

Unitarity-based methods and Integration-by-parts identities for Feynman integrals

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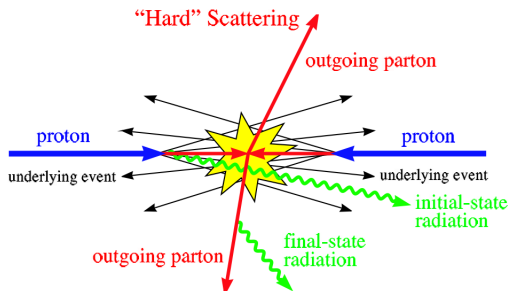


Università degli studi di Padova

19 December 2017

- 1 Scattering amplitudes
 - Phenomenology and theory
- 2 Feynman integrals
 - Integration-by-parts-identities
 - Current status
 - Open problems
 - Limitations and difficulties
- 3 A possible solution: PARSIVAL
 - Unitarity-based methods
 - Unitarity and Integration-by-parts-identities
 - New algorithm
 - PARSIVAL
 - Applications and Tests
- 4 Conclusions and Outlook

High-energy Particle Physics

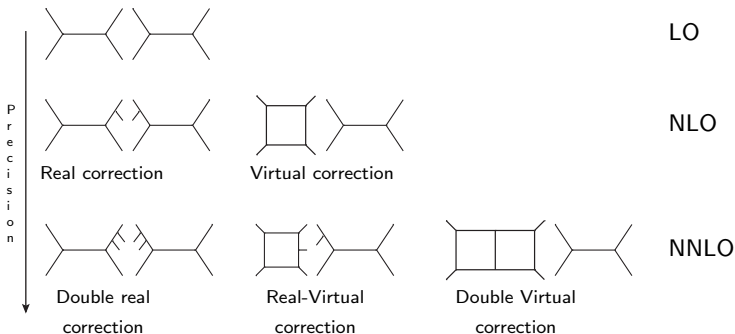
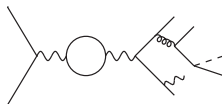
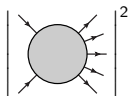


- Discovery of *Higgs boson* \implies Success of the *Standard Model* (SM)
- Evidences suggest that SM is not a complete theory (neutrino mass, Dark Matter etc.)
- New physics effects carried by massive particles \implies Precision measurements of *cross sections*
- High energies vs. High precisions
 - Direct production of new particles (+ jets) \implies many legs;
 - Indirect searches involving high precision measurements \implies many loops;

Scattering Amplitudes and Feynman diagrams

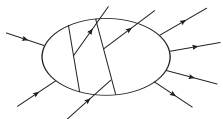
Cross section \leftarrow Scattering amplitudes \xleftarrow{QFT} Feynman diagrams

$$\frac{d\sigma}{d\Omega}$$



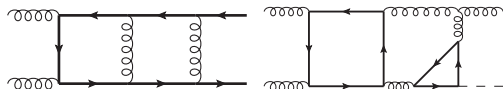
This thesis is focused on the *evaluation of virtual corrections*

Evaluating Feynman Diagrams

$$\mathcal{F}(s) = \text{Diagram} = \int d^d k_1 \cdots d^d k_l \frac{\mathcal{N}(k_1, \dots, k_l, \bar{p})}{\mathcal{D}(k_1, \dots, k_l, \bar{p})}$$


Feynman Diagrams Evaluation: Problems

- Complexity of calculation depends on
 - number of legs;
 - number of free scale parameters;
 - number of loops;
- Evaluate a large number of Feynman integrals;
- Virtual corrections are the realm of integral calculus.



Evaluating Feynman Diagrams

Feynman Integrals Evaluation: Solution

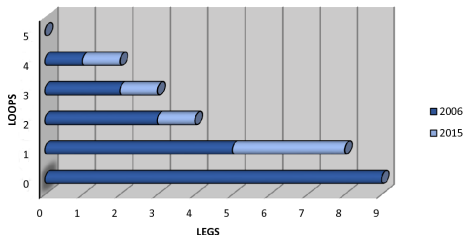
- Passage from Integrals to Algebraic problems
- Automate the calculation

State of Automation

Tree-level \implies solved

1-loop \implies solved

2-loop \implies open problem



Automation of 1-loop diagram evaluation improved using *unitarity-based methods*.

Current evaluation approach

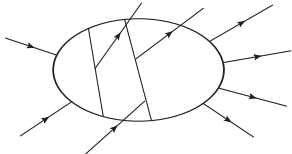
1. Amplitude decomposition
 - IBPs and generalized unitarity
⇒ Decomposition of Feynman diagrams
in terms of *Master Integrals* (MI)
2. Master integrals evaluation
 - Differential and difference equations
 - Direct integration

The core of this procedure are the
Integration-by-parts identities (IBPs)

[K. G. Chetyrkin and F.V. Tkachov (1981)]
[S. Laporta (2001)]
[A. V. Kotikov (1991)]
[Remiddi (1997)]
[Gehrmann Remiddi (1999-2000)]
[Bonciani Mastrolia Remiddi (2003)]
[Bern Dixon Dunbar Kosower (1996)]
[Britto Cachazo Feng (2004)]
[Ossola Papadopoulos Pittau (2006)]

Tensor Decomposition and Feynman Integrals

Factorize the tensorial dependence of Feynman diagrams \implies Preliminary decomposition in *Feynman Integrals*

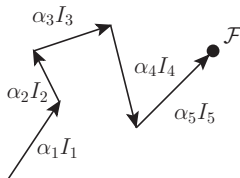
$$\mathcal{F}(\mathbf{s}) = \text{Diagram} = \sum_i \alpha_i I_i(\mathbf{s})$$


Feynman integrals

$$I(\mathbf{s}) = \int d^d k_1 \cdots d^d k_l \frac{\mathcal{N}(k_1, \dots, k_l, \bar{\mathbf{p}})}{D_1^{a_1} \cdots D_t^{a_t}}$$

Observation

- Integrals $I_i(\mathbf{s})$ are related by symmetries
- Evaluate the smallest set of Feynman integrals required by the full calculation



Integration-by-parts identity

$$\int d^d k_1 \cdots d^d k_l \frac{\partial}{\partial k_i^\mu} \left[q_j^\mu \frac{\mathcal{N}(k_1, \dots, k_l, \mathbf{s})}{D_1^{a_1} \cdots D_t^{a_t}} \right] = 0$$

$$q_j^\mu = \{k_j^\mu, p_j^\mu\}$$

[K. G. Chetyrkin and F.V. Tkachov (1981)]

- Comes from:
 - Dimensional regularization;
 - Invariance of Feynman integrals with respect to shifts of loop momenta $k_i \rightarrow A_{ij}k_j + B_{ij}p_j$ ($|\det(A)| = 1$);
- Generates a linear system of equations;
- IBP generation is performed by computer programs (REDUZE, KIRA, FIRE, etc.)

[S. Laporta (2001)]

[A.V. Smirnov (2008)]

[A. von Manteuffel and C. Studerus (2012)]

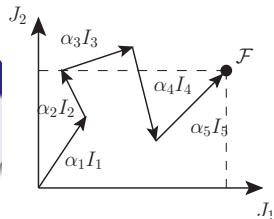
[P. Maierhoefer, J. Usovitsch and P. Uwer (2017)]

Reduction of Feynman integrals

$$\mathcal{F}(\mathbf{s}) = \sum_i \beta_i J_i(\mathbf{s}) \implies \text{Diagram 1} = \sum_i \beta_i \text{Diagram 2}$$

Reduction:

Decomposition of a Feynman integral into a basis of *master integrals*



- Generation of IBPs \implies Identification of MIs, $\bar{J}(\mathbf{s})$;
- MIs \implies Reduction of Feynman integrals

Automating IBP reduction



Issue

The current IBP reduction strategy is *not optimized* for NNLO calculations:

- $2 \rightarrow 3$ NNLO corrections;
- $2 \rightarrow 2$ NNLO corrections with more than four scale parameters;

Idea

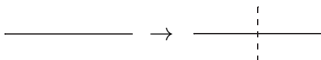
Exploit *Unitarity of Scattering Amplitudes* in order to develop an algorithm aimed at ameliorating the current IBP reduction strategy

Unitarity of the S matrix

$$S^+ S = \mathbb{I} \implies i(T^+ - T) = T^+ T$$

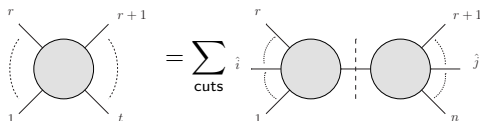
Consequences:

- Generalized Unitarity Cuts

$$\frac{1}{q^2 - m^2 + i\epsilon} \rightarrow 2\pi i \delta(q^2 - m^2),$$


[Z. Bern, L. Dixon, D.C. Dunbar, D.A. Kosower (1994)]






- BCFW tree-level recurrence relation





[R. Britto, F. Cachazo and B. Feng, (2005)]




[R. Britto, F. Cachazo, B. Feng, and E. Witten (2005)]





Unitarity-Based Methods and Reduction






4-point function reduction:  = C_4  + C_3  + C_2  + C_1 

Multiple cuts on
4-point function:

 = C_4 

 = C_4  + C_3 

 = C_4  + C_3  + C_2 

 = C_4  + C_3  + C_2  + C_1 

Generalized unitarity

Use of unitarity cuts \implies Reconstruction of
Feynman integrals
reductions

New idea: Partial fractioning for Feynman integrals

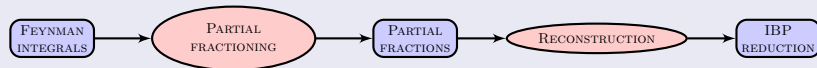
Observation

Partial fractioning \sim Unitarity cuts \implies IBP Reconstruction

Goal of our work

Develop a *new reduction strategy inspired by unitarity of Scattering Amplitudes*
 \implies *Partial fractioning of integrands of Feynman integrals*

Algorithm



Result 1: Partial Fractioning for Feynman integrals

Result 1: We found a **symbolical expression** of the partial fractioning identity

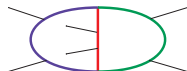
Generalization of Partial Fractioning

$$\frac{1}{D_1^{a_1} D_2^{a_2} \dots D_t^{a_t}} = \sum_{p=1}^t \sum_{i=0}^{a_p-1} \sum_{\{\hat{j}_1, \dots, \hat{j}_p, \dots, \hat{j}_t\}=0}^{\hat{j}_1+\dots+\hat{j}_p+\dots+\hat{j}_t=i} \frac{\prod_{k=1, k \neq p}^t (-1)^{j_k} \binom{a_k+j_k-1}{a_k-1}}{(D_1 - D_p)^{a_1+j_1} \dots D_p^{a_p-i} \dots (D_p - D_t)^{a_t+j_t}}$$

Properties of partial fractions:

- $D_j - D_i$ is *linear in loop momenta*;
- the j -th propagator is represented by a dashed line

$$D_j \rightarrow D_j - D_i \Rightarrow \text{—————} \rightarrow \text{-----}$$



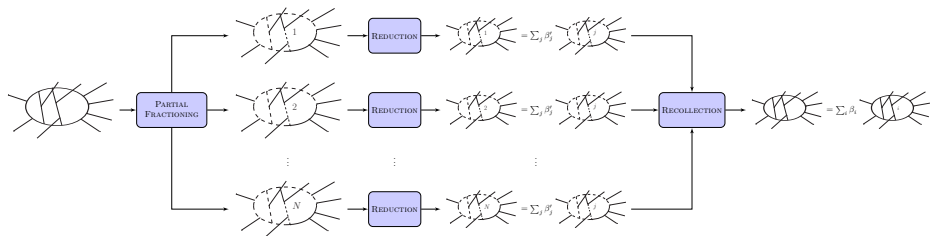
- Integrals of partial fractions dubbed *linearized Feynman integrals*;
- Multi-loop extension \Rightarrow Partial fractioning for every *Branch*.

A diagram showing the partial fractioning of a multi-loop integral. On the left, a loop with several internal lines is shown. This is equated to a sum over \bar{i} of $B_{\bar{i}}$, where each $B_{\bar{i}}$ is a diagram with a dashed line representing the branch \bar{i} .

New idea: Developments of a novel algorithm

Observations

- degree of partial fractions denominators < degree of starting integrand denominators
- The IBP reductions of linearized Feynman integrals can be parallelised



Recollection

Combining Master Integrals found in the reductions of linearized Feynman Integrals

Result 2: Reconstruction of IBP Reduction

IBP Reduction of 1-loop Feynman Integral:

$$\text{Sunset} = C_4 \text{Box} + C_3 \text{Triangle} + C_2 \text{Bubble} + C_1 \text{Line}$$

IBP Reconstruction of 1-loop Feynman Integral with our algorithm:

$$\text{Sunset} = C_4 \text{Box} + C_3 \text{Triangle} + C_2 \text{Bubble} + C_1 \text{Line} + D_1 \text{Triangle} + D_2 \text{Bubble}$$

Spurious terms appear

Linearized Feynman Integrals with no quadratic propagators. In general $D_1 \neq 0$ and $D_2 \neq 0$.

Result 2: Classification of linearized Feynman integrals lead us to the identification of vanishing integrals

Classification

We proven that every *Linearized Feynman Integral made solely of linear denominators vanishes.*

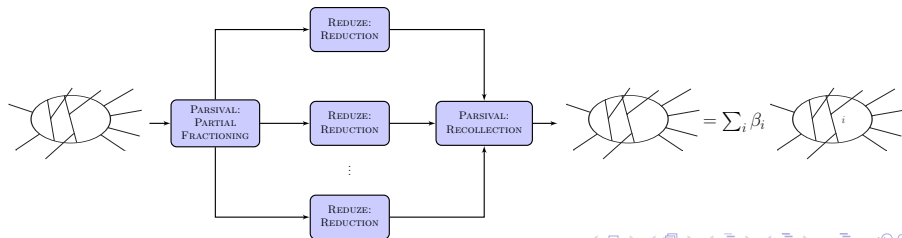
$$\text{Linearized Sunset} = 0$$

PARSIVAL: PARTIAL fraction-based Strategy for Feynman Integral eVALuation

Result 3: We implemented our new algorithm in a MATHEMATICA code, called PARSIVAL.

It automates:

- the decomposition Feynman integrals in linearized Feynman Integrals;
- the parallelisation of IBP reductions of linearized Feynman Integrals;
- the reconstruction of the IBPs of the starting Feynman integral

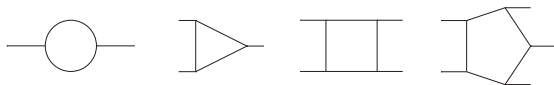


Strength of PARSIVAL

- **Generality:** partial fractioning can be applied to any multi-loop Feynman integral;
- **Flexibility:** it can be interfaced with any available IBP generator.

Tests

- 1-loop Feynman diagrams



- 2-loop Feynman diagrams



PARSIVAL: RUNNING TESTS

Massless planar box 2-loop diagram

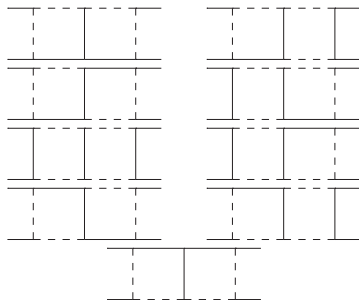
Test with existing routine: REDUZE

Test with our code: PARSIVAL+REDUZE



RUNNING TIME

20399s



RUNNING TIMES

6936s 14080s

5897s 5448s

5497s 16150s

15001s 9842s

6275s

CONCLUSIONS AND OUTLINES

Conclusions:

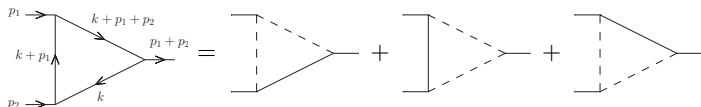
- Novel algorithm
 - The improvement of Feynman integrals evaluation passes through unitarity-based methods;
 - BCFW tree-level recurrence inspired the use of partial fractioning;
 - Partial fractions shares properties with unitarity cuts;
 - From IBP reconstruction we get the right IBP reduction;
 - The procedure exposed is completely general, applicable to any Feynman integral.
- Automation
 - Implementation of the novel algorithm in `PARSIVAL`;
 - 1-loop and simple 2-loop tests shown an improvement of IBP generation time;

Outlook:

- Possibility to improve `PARSIVAL` (e.g. adding symmetries of the diagram);
- Improvements on our new algorithm passes through the understanding of *linearized Feynman integrals*;
- Use of `PARSIVAL` in the evaluation of $2 \rightarrow 2$ and $2 \rightarrow 3$ 2-loop Feynman integrals.

Thank you for the attention!

Example: Partial fractioning for QED 1-loop vertex



$$\int \frac{d^d k}{D_1 D_2 D_3} = \int \frac{d^d k}{D_1 (D_2 - D_1) (D_3 - D_1)} + \int \frac{d^d k}{(D_1 - D_2) D_2 (D_3 - D_2)} + \int \frac{d^d k}{(D_1 - D_3) (D_2 - D_3) D_3}$$

Example

$$D_1 = k^2 - m_1^2, \quad D_2 = (k + p_1)^2 - m_2^2, \quad D_3 = (k + p_1 + p_2)^2 - m_3^2$$

$$D_2 - D_1 = p_1 \cdot (2k + p_1) - (m_2^2 - m_1^2)$$

$$D_3 - D_1 = (p_1 + p_2) \cdot (2k + p_1 + p_2) - (m_3^2 - m_1^2)$$

$$D_3 - D_2 = p_2 \cdot (2k + 2p_1 + p_2) - (m_3^2 - m_2^2)$$

Action of generalized unitarity

1-loop bubble diagram:

$$\text{---} \bigcirc \text{---} = \int d^d k \frac{1}{D_1 D_2}$$

Single cut of 1-loop bubble diagram:

$$\text{---} \overset{\uparrow}{\bigcirc} \text{---} = \int d^d k \frac{\delta(D_1)}{D_2|_{D_1=0}} = \int d^d k \frac{\delta(D_1)}{D_2 - D_1}$$

Cuts generates *differences between denominators*

Using a little abuse of notation:

Partial fractioning of 1-loop bubble diagram:

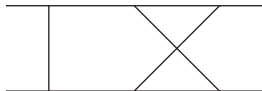
$$\begin{aligned} \text{---} \bigcirc \text{---} &= \int d^d k \frac{1}{D_1(D_2 - D_1)} + \int d^d k \frac{1}{(D_1 - D_2)D_2} \\ &= \int d^d k \frac{1}{D_1} \text{---} \overset{\uparrow}{\bigcirc} \text{---} + \int d^d k \frac{1}{D_2} \text{---} \underset{\downarrow}{\bigcirc} \text{---} \end{aligned}$$

PARSIVAL: RUNNING TESTS

Massless non-planar box 2-loop diagram

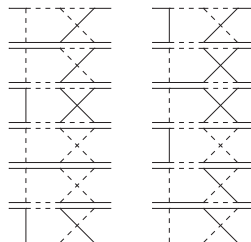
Test with existing routine: REDUZE

Test with our code: PARSIVAL+REDUZE



RUNNING TIME

37235s



RUNNING TIMES

3963s 2627s

11023s 6589s

17967s 7124s

4924s 12570s

4929s 18957s

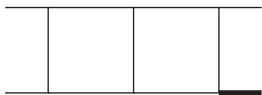
4154s 7850s

PARSIVAL: RUNNING TESTS

Massless planar box 2-loop diagram - 1 massive external leg

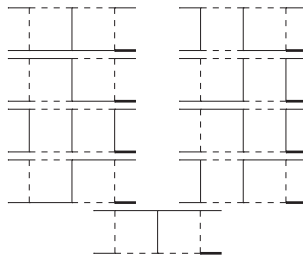
Test with existing routine: REDUZE

Test with our code: PARSIVAL+REDUZE



RUNNING TIME

38722s



RUNNING TIMES

5608s 11795s

6890s 61770s

79043s 73513s

44039s 54455s

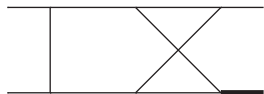
74978s

PARSIVAL: RUNNING TESTS

Massless non-planar box 2-loop diagram - 1 massive external leg

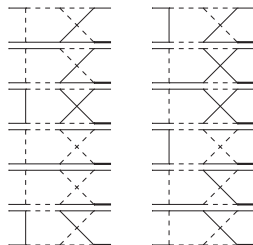
Test with existing routine: REDUZE

Test with our code: PARSIVAL+REDUZE



RUNNING TIME

107277s



RUNNING TIMES

50571s	63103s
83474s	104584s
103021s	111266s
91578s	75797s
84110s	94146s
56200s	58803s