

Giant K factors

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Work performed with Mathieu Rubin and Sebastian Sapeta, [arXiv:1006.2144](https://arxiv.org/abs/1006.2144)

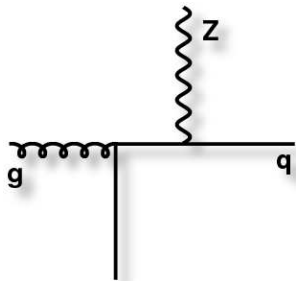
Kick-off meeting of LHCPHenoNet Initial Training Network
Valencia, Spain, 1–4 February 2011

Many searches for New Physics (eg. SUSY) rely on
Leading Order predictions for backgrounds.

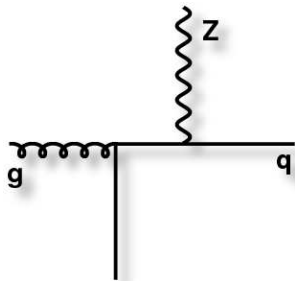
eg. Z+4jet background to gluino pair production
with NLO technology rapidly becoming mature for such cases

LO often considered good to within a factor of 2
NLO to within 10-20%

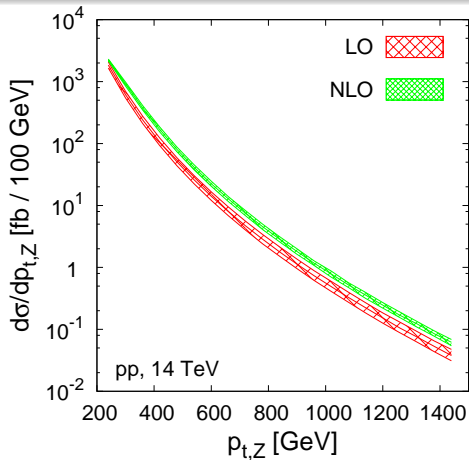
This talk is about cases where such “rules of thumb” fail
(spectacularly)



Use MCFM to examine various properties of such events at LO and NLO.

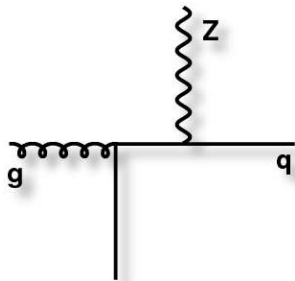


Use MCFM to examine various properties of such events at LO and NLO.

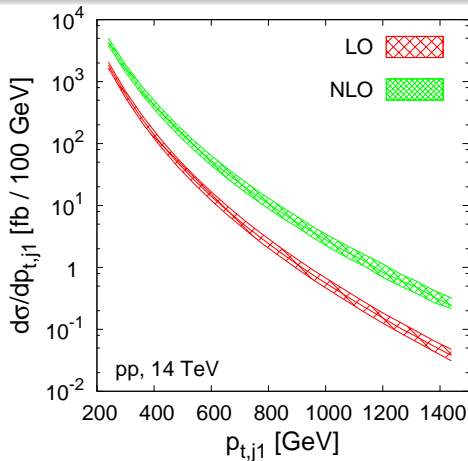


p_t spectrum of **Z boson** gets K -factor of 1.5

Fairly standard kind of occurrence

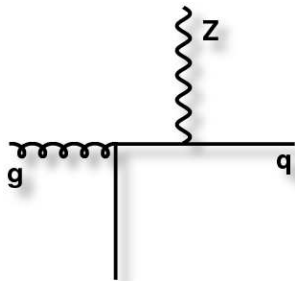


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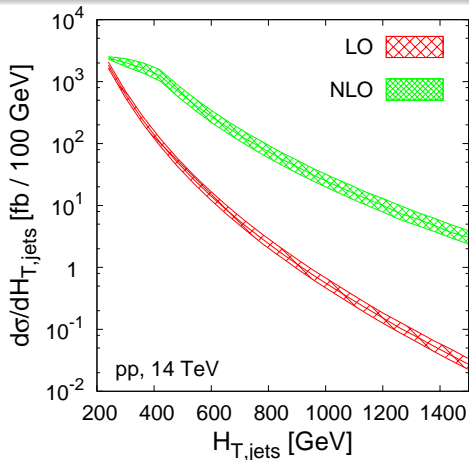


p_t spectrum of **leading jet** gets K -factor of 5–10

related issues in Butterworth, Davison, Rubin & GPS '08
Bauer & Lange '09; Denner, Dittmaier, Kasprzik & Much '09

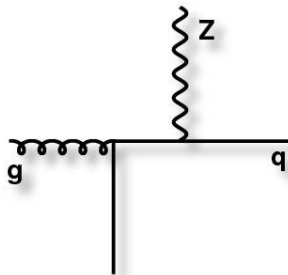


Use MCFM to examine various properties of such events at LO and NLO.

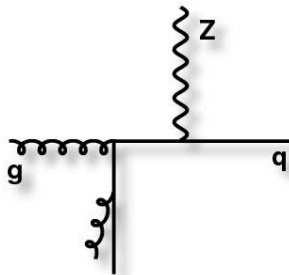


$$H_{T,jets} \equiv \sum_{i \in jets} p_{t,i} \text{ gets } K\text{-factor of up to 100}$$

Such things are not supposed to happen with $\alpha_s = 0.1!$

Leading Order

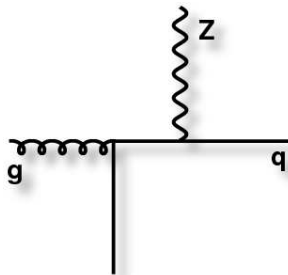
$$\alpha_s \alpha_{EW}$$

Next-to-Leading Order

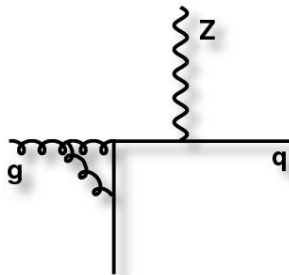
$$\alpha_s^2 \alpha_{EW}$$

LHC probes scales well above EW scale, $\sqrt{s} \gg M_Z$.
EW bosons are **light**. New log-enhanced topologies appear.

$H_{T,jets}$ is extreme, because at LO $H_{T,jets} \simeq p_{t,jet 1}$; NLO: $H_{T,jets} \simeq 2p_{t,jet 1}$

Leading Order

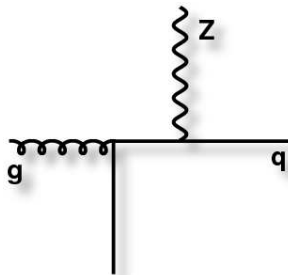
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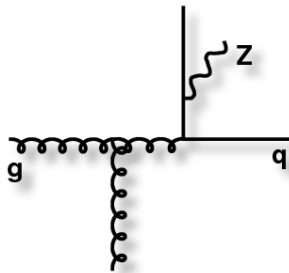
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Leading Order

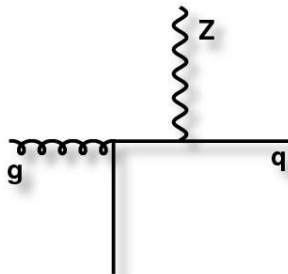
$$\alpha_s \alpha_{EW}$$

Next-to-Leading Order

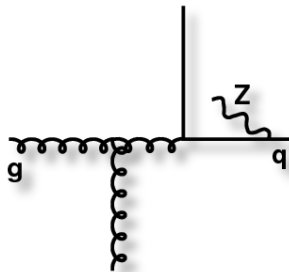
$$\alpha_s^2 \alpha_{EW} \ln^2 \frac{p_t}{M_Z}$$

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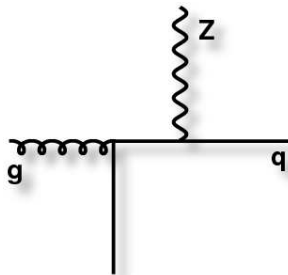
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Next-to-Leading Order

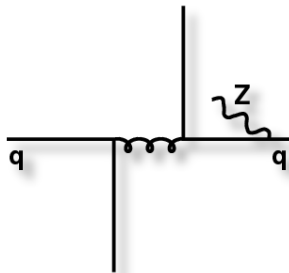
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Leading Order

$$\alpha_s \alpha_{EW}$$

Next-to-Leading Order

$$\alpha_s^2 \alpha_{EW} \ln^2 \frac{p_t}{M_Z}$$

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$H_{T,jets}$ is extreme, because at LO $H_{T,jets} \simeq p_{t,jet 1}$; NLO: $H_{T,jets} \simeq 2p_{t,jet 1}$

We calculated $Z+\text{jet@NLO}$.

But giant K -factors are dominated by “LO” $Z+2\text{-parton}$ piece of $Z+\text{jet@NLO}$.

We know LO calculations aren't reliable.

We'd ideally want to combine $Z+\text{jet@NLO}$ with $Z+2\text{-jet@NLO}$

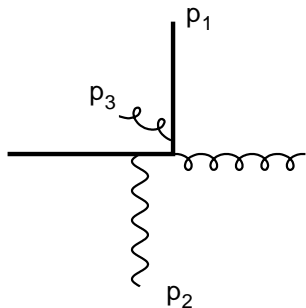
without double counting
without having to do full $Z+\text{jet@NNLO}$ calculation

First try something simpler:

Take the “leading” process
[Z + jet @ LO]

and add in process with one extra jet.
[i.e. include Z + 2 jets @ LO]

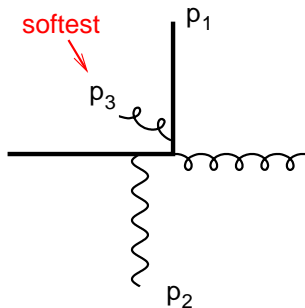
approximate the 1-loop Z+jet term, by requiring
cancellation of all divergences
[those from singly unresolved limit of Z + 2 jets]



Z + 2 partons

$$|M^2(p_1, p_2, p_3)|$$

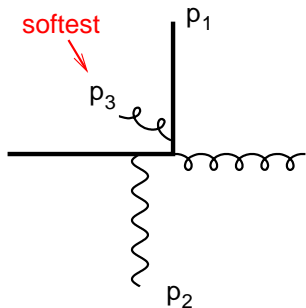
- ▶ Identify softest or most collinear parton [with help of a jet algorithm]
- ▶ “Loop” it \equiv remove it from event, reshuffle other momenta;
 weight of looped event is $(-1) \times$ weight of tree-level event



Z + 2 partons

$$|M^2(p_1, p_2, p_3)|$$

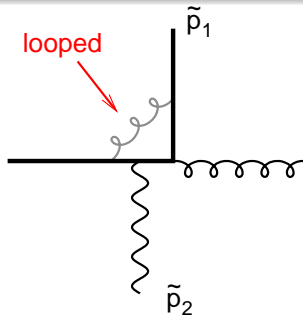
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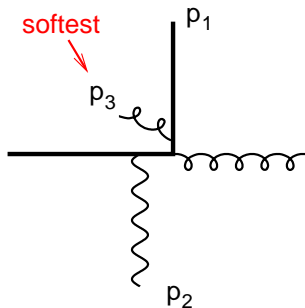
+



Z + 1 parton + 1 sim. loop

$$-|M^2(p_1, p_2, p_3)|$$

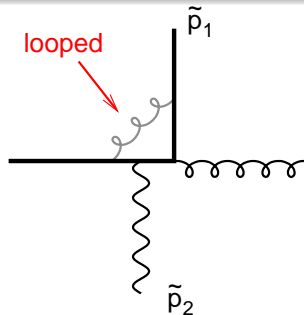
- ▶ Identify softest or most collinear parton [with help of a jet algorithm]
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Z + 2 partons

$$|M^2(p_1, p_2, p_3)|$$

+



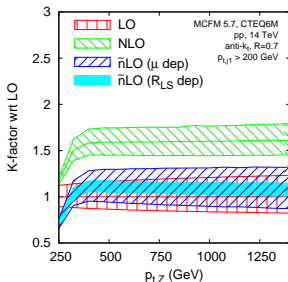
Z + 1 parton + 1 sim. loop

$$-|M^2(p_1, p_2, p_3)|$$

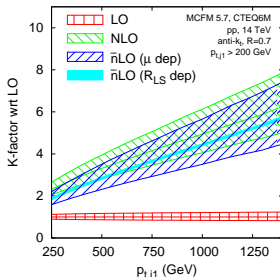
- ▶ Identify softest or most collinear parton [with help of a jet algorithm]
- ▶ “Loop” it \equiv remove it from event, reshuffle other momenta; weight of looped event is $(-1) \times$ weight of tree-level event

This cancels all the “single-unresolved” divergences in the Z+2 events

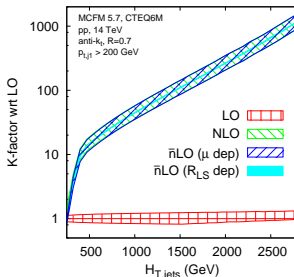
p_t of Z-boson



p_t of jet 1



$H_{T,jets} = \sum_{jets} p_{t,j}$



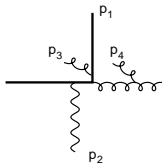
When the K -factors are large, \bar{n} LO agrees well with NLO

MLM matching also does a similar job
 cf. de Aquino, Hagiwara, Li & Maltoni '11

Differences between LoopSim and MLM/CKKW matching:

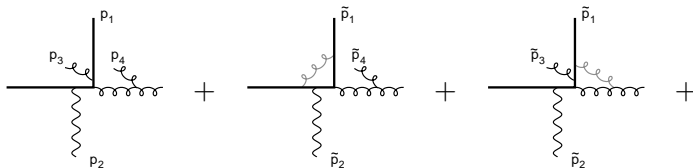
1. Does not rely on shower (✓: simplicity; ✗: not easily integrated with shower MCs)
2. Does not need arbitrary separation of $Z+1/Z+2$ /etc. samples with (hard-to-choose) momentum cutoff
3. Can easily be extended beyond LO matching

add tree-level Z+3,
cancel divergences in single + doubly unresolved limits: **n̄n̄LO**



$$|M^2(p_1, p_2, p_3, p_4)|$$

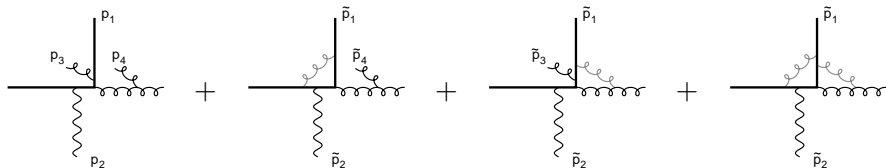
add tree-level Z+3,
cancel divergences in single + doubly unresolved limits: **n̄n̄LO**



$$|M^2(p_1, p_2, p_3, p_4)| \quad - |M^2(p_1, p_2, p_3, p_4)| \quad - |M^2(p_1, p_2, p_3, p_4)|$$

provides approximation of 1-loop

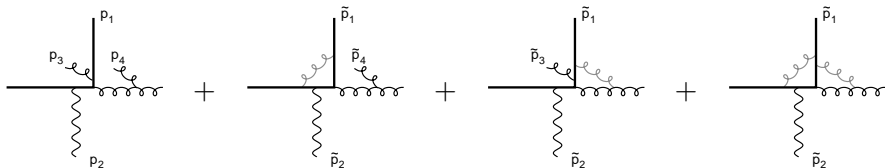
add tree-level Z+3,
cancel divergences in single + doubly unresolved limits: **n̄n̄LO**



$$|M^2(p_1, p_2, p_3, p_4)| \quad - |M^2(p_1, p_2, p_3, p_4)| \quad - |M^2(p_1, p_2, p_3, p_4)| \quad + |M^2(p_1, p_2, p_3, p_4)|$$

provides approximation of 1-loop and 2-loop contributions; total sums to zero

add tree-level Z+3,
cancel divergences in single + doubly unresolved limits: **n̄n̄LO**



$$|M^2(p_1, p_2, p_3, p_4)| - |M^2(p_1, p_2, p_3, p_4)| - |M^2(p_1, p_2, p_3, p_4)| + |M^2(p_1, p_2, p_3, p_4)|$$

provides approximation of 1-loop and 2-loop contributions; total sums to zero

add in (exact Z+2 @ 1-loop) – (approximate Z+2 @ 1-loop)
+ extra simulated 2-loop piece to cancel new Z+2@1-loop divergences

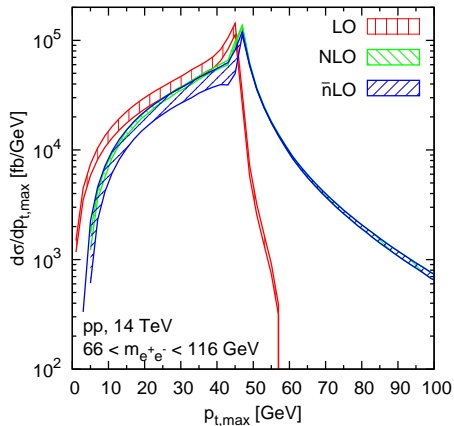
This is n̄NLO

allows combination of Z+1@NLO with Z+2@NLO

Testing NLO Merging, in 3 processes

1. $Z@NLO$ with $Z+j@NLO$
2. $Z+j@NLO$ with $Z+2j@NLO$
3. $2j@NLO$ with $3j@NLO$

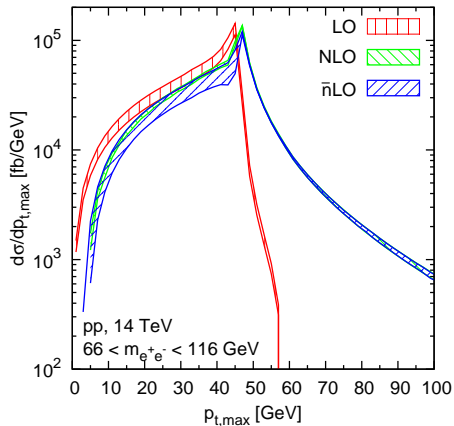
\bar{n} LO v. NLO



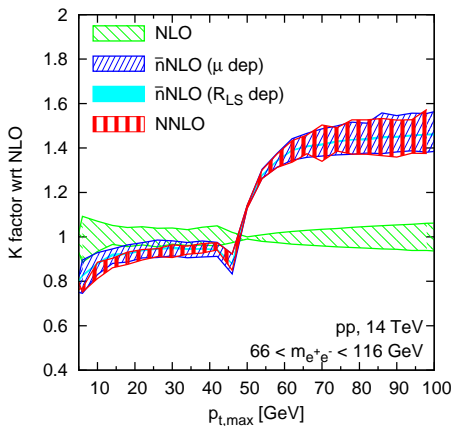
Z (i.e. DY) with Z+j from MCFM & LoopSim

For $p_{t,ell} \gtrsim \frac{1}{2}M_Z + \Gamma_Z$ (giant K -factor!) it had to work
For $p_{t,\ell} \lesssim \frac{1}{2}M_Z + \Gamma_Z$ it's remarkable that it still works

\bar{n} LO v. NLO



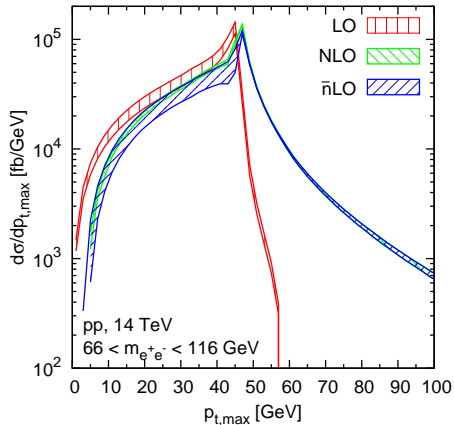
\bar{n} NLO v. NNLO



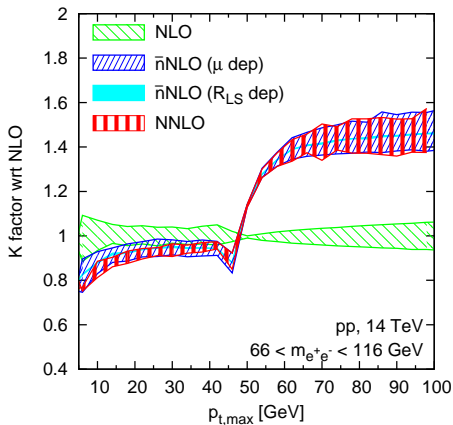
NNLO from DYNNOLO, Z (i.e. DY) with Z+j from MCFM & LoopSim

For $p_{t,ell} \gtrsim \frac{1}{2}M_Z + \Gamma_Z$ (giant K -factor!) it had to work
 For $p_{t,ell} \lesssim \frac{1}{2}M_Z + \Gamma_Z$ it's remarkable that it still works

\bar{n} LO v. NLO



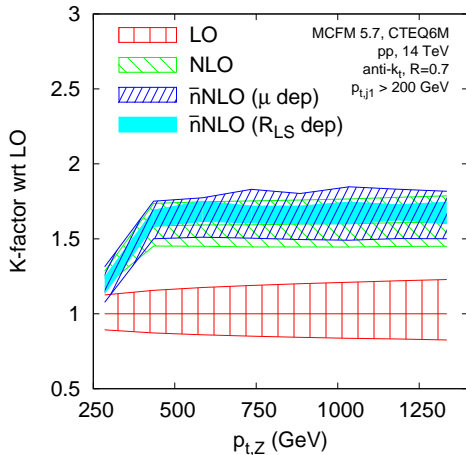
\bar{n} NLO v. NNLO



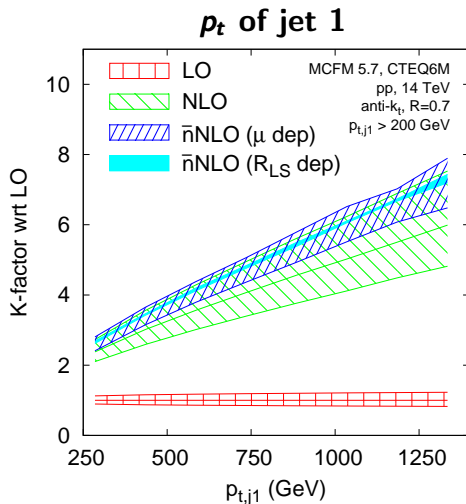
NNLO from DYNNOLO, Z (i.e. DY) with Z+j from MCFM & LoopSim

For $p_{t,ell} \gtrsim \frac{1}{2}M_Z + \Gamma_Z$ (giant K -factor!) it had to work
 For $p_{t,l} \lesssim \frac{1}{2}M_Z + \Gamma_Z$ it's remarkable that it still works

p_{tZ} of Z-boson

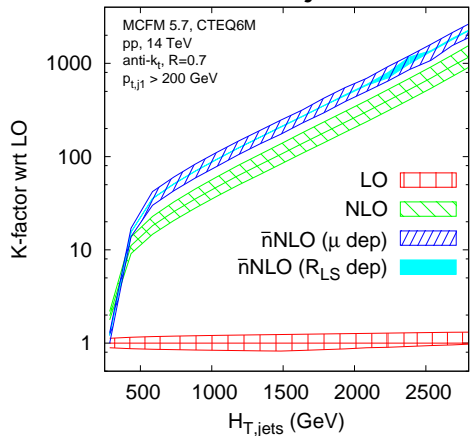


- ▶ p_{tZ} distribution didn't have giant K -factors.
- ▶ \bar{n} NLO brings no benefit
 To get improvement you would need exact 2-loop terms



- ▶ p_{tj} distribution seems to converge at \bar{n} NLO
- ▶ scale uncertainties reduced by \sim factor 2

$$H_{T,jets} = \sum_{jets} p_{t,j}$$



- ▶ Significant further enhancement for $H_{T,jets}$
- ▶ \bar{n} NLO brings clear message:

$H_{T,jets}$ is not a good observable!

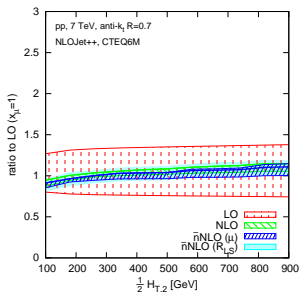
H_T (effective mass) type observables are widely used in searches

- ▶ H_T has a steeply falling distribution (like p_{tj} , p_{tZ})
- ▶ At each order (NLO, NNLO), an extra (soft) jet contributes to the H_T sum e.g. from ISR
- ▶ That shifts H_T up, which translates to a substantial increase in the cross section

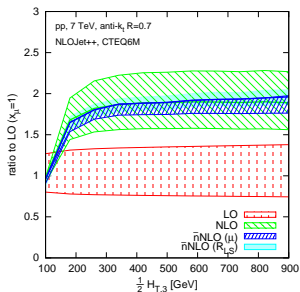
We can test this hypothesis for plain jet events, using a truncated sum,

$$H_{T,n} = \sum_{i=1}^n p_{t,\text{jet } i}$$

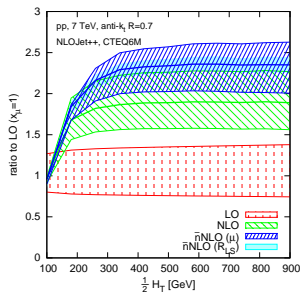
$H_{T,2}$



$H_{T,3}$



$H_{T,\infty}$



A clear message:

for a process with n objects at lowest order, use $H_{T,n}$

Do you know what gets used in your experiment's searches?

Many writers of LHC SUSY proceedings didn't...

Be aware that giant K -factors exist

Always look one order beyond the leading order, for example with
MLM/CKKW matching

New tool to get good predictions in such cases: **LoopSim**

Basically an “operator” to generate approximations to unknown loops

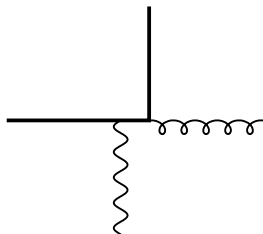
Combine $Z+j@NLO$, $Z+2j@NLO$ to get “ $\bar{n}NLO$ ” $Z+jet$

It sometimes works even beyond “giant- K -factor” regions

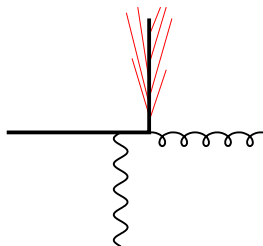
Watch out for H_T

Even for simple processes, it converges very poorly
unless you define it carefully (limit number of objects in sum)

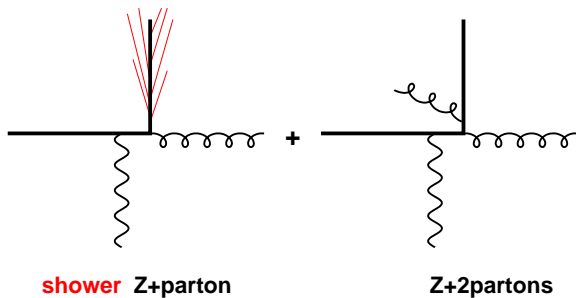
EXTRAS

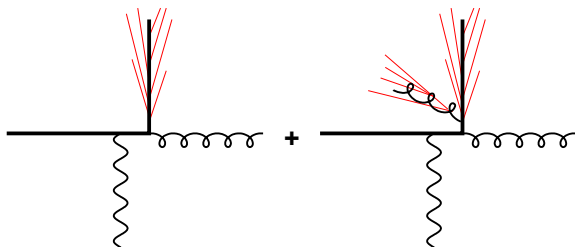


Z+parton



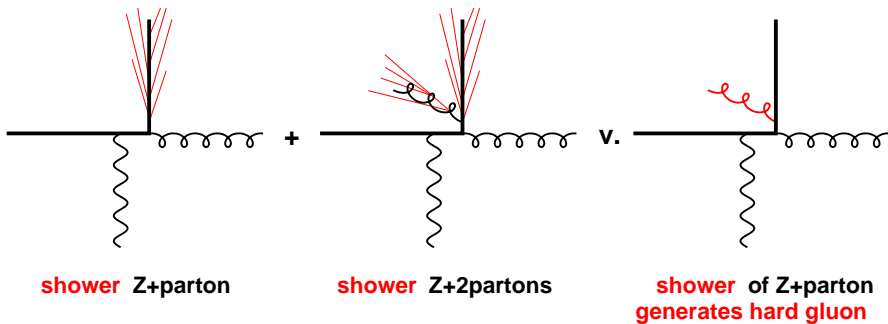
shower Z+parton

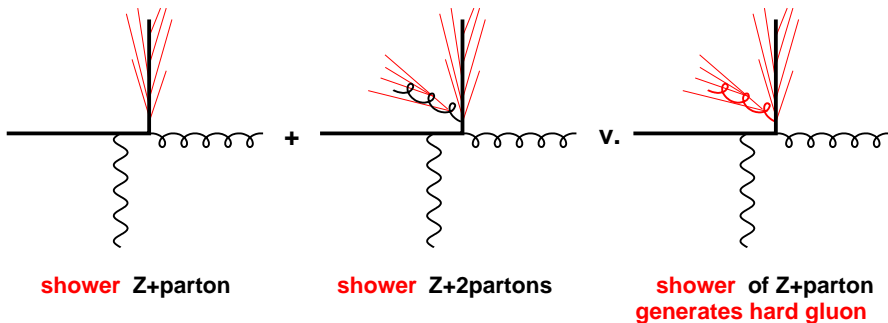


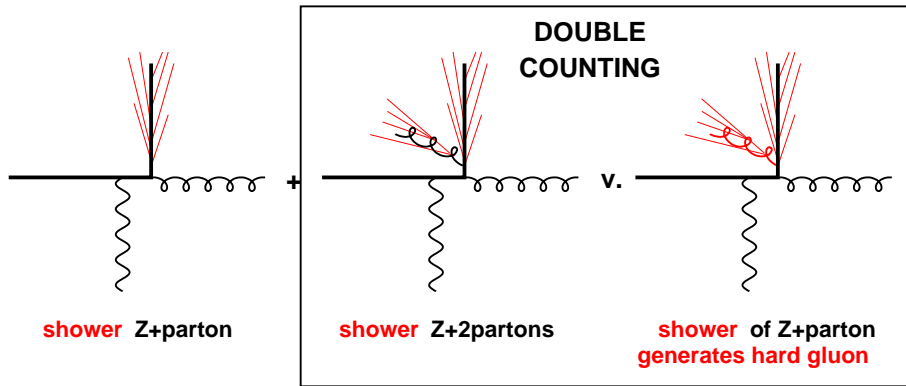


shower Z+parton

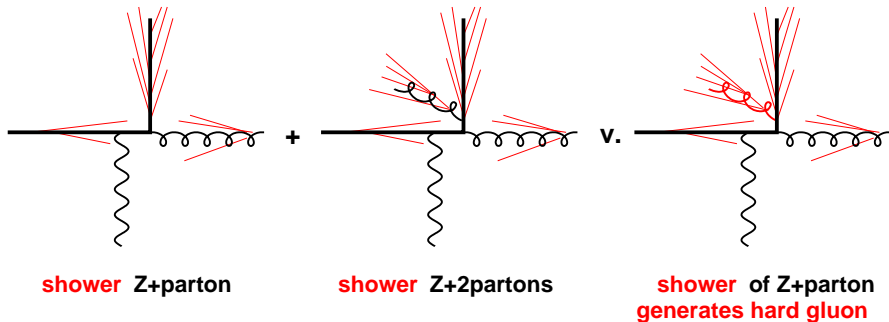
shower Z+2partons







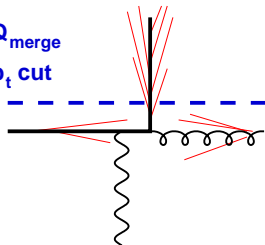
Z + parton implicitly includes part of Z + 2 partons
It's just that the 2nd parton isn't always explicitly "visible"



- ▶ MLM merging relies on parton shower to help figure out what fraction of $Z + \text{parton}$ is really $Z + 2 \text{ partons}$.
- ▶ Our aim is to do that without the parton shower

ACCEPT

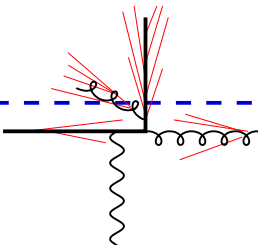
Q_{merge}
 $p_t \text{ cut}$



shower $Z+\text{parton}$

ACCEPT

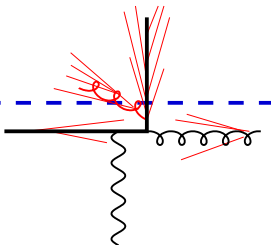
+



shower $Z+2\text{partons}$

v.

REJECT



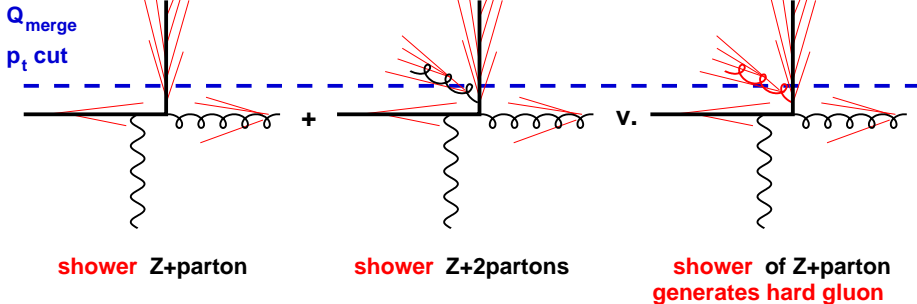
shower of $Z+\text{parton}$
generates hard gluon

- ▶ MLM merging relies on parton shower to help figure out what fraction of $Z + \text{parton}$ is really $Z + 2 \text{ partons}$.
- ▶ Our aim is to do that without the parton shower

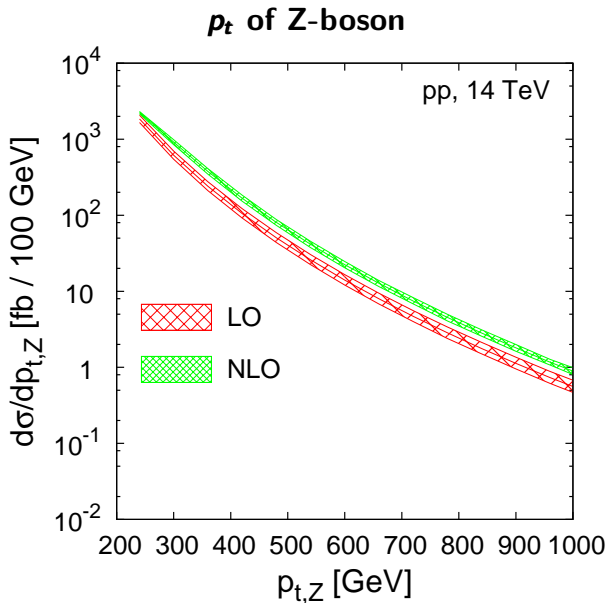
ACCEPT

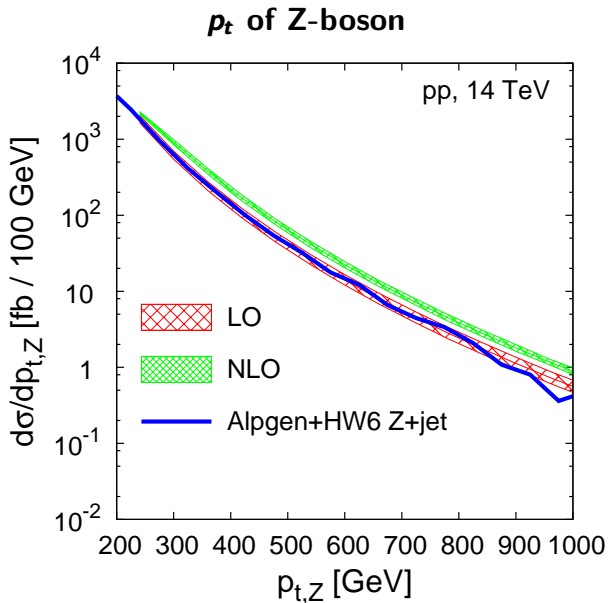
ACCEPT

REJECT

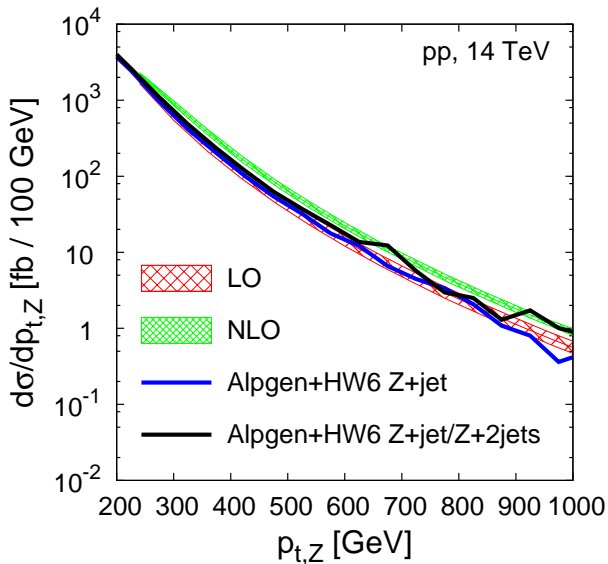


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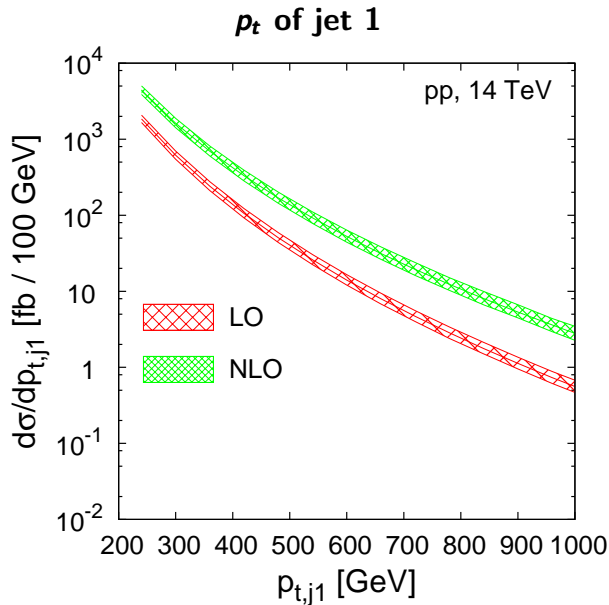


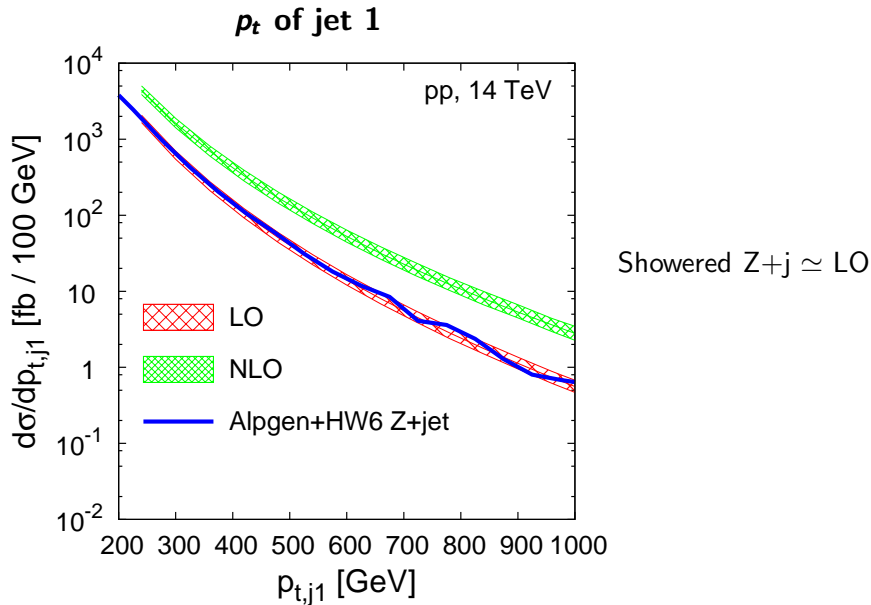


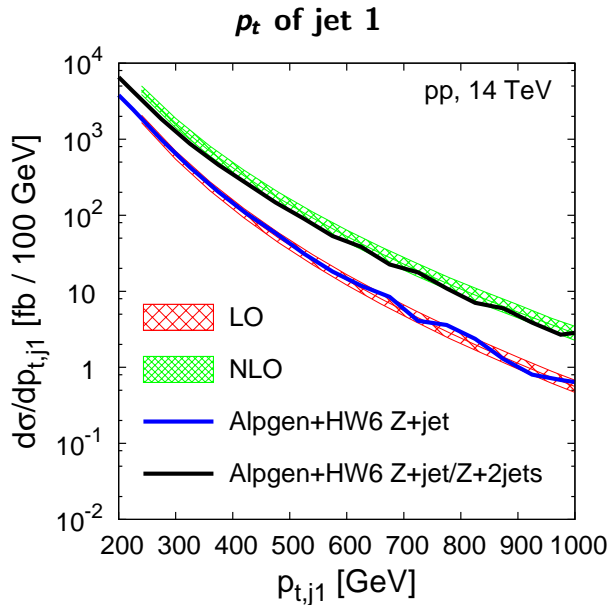
p_t of Z-boson



All predictions similar and stable

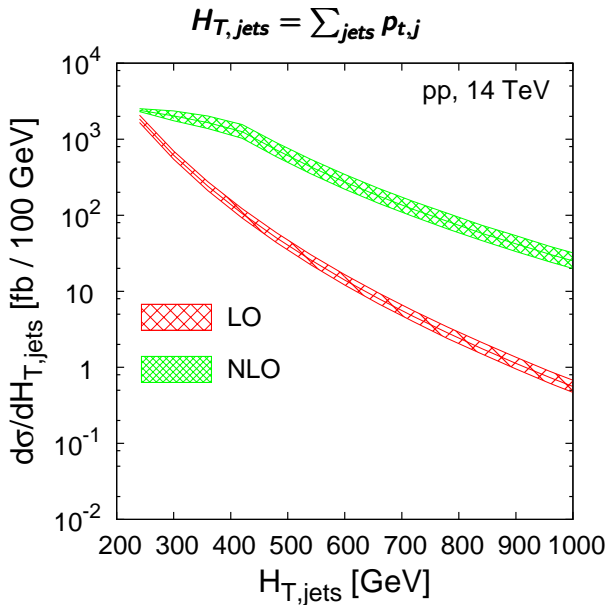


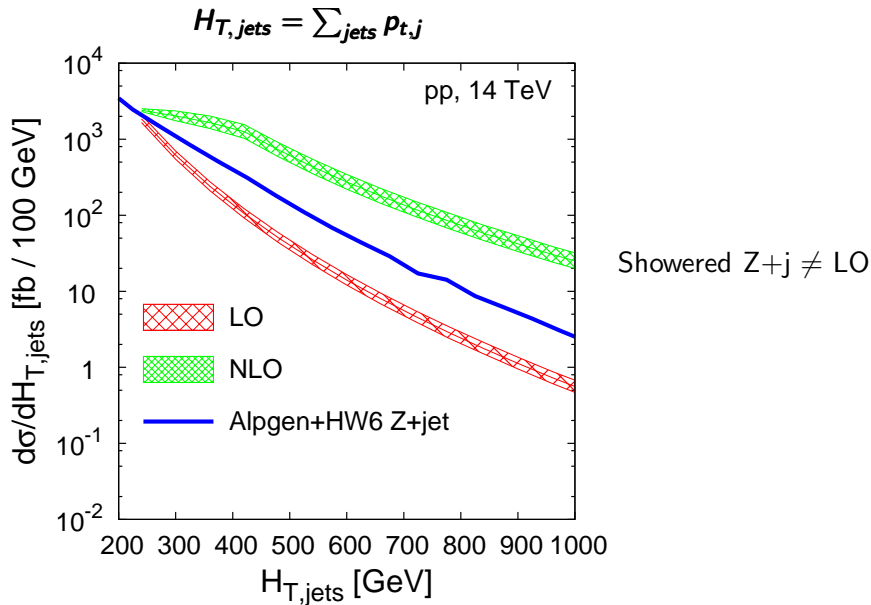


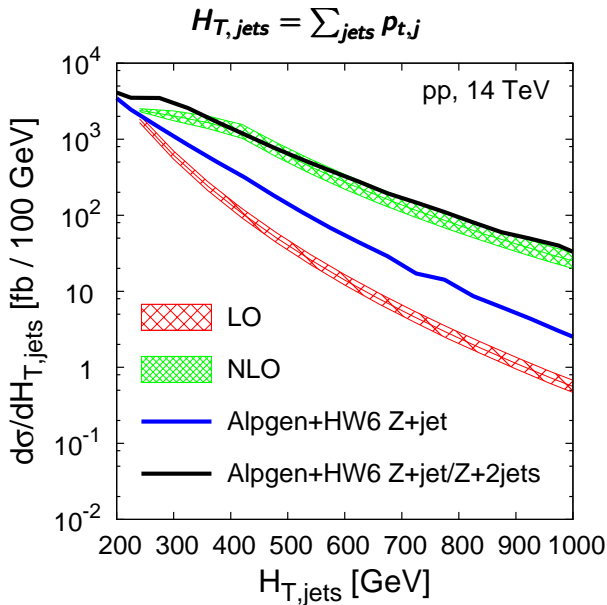


Showered Z+j \simeq LO

Showered Z+j/Z+2j
 \simeq NLO







Showered Z+j \neq LO

Showered Z+j/Z+2j
 \simeq NLO