

Update on top CPV couplings and EFT analysis

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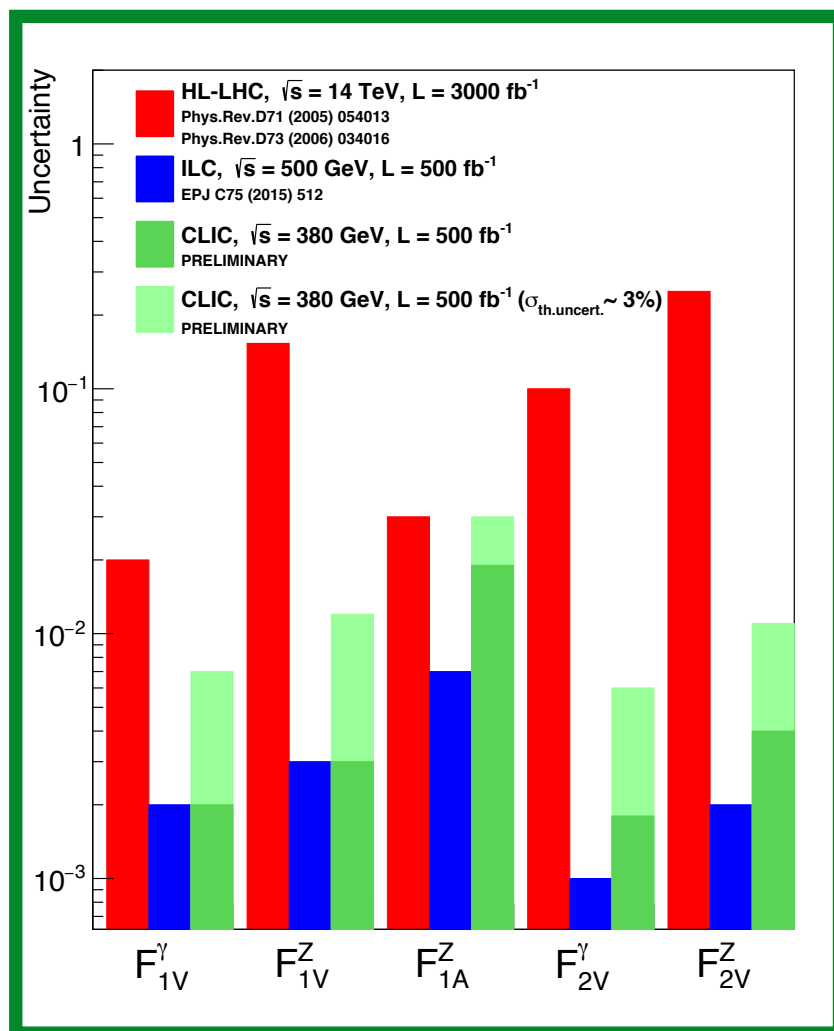
CLICdp WG Analysis Meeting 26/09/17

Top quark electroweak couplings

$$\Gamma_{\mu}^{ttX}(k^2) = -ie \left\{ \gamma_{\mu} \left(F_{1V}^X(k^2) + \gamma_5 F_{1A}^X(k^2) \right) + \frac{\sigma_{\mu\nu} k^{\nu}}{2m_t} \left(i F_{2V}^X(k^2) + \gamma_5 F_{2A}^X(k^2) \right) \right\}$$

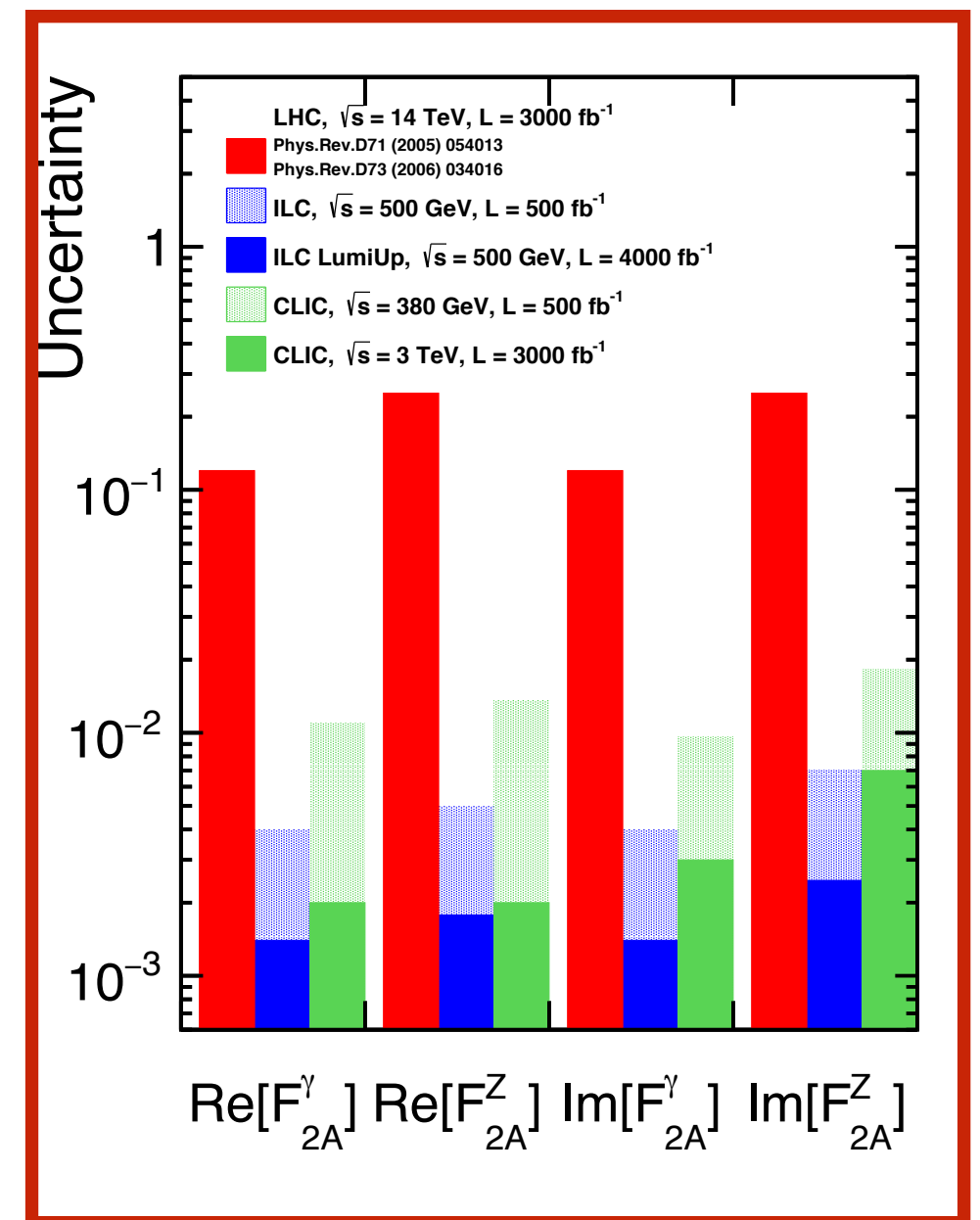
CP-conserving couplings

CP-violating couplings

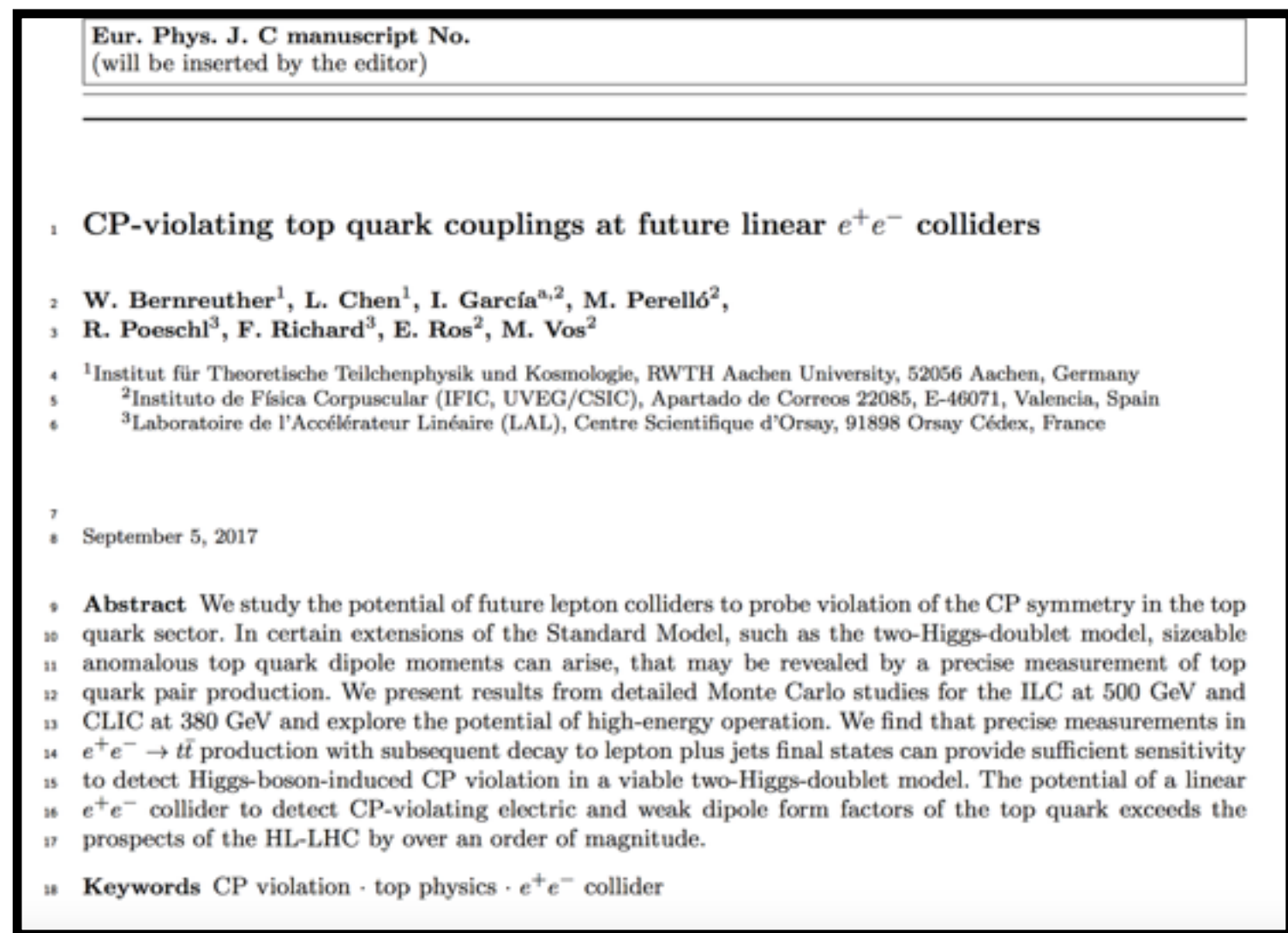


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s10052-015-3746-5

Future e+e- colliders can measure CP-conserving top quark electroweak couplings with a precision that exceeds that of the HL-LHC



CPV paper status



Useful comments from CLIC people (see all the discussion in CDS):

- **Lucie Linssen**

- low-energy stage -> initial stage when referring to 380GeV
- changed the phrasing of the positron polarization to a positive variant
- ...

- **Matthias Artur Weber**

- Introduction of the CP violation as charge conjugation and parity (CP) violation.
- ...

- **Jan Henryk Kalinowski**

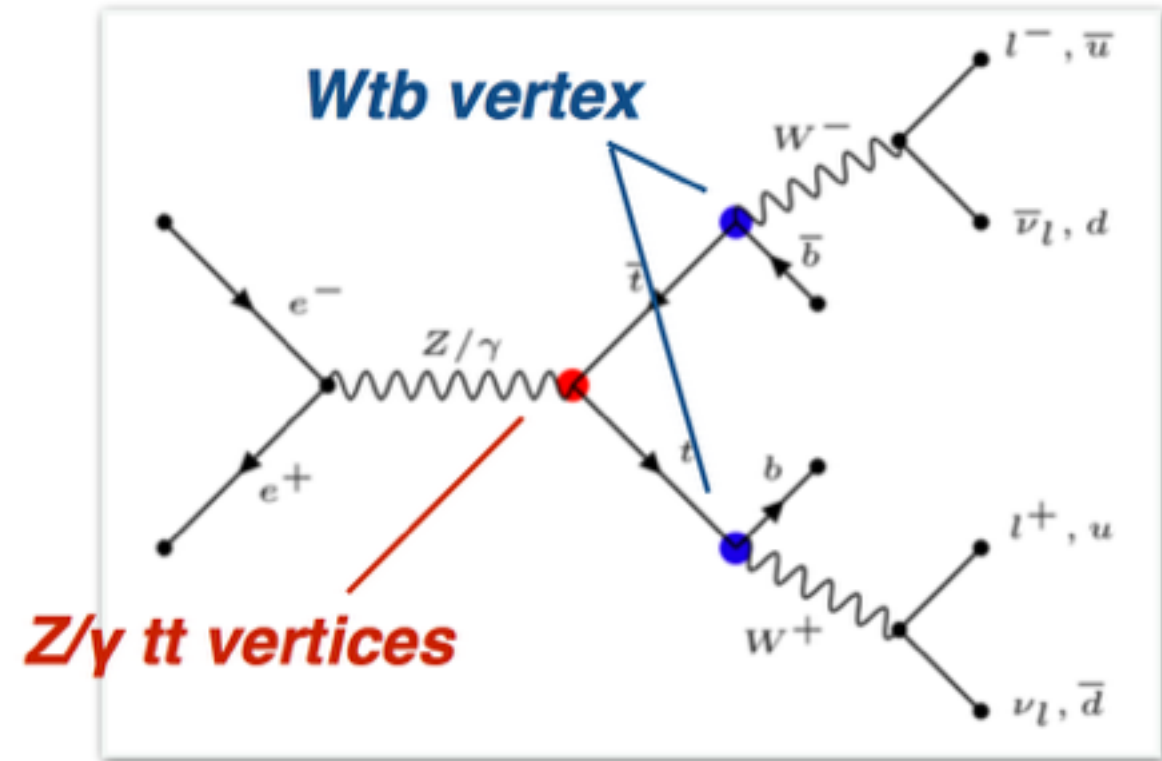
- some redefinitions and writing suggestions
- ...

**thanks for all the
suggestions!!**

Change to an EFT approach

$$\begin{aligned}
 O_{\varphi q}^1 &\equiv \frac{y_t^2}{2} \bar{q} \gamma^\mu q \varphi^\dagger i \overleftrightarrow{D}_\mu \varphi \\
 O_{\varphi q}^3 &\equiv \frac{y_t^2}{2} \bar{q} \tau^I \gamma^\mu q \varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi \\
 O_{\varphi u} &\equiv \frac{y_t^2}{2} \bar{u} \gamma^\mu u \varphi^\dagger i \overleftrightarrow{D}_\mu \varphi \\
 O_{\varphi ud} &\equiv \frac{y_t^2}{2} \bar{u} \gamma^\mu d \varphi^T \epsilon i D_\mu \varphi
 \end{aligned}$$

$$\begin{aligned}
 O_{uG} &\equiv y_t g_s \bar{q} T^A \sigma^{\mu\nu} u \epsilon \varphi^* G_{\mu\nu}^A \\
 O_{uW} &\equiv y_t g_W \bar{q} \tau^I \sigma^{\mu\nu} u \epsilon \varphi^* W_{\mu\nu}^I \\
 O_{dW} &\equiv y_t g_W \bar{q} \tau^I \sigma^{\mu\nu} d \epsilon \varphi^* W_{\mu\nu}^I \\
 O_{uB} &\equiv y_t g_Y \bar{q} \sigma^{\mu\nu} u \epsilon \varphi^* B_{\mu\nu}
 \end{aligned}$$

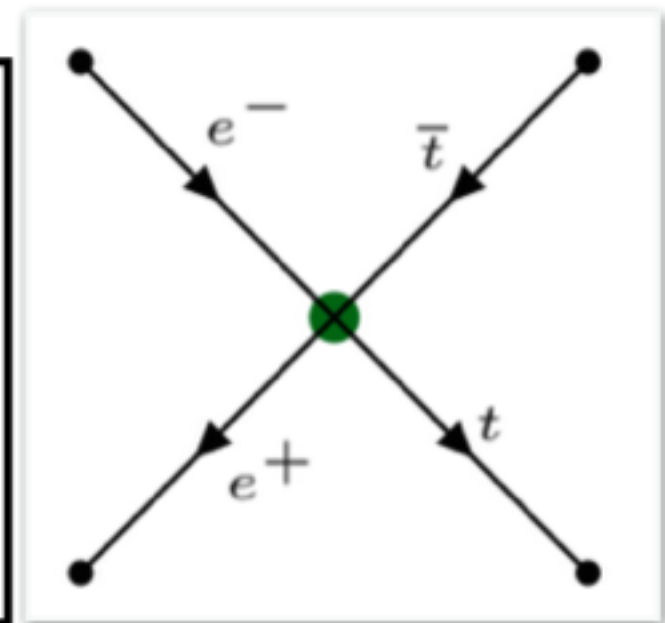


$$\begin{aligned}
 O_{lq}^1 &\equiv \bar{q} \gamma_\mu q \bar{l} \gamma^\mu l \\
 O_{lq}^3 &\equiv \bar{q} \tau^I \gamma_\mu q \bar{l} \tau^I \gamma^\mu l \\
 O_{lu} &\equiv \bar{u} \gamma_\mu u \bar{l} \gamma^\mu l \\
 O_{eq} &\equiv \bar{q} \gamma_\mu q \bar{e} \gamma^\mu e \\
 O_{eu} &\equiv \bar{u} \gamma_\mu u \bar{e} \gamma^\mu e
 \end{aligned}$$

Contact interactions

$$O_{lequ}^T \equiv \bar{q} \sigma^{\mu\nu} u \epsilon \bar{l} \sigma_{\mu\nu} e$$

$$\begin{aligned}
 O_{lequ}^S &\equiv \bar{q} u \epsilon \bar{l} e \\
 O_{ledq} &\equiv \bar{d} q \bar{l} e
 \end{aligned}$$



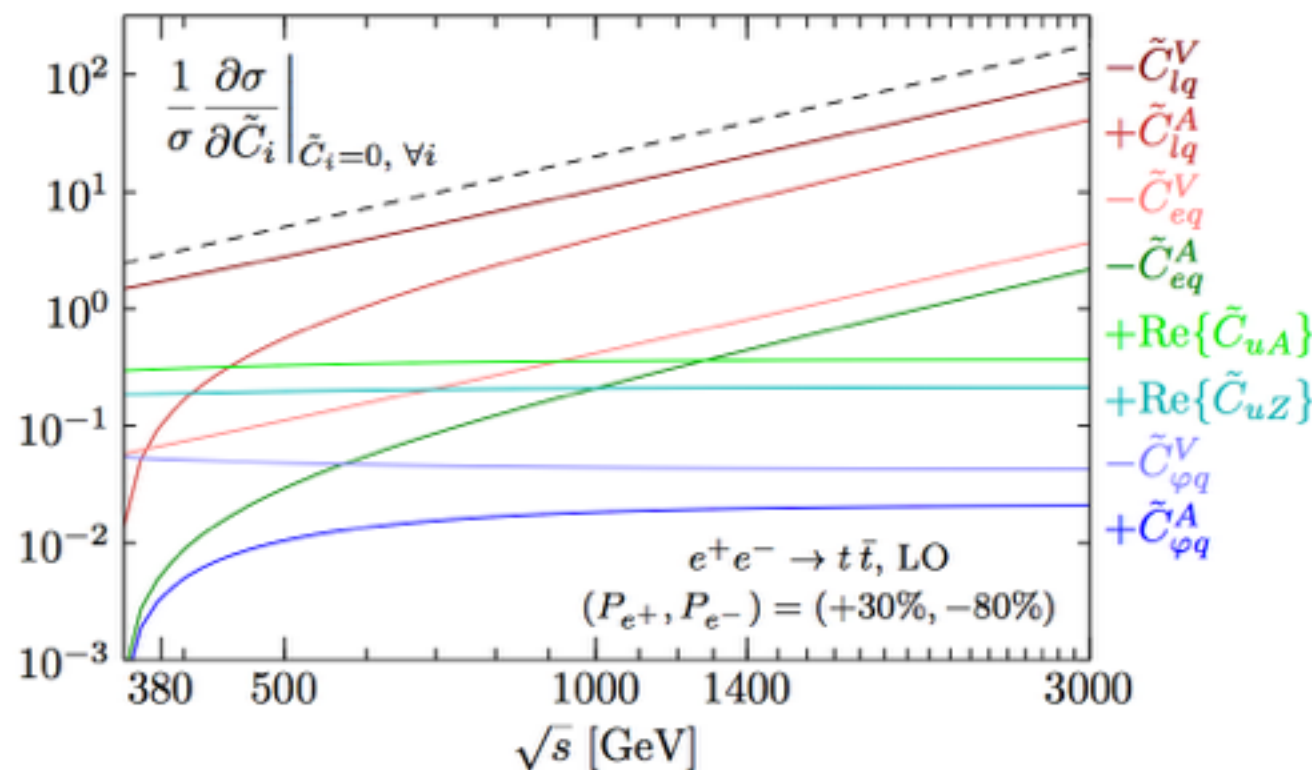
Typical procedure...

Observables sensitivity

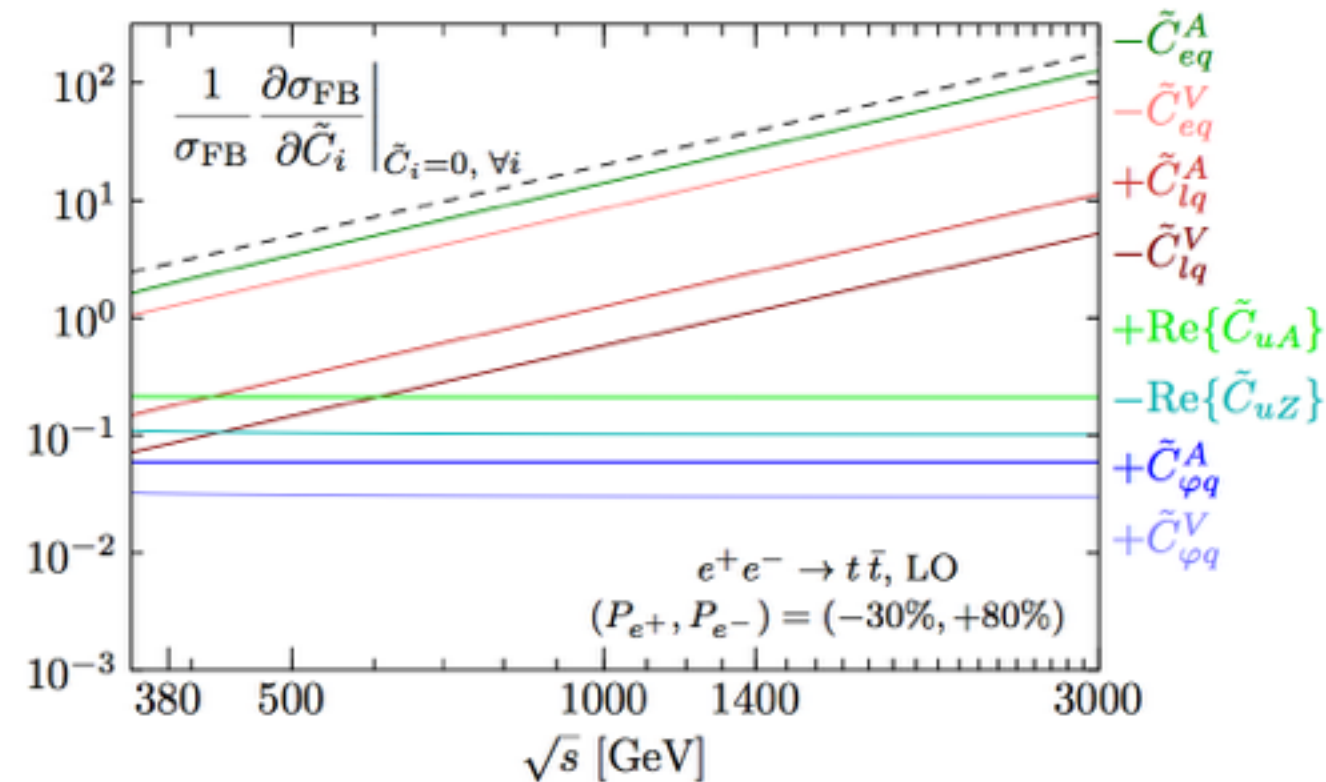
$$e^+e^- \rightarrow t\bar{t}, \text{ LO}$$

Durieux, Perelló, Vos, Zhang, to be published

Cross-section



Forward-backward asymmetry



Sensitivity:

Relative change in cross-section due to non-zero operator coefficient
 $\Delta\sigma(C) / \sigma / \Delta C$

(multi-) TeV operation provides better sensitivity to contact-interaction operators.

Full-simulation

Studies included in I. Garcia thesis

ILC@500GeV L=500fb⁻¹ [arXiv:1505.06020]

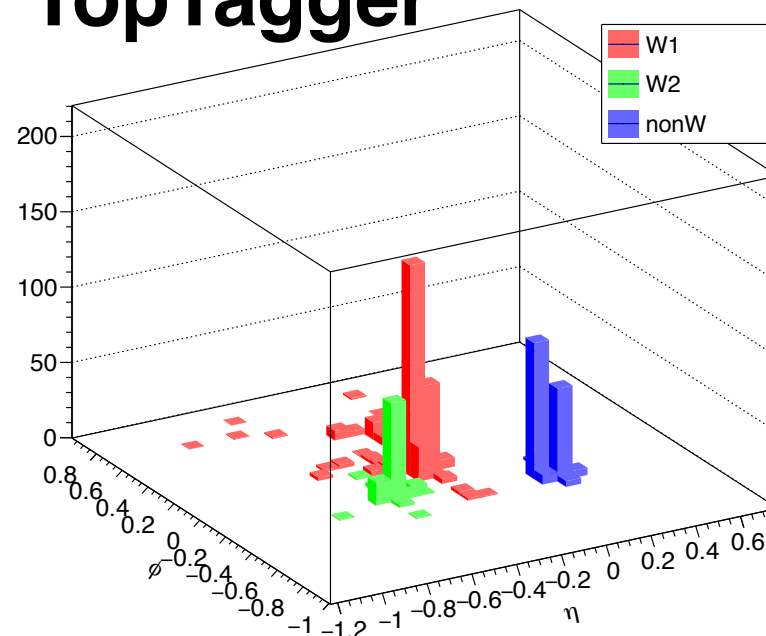
$\mathcal{P}_{e^-}, \mathcal{P}_{e^+}$	$(\delta\sigma/\sigma)_{\text{stat.}} (\%)$	$(\delta A_{\text{FB}}^t/A_{\text{FB}}^t)_{\text{stat.}} (\%)$
-0.8, +0.3	0.47	1.8
+0.8, -0.3	0.63	1.3

CLIC@380GeV L=500fb⁻¹

$\mathcal{P}_{e^-}, \mathcal{P}_{e^+}$	$(\delta\sigma/\sigma)_{\text{stat.}} (\%)$	$(\delta A_{\text{FB}}^t/A_{\text{FB}}^t)_{\text{stat.}} (\%)$
-0.8, 0	0.47	3.8
+0.8, 0	0.83	4.6

At higher energies...

TopTagger



CLIC@1.4TeV L=1500fb⁻¹

Results	$\mathcal{P}(e^-) = -80\%$	$\mathcal{P}(e^-) = +80\%$
ΔA_{FB}^t	0.011	0.015
$\frac{\Delta A_{\text{FB}}^t}{A_{\text{FB}}^t}$	0.025	0.028

Numbers not updated!!

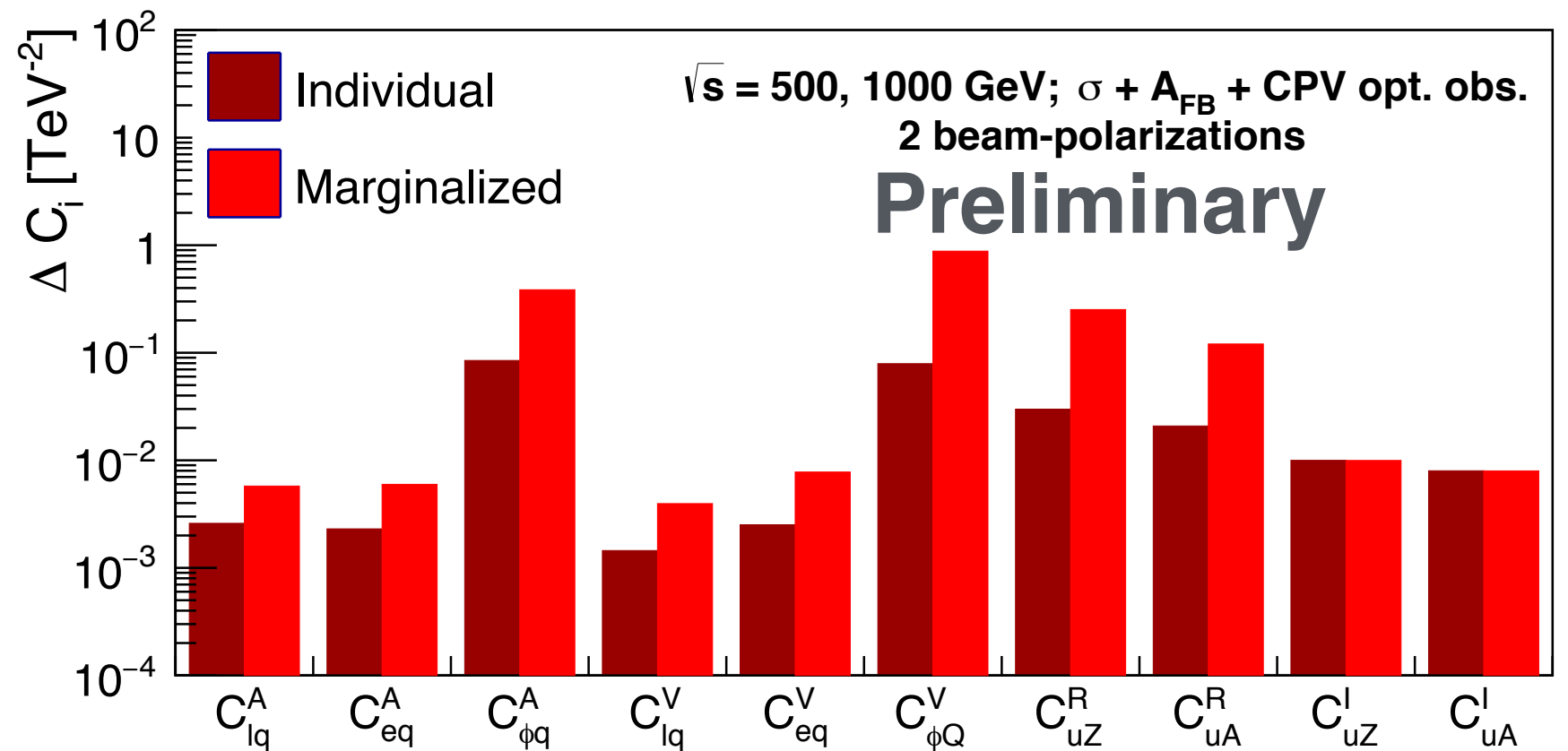
For signal only, w/o scaling or bkg. subtraction

Thanks to Rickard

Fast-simulation (luminosity scaling) for **ILC@1TeV** and **CLIC@3TeV**

M.P. @TopLC2017

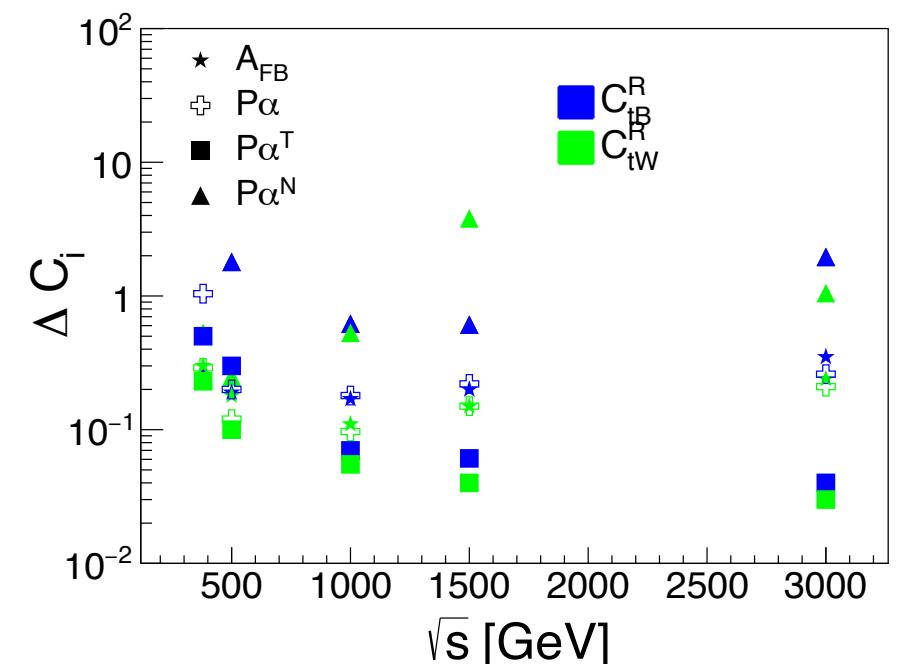
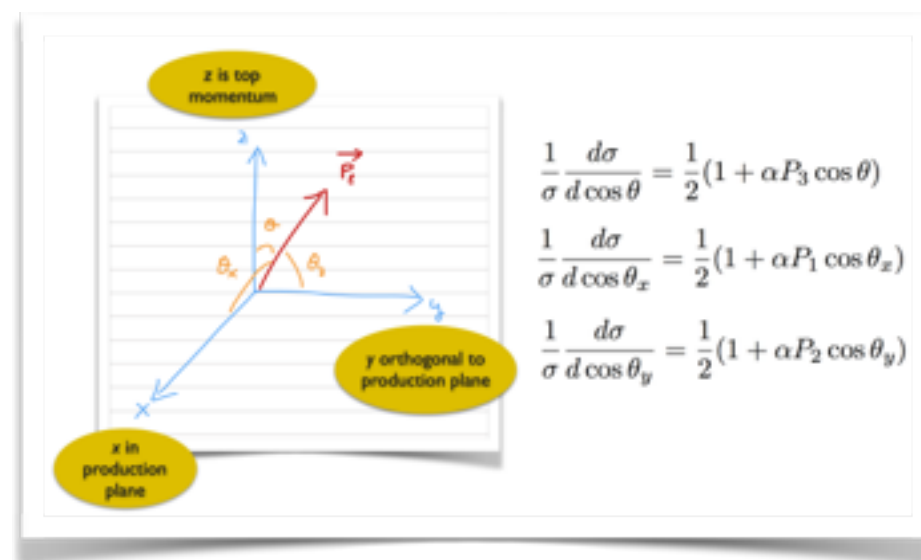
**Preliminary
status...**



**Looking for extra observables
to improve the fit**

i.e. top
polarisation in
different axes

motivation from JA Aguilar



G. Durieux @TopLC 2017:

<https://indico.cern.ch/event/595651/contributions/2573918/attachments/1473086/2280215/durieux-top-lc-2017.pdf>

Statistically optimal observables

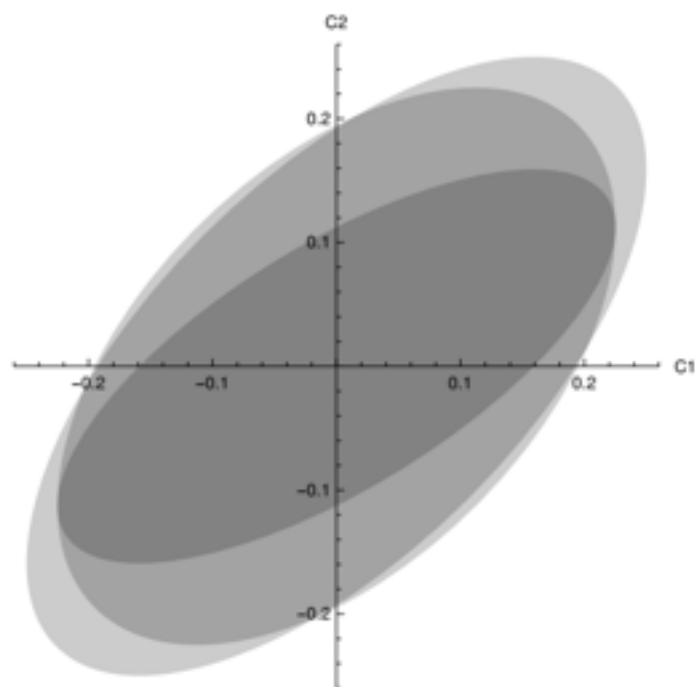
[Atwood, Soni '92]

[Diehl, Nachtmann '94]

minimize the one-sigma ellipsoid in EFT parameter space.

(joint efficient set of estimators, saturating the Rao-Cramér-Fréchet bound: $V^{-1} = I$)

For small C_i , with a phase-space distribution $\sigma(\Phi) = \sigma_0(\Phi) + \sum_i C_i \sigma_i(\Phi)$,
the statistically optimal set of observables is: $O_i(\Phi) = \sigma_i(\Phi)/\sigma_0(\Phi)$.



e.g. $\sigma(\phi) = 1 + \cos(\phi) + C_1 \sin(\phi) + C_2 \sin(2\phi)$

1. asymmetries: $O_i \sim \text{sign}\{\sin(i\phi)\}$

2. moments: $O_i \sim \sin(i\phi)$

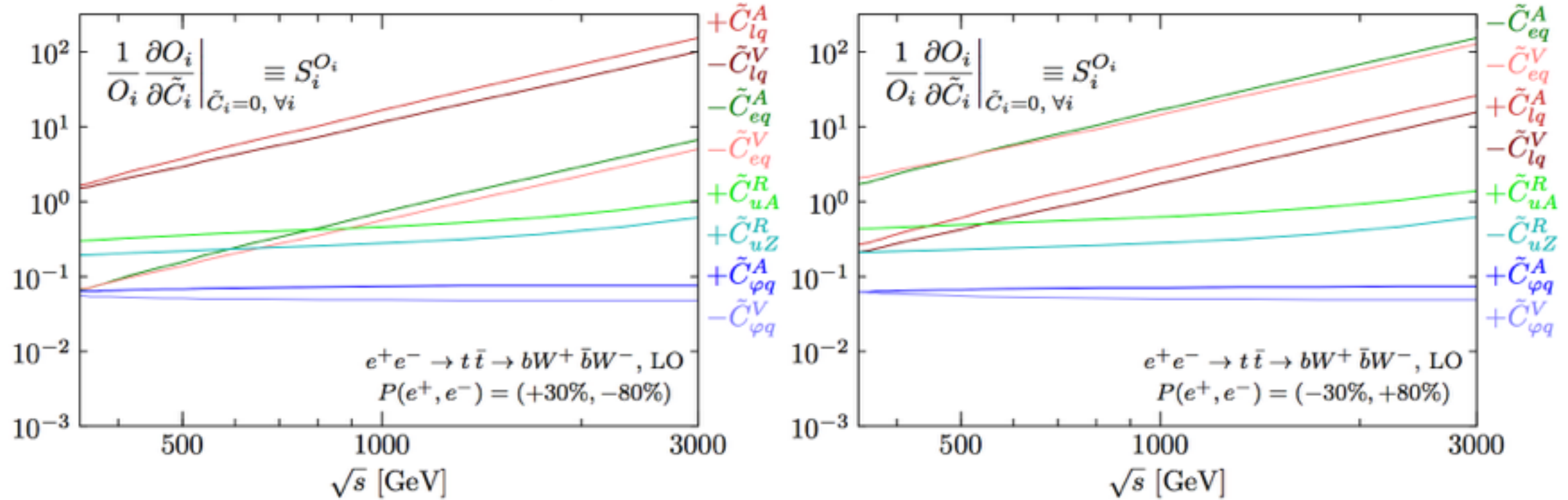
3. statistically optimal: $O_i \sim \frac{\sin(i\phi)}{1 + \cos \phi}$

\Rightarrow area ratios 1.9 : 1.7 : 1

Previous applications in $e^+e^- \rightarrow t\bar{t}$:

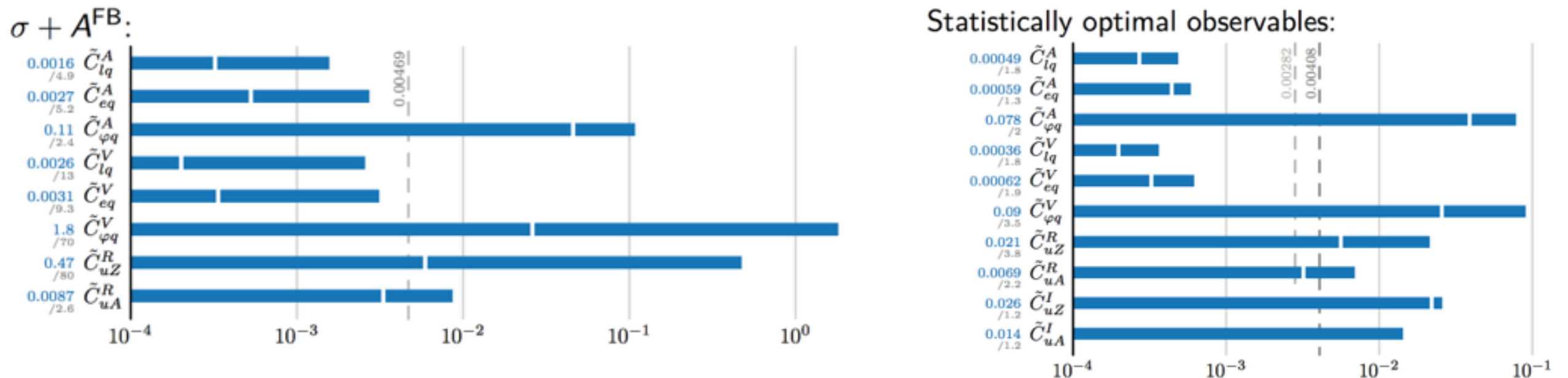
[Grzadkowski, Hioki '00] [Janot '15] [Khiem et al '15]

Statistically optimal observable sensitivities



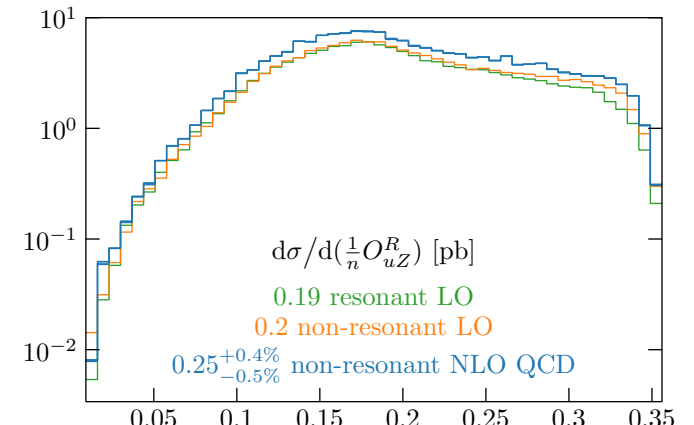
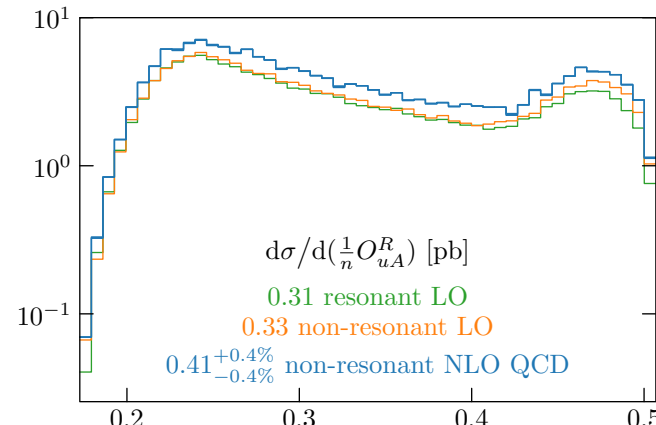
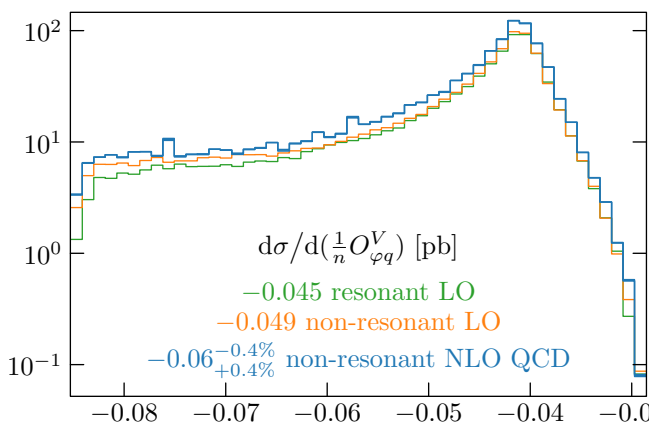
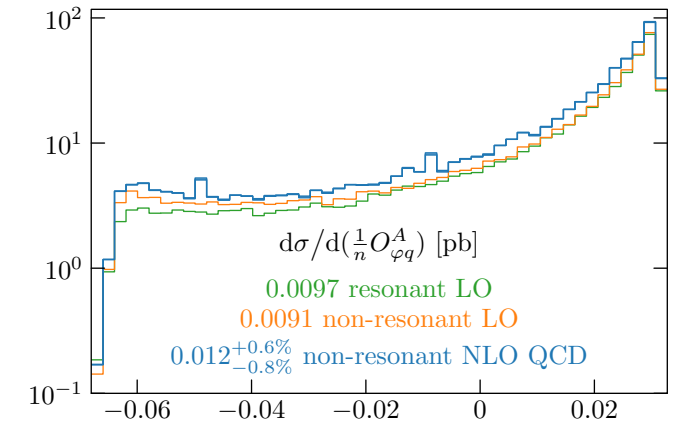
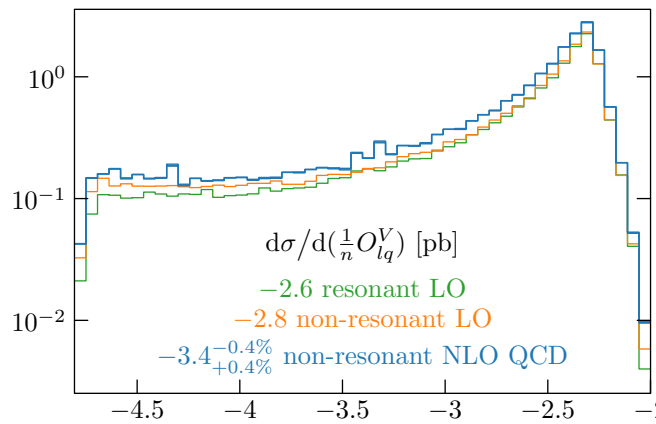
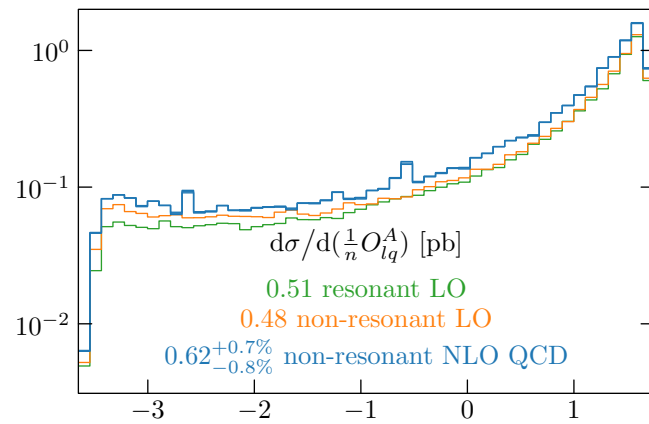
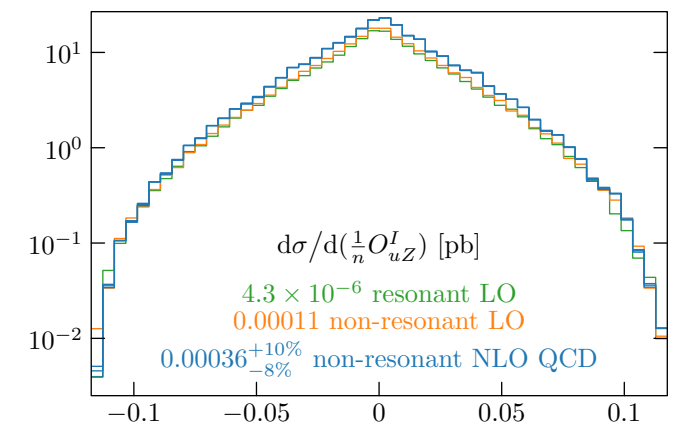
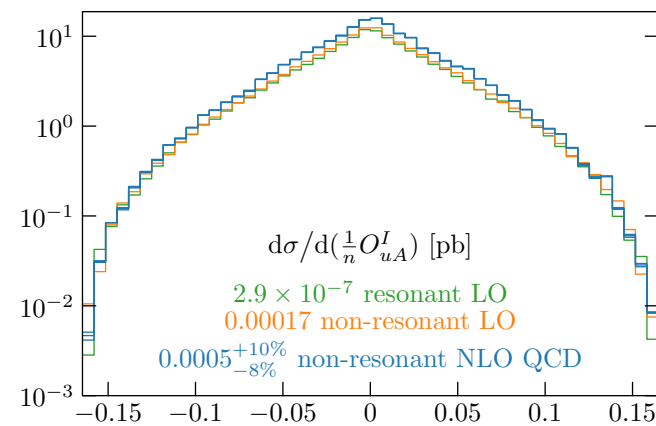
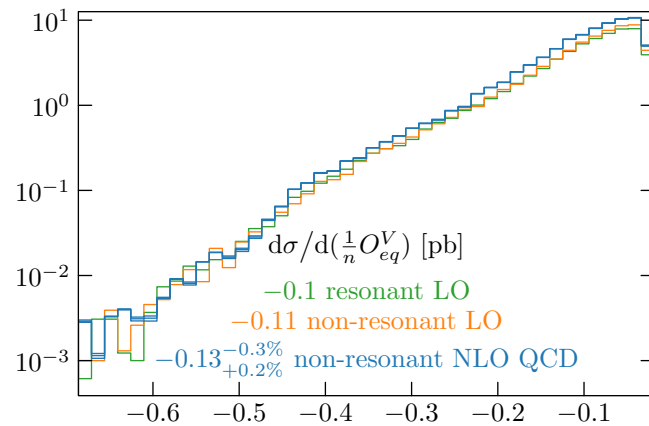
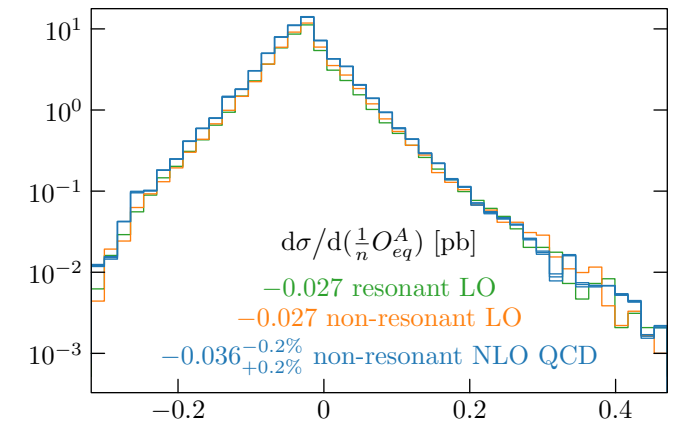
Optimal observables provide much better global limits

Comparison in the global limits (500GeV + 1TeV for 2 pols.):



How the optimal observables look like at 500 GeV (e^-, e^+) = (-0.8, 0.3)

10 observables in total (one per parameter)



Two approaches for the optimal observables systematic studies

Optimal observables

$$O_i = (\sum \sigma_i / \sigma_0)$$

- Depends strongly on the cross-section
- Errors on the shape and normalization

Now in progress

Definition based on the mean of normalized distributions

$$O_i = 1/n (\sum \sigma_i / \sigma_0)$$

- Sensitivities are calculated using real definition
- Errors only on the shape

Results in next slides

Starting reconstruction at CLIC@380 and ILC@500

(Same samples that Nacho used in his studies)

Signal selection:

Same cuts used in previous studies which reduce background.

- **Hadronic top in the range: $120 < m_t < 230$**
- **Hadronic W: $50 < m_W < 110$**
- **only 1 lepton per event**
- **2 btags ($b_{tag1} > 0.8$ and $b_{tag2} > 0.5$)**

Statistical uncertainties:

statistical uncertainty of the distribution mean [%]	lqA	eqA	pqA	lqV	eqV	pqV	ReuZ	ReuA	ImuZ*	ImuA*	*Absolute uncertainty
380 (e-,e+) = (-0.8, 0)	4	6	4	0,1	0,6	0,1	0,3	0,1	1E-3	2E-3	
380 (e-,e+) = (0.8, 0)	6	4	4	0,5	0,2	0,4	0,3	0,2	2E-3	2E-3	
500 (e-,e+) = (-0.8, 0.3)	2	10	2	0,2	5	0,3	0,3	0,2	2E-3	4E-3	
500 (e-,e+) = (0.8, -0.3)	8	2	2	2	0,5	0,9	0,9	0,3	4E-3	7E-3	

Starting reconstruction at CLIC@380 and ILC@500

(Same samples that Nacho used in his studies)

Reconstruction effects

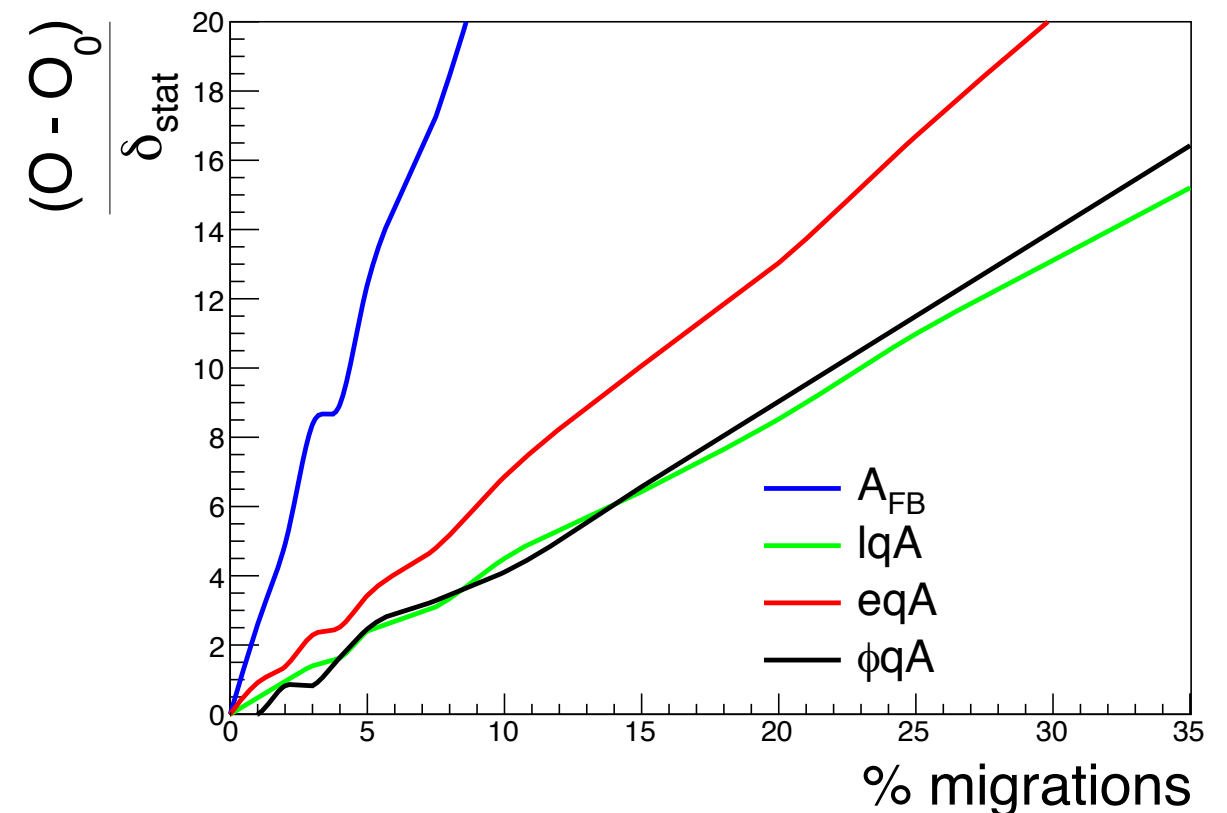
Need of a quality cut
(mainly for reducing
migrations)

numbers are for CLIC380 (I. Garcia's analysis)

$$\chi^2 = \left(\frac{M_t - 172.5}{\sigma_{M_{top}}} \right)^2 + \left(\frac{E_t - 190}{\sigma_{E_{top}}} \right)^2 + \left(\frac{E_b^* - 68}{\sigma_{E_b^*}} \right)^2 + \left(\frac{\cos\theta_{bW} + (-0.6)}{\sigma_{\cos\theta_{bW}}} \right)^2$$

	efficiency	quality cut chi2 < X	efficiency after quality cut
380L	37%	5	18%
380R		40	30,4%
500L	34,4%	50	29,4%
500R	35%	50	30,1%

380 GeV (e-,e+) = (-0.8, 0)



Similar behaviour we observed in the asymmetry.

Systematic uncertainties

(of the distributions means)

selection effect (impact in $\# \sigma$)	lqA	eqA	pqA	lqV	eqV	pqV	ReuZ	ReuA	ImuZ	ImuA
380 (e-,e+) = (-0.8, 0)	2	3	2	1	3	0	2	1	0,2	0,2
380 (e-,e+) = (0.8, 0)	3	0,5	0,1	0,1	0	0	0,5	0	0,4	0,3
500 (e-,e+) = (-0.8, 0.3)	0,8	0,6	0,6	0,4	0,3	0,1	2	2	0,5	0,5
500 (e-,e+) = (0.8, -0.3)	0,4	2	2	0,7	0,8	0,2	0,2	1	1	1

Selection
biases around
 1σ in almost all
cases

reconstruction effect (impact in $\# \sigma$)	lqA	eqA	pqA	lqV	eqV	pqV	ReuZ	ReuA	ImuZ	ImuA
380 (e-,e+) = (-0.8, 0)	2	3	1	2	1	2	1	1	0,1	0,1
380 (e-,e+) = (0.8, 0)	2	2	2	1	2	1	1	2	0,2	0,3
500 (e-,e+) = (-0.8, 0.3)	2	0,3	2	1	0,3	1	0,5	2	0,4	0,5
500 (e-,e+) = (0.8, -0.3)	0,5	2	2	0,2	2	1	2	2	0,4	0,5

Reconstruction
biases around
 1σ in almost all
cases

Beam structure effects (of the distributions means)

Using WHIZARD 2.3.1 for MC generation:

Beamstrahlung switching on/off CIRCE1:

selection effect (impact in $\# \sigma$)	lqA	eqA	pqA	lqV	eqV	pqV	ReuZ	ReuA	ImuZ	ImuA
380 (e-,e+) = (-0.8, 0)	0,4	0,4	0,3	0,3	0,3	0,2	0	0,3	0	0
380 (e-,e+) = (0.8, 0)	0,5	1	1	0	1	0,8	1	0,9	0	0
500 (e-,e+) = (-0.8, 0.3)	1	2	1	1	0,4	1	0,6	1	0	0
500 (e-,e+) = (0.8, -0.3)	0,4	1	1	0,5	1	0,9	1	1	0	0

Parameters variation give
rise to much smaller effects

Switching on/off ISR (using parameters by default):

selection effect (impact in $\# \sigma$)	lqA	eqA	pqA	lqV	eqV	pqV	ReuZ	ReuA	ImuZ	ImuA
380 (e-,e+) = (-0.8, 0)	1	1	1	1	0,1	1	0,7	1	0	0
380 (e-,e+) = (0.8, 0)	1	1	1	0,5	1	0,6	0,7	1	0	0
500 (e-,e+) = (-0.8, 0.3)	2	2	2	3	0,7	3	2	2	0	0
500 (e-,e+) = (0.8, -0.3)	1	2	2	0,7	3	2	3	2	0	0

Low impact of the beam structure

Conclusions

- Optimal observables for global EFT fit are found to be robust

Future work

- EFT paper in preparation
- Complet systematic studies ongoing
- Move to CLIC@1400 (collaboration with Rickard)