

Mass Measurements with Ion Traps for Nuclear Physics

Daniel Rodríguez
Departamento de Física Atómica Molecular y Nuclear
Universidad de Granada
danielrodriguez@ugr.es

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Outline

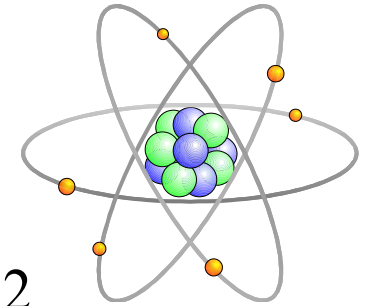
1. General remarks
 - Importance of the mass
 - Physics goals & mass uncertainty
2. The Penning trap
 - Fundamentals
 - Techniques for exotic nuclei
 - Existing facilities & physics survey
 - Mass uncertainty & half-life
 - Recent highlights
3. The project TRAPSENSOR in Granada
 - The technique
 - Implementation
 - First results

GENERAL REMARKS

Importance of the mass

The nuclear binding energy (B)

The nuclear binding energy reflects all forces acting in a nucleus



$$B(N, Z) = \left\{ Zm_p + Nm_n - \underbrace{M(N, Z)}_{\text{Nuclear mass}} \right\} \cdot c^2$$

$$B(N, Z) = \left\{ Zm_p + Nm_n - \left[\underbrace{m(^A X)}_{\text{Atomic mass}} - Zm_e - \sum_i B_i^e \right] \right\} \cdot c^2$$

$$ME(N, Z) = m(^A X) - A \cdot u$$

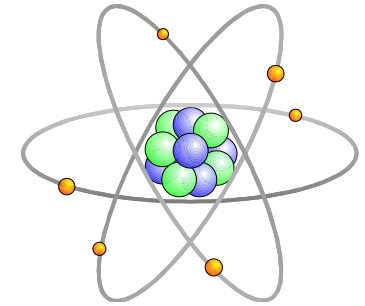
Mass excess (keV)

$$1 \text{ keV} \rightarrow \frac{\delta m}{m} \approx 10^{-8}$$

Importance of the mass

Nucleon separation energies (S_n , S_p , S_{2n} , S_{2p})

Nuclear Structure & astrophysics



$$S_{2n} = B(N, Z) - B(N - 2, Z)$$

$$S_{2p} = B(N, Z) - B(N, Z - 2)$$

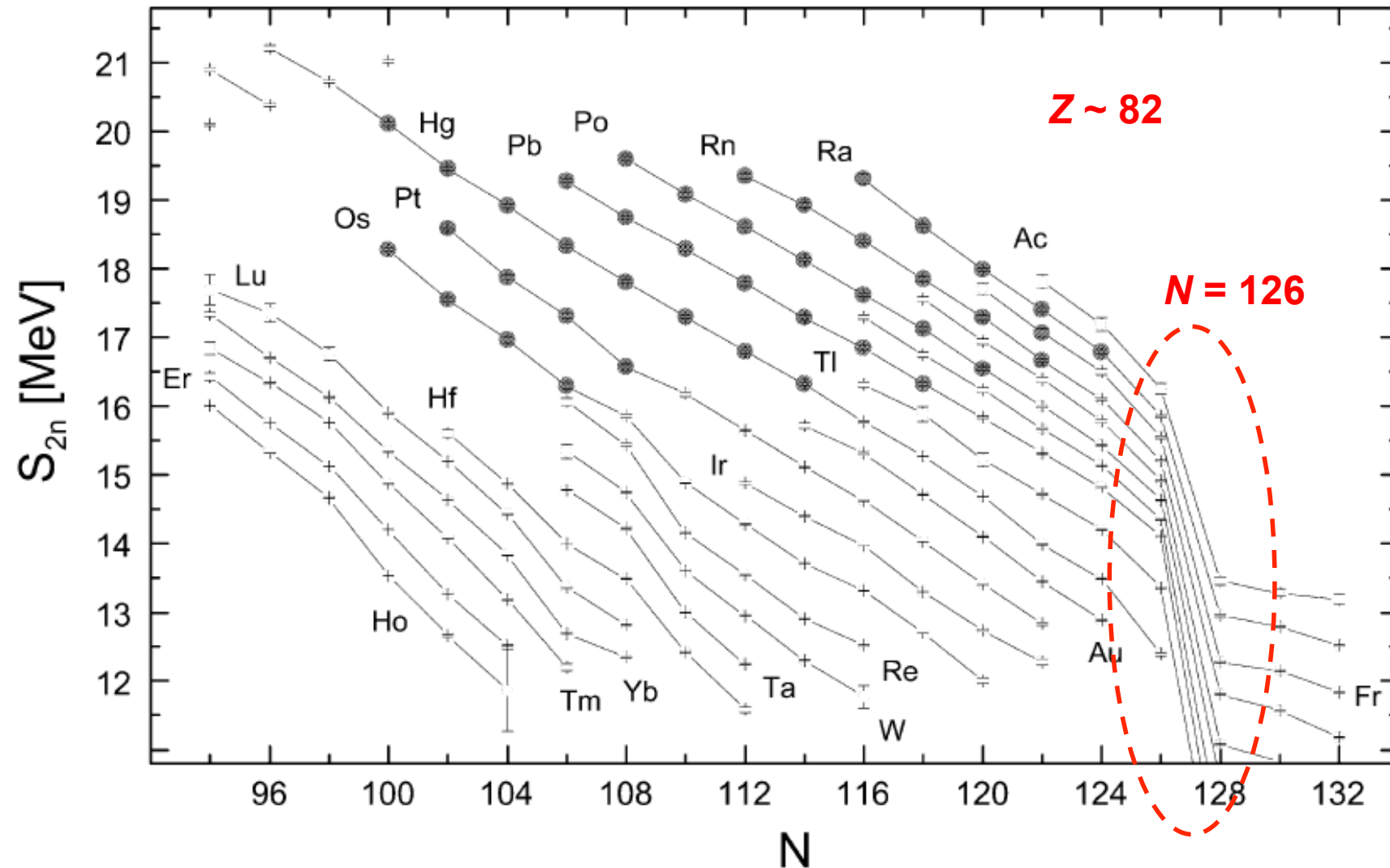
$$S_n = B(N, Z) - B(N - 1, Z)$$

$$S_p = B(N, Z) - B(N, Z - 1)$$

Other quantities are also important for nuclear structure,
all of them obtained from nuclear binding energies

Importance of the mass

Nucleon separation energies (S_n , S_p , S_{2n} , S_{2p})



S. Schwarz *et al.*, Nucl. Phys. A 693 (2001) 533

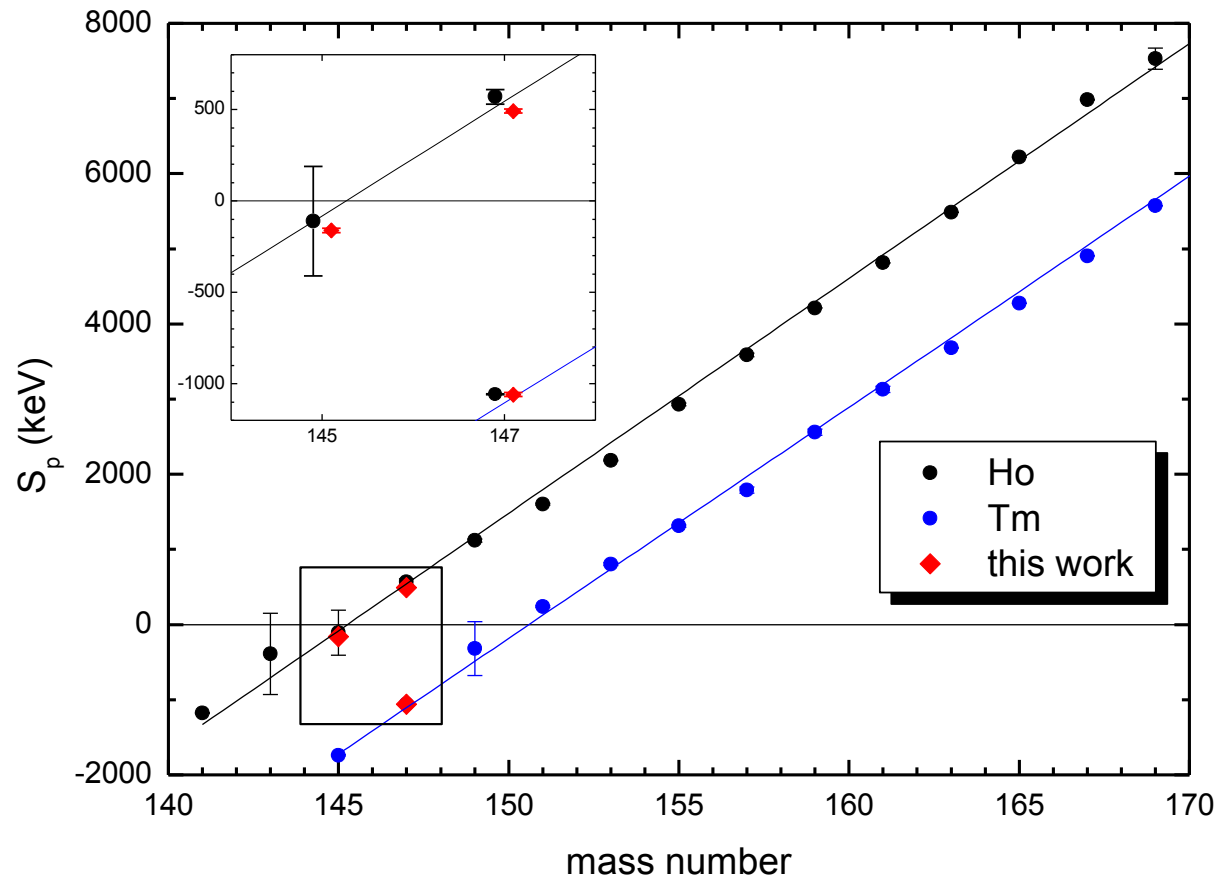
Importance of the mass

Nucleon separation energies (S_n , S_p , S_{2n} , S_{2p})

Location of the proton-drip line

146 69 Tm 77	147 69 Tm 78	148 69 Tm 79	149 69 Tm 80	150 69 Tm 81	151 69 Tm 82
72 ms (10 ⁺) E _{ex} 71 (6) p ⁺ =2 p ⁺ =164%	360 ns 30 ⁺ E _{ex} 63 (5) p ⁺ =100% p ⁺ =100%	500 ns 112 ⁺ M ⁺ 30700 (3000) p ⁺ =15 (51%) p ⁺ =100%	700 ms (10 ⁺) M ⁺ 392700 (4000) p ⁺ =100%	52 ms (10 ⁺) E _{ex} 40 (4) p ⁺ =100% p ⁺ =100%	34 s (1 ⁺) M ⁺ 440400 (3000) p ⁺ =100%
145 68 Er 77	146 68 Er 78	147 68 Er 79	148 68 Er 80	149 68 Er 81	150 68 Er 82
900 ms 1/2 ⁺ M ⁺ 396900 (4000) p ⁺ =100% p ⁺ =2	1.7 s 0 ⁺ M ⁺ 447100 (3000) p ⁺ =100% p ⁺ =2	25 s (112 ⁺) E _{ex} 100 (50) p ⁺ =100% p ⁺ =2	4.6 s 0 ⁺ M ⁺ 516500 (2000) p ⁺ =100% p ⁺ =100%	8.9 s (112 ⁺) E _{ex} 741.8 (8.2) p ⁺ =100% p ⁺ =100%	4 s (10 ⁺) M ⁺ 53742 (20) p ⁺ =100% p ⁺ =2
144 67 Ho 77	145 67 Ho 78	146 67 Ho 79	147 67 Ho 80	148 67 Ho 81	149 67 Ho 82
700 ms M ⁺ 452000 (3000) p ⁺ =100% p ⁺ =2	100 ms 50 ⁺ E _{ex} 100 (100) p ⁺ =2 p ⁺ =100%	24 s (112 ⁺) M ⁺ 48100 (200) p ⁺ =100% p ⁺ =2	3.6 s (10 ⁺) M ⁺ 515700 (2000) p ⁺ =100% p ⁺ =2	5.8 s (112 ⁺) E _{ex} 56837 (28) p ⁺ =100% p ⁺ =2	22 ms (10 ⁺) E _{ex} 969 (6) p ⁺ =100% p ⁺ =2
143 66 Dy 77	144 66 Dy 78	145 66 Dy 79	146 66 Dy 80	147 66 Dy 81	148 66 Dy 82
3.8 s (112 ⁺) E _{ex} 357.7 (8.3) p ⁺ =100% p ⁺ =2	9.1 s 0 ⁺ M ⁺ 56500 (30) p ⁺ =100% p ⁺ =2	14 s (112 ⁺) E _{ex} 112 (8.2) p ⁺ =100% p ⁺ =2	9.6 s (10 ⁺) M ⁺ 6090 (50) p ⁺ =100% p ⁺ =2	150 ms 10 ⁺ E _{ex} 258.7 (8.5) p ⁺ =100% p ⁺ =2	33.2 s 0 ⁺ M ⁺ 5284 (27) p ⁺ =100% p ⁺ =2
142 65 Tb 77	143 65 Tb 78	144 65 Tb 79	145 65 Tb 80	146 65 Tb 81	147 65 Tb 82
300 ms (5 ⁺) E _{ex} 285.2 (1.6) p ⁺ =100% p ⁺ =0.5%	597 ms 1 ⁺ M ⁺ 12000 (1000) p ⁺ =100% p ⁺ =0.022 (11%)	<21 s 5/2 ⁺ E _{ex} 94 (100) p ⁺ =100% p ⁺ =2	12 s (112 ⁺) M ⁺ 10430 (80) p ⁺ =100% p ⁺ =2	28 ms (5 ⁺) E _{ex} 62 (5) p ⁺ =100% p ⁺ =2	426 ms (10 ⁺) E _{ex} 1200 (20) p ⁺ =100% p ⁺ =2

One-proton separation energy for odd-A (even-N) isotopes of Ho (Z=67) and Tm (Z=69):



C. Rauth *et al.*, PRL 100 (2008) 012501

Importance of the mass

Astrophysics

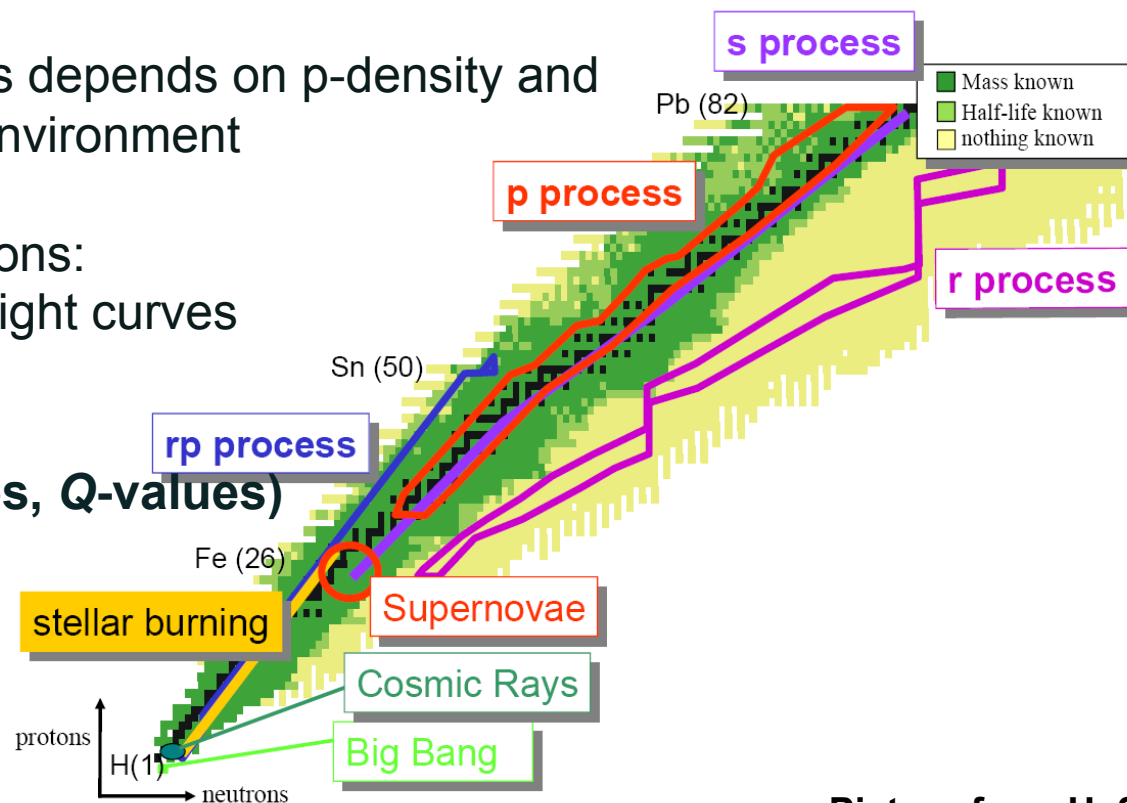
- Nucleosynthesis e.g. on accreting neutron stars
- Explosive hydrogen burning (X-ray bursts)
- Steady-state burning

Evolution of the process depends on p-density and temperature in stellar environment

Astrophysical observations:
elemental abundance, light curves

Required nuclear data:

- **Masses (sep. energies, Q-values)**
- β decay half-lives
- Reaction rates

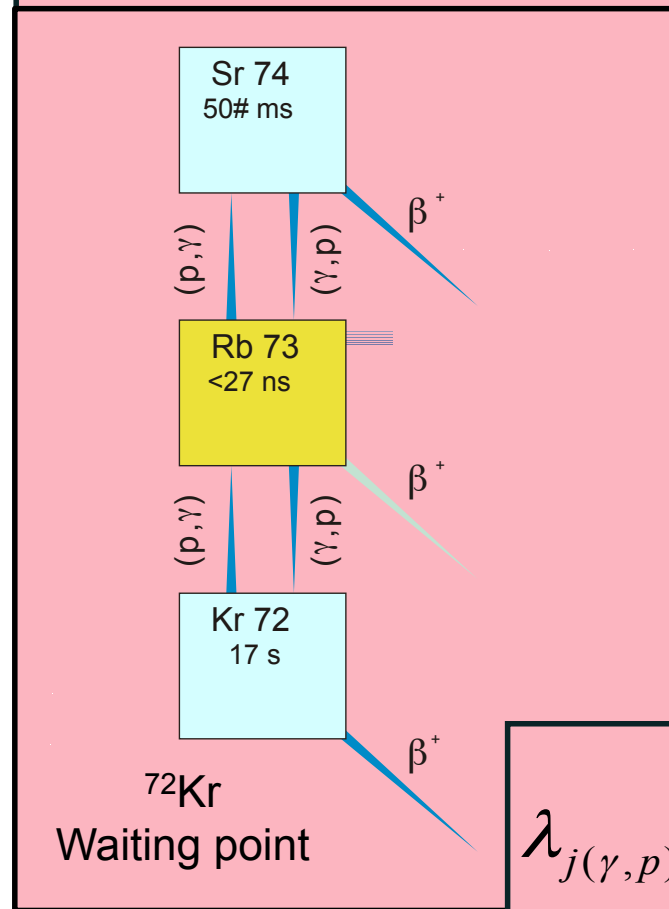


Picture from H. Schatz

Importance of the mass

Nucleon separation energies (S_n , S_p , S_{2n} , S_{2p})

$$Q_{i(p,\gamma)j} = S_j - S_i \quad \text{structure \& astrophysics}$$



$$B(N, Z) - B(N - 2, Z)$$

$$B(N, Z) - B(N, Z - 2)$$

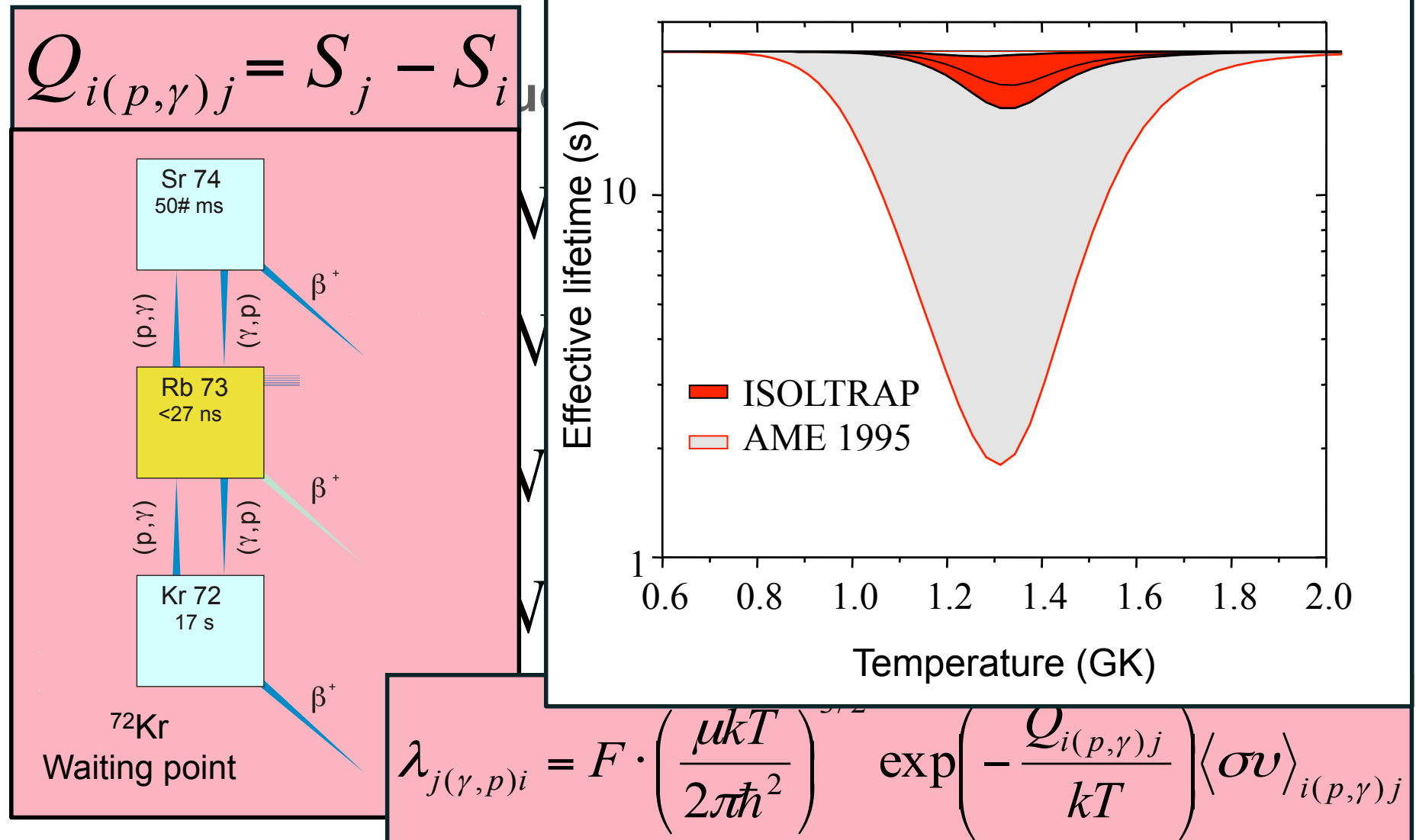
$$B(N, Z) - B(N - 1, Z)$$

$$B(N, Z) - B(N - 1, Z)$$

$$\lambda_{j(\gamma,p)i} = F \cdot \left(\frac{\mu kT}{2\pi\hbar^2} \right)^{3/2} \exp\left(-\frac{Q_{i(p,\gamma)j}}{kT} \right) \langle \sigma v \rangle_{i(p,\gamma)j}$$

Importance of the mass

Nucleon separation energies (S_n , S_p , S_{2n} , S_{2p})



D. Rodríguez *et al.*, PRL 93 (2004) 161104

Importance of the mass

Q-values from β -decay

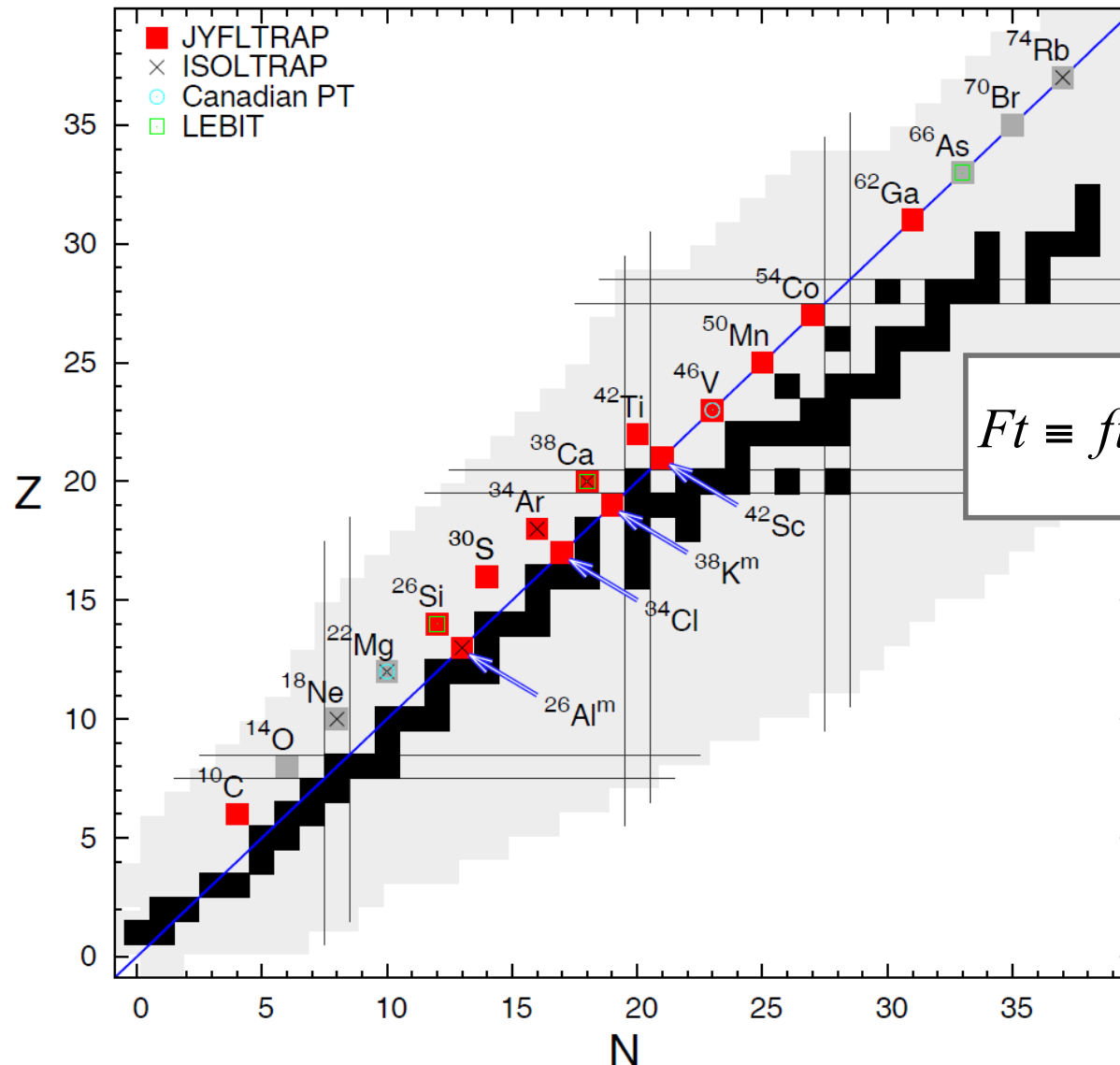
Fundamental symmetries & Neutrino-related physics

$$Q_{EC} = ME(Z, N) - ME(Z - 1, N + 1)$$

$$Q_{\beta^+} = ME(Z, N) - ME(Z - 1, N + 1) - 2m_e$$

$$Q_{\beta^-} = ME(Z, N) - ME(Z + 1, N - 1)$$

Importance of the mass Q-values from β -decay



CVC hypothesis
Superallowed
 β emitters
 $0^+ \rightarrow 0^+$

$$Ft \equiv ft(1 + \delta_R)(1 - \delta_C) = \frac{K}{2G_V^2(1 + \Delta_R^V)}$$

$f(Q^5)$ statistical rate function
 t ($t_{1/2}$, $1/R$) partial half-life

$$V_{ud}^2 = G_V^2 / G_\mu^2$$

$$|V_{ud}^2| + |V_{us}^2| + |V_{ub}^2| = 1 - \Delta$$

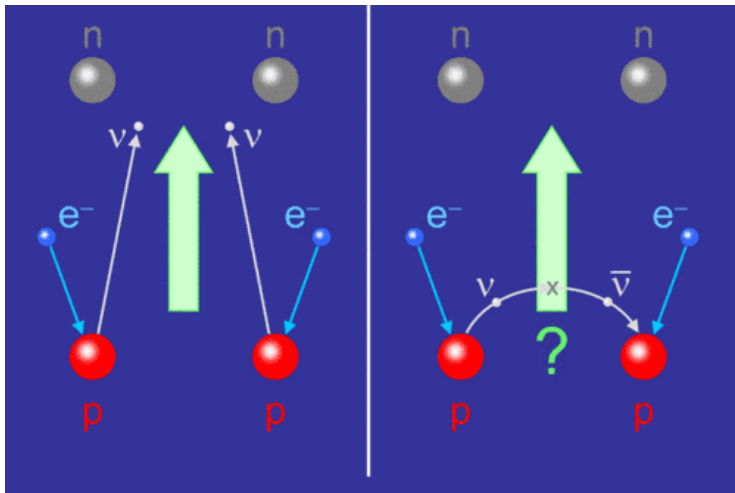
Importance of the mass Q-values from EC-decay

Is the neutrino a Majorana or Dirac particle?

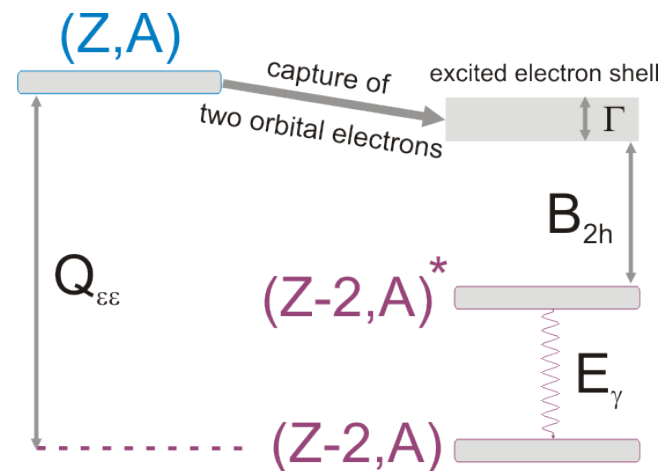
$2\nu 2EC$ ($T_{1/2} > 10^{24}$ y)

$0\nu 2EC$ ($T_{1/2} > 10^{30}$ y)

$$\frac{1}{T_{1/2}} = C \times m_\nu^2 \times |M|^2 \times |\Psi_{1e}|^2 \times |\Psi_{2e}|^2 \times \frac{\Gamma}{(Q - B_{2h} - E_\gamma^2)^2 + \frac{1}{4}\Gamma^2}$$



$0\nu 2EC$ might be resonantly enhanced ($T_{1/2} \sim 10^{25}$ y)



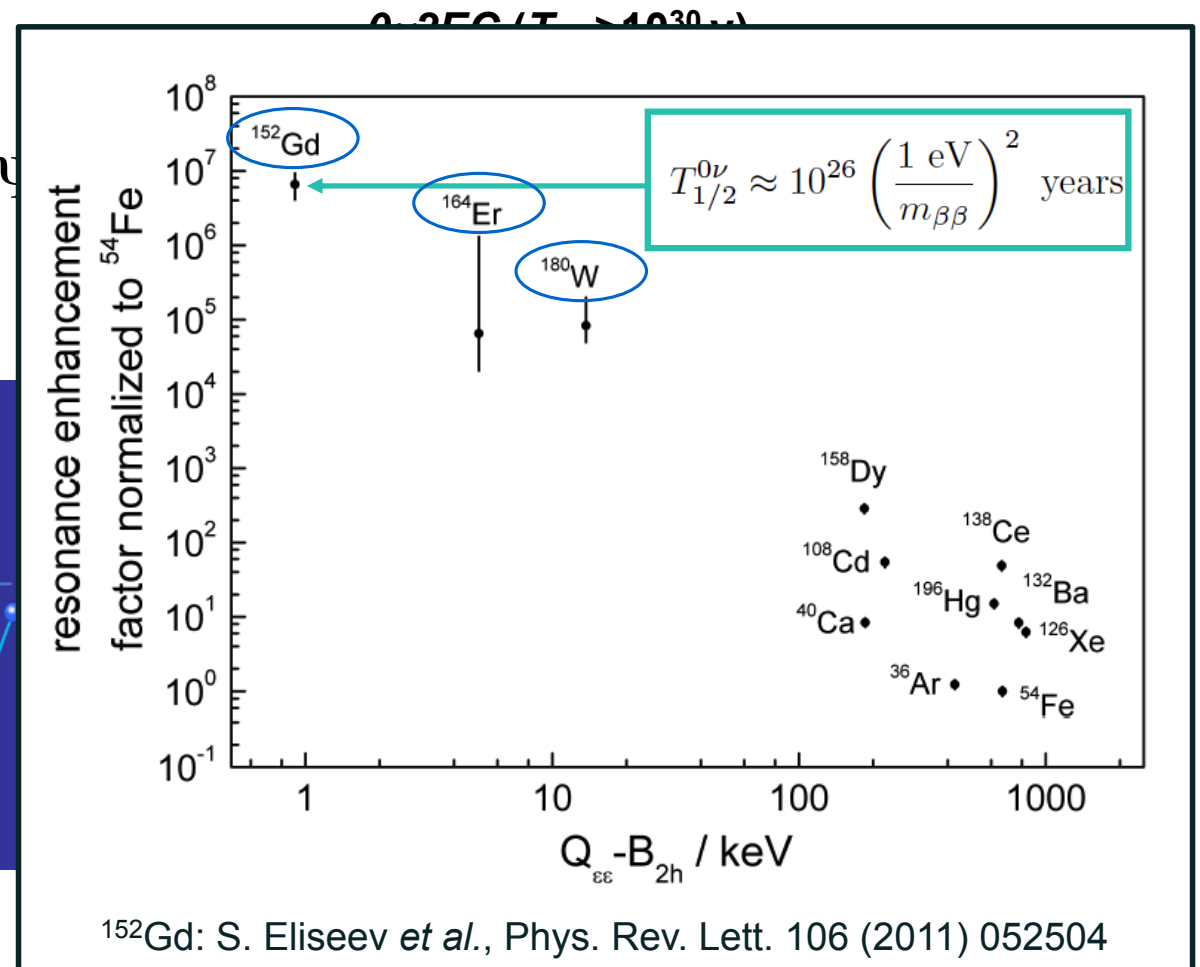
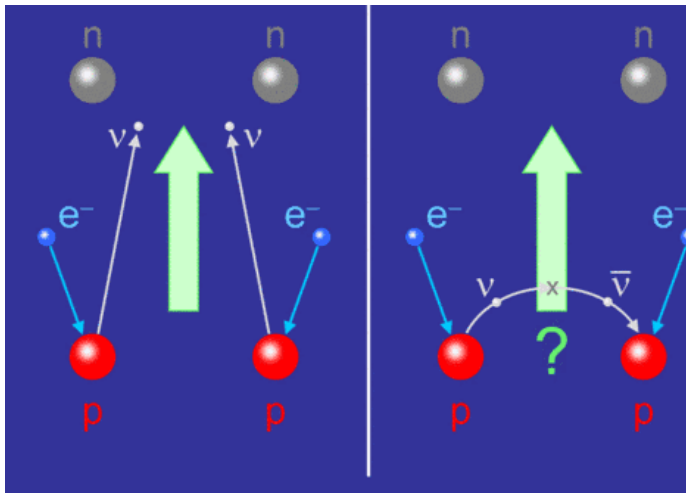
Search for nuclides with $D = (Q_{ee} - B_{2h} - E_g) < 1$ keV

Importance of the mass Q-values from EC-decay

Is the neutrino a Majorana or Dirac particle?

$2\nu 2EC$ ($T_{1/2} > 10^{24}$ y)

$$\frac{1}{T_{1/2}} = C \times m_\nu^2 \times |M|^2 \times |\psi|^2$$

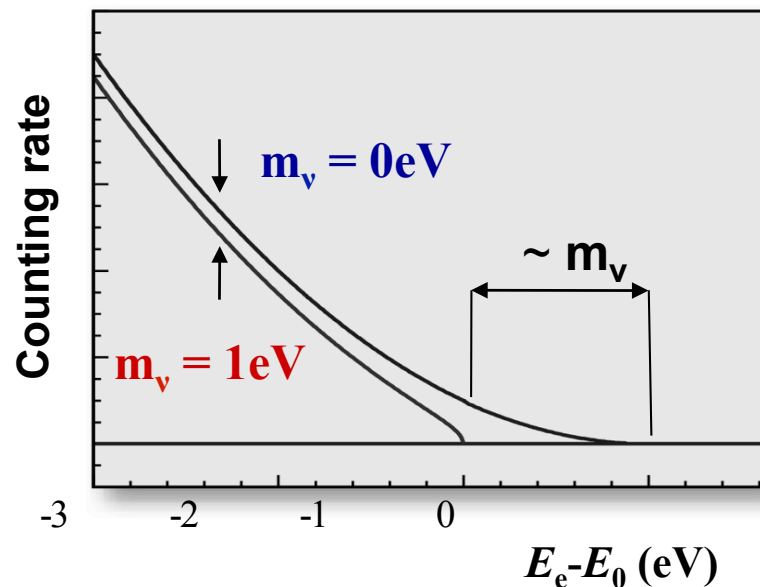


Importance of the mass

Q-values from β -decay

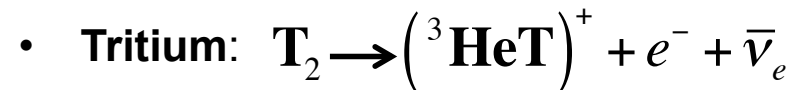
- Direct measurement of the mass by analyzing the end of the β -decay spectrum independent whether the neutrino is a Majorana or a Dirac particle

$$\frac{dN}{dE} = k \times F(E, Z) \times p \times E_{total} \times (E_0 - E_e)^2 \times \left[(E_0 - E_e)^2 - m_\nu^2 \right]^{1/2}$$

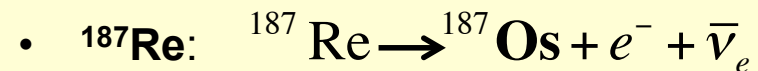


$$F(\delta E) = \int_{-\delta E}^0 N(E_\beta) dE \approx 2 \left(\frac{\delta E}{E_0} \right)^3$$

- The best candidates will be decaying nuclei with low *end-point energies*



$$Q = 18.6 \text{ keV} \quad T_{1/2} = 12.3 \text{ y}$$

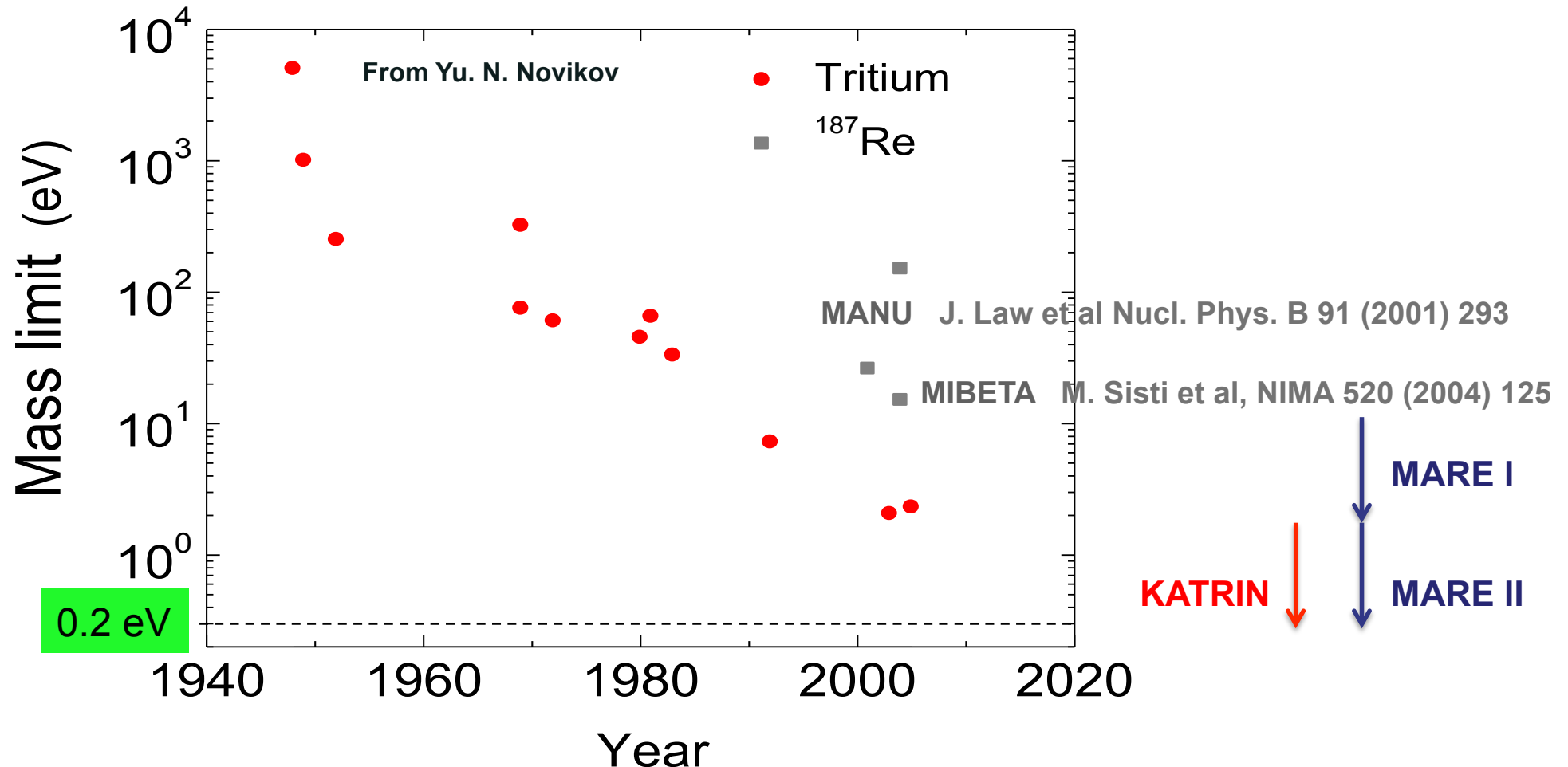


$$Q = 2.5 \text{ keV} \quad T_{1/2} = 43.2 \times 10^9 \text{ y}$$

Q has to be unambiguously determined (with $\delta m/m$ of at least 10^{-11}) from the mass difference using a Penning trap

History of antineutrino mass measurements

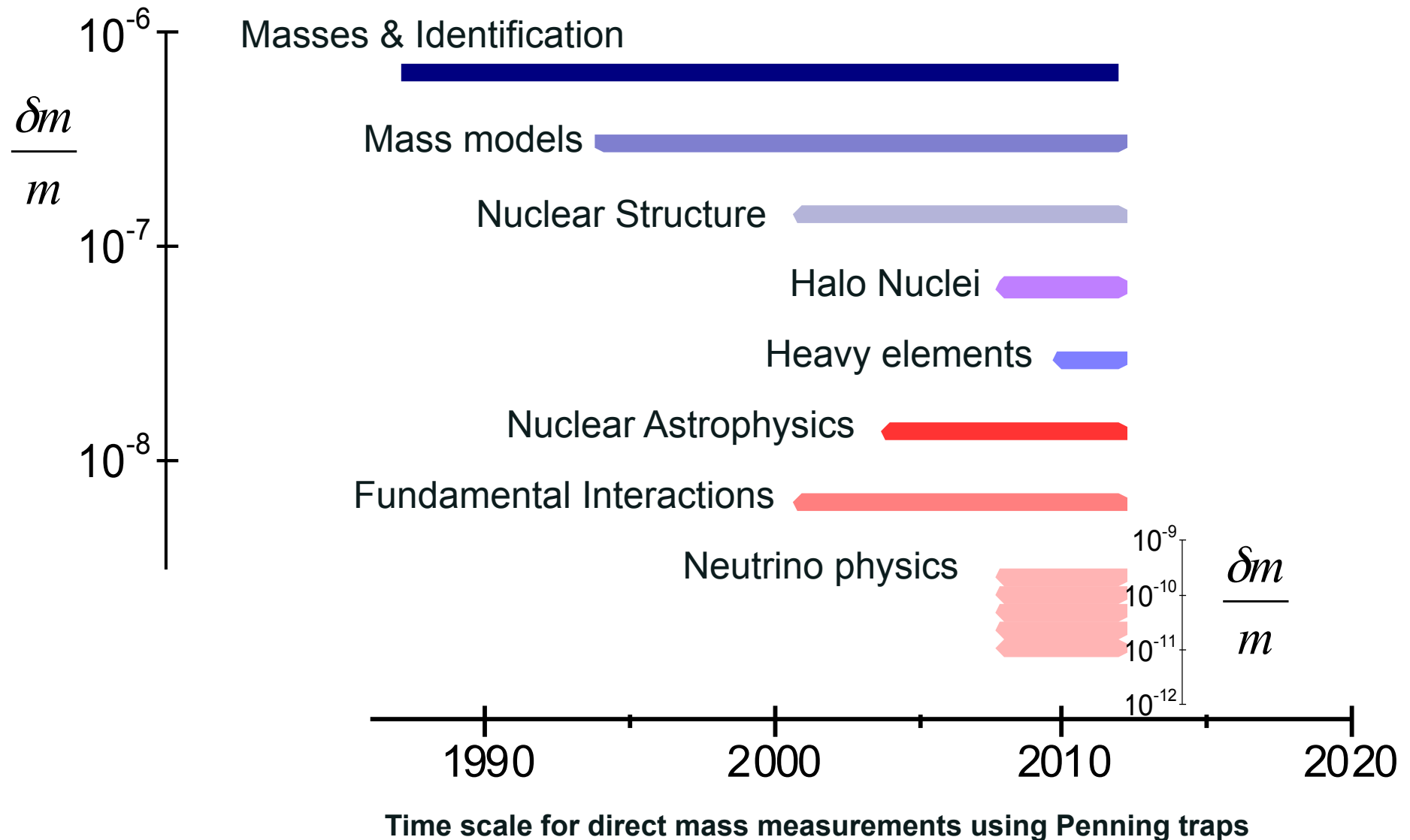
Tritium and ^{187}Re



- There are two International collaborations aiming at measuring with a sensitivity of about $0.2 \text{ eV}/c^2$: KATRIN (KARlsruhe TRItium Neutrino experiment) and MARE (Micro calorimeter Array for a Rhenium Experiment)

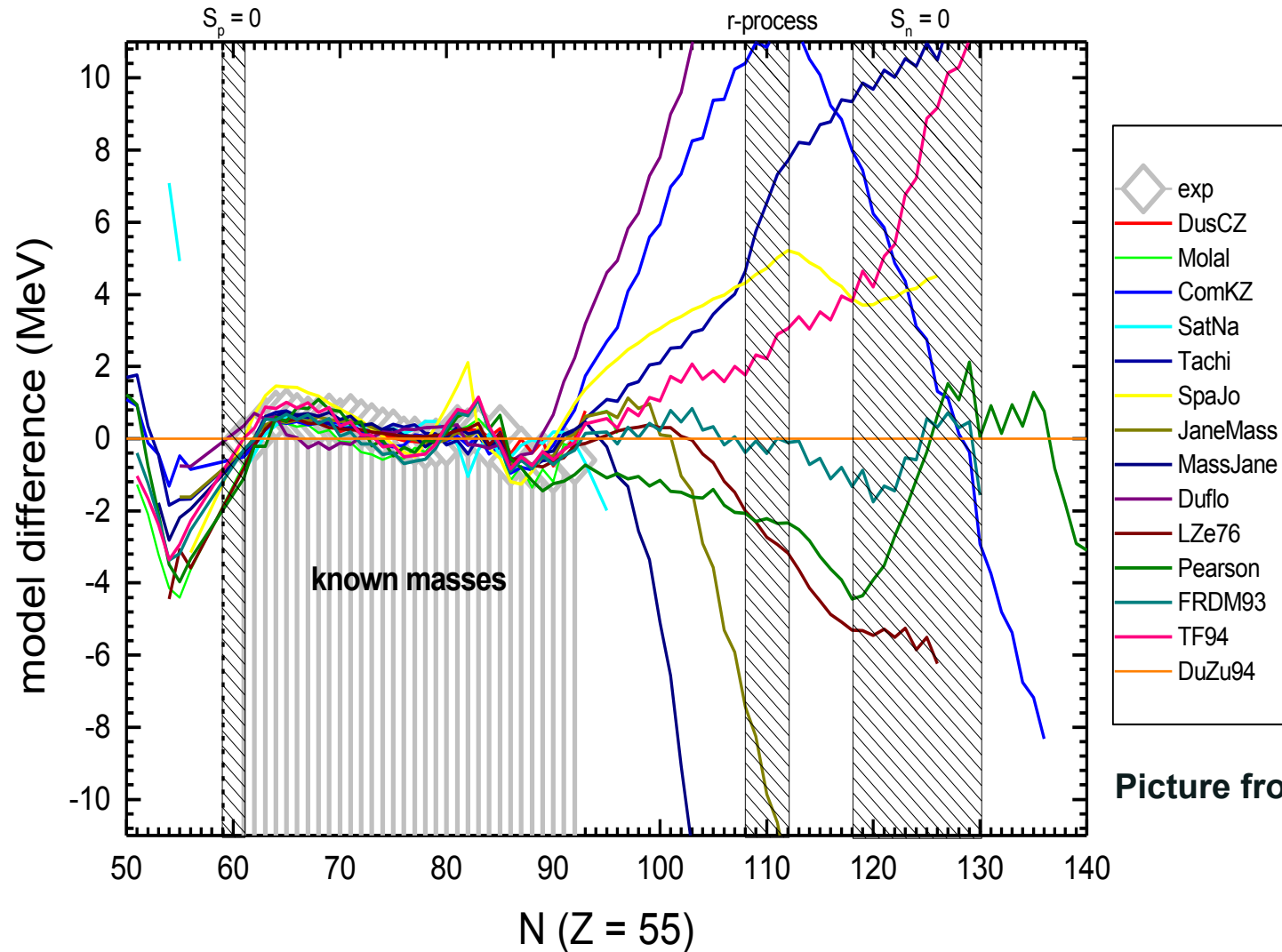
Physics goals

Mass uncertainty required



Experimental results and models

Discrepancies between the models

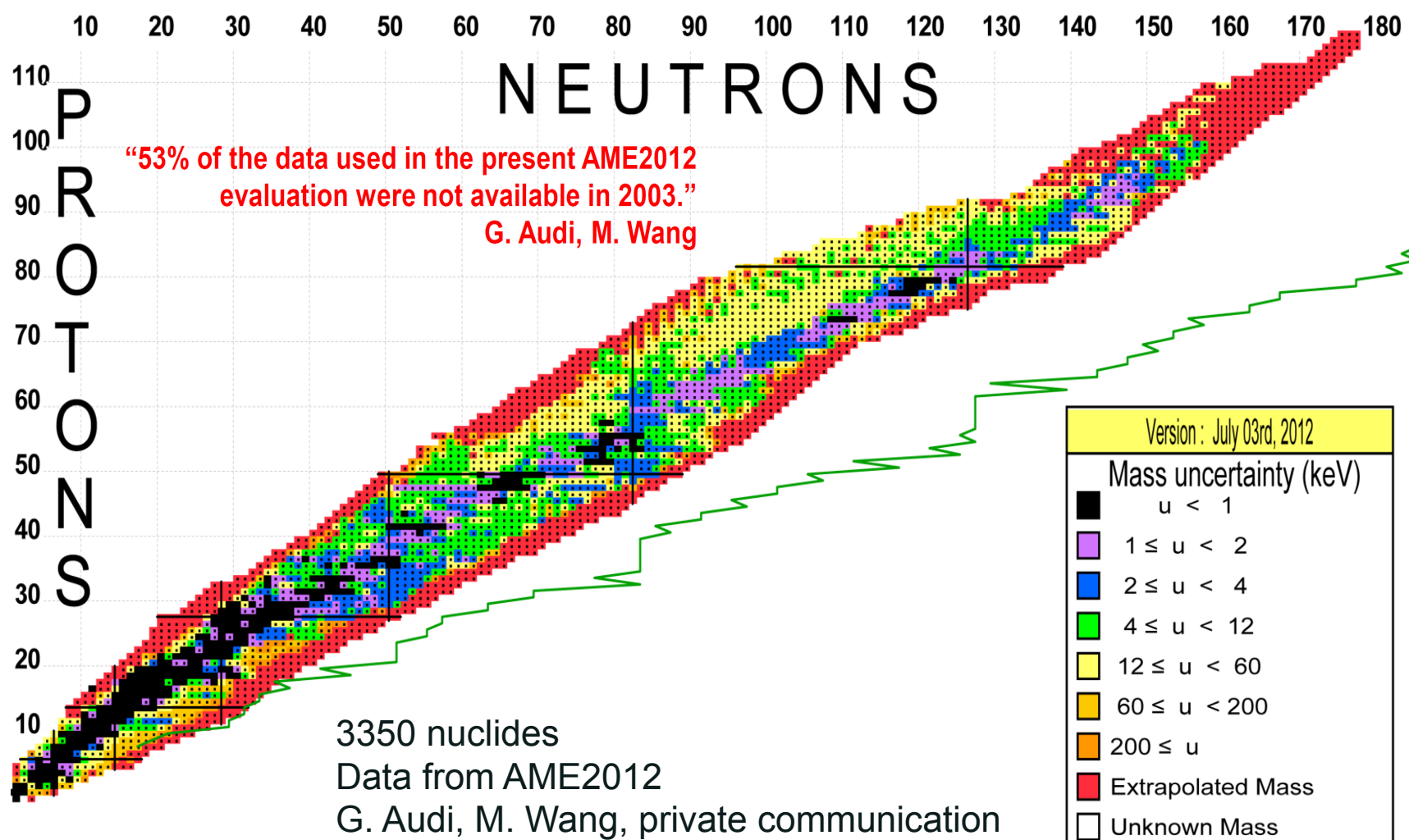


Picture from D. Lunney

Particularly relevant in the region of Superheavy Elements

Mass measurements in the last decade

The 2012 Atomic Mass Evaluation

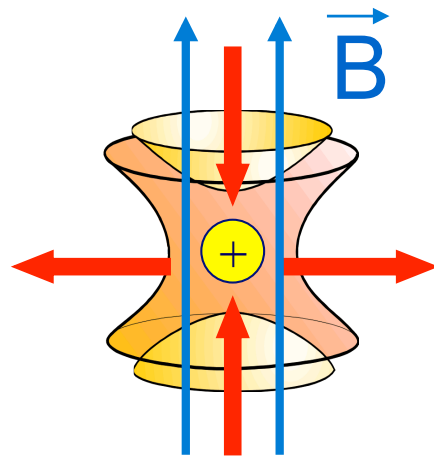


THE PENNING TRAP

Penning traps

Fundamentals

Magnetic + Electrostatic field

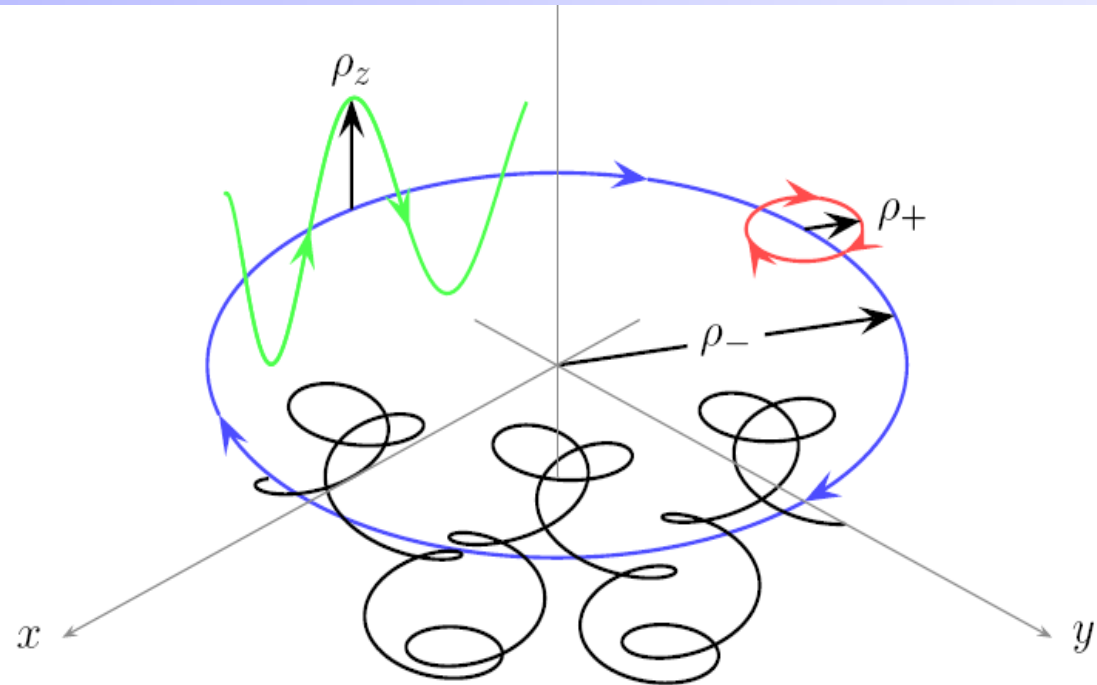


Axial motion

$$\omega_z = \sqrt{\frac{q\Phi_0}{md^2}}$$

Reduced cyclotron motion

Magnetron motion



$$\left. \begin{aligned} \omega_+ &= \frac{1}{2} \left[\omega_c + \sqrt{(\omega_c^2 - 2\omega_z^2)} \right] \\ \omega_- &= \frac{1}{2} \left[\omega_c - \sqrt{(\omega_c^2 - 2\omega_z^2)} \right] \end{aligned} \right\}$$

$$\omega_c = \frac{q}{m} B$$

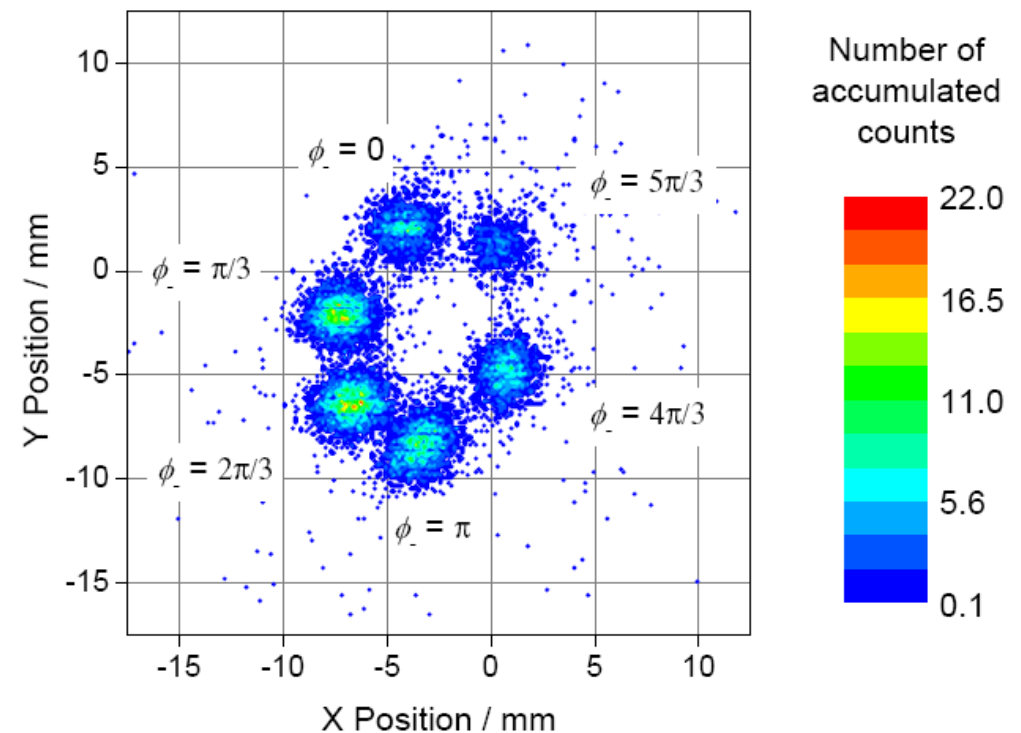
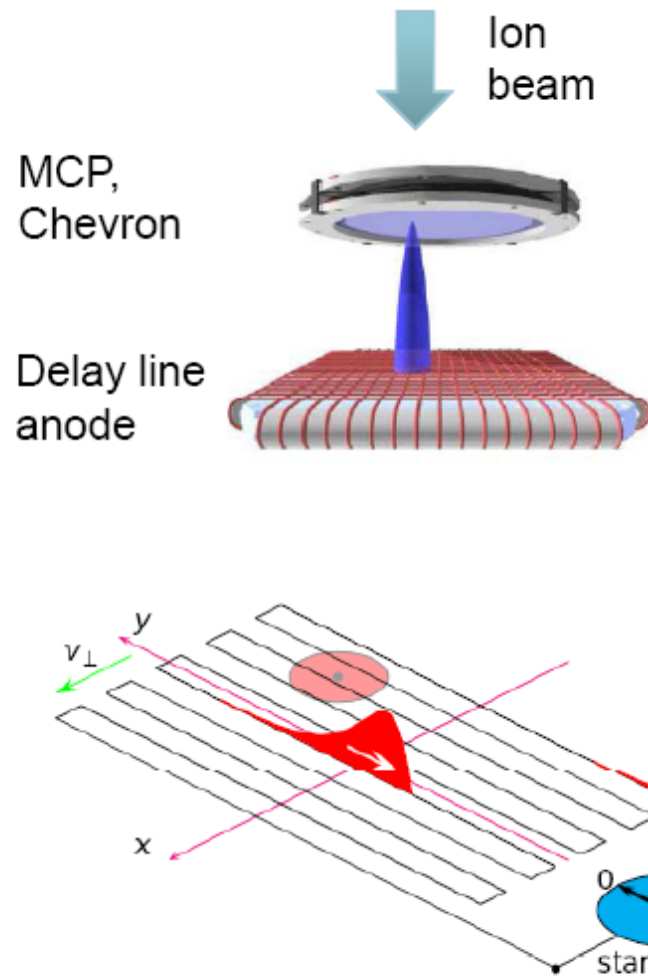
$$\omega_c = \omega_+ + \omega_-$$

$$\omega_+ > \omega_z > \omega_-$$

$$\omega_+^2 + \omega_z^2 + \omega_-^2 = \omega_c^2$$

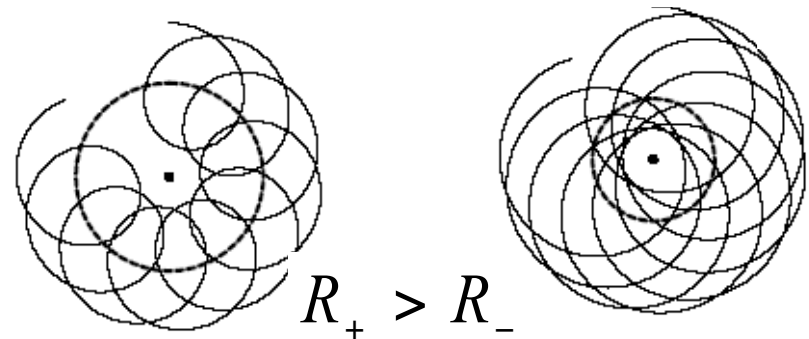
Penning traps

Fundamentals



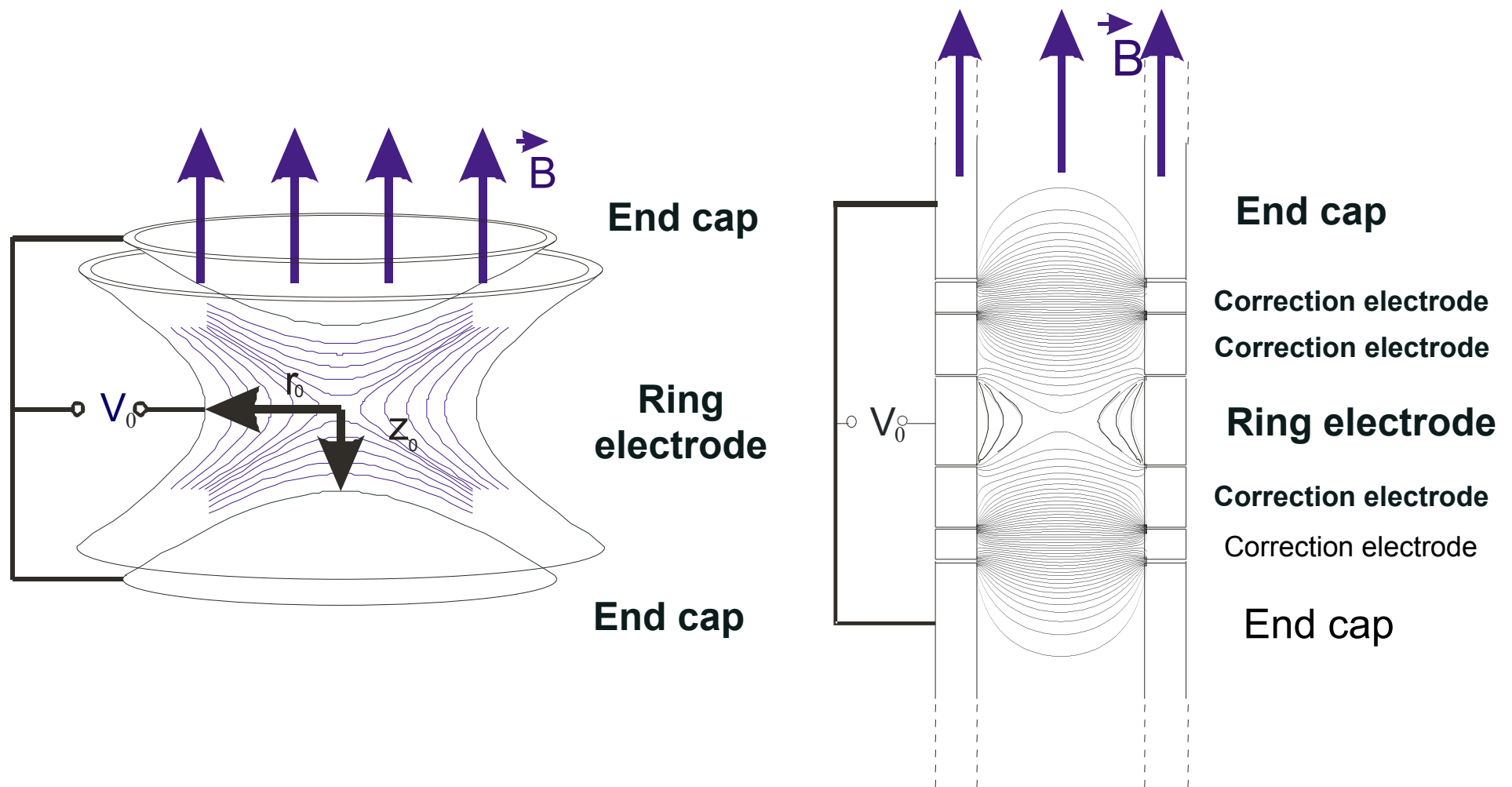
G. Eitel *et al.*, Nucl. Instrum. Methods A, 606, (2009) 475

$$R_+ < R_-$$



Penning traps

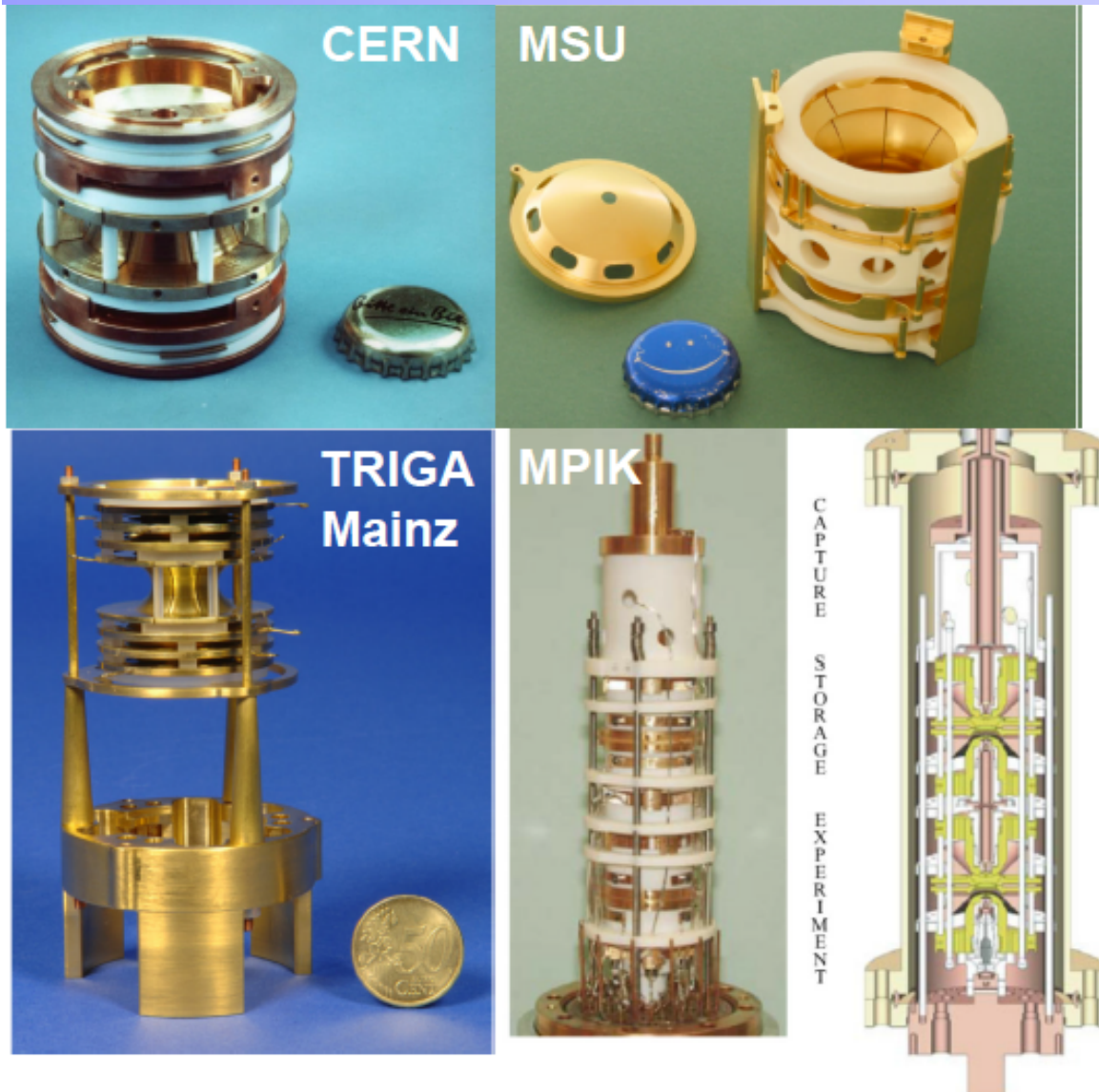
Different geometries



L.S. Brown and G. Gabrielse, Rev. Mod. Phys. 58 (1986) 233

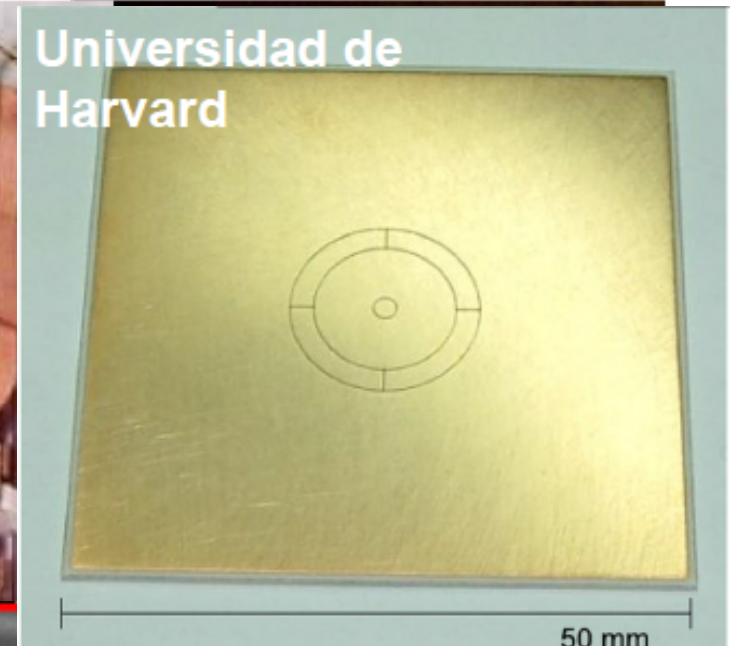
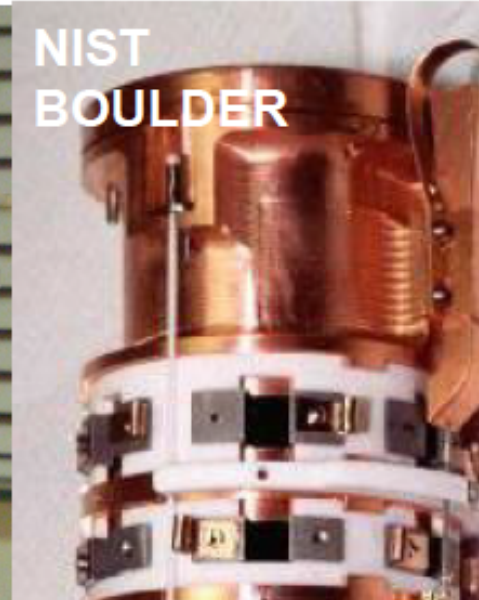
Penning traps

Different geometries



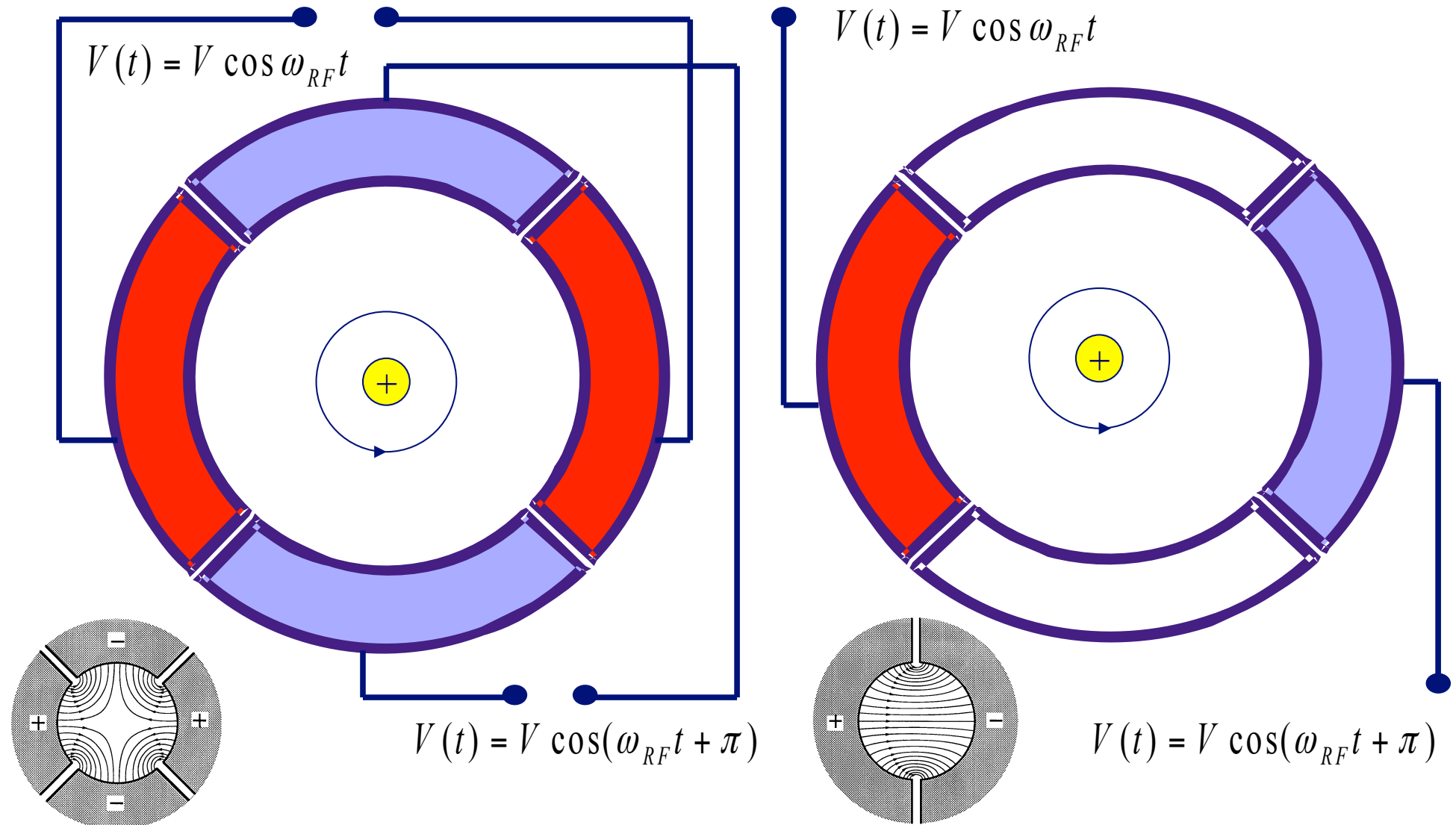
Penning traps

Different geometries



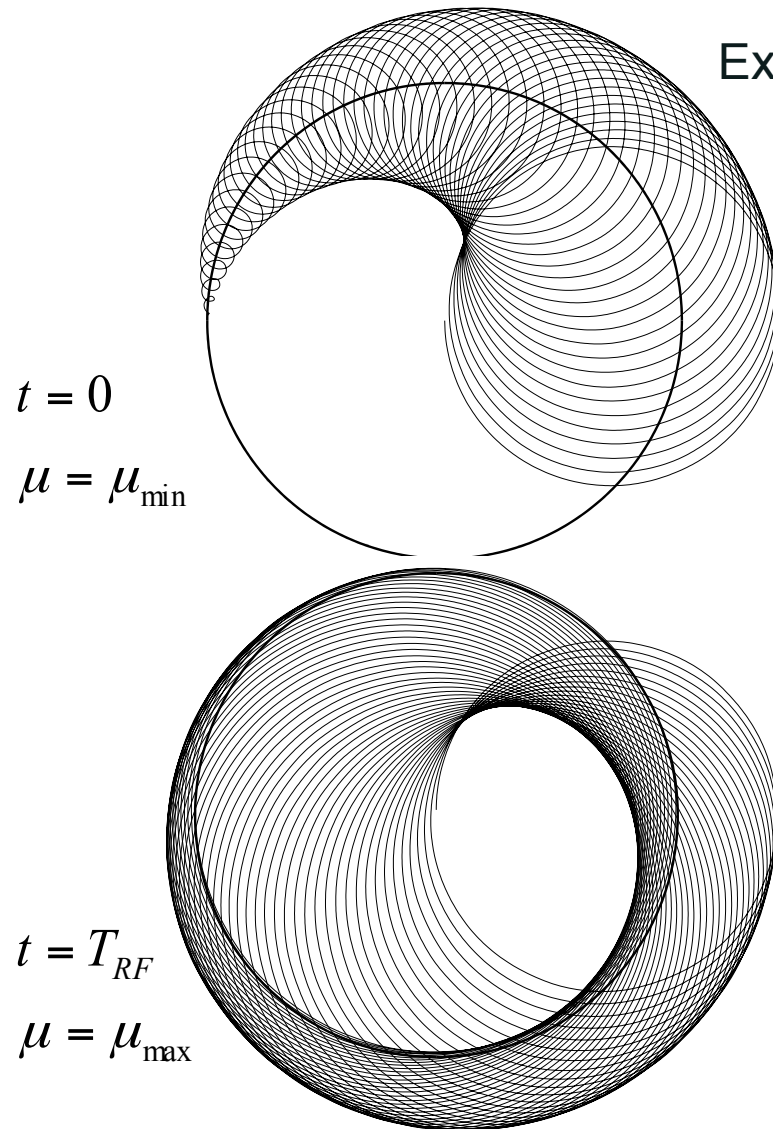
Penning-trap techniques

External Fields

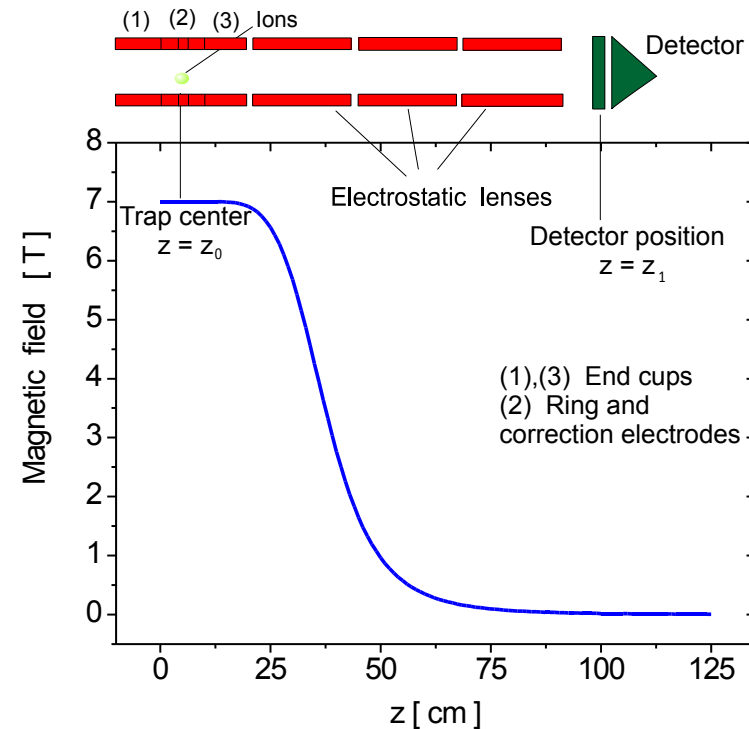


Penning-trap techniques

Time-of-Flight Ion Cyclotron Resonance (TOF-ICR)



External RF quadrupole field with $\omega_{RF} = \omega_c (= \omega_+ + \omega_-)$



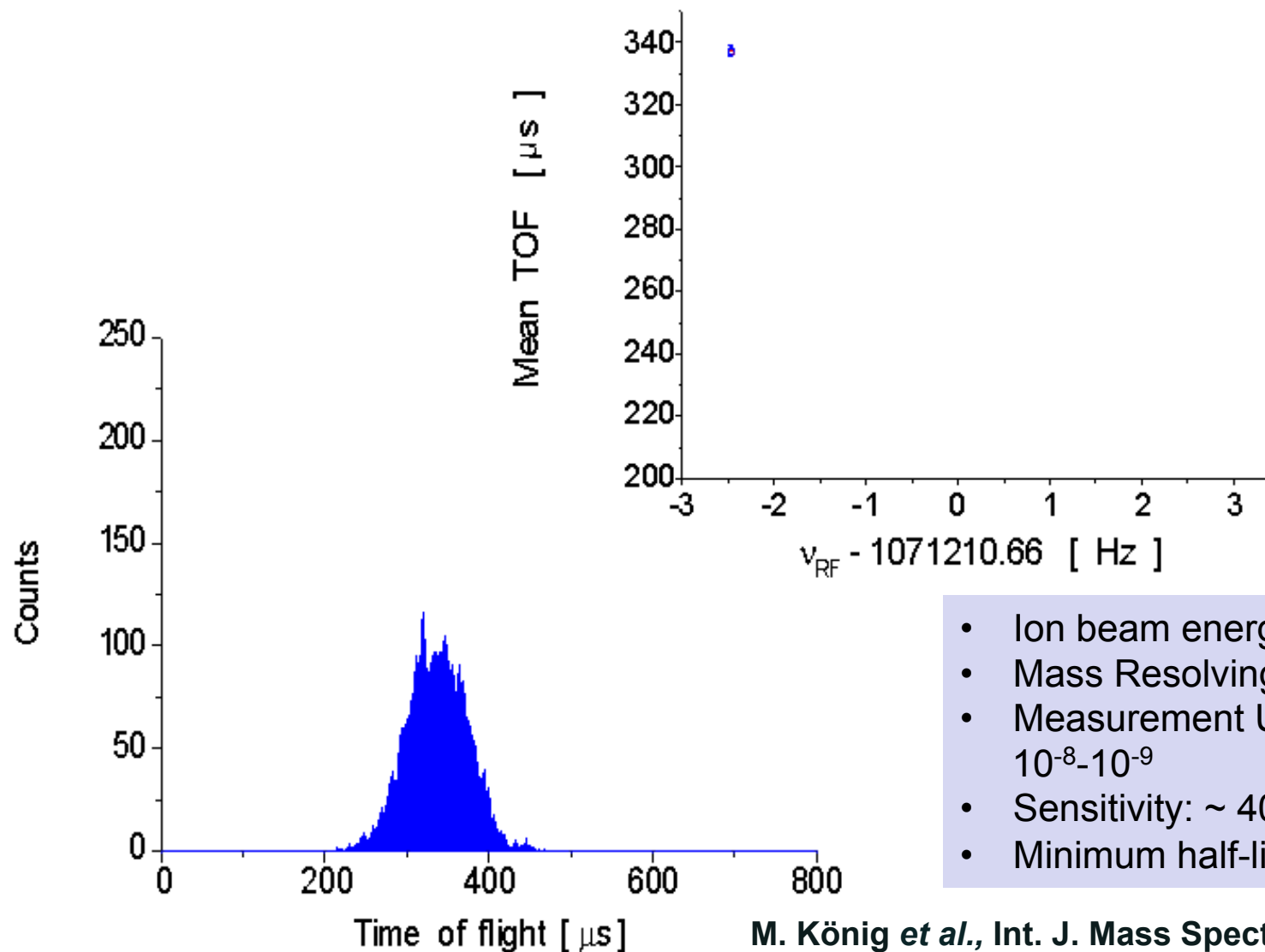
$$F_z = \mu \frac{\partial B}{\partial z}$$

$$TOF_{\min} \Leftrightarrow \mu_{\max} \Leftrightarrow \omega_{RF} = \omega_c$$

G. Gräff et al., Z. Phys. A 297 (1980) 35

Penning-trap techniques

Time-of-Flight Ion Cyclotron Resonance (TOF-ICR)

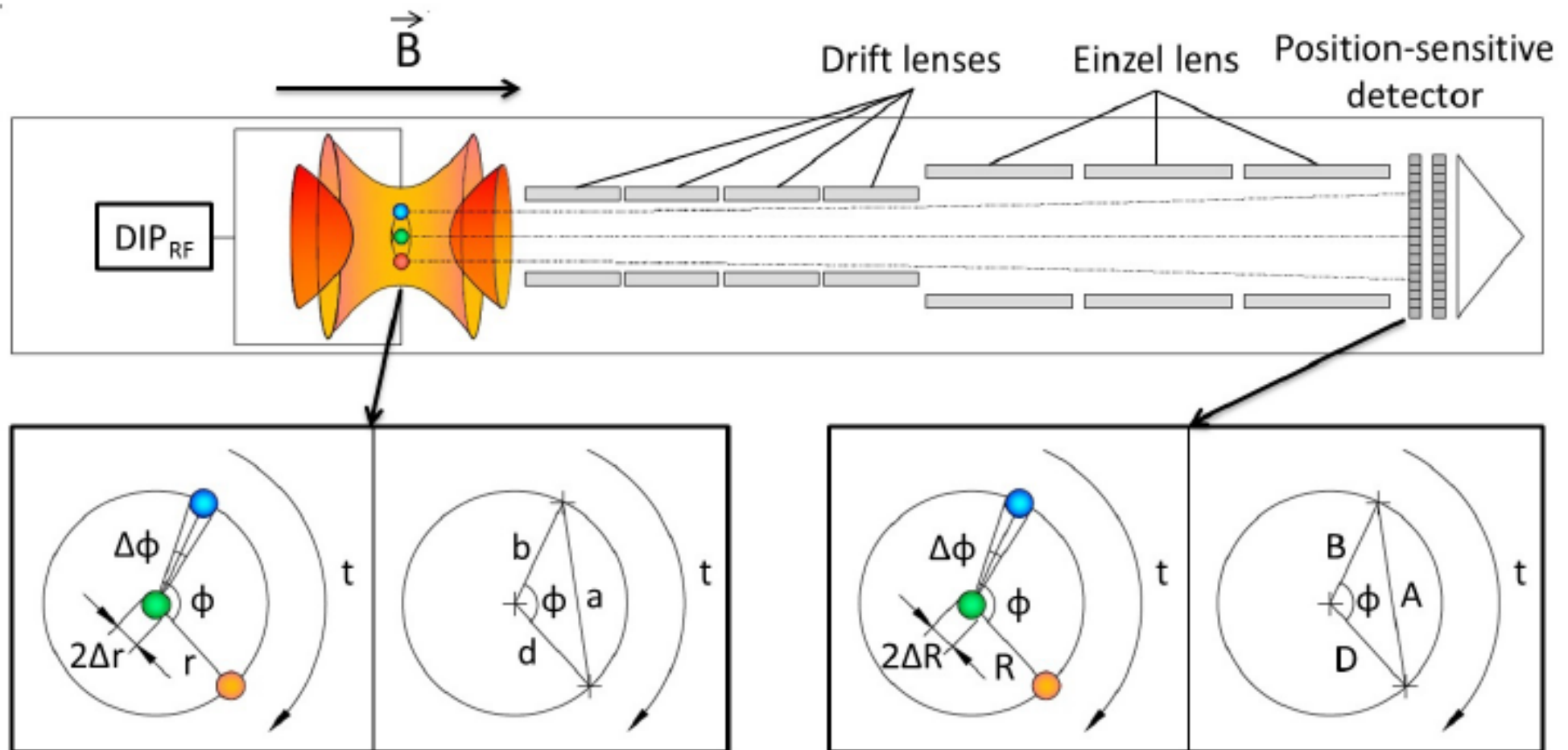


- Ion beam energy: ~ 1 eV
- Mass Resolving Power: 1×10^6
- Measurement Uncertainty: 10^{-8} - 10^{-9}
- Sensitivity: ~ 40 ions
- Minimum half-life: 8 ms

M. König *et al.*, Int. J. Mass Spectrom. 142 (1995) 95

Penning-trap techniques

Phase-Imaging Ion Cyclotron Resonance (PI-ICR)



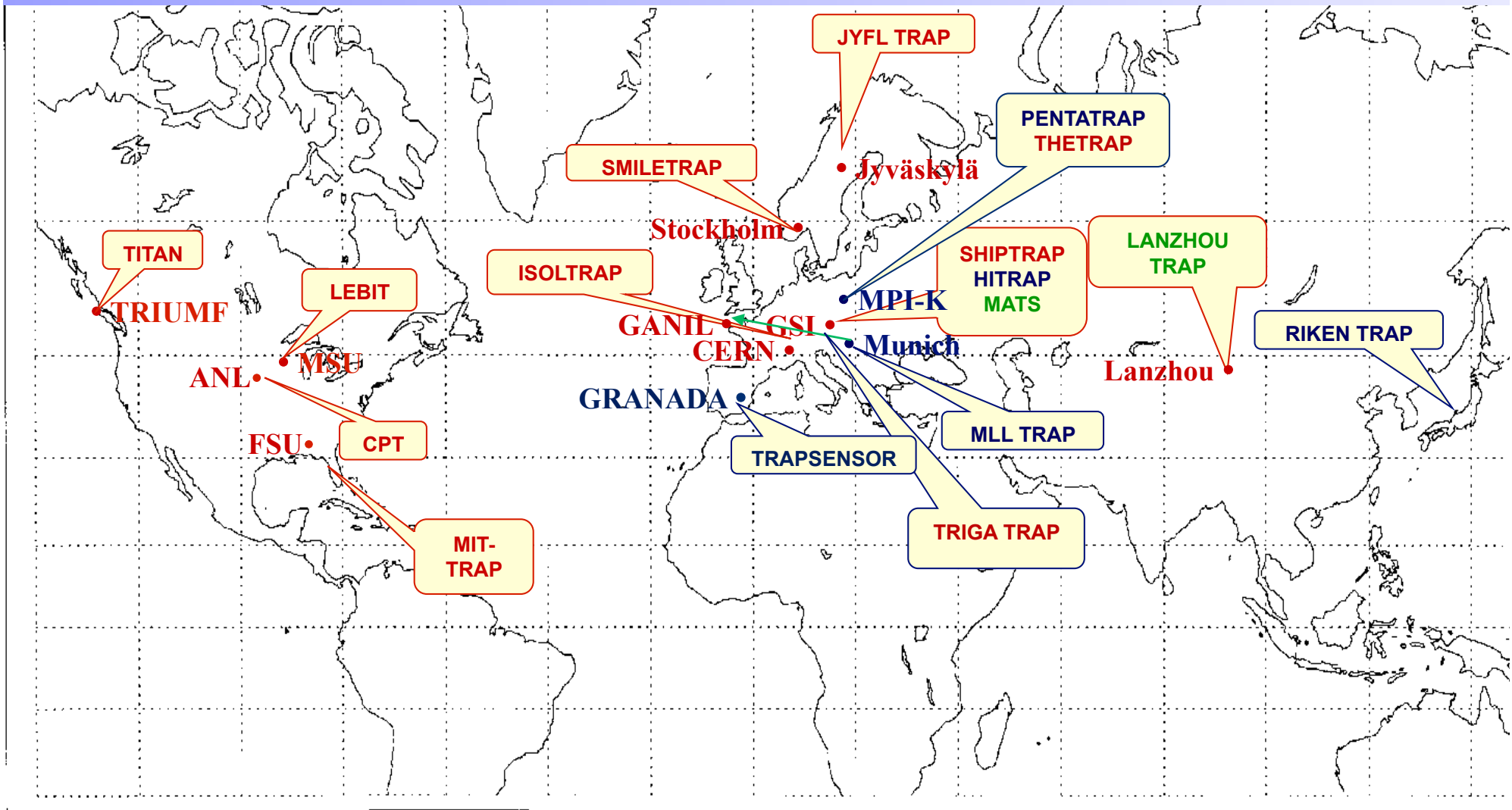
- Ion beam energy: ~ 1 eV
- Mass Resolving Power: 1×10^7
- Measurement Uncertainty: 10^{-9} - 10^{-10}

- Sensitivity: ~ 10 -20 ions
- Minimum half-life: ~ 1 ms

S. Eliseev *et al.*, PRL 110, (2013) 082501

Existing Penning-trap facilities

Mass spectrometry of atomic nuclei

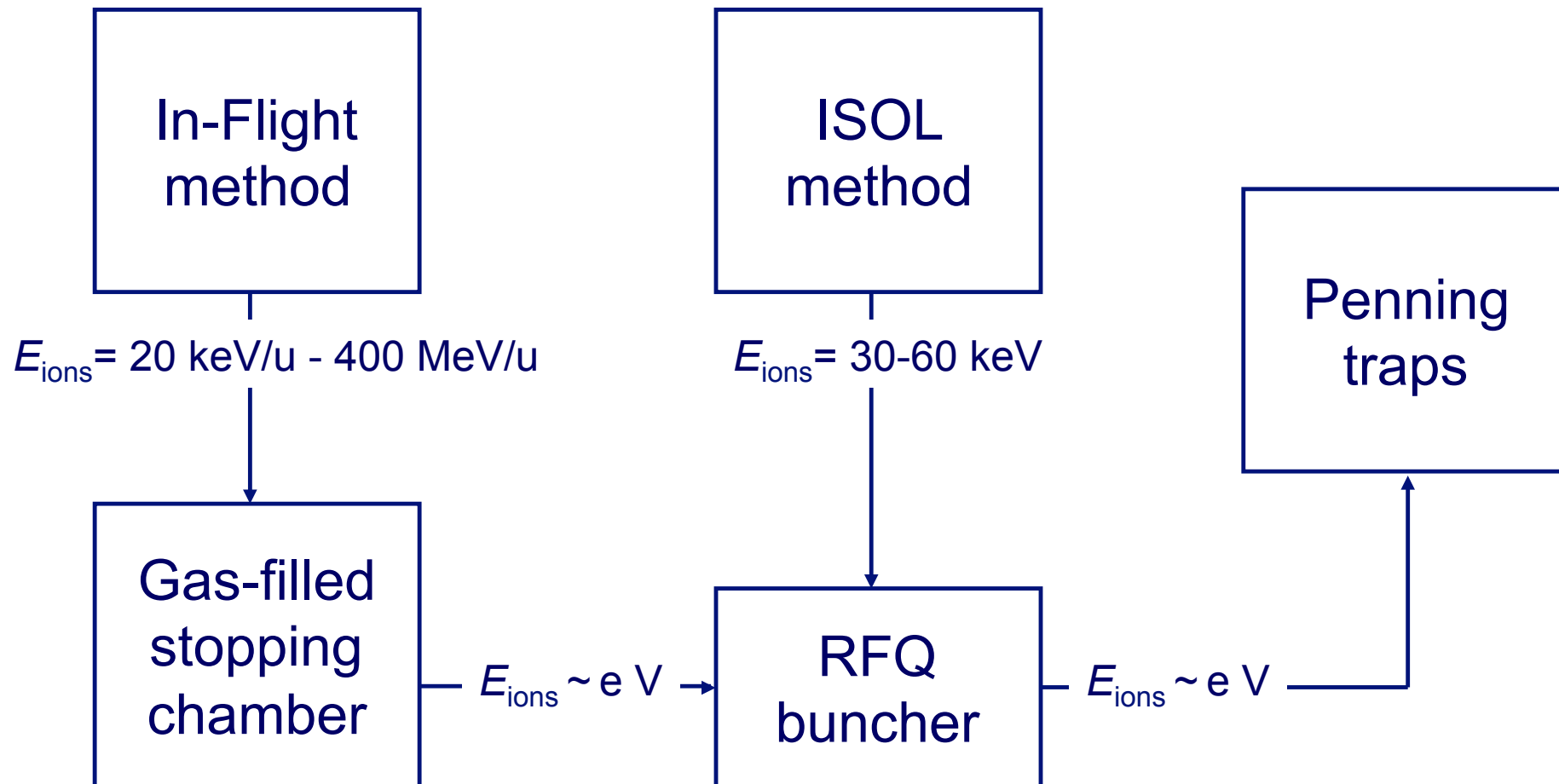


• Running facilities

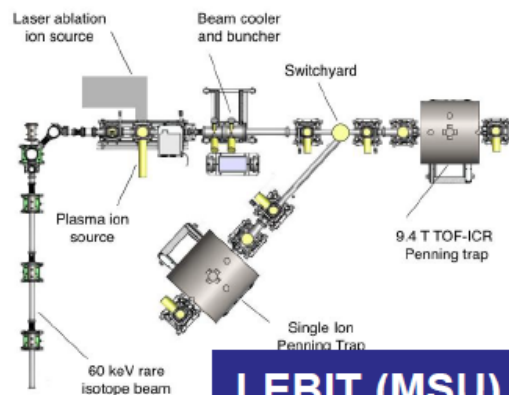
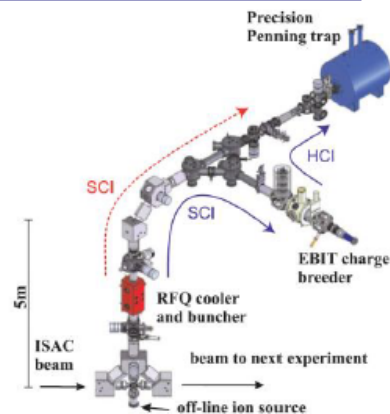
• Under construction

• Future facilities

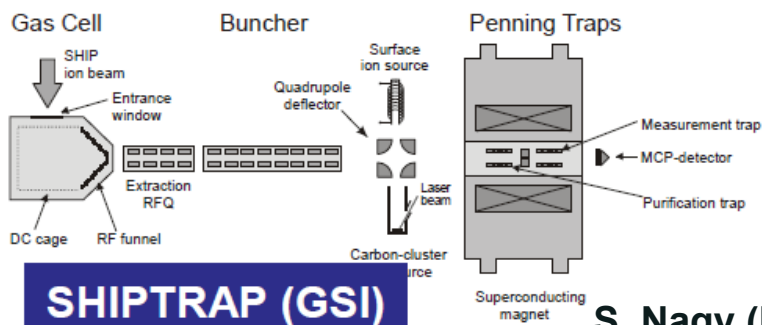
Coupling Penning traps to RIB facilities



TITAN (TRIUMF)

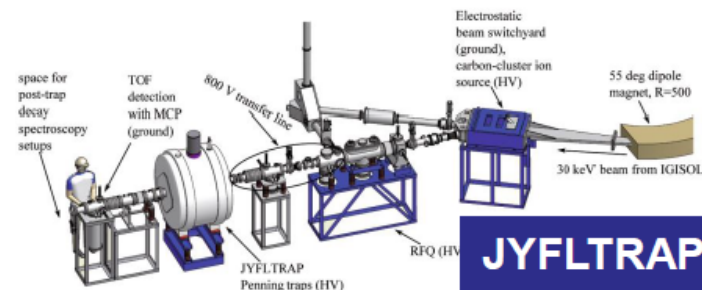


LEBIT (MSU)

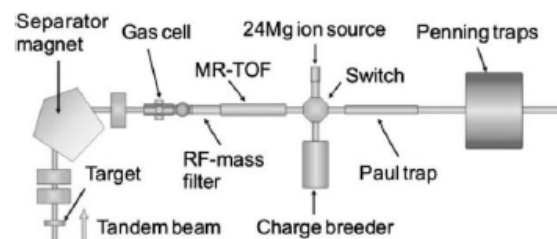


SHIPTRAP (GSI)

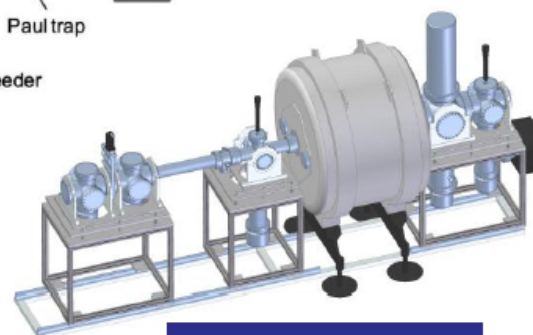
Growing facilities



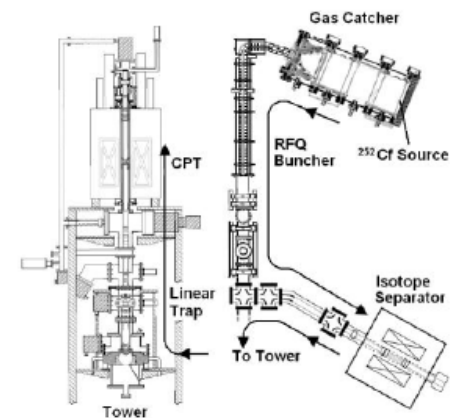
JYFLTRAP



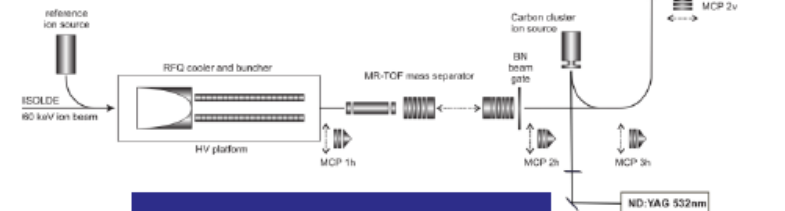
MLLTRAP



TRIGA-TRAP



CPT (ARGONNE)



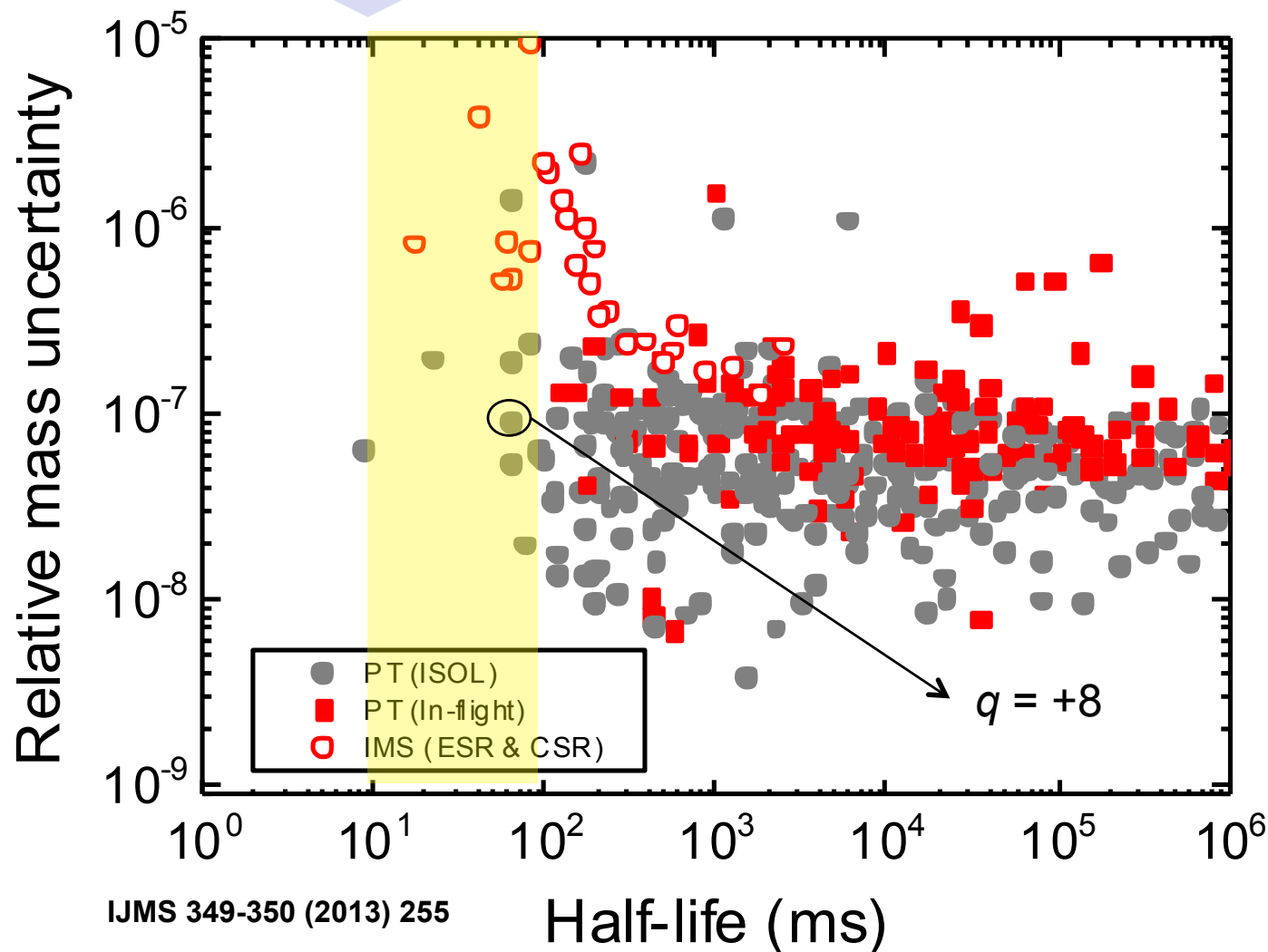
ISOLTRAP(CERN)

S. Nagy (EMIS2012)

Physics survey

Penning-trap mass measurements at RIB

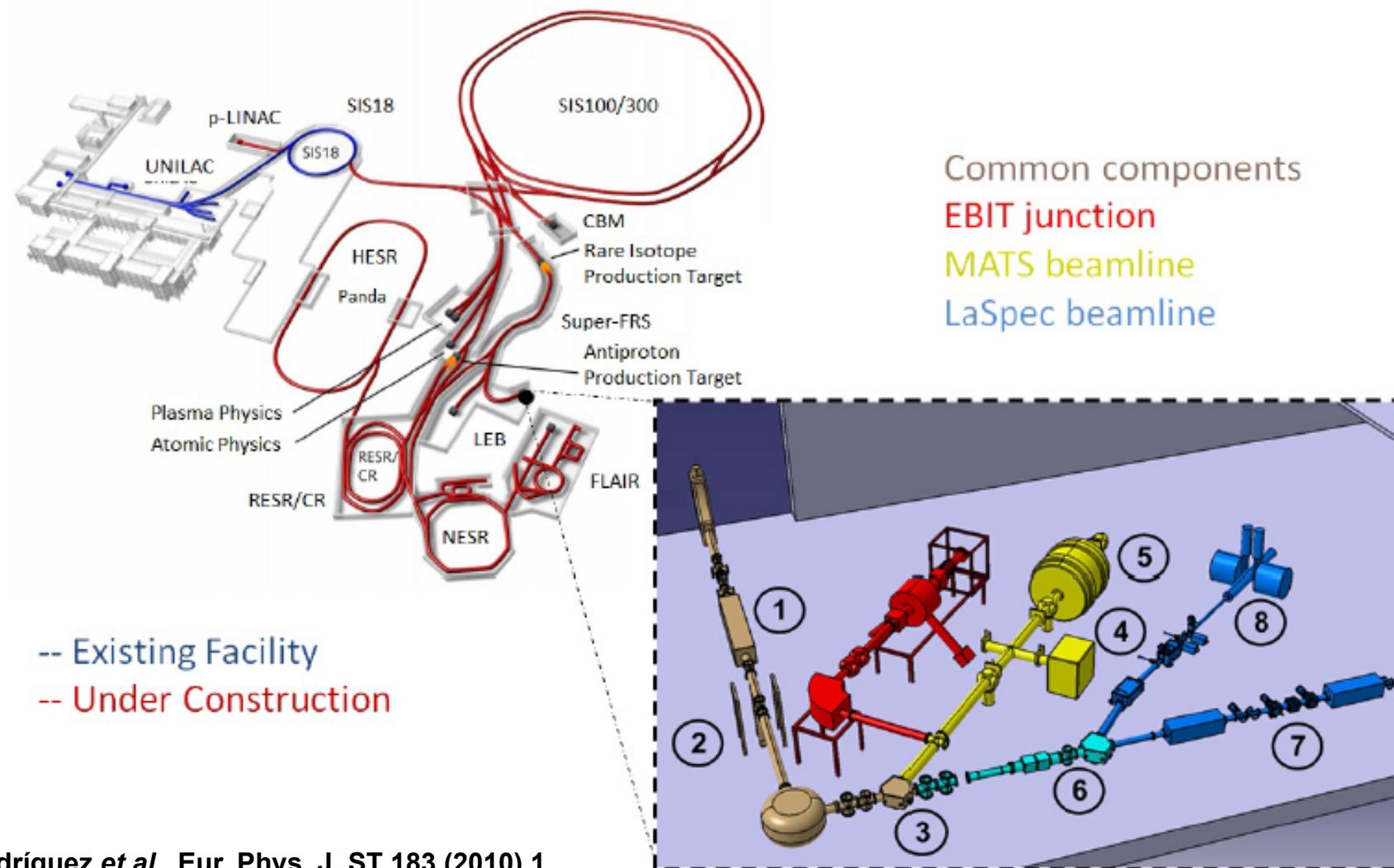
Practicable with higher production rates → MATS at FAIR



Most of the PT measurements performed with the TOF-ICR technique on 1+ or 2+ ions

Physics survey

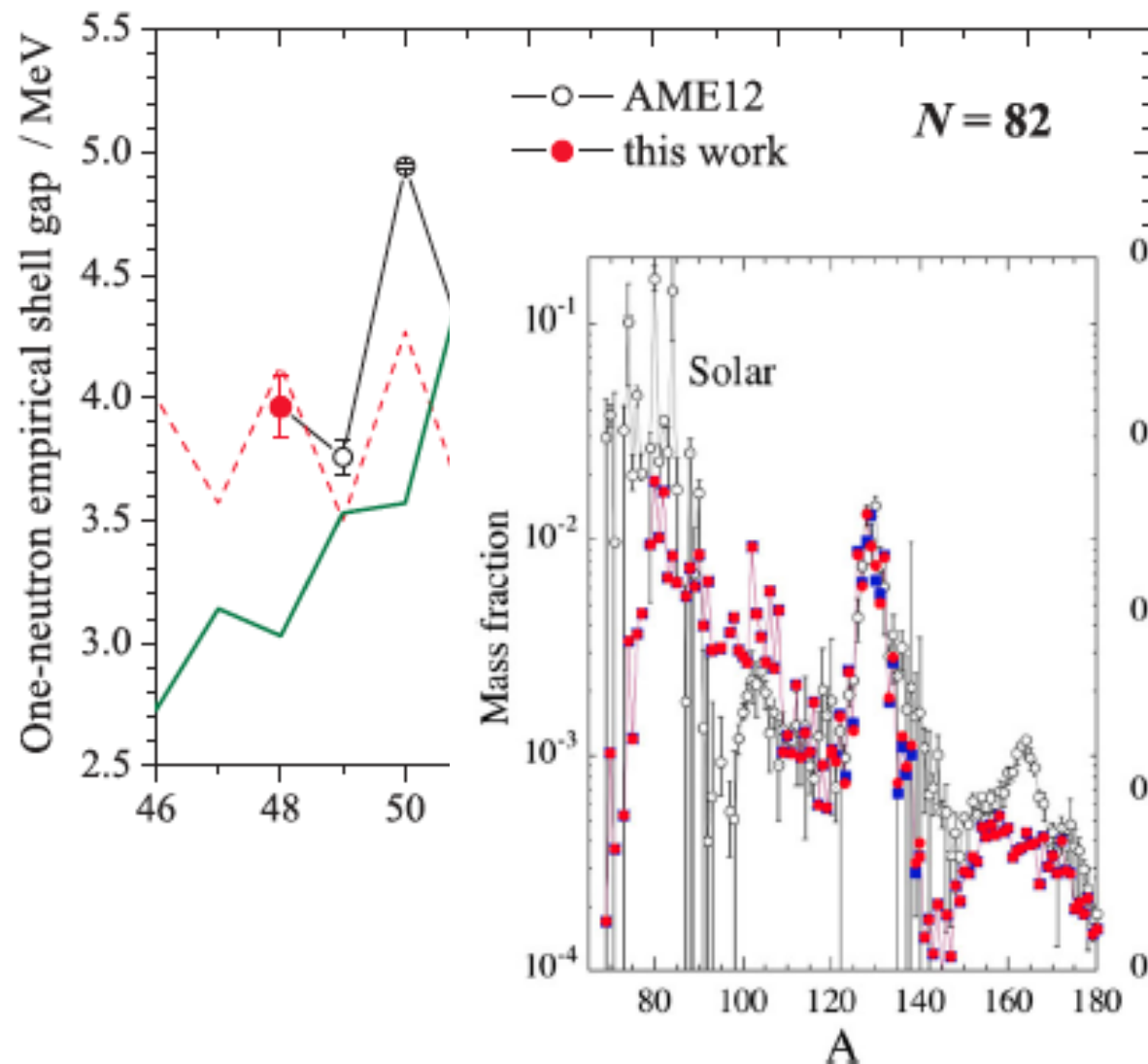
Penning-trap mass measurements at RIB



D. Rodríguez *et al.*, Eur. Phys. J. ST 183 (2010) 1

Recent highlights

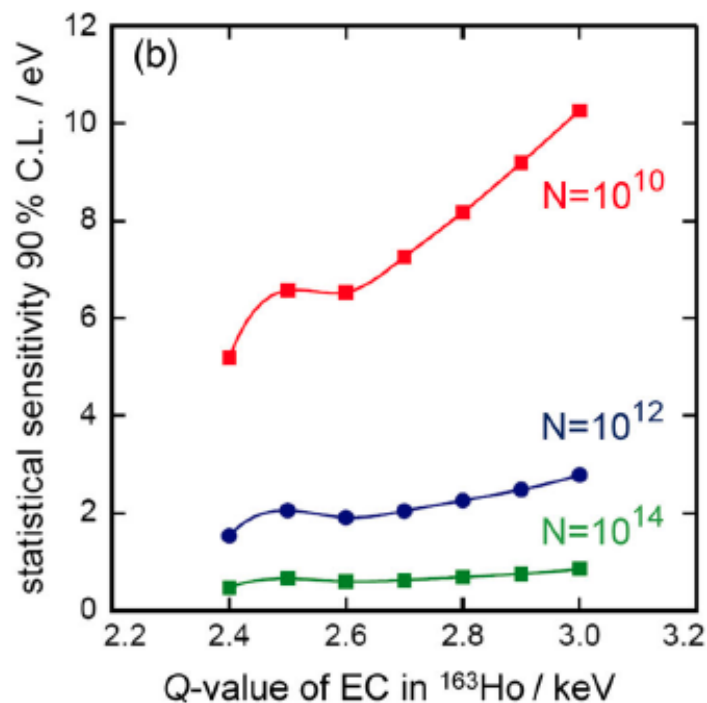
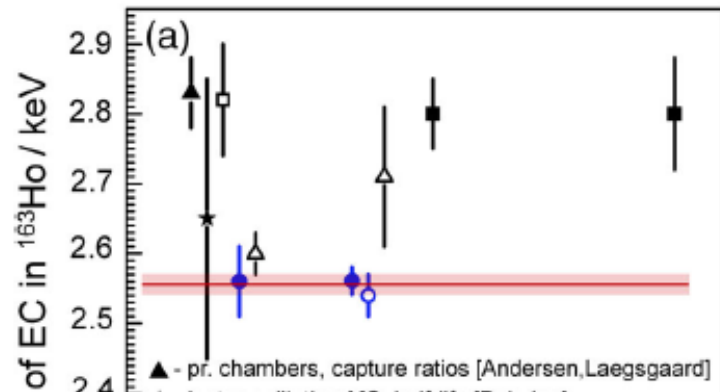
Nuclear (Penning-trap) astrophysics at ISOLDE (2015)



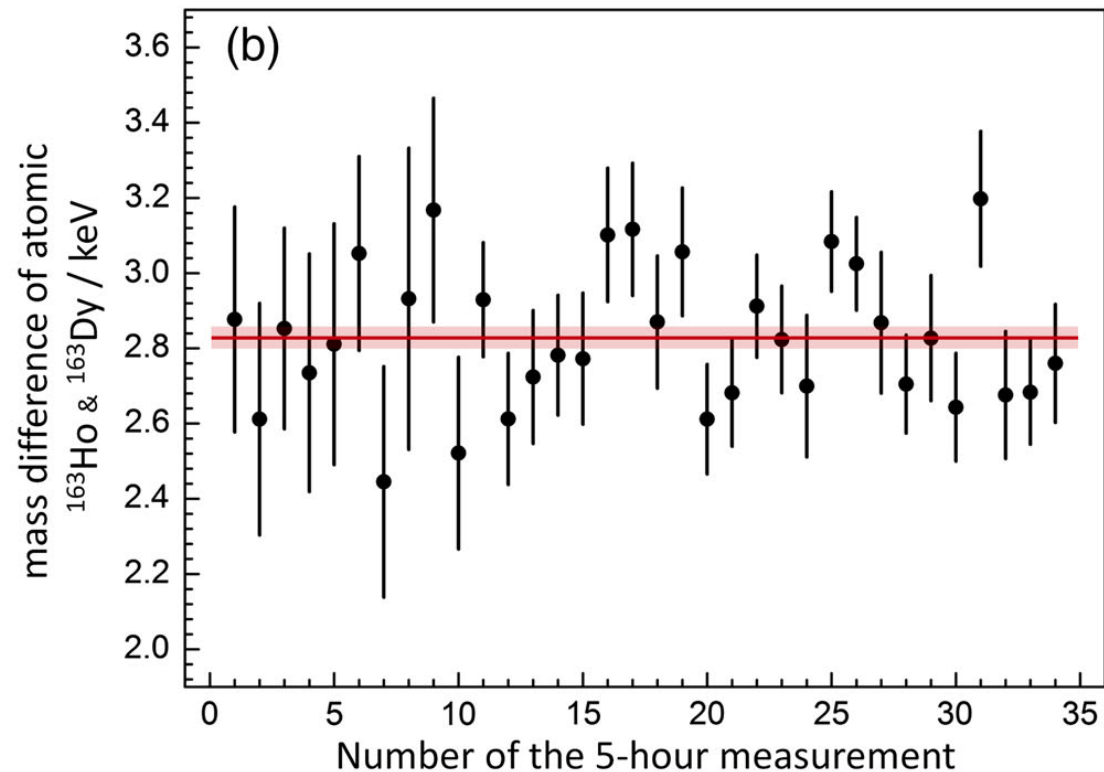
Pictures from
S. Eliseev *et al.*, PRL 115, (2015) 232501

Recent highlights

Penning-trap Q-value for neutrino physics at GSI (2015)



$$\Delta m = 2833(30_{\text{stat}})(15_{\text{sys}}) \text{ eV}/c^2$$



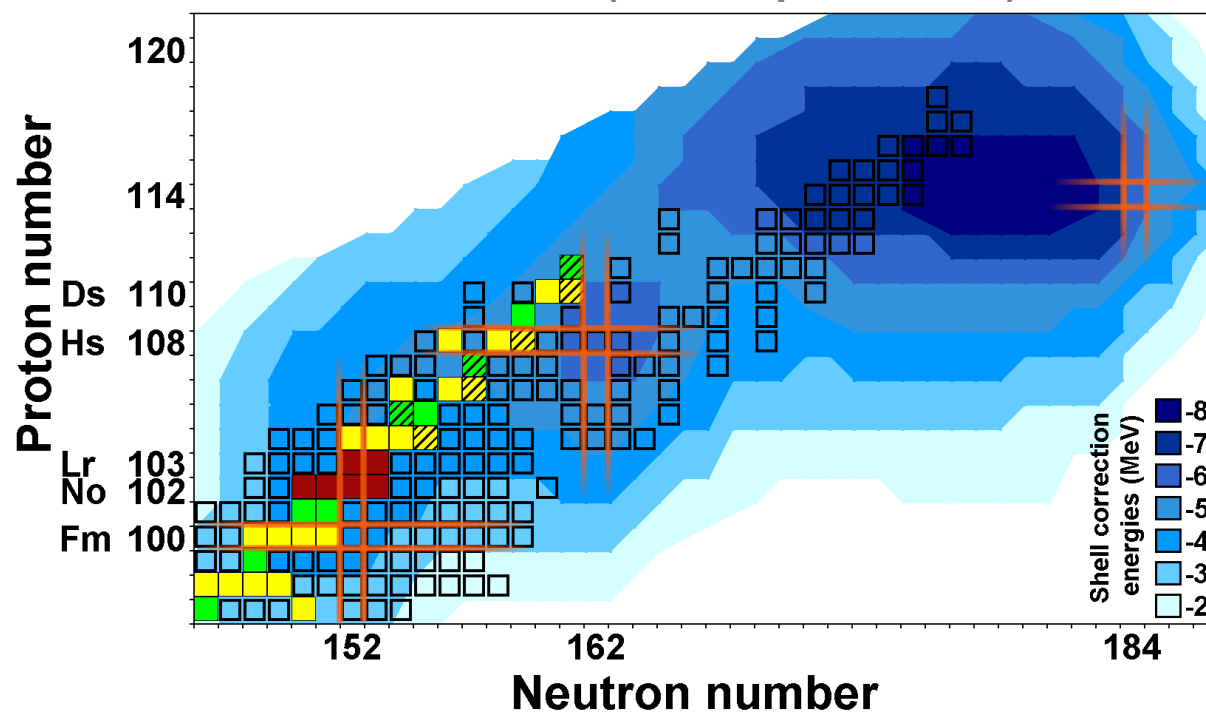
Pictures from S. Eliseev *et al.*, PRL 115, (2015) 062501

THE PROJECT TRAPSENSOR IN GRANADA

The TRAPSENSOR project

The technique

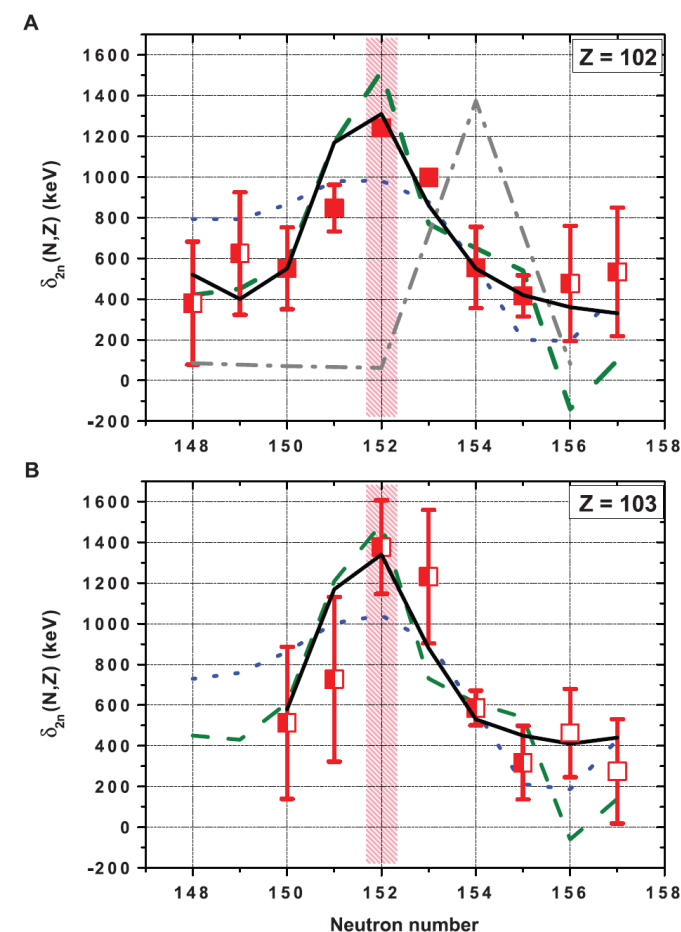
Nuclear Structure (minute production)



- The isotope with lowest production rate ever measured in a Penning trap (^{256}Lr , 60(18) nb)
→ 48 ions detected in 93 hours !!!

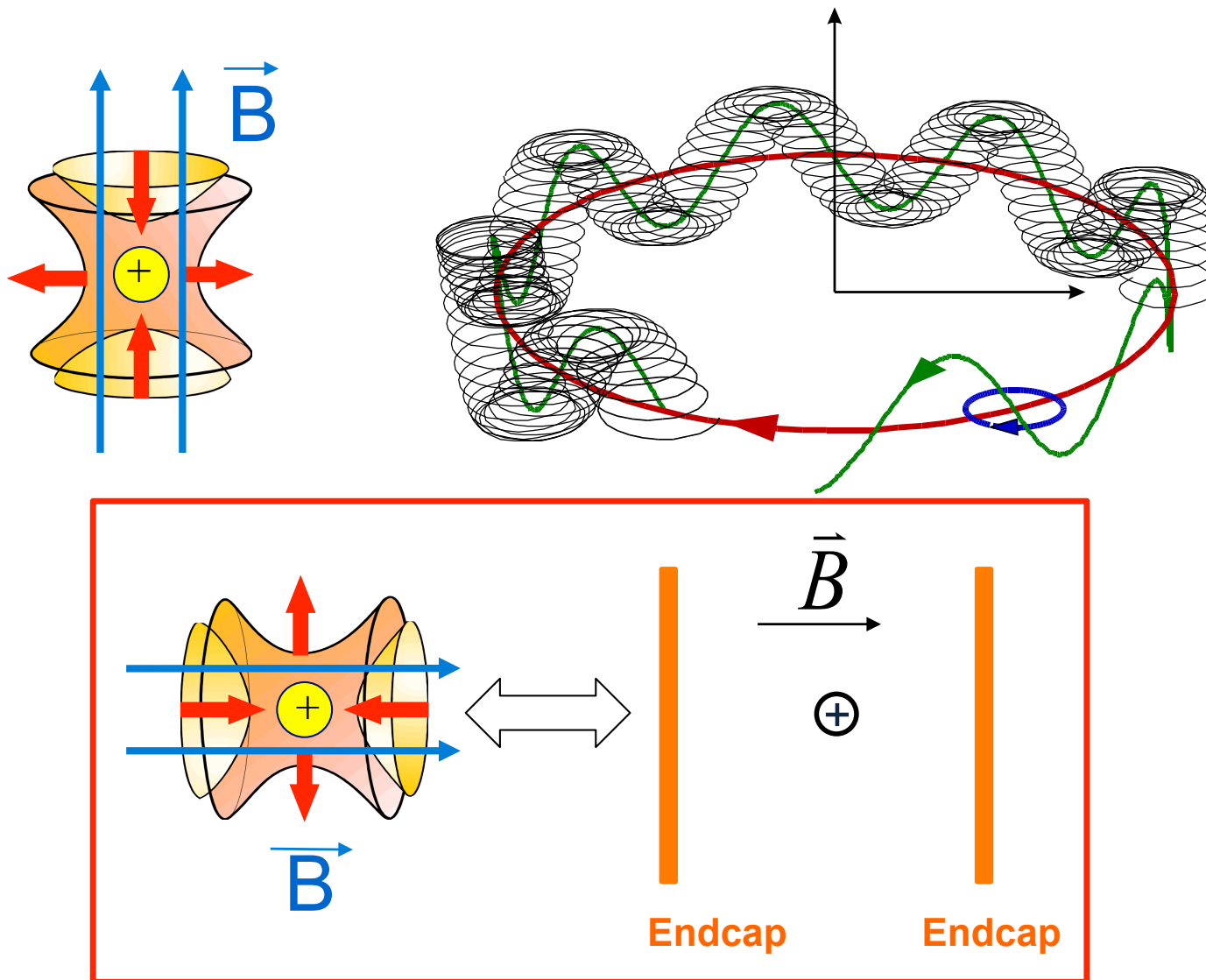
E. Minaya Ramirez *et al.*, Science 337 (2012) 1207

$$\begin{aligned}\delta_{2n}(N,Z) &= S_{2n}(N,Z) - S_{2n}(N+2,Z) \\ &= -2M_{\text{exc}}(N,Z) + M_{\text{exc}}(N-2,Z) \\ &\quad + M_{\text{exc}}(N+2,Z),\end{aligned}$$



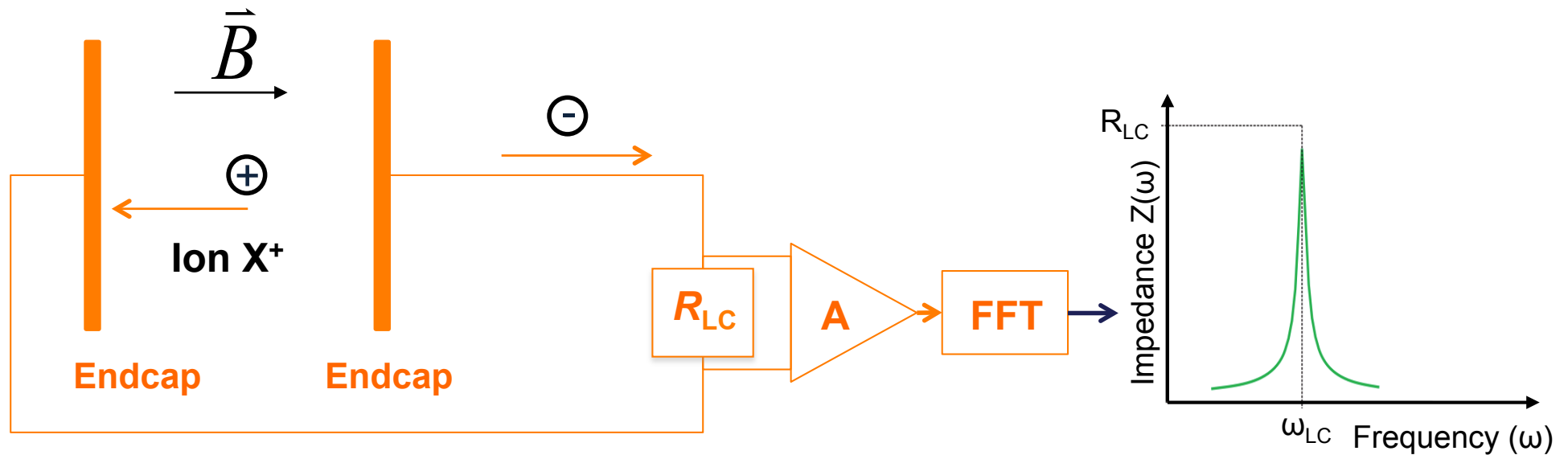
The TRAPSENSOR project

The technique



The TRAPSENSOR project

The technique



Induced image-current $\rightarrow$ Oscillation frequency $\rightarrow m/q$

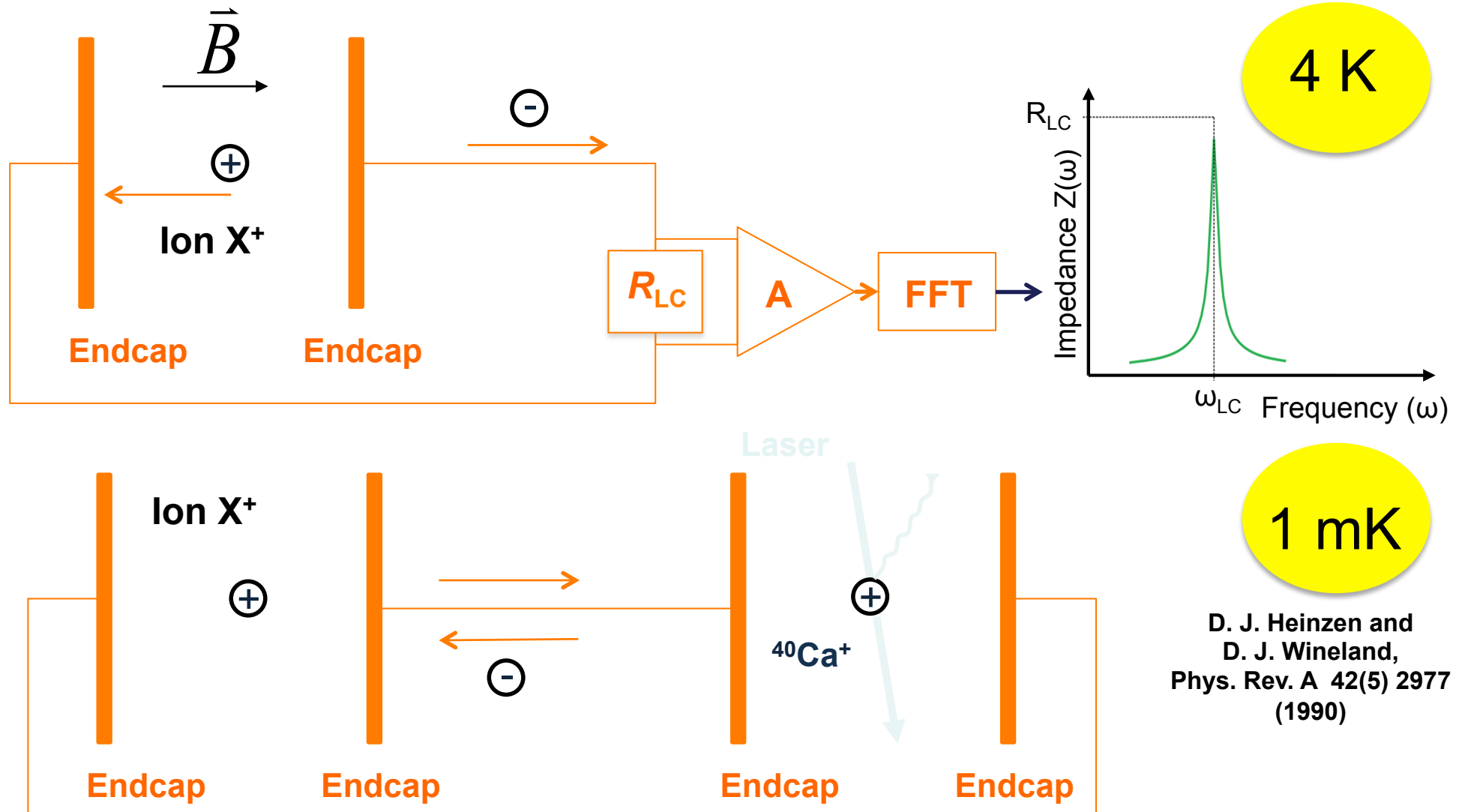
$$\omega_+^2 + \omega_z^2 + \omega_-^2 = \omega_c^2$$

L. S. Brown and G. Gabrielse,
Rev. Mod. Phys. 58, 233 (1986)

It has been demonstrated for low or medium m/q

The TRAPSENSOR project

The technique



D. J. Heinzen and
D. J. Wineland,
Phys. Rev. A 42(5) 2977
(1990)

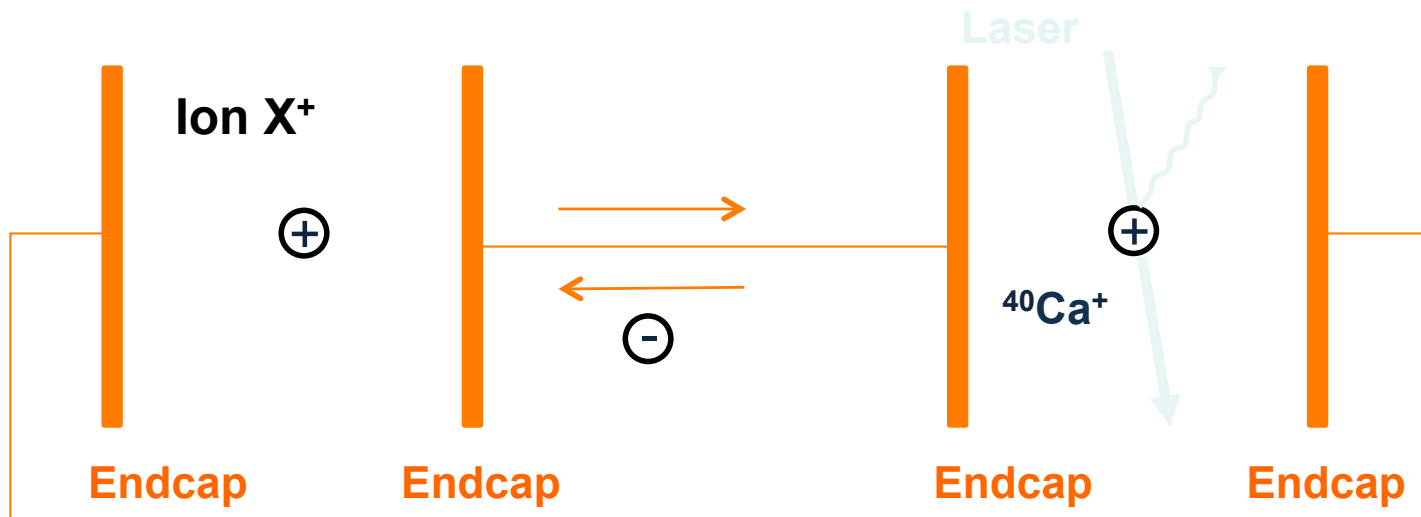
The TRAPSENSOR project

The technique

1. Ion X^+ is probed with a radiofrequency field and will be detected in the optical regime.
2. Ion X^+ can be any ion regardless of its mass, charge and polarity.



MICRO-TRAPS + LASERS + EXTERNAL ION SOURCE + PENNING TRAP BEAMLINE

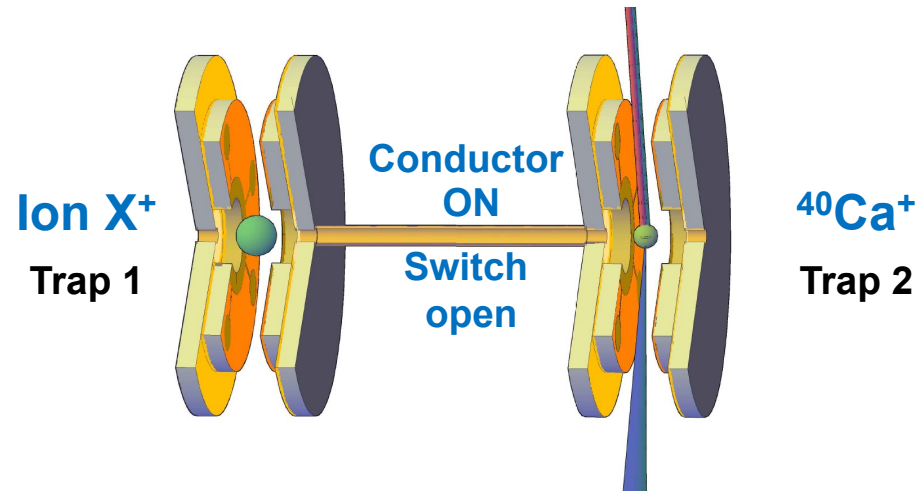


D. J. Heinzen and
D. J. Wineland,
Phys. Rev. A 42(5) 2977
(1990)

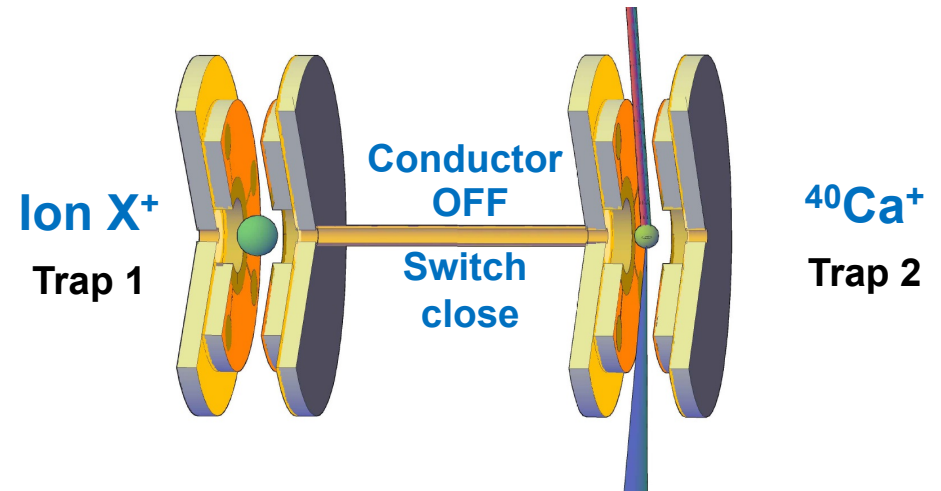
The project TRAPSENSOR

The technique

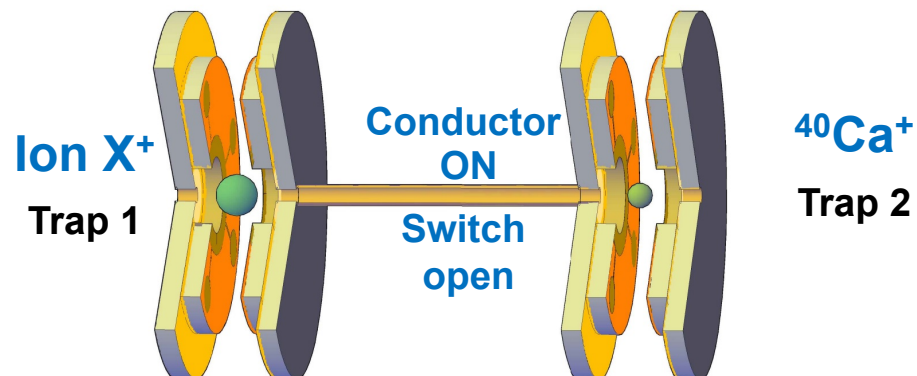
1. COOLING



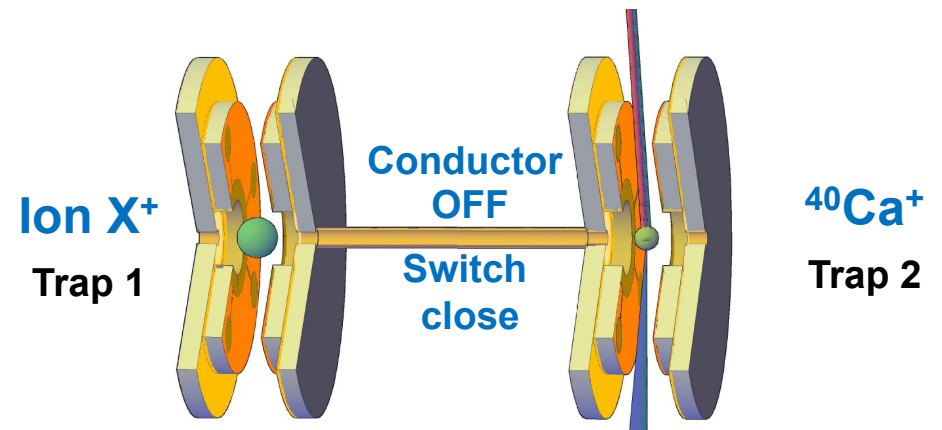
2. PROBING ION X^+



3. ION-ION INTERACTION

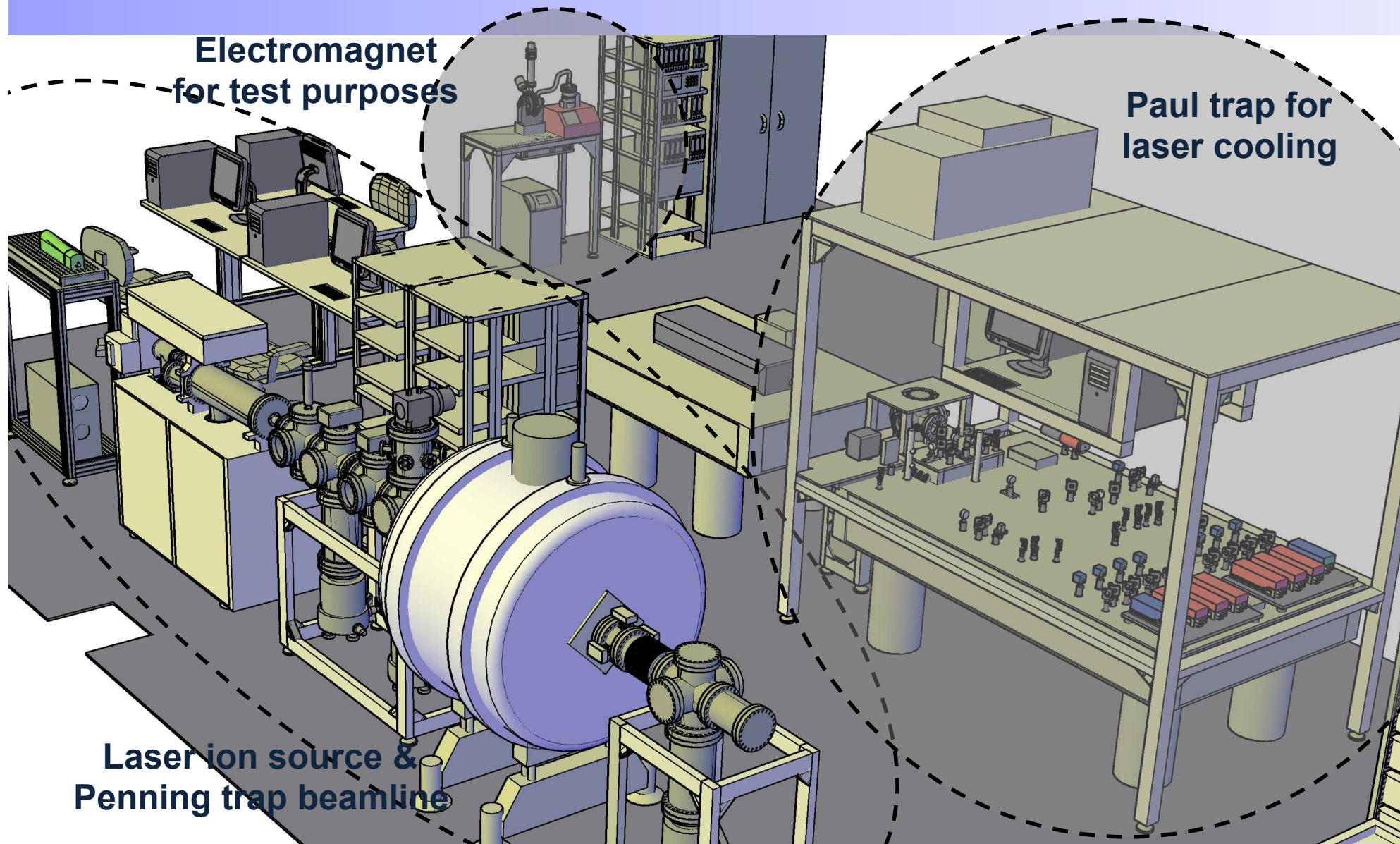


4. INTERROGATION



The project TRAPSENSOR

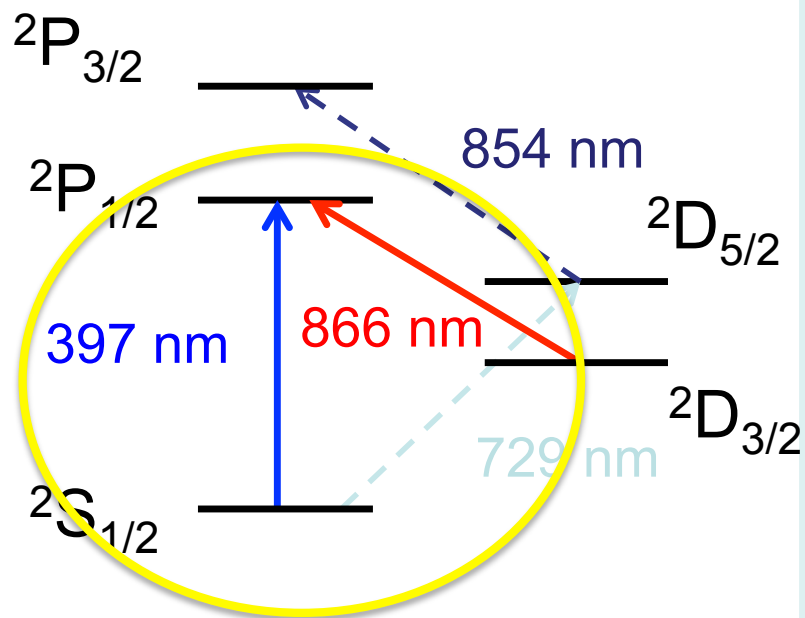
An overview of a Penning-trap beamline



The project TRAPSENSOR

Laser system

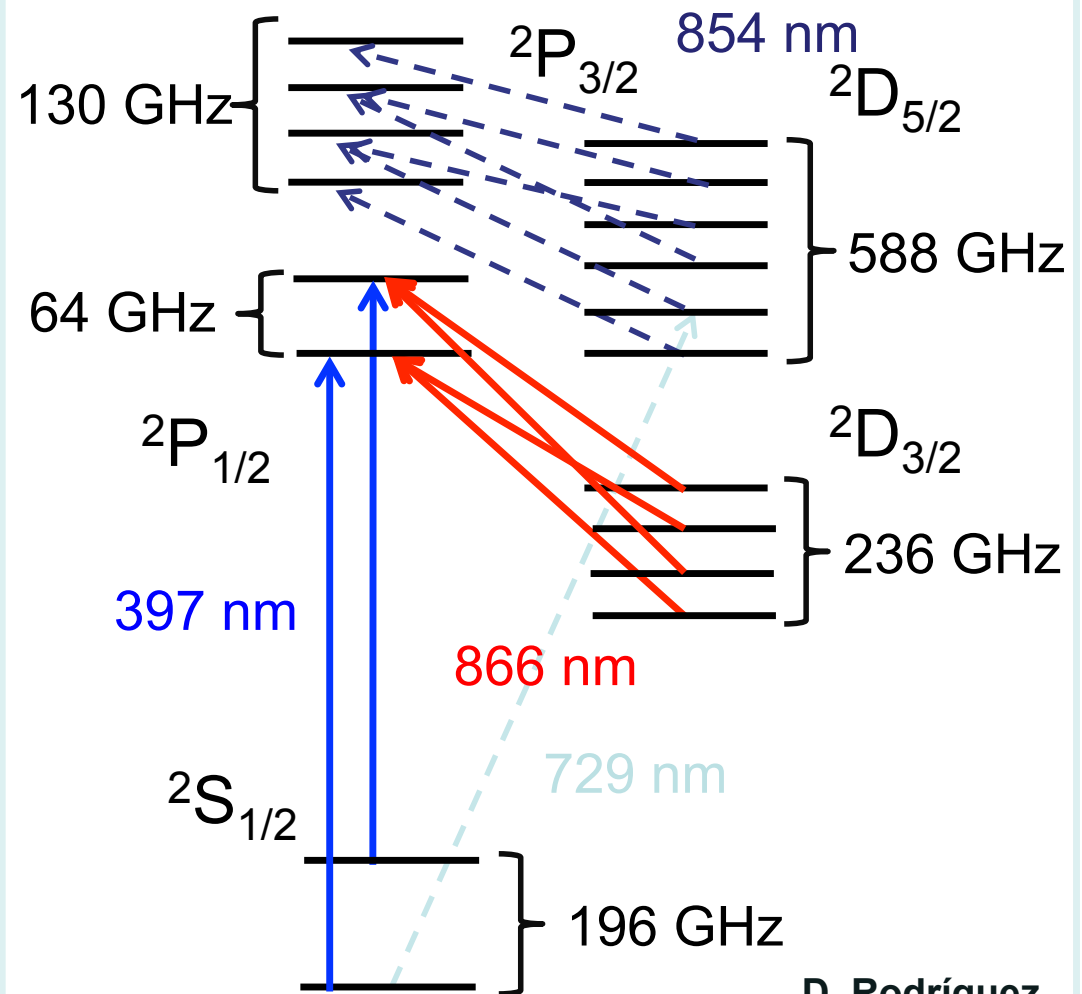
1. Level scheme for a $^{40}\text{Ca}^+$ ion in a Paul trap



**Doppler
cooling**

“Curso FNEXP” Valencia • January 2016

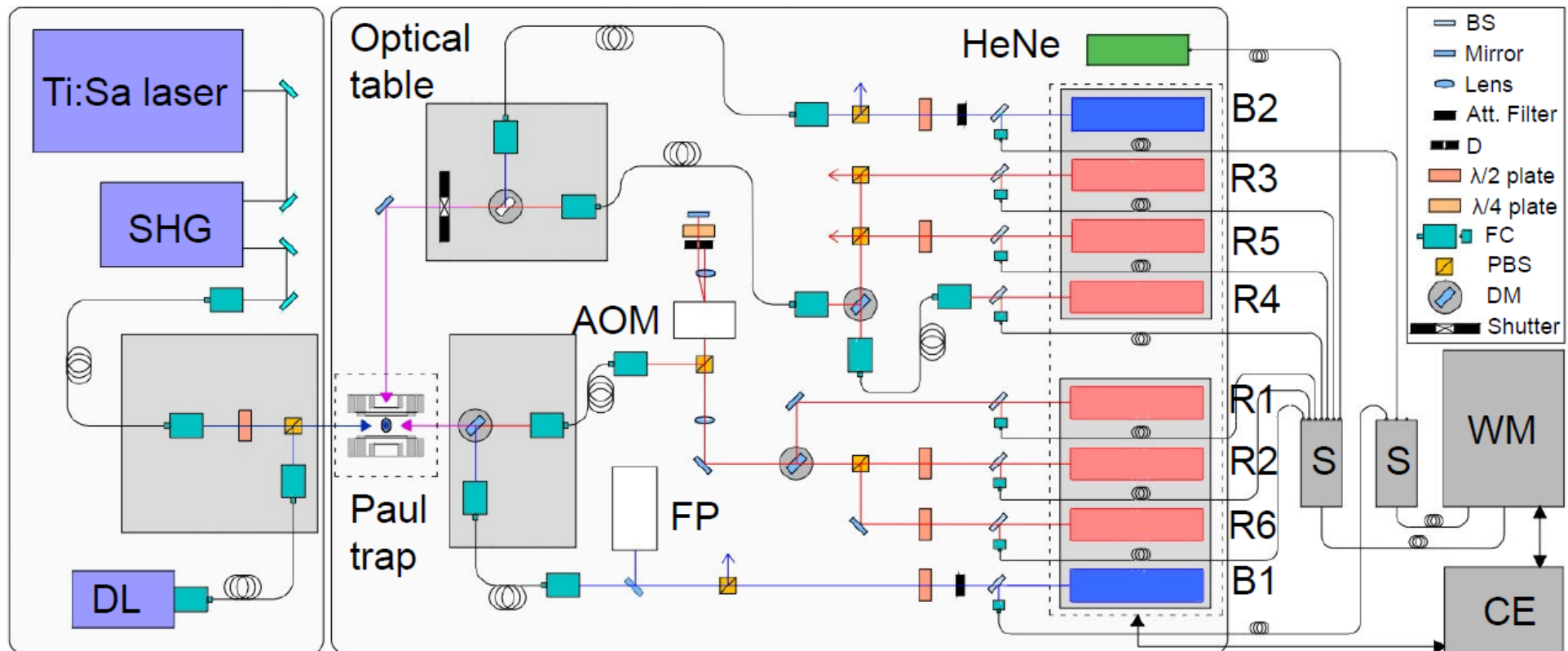
2. Level scheme for a $^{40}\text{Ca}^+$ ion in a 7-T magnetic field



D. Rodríguez

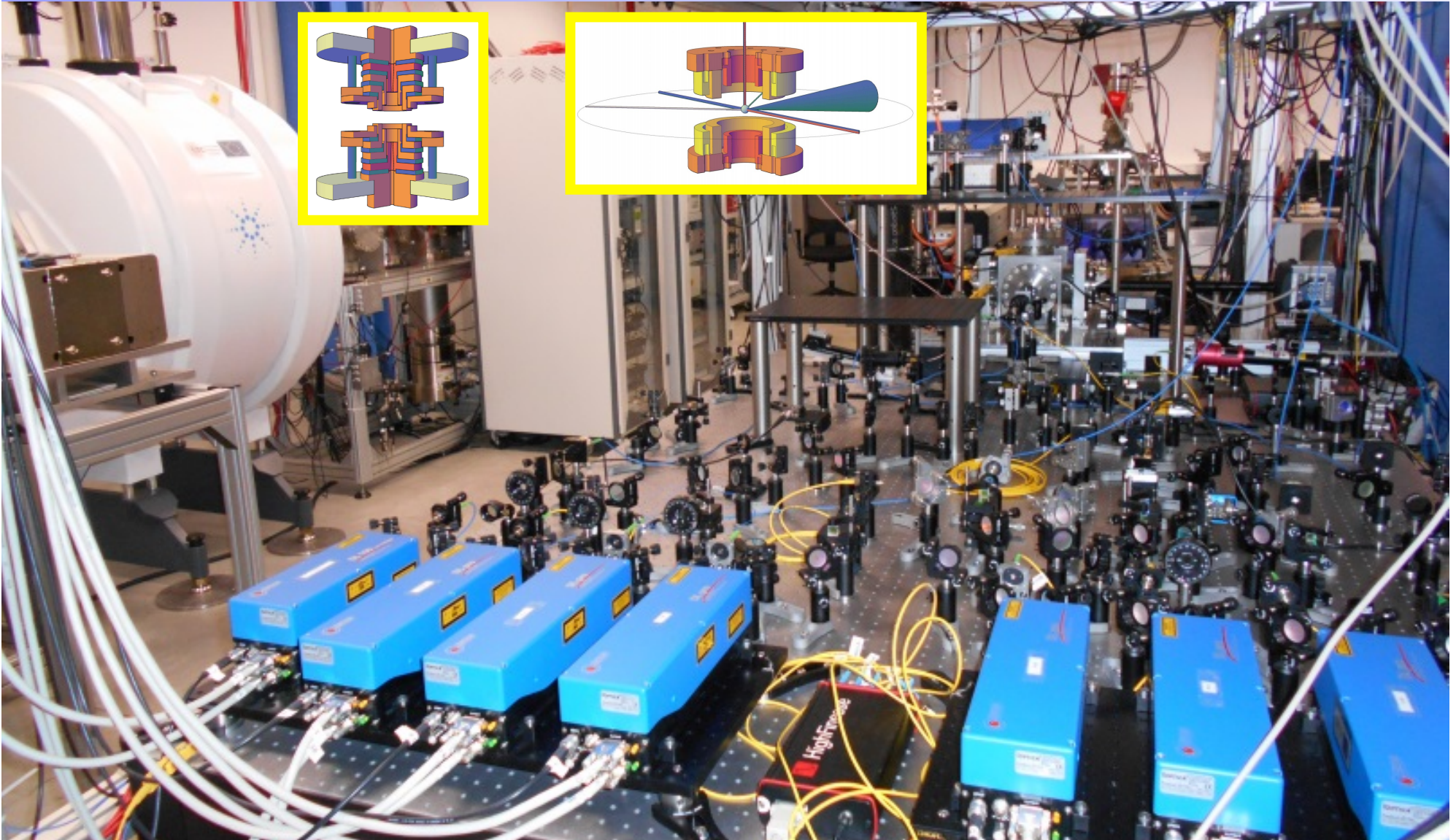
The project TRAPSENSOR

Laser system



The project TRAPSENSOR

Laser system

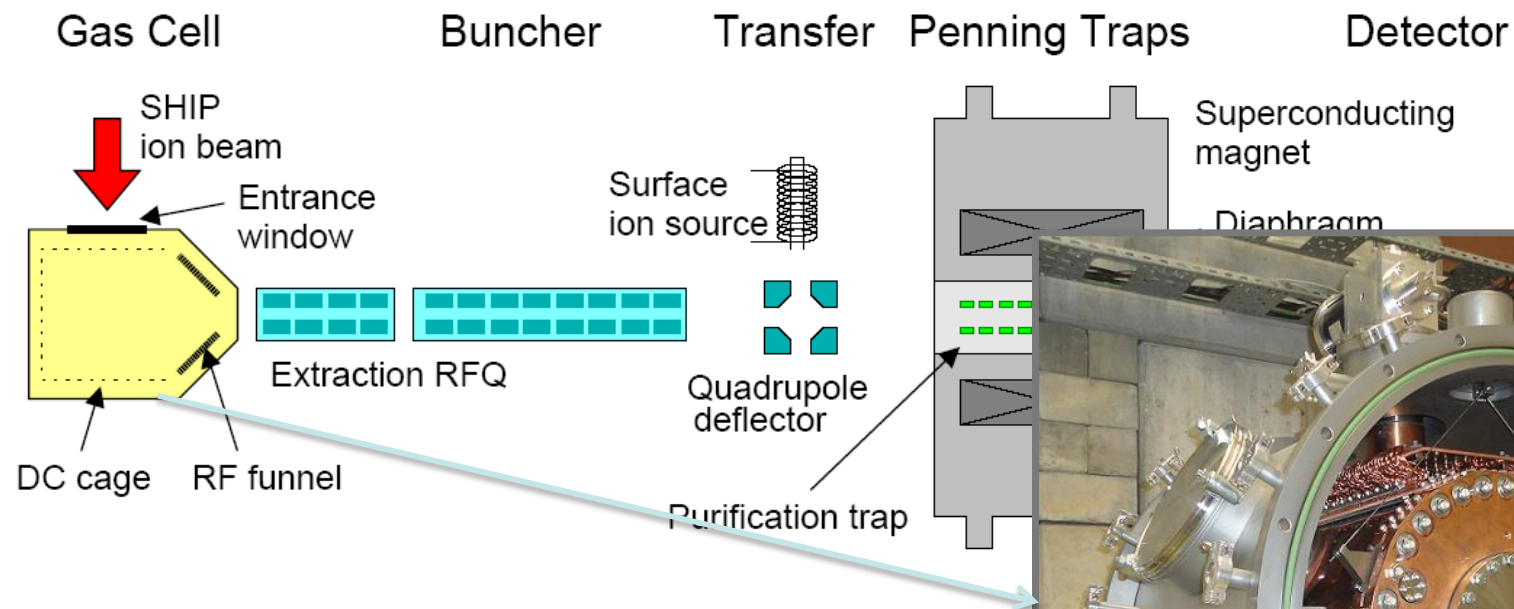


The project TRAPSENSOR SHIPTRAP beamline

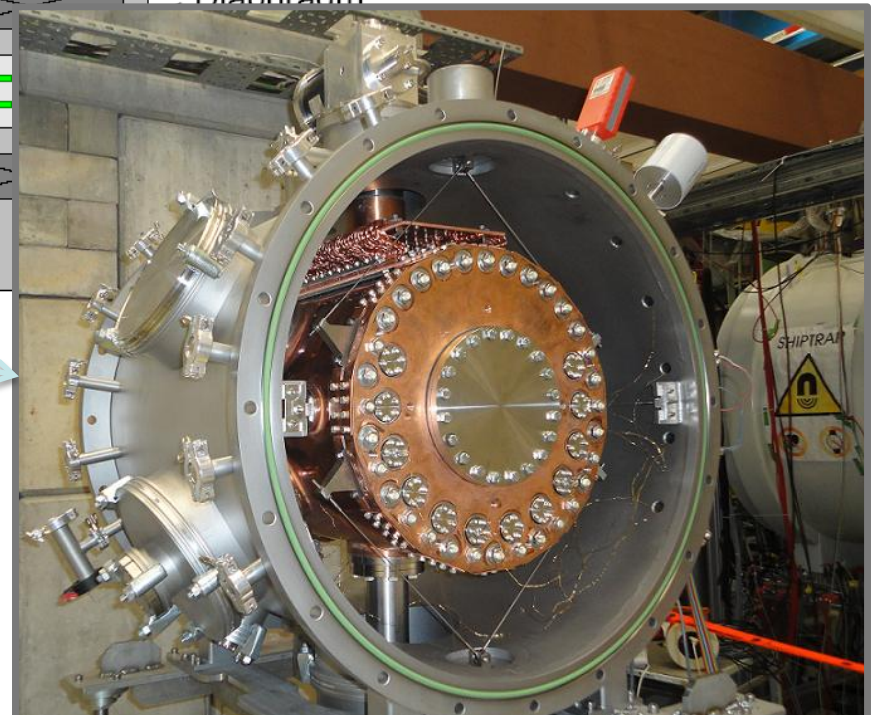
$\approx 50 \text{ MeV}$



$\approx 1 \text{ eV}$

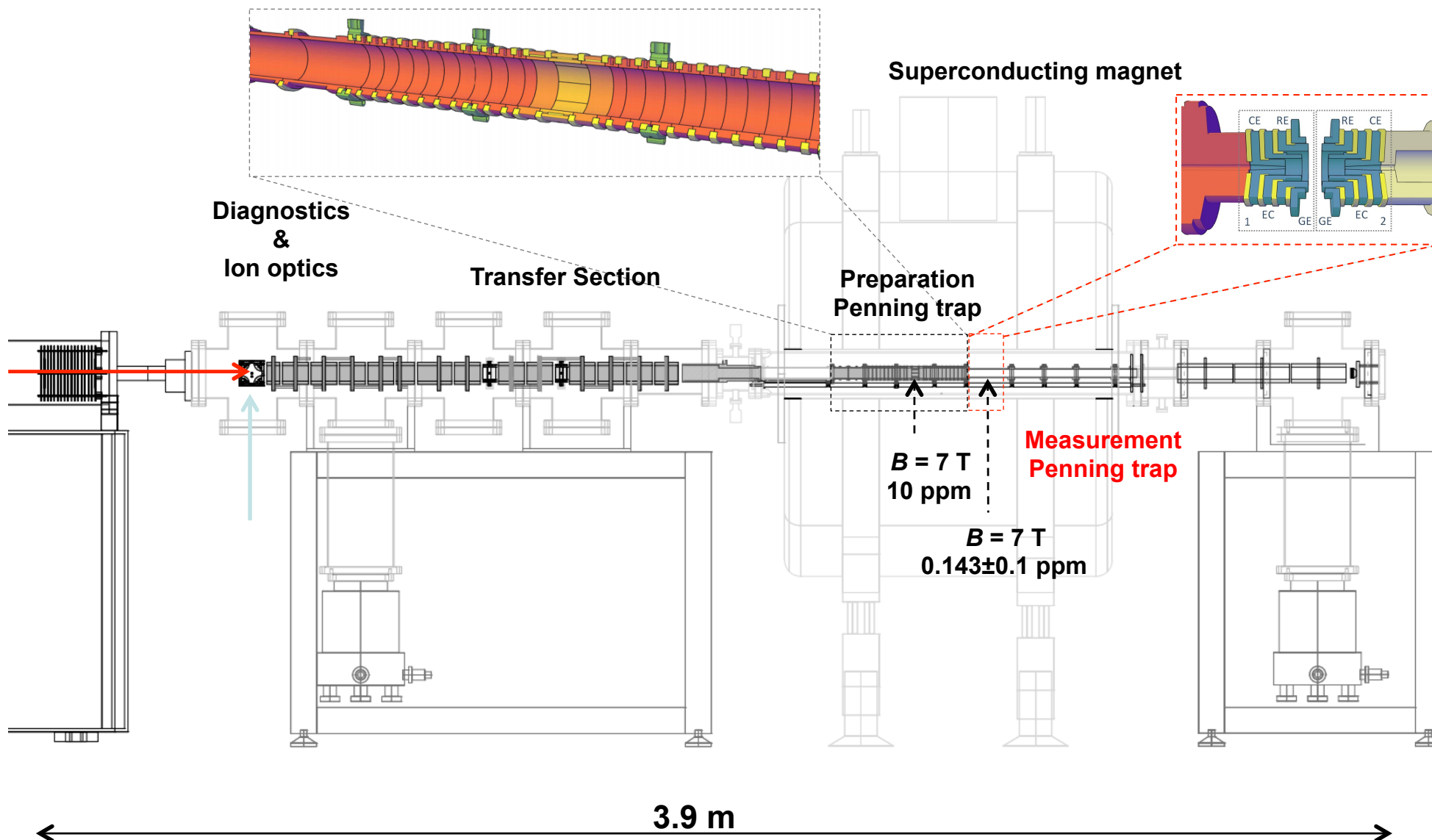


Several developments are on-going



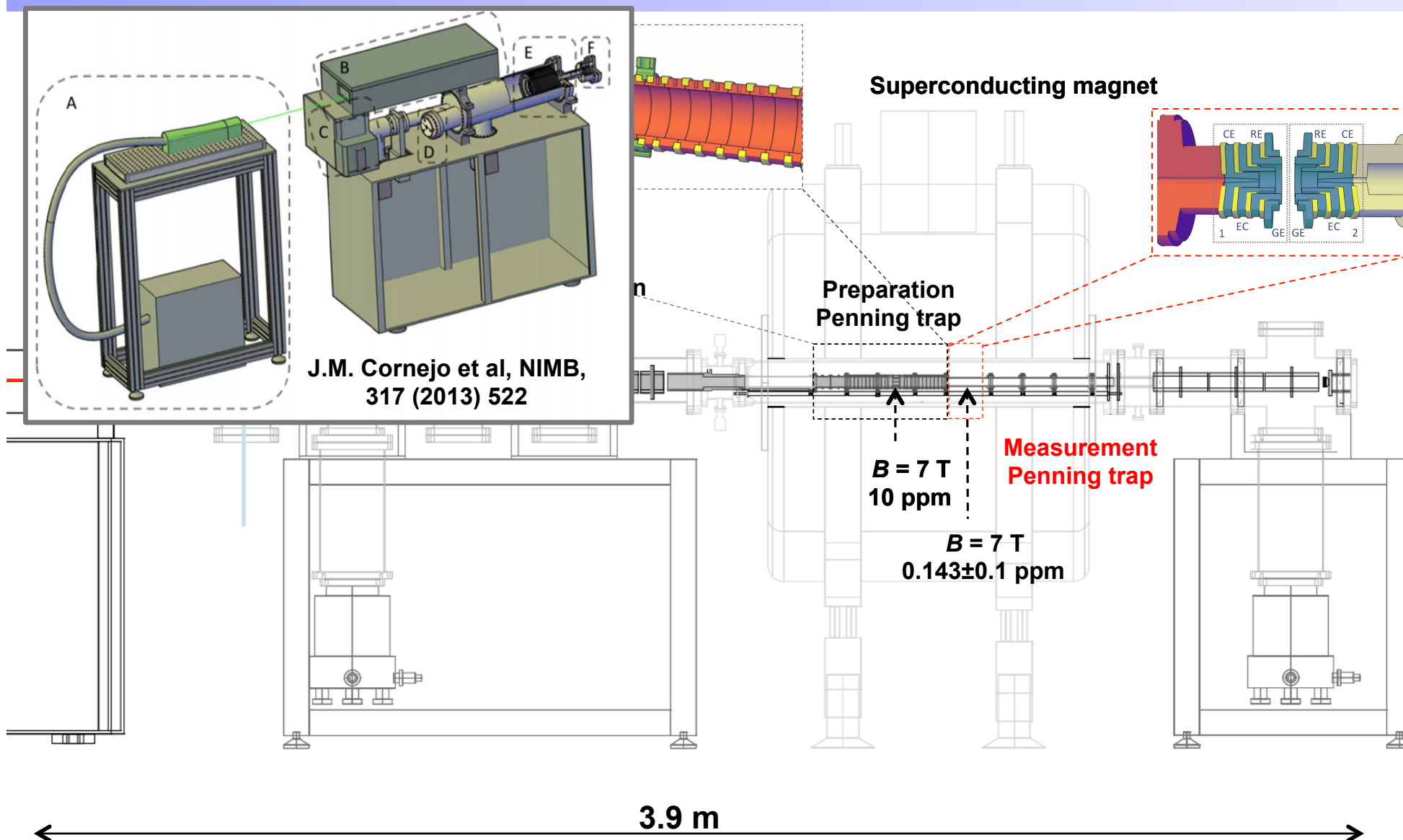
The project TRAPSENSOR

The Penning-traps beamline



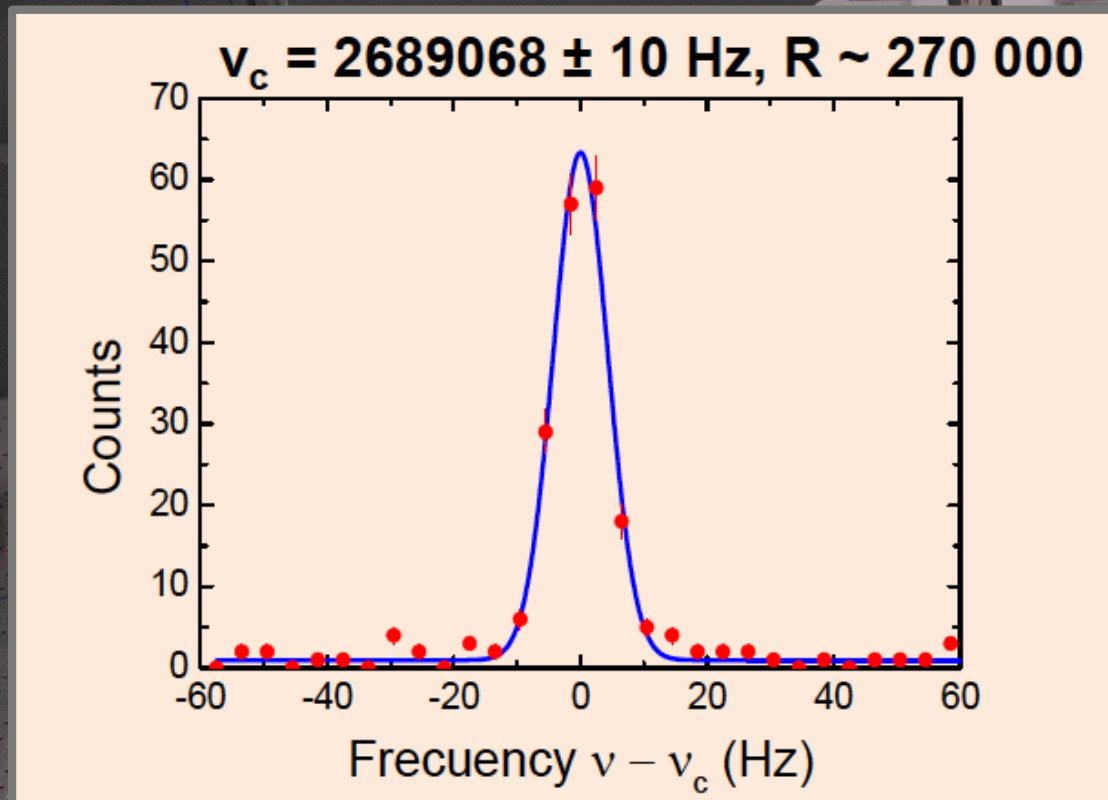
The project TRAPSENSOR

The Penning-traps beamline



The project TRAPSENSOR

The preparation Penning trap

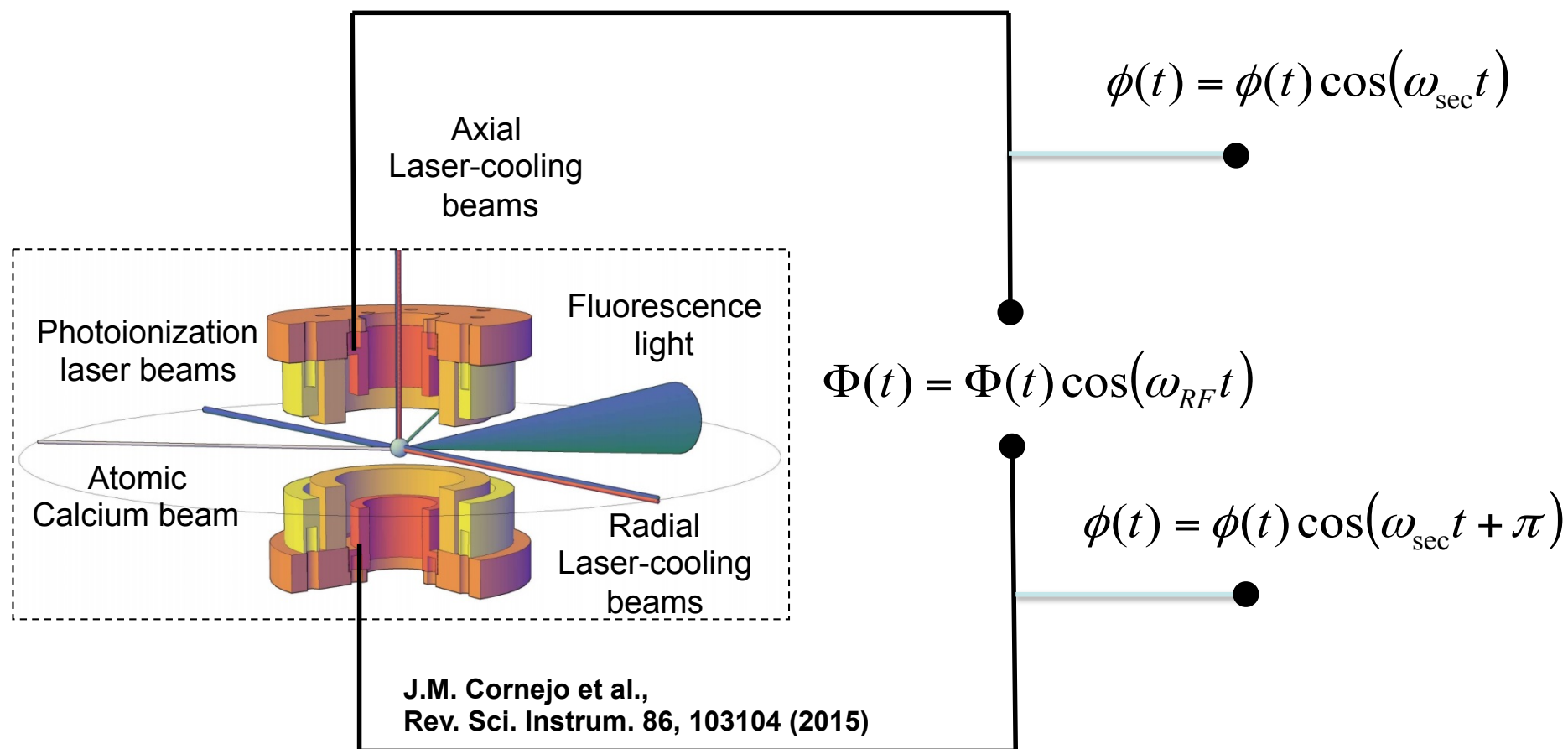


J.M. Cornejo and D. Rodríguez,
NIMB, in press, corrected proofs



The project TRAPSENSOR

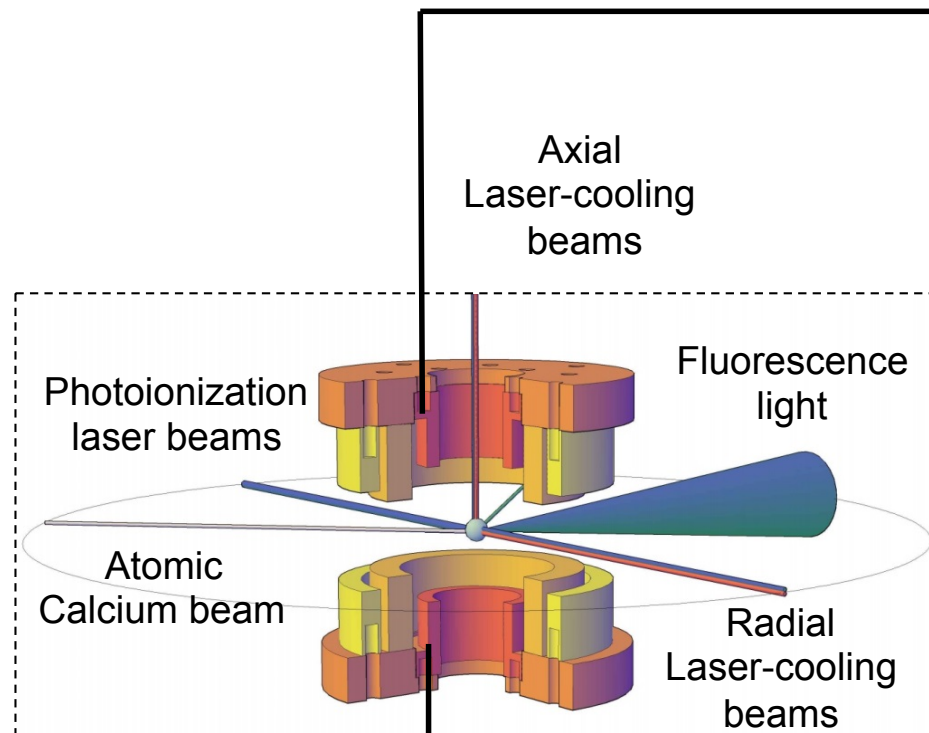
First experiment (laser cooling)



**Measuring the sensitivity of the laser-cooled $^{40}\text{Ca}^+$ ion
using an external RF field**

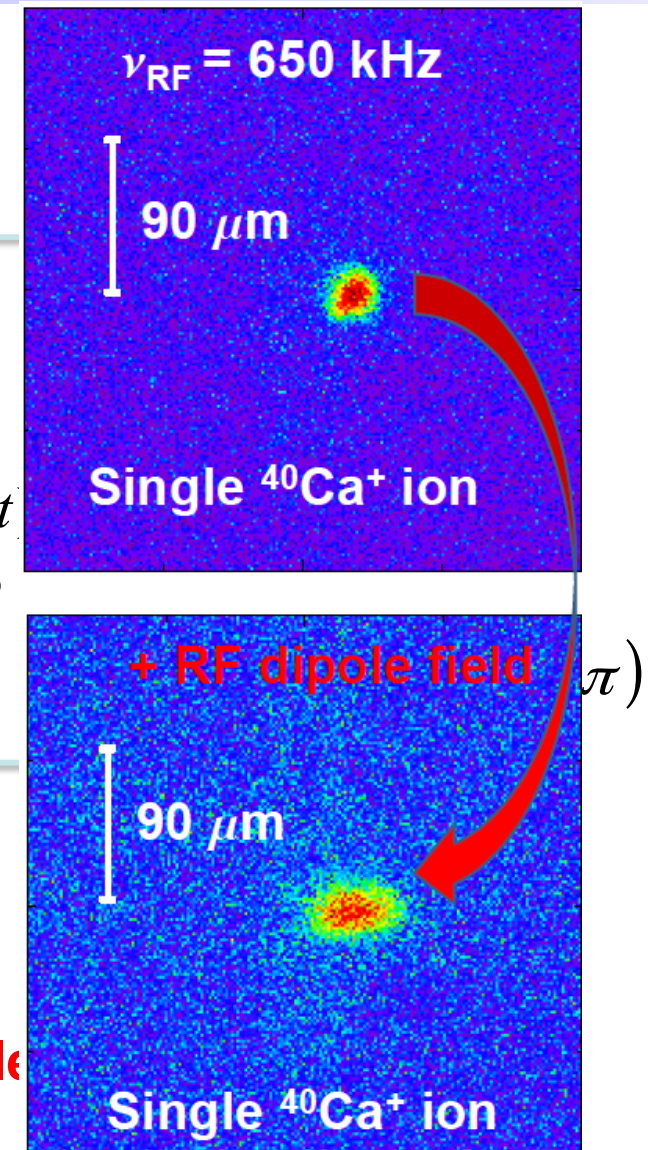
The project TRAPSENSOR

First experiment (laser cooling)



J.M. Cornejo et al.,
Rev. Sci. Instrum. 86, 103104 (2015)

$$\Phi(t) = \Phi(t)$$



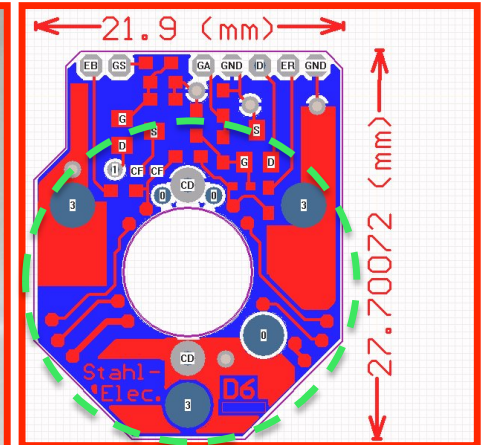
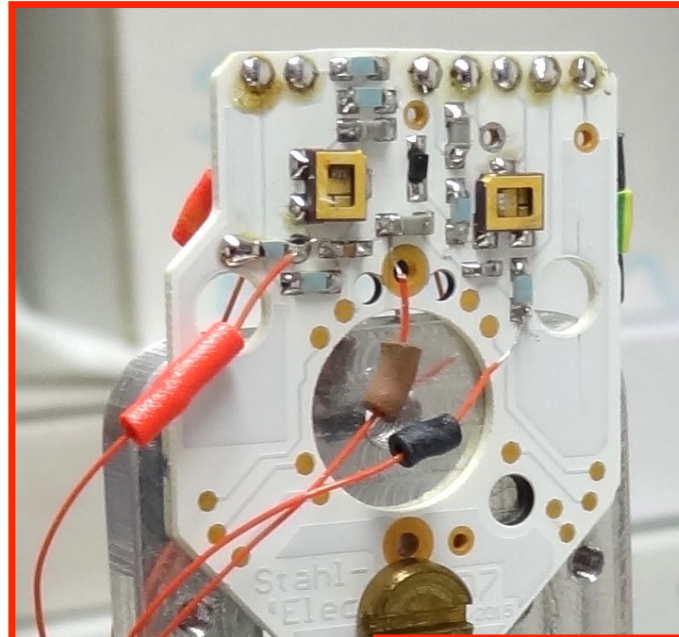
**Measuring the sensitivity of the laser-cooled ion
using an external RF field**

The project TRAPSENSOR

On-going work

I. Micro-trap operation at room temperature

- The final geometry was decided in September 2014.
- The electrodes have been fabricated by TRANSLUME (May 2015).
- Electrodes are now in the US for coating and electroplating
- The electronics has been built by Stahl electronics.
- A new vacuum chamber is ready for the commissioning of the system.



II. Laser cooling in the 7-T magnetic field

- A new trap has been built to study laser cooling in the 7-T magnetic field.
- We have all the lasers and we have recently got the microwave generator to drive the EOM to generate the remaining 4 laser beams.

