^{\dagger}(i\overset{\leftrightarrow}{D_{\mu}}\Phi)(\overline{Q}_{i}\gamma^{\mu}T_{a}Q_{j}) \\ \mathcal{O}_{\Phi u,ij}^{(1)} &= \Phi^{\dagger}(i\overset{\leftrightarrow}{D_{\mu}}\Phi)(\overline{u}_{R_{i}}\gamma^{\mu}u_{R_{j}}) \ , \quad \mathcal{O}_{\Phi d,ij}^{(1)} &= \Phi^{\dagger}(i\overset{\leftrightarrow}{D_{\mu}}\Phi)(\overline{d}_{R_{i}}\gamma^{\mu}d_{R_{j}}) \\ \mathcal{O}_{\Phi e,ij}^{(1)} &= \Phi^{\dagger}(\overset{\leftrightarrow}{D_{\mu}}\Phi)(\overline{e}_{R_{i}}\gamma^{\mu}e_{R_{j}}) \ , \quad \mathcal{O}_{\Phi ud}^{(1)} &= \tilde{\Phi}^{\dagger}(i\overset{\leftrightarrow}{D_{\mu}}\Phi)(\overline{u}_{R}\gamma^{\mu}d_{R} + \mathrm{h.c.}) \end{split}$$

The fermionic operators

Impact of the dipole operators on the TGV



 The impact of the dipoles is minimum.

Neverthless the TGV can impose strong bounds on the dipole operators! How strong????

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How these bounds are comparable with other data set?

LHC vs EWPD



 The bounds from EWPD are weaker than those from LHC EWDBD.

 The contributions from dipole operators to EWDBD grow as s.

 DY results totally resolve the light-quark dipole couplings.

> ↓ stronger constrains

Summary

- Constraints derived on all the Wilson coefficients of those non-dipole operators is robust under the inclusion of the light-quark dipole operators.
- Analyses of LHC data improves over EWPD.
- The improvements driven both by the growth of the dipole contribution with energy.
- LHC also "sees" \boldsymbol{Z} and γ couplings with similar weight.

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LHC as a precision machine!!!

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LHC as a precision machine!!!

We should never forget that energy can provide a new discovery...

Thank you very much!!!