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Top-down approach to EFTs

José Santiago



Based on:

F. del Aguila, Z. Kunszt, J.S. [1602.00126]

C. Anastasiou, A. Carmona, A. Lazopoulos, J.S. (in progress)

Effective theories: bottom-up

- Effective Lagrangians: model-independent description of new physics in the presence of a mass gap
- Map experimental (pseudo)observables to the Wilson coefficients in the effective Lagrangian to obtain all the experimental information in a model independent way
- Truly global fit to new physics now possible (EWPD plus LHC data -Higgs and otherwise-)

Ciuchini, Franco, Mishima, Silvestrini ('13); Blas, Chala, J.S. ('13, '15); Pomarol, Riva ('14); Falkowski, Riva ('15); Buckley, Englert, Ferrando, Miller, Moore, Russell, White ('15); Berthirer, Trott ('15), ...

- Efforts to extend to NLO already on the way

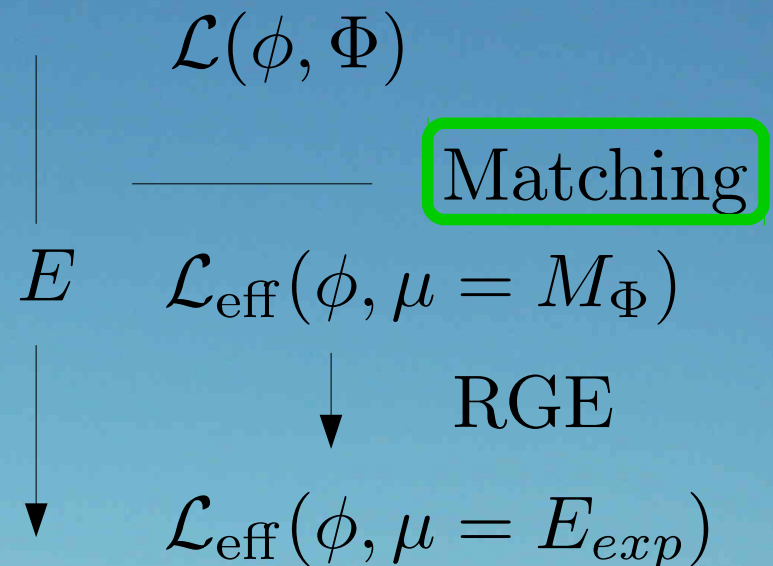
Ghezzi, Gomez-Ambrosio, Passarino, Uccirati ('15), Hartmann, Trott ('15), David, Passarino ('15), Boggia, Gomez-Ambrosio, Passarino ('16) ...



Effective theories: top-down

- A complementary approach is to consider specific UV completions

- Correlations among Wilson coefficients in specific models (eventually observable in data)
- Better control over range of validity
- Give up model-independence? Not if all models considered



- Goal: Generate a UV/IR dictionary (map all possible UV completions of the SM to the Wilson coefficients of the SM effective Lagrangian)

Outline

- UV/IR tree-level dictionary
- UV/IR one-loop dictionary
 - Effective Lagrangian at one loop: functional methods and matching
 - Matchmaker: automated one-loop matching in effective theories
 - Consistent one-loop calculations with EFTs from the top-down
- Conclusions and outlook



Tree-level dictionary (non-mixed contributions)

New Quarks:
F. Aguila, M. Perez-Victoria, J.S., JHEP (00)

$Q^{(m)}$	U	D	$\begin{pmatrix} U \\ D \end{pmatrix}$	$\begin{pmatrix} X \\ U \end{pmatrix}$	$\begin{pmatrix} D \\ Y \end{pmatrix}$	$\begin{pmatrix} X \\ U \\ D \end{pmatrix}$	$\begin{pmatrix} U \\ D \\ Y \end{pmatrix}$
isospin	0	0	1/2	1/2	1/2	1	1
hypercharge	2/3	-1/3	1/6	7/6	-5/6	2/3	-1/3

New Leptons:
F. Aguila, J. Blas, M. Perez-Victoria, PRD (08)

Leptons	N	E	$\begin{pmatrix} N \\ E^- \end{pmatrix}$	$\begin{pmatrix} E^- \\ E^{--} \end{pmatrix}$	$\begin{pmatrix} E^+ \\ N \\ E^- \end{pmatrix}$	$\begin{pmatrix} N \\ E^- \\ E^{--} \end{pmatrix}$
Notation			Δ_1	Δ_3	Σ_0	Σ_1
$SU(2)_L \otimes U(1)_Y$	1_0	1_{-1}	$2_{-(1/2)}$	$2_{-(3/2)}$	3_0	3_{-1}
Spinor	Dirac or Majorana	Dirac	Dirac	Dirac	Dirac or Majorana	Dirac

New Vectors:
F. Aguila, J. Blas, M. Perez-Victoria, JHEP (10)

Vector	\mathcal{B}_μ	\mathcal{B}_μ^1	\mathcal{W}_μ	\mathcal{W}_μ^1	\mathcal{G}_μ	\mathcal{G}_μ^1	\mathcal{H}_μ	\mathcal{L}_μ
Irrep	$(1, 1)_0$	$(1, 1)_1$	$(1, \text{Adj})_0$	$(1, \text{Adj})_1$	$(\text{Adj}, 1)_0$	$(\text{Adj}, 1)_1$	$(\text{Adj}, \text{Adj})_0$	$(1, 2)_{-\frac{3}{2}}$
Vector	\mathcal{U}_μ^2	\mathcal{U}_μ^5	\mathcal{Q}_μ^1	\mathcal{Q}_μ^5	\mathcal{X}_μ	\mathcal{Y}_μ^1	\mathcal{Y}_μ^5	
Irrep	$(3, 1)_{\frac{2}{3}}$	$(3, 1)_{\frac{5}{3}}$	$(3, 2)_{\frac{1}{6}}$	$(3, 2)_{-\frac{5}{6}}$	$(3, \text{Adj})_{\frac{2}{3}}$	$(\bar{6}, 2)_{\frac{1}{6}}$	$(\bar{6}, 2)_{-\frac{5}{6}}$	

New Scalars:
J. Blas, M. Chala, M. Perez-Victoria, J.S., JHEP (15)

Colorless Scalars	\mathcal{S}	\mathcal{S}_1	\mathcal{S}_2	φ	Ξ_0	Ξ_1	Θ_1	Θ_3
Irrep	$(1, 1)_0$	$(1, 1)_1$	$(1, 1)_2$	$(1, 2)_{\frac{1}{2}}$	$(1, 3)_0$	$(1, 3)_1$	$(1, 4)_{\frac{1}{2}}$	$(1, 4)_{\frac{3}{2}}$
Colored Scalars	ω_1	ω_2	ω_4	Π_1	Π_7	ζ		
Irrep	$(3, 1)_{-\frac{1}{3}}$	$(3, 1)_{\frac{2}{3}}$	$(3, 1)_{-\frac{4}{3}}$	$(3, 2)_{\frac{1}{6}}$	$(3, 2)_{\frac{7}{6}}$	$(3, 3)_{-\frac{1}{3}}$		
Colored Scalars	Ω_1	Ω_2	Ω_4	Υ	Φ			
Irrep	$(6, 1)_{\frac{1}{3}}$	$(6, 1)_{-\frac{2}{3}}$	$(6, 1)_{\frac{4}{3}}$	$(6, 3)_{\frac{1}{3}}$	$(8, 2)_{\frac{1}{2}}$			

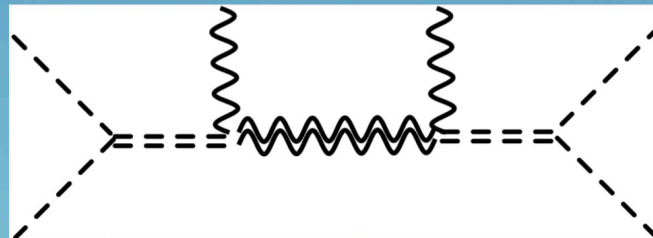
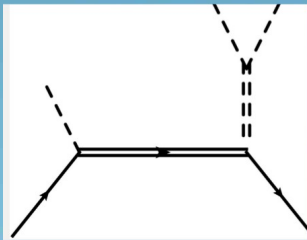
Tree-level dictionary (mixed contributions)

Mixed contributions: J. Blas, M. Chala, J.C. Criado, M. Perez-Victoria, J.S., to appear soon

- Dimensionful couplings imply that particles with different spins can simultaneously contribute to \mathcal{L}_6 at tree level

$$\kappa\phi_1\phi_2\phi_3 + \kappa'V^\mu D_\mu\phi + \kappa''V^\mu V'_\mu\phi + \dots$$

- We are currently classifying and computing all possible contributions



- Only a subset of the representations in the previous list contributes
- With this, the tree-level, dimension 6 UV/IR dictionary is complete: we can map arbitrary UV extensions to the SM EFT

One-loop UV/IR dictionary

- Many contributions to the effective Lagrangian can be only generated at the quantum level
- Even contributions that can potentially arise at tree-level only appear at loop level in specific models
- The dictionary should be extended to one loop if we want to account for these cases
- The number of possibilities increases dramatically: automation seems compulsory



Functional methods and matching

- Interesting progress has been recently made using functional methods Henning, Lu, Murayama ('14); Gaillard ('86); Cheyette ('86)
- There has been a great deal of developments in the last year: Henning, Lu, Murayama ('14); Drozd, Ellis, Quevillon, You ('15), ...
 - Initial attempts were not complete in the case of linear couplings to heavy states F. Aguila, Z. Kunszt, J.S. ('16); Bilenky, Santamaría ('94)
 - The missing terms are local and can only be recovered by matching which can be performed:
 - diagrammatically Anastasiou, Carmona, Lazopoulos, J.S.
 - by functional methods Henning, Lu, Murayama ('16); Ellis, Quevillon, You, Zhang ('16)



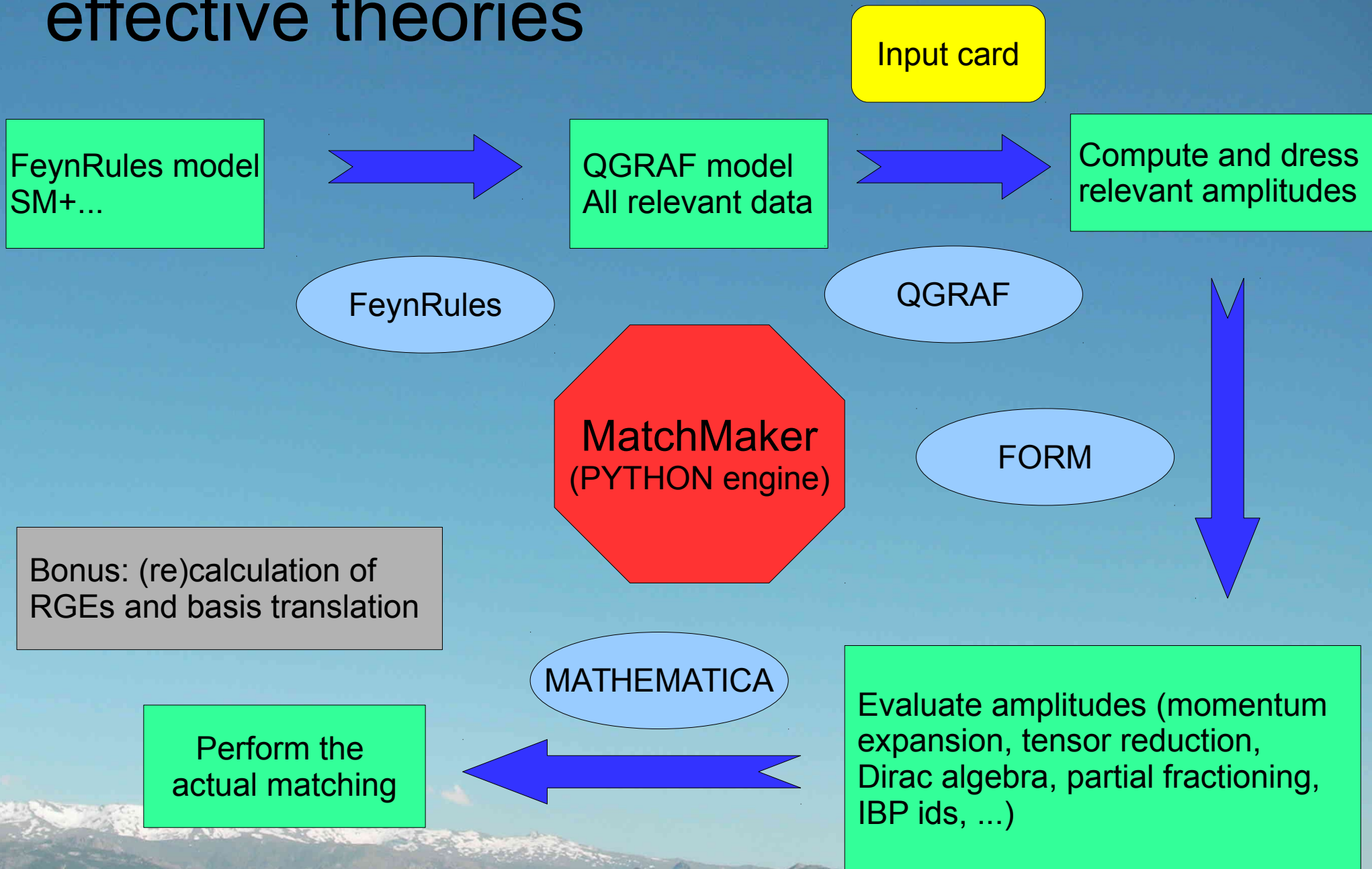
MatchMaker: automated matching in effective theories

Anastasiou, Carmona, Lazopoulos, J.S., in progress

- We are developing an automated tool to perform tree-level and one-loop matching of arbitrary theories into arbitrary effective Lagrangians
- Based on standard, well-tested tools (FeynRules, QGRAF, FORM, Mathematica, Python)
- Flexible (from full matching to specific operators), fully automated and general
- Unified treatment (effective theory just another model)
- Off-shell matching with (initially) massless particles in the effective theory (e.g. unbroken phase of the SM)



MatchMaker: automated matching in effective theories



How to use EFTs (from the top-down) at one loop

F. Aguila, Z. Kunszt, J.S., ('16)

- Sample result: T parameter from charge 2/3 vector-like quark singlet

$$\mathcal{L}_T = \bar{T}(i\not{D} - M)T - \left[\lambda_T \bar{q}_L \tilde{\phi} T_R + \text{h.c.} \right]$$

- Computed in the physical basis (full model)

$$\Delta\hat{T} = \frac{N_C}{32\pi^2} \frac{v^2}{M^2} \left[|\lambda_T|^4 + 2\lambda_t^2 |\lambda_T|^2 \left(\log \frac{M^2}{m_t^2} - 1 \right) \right] \quad \text{Carena, Ponton, J.S., Wagner ('06)}$$

- Computed in an EFT approach (3 steps)
 - Matching at M
 - Running to m_t
 - Matching at m_t

How to use EFTs (from the top-down) at one loop

- Sample result: T parameter from charge 2/3 vector-like quark singlet
 - Matching at M: off-shell (3 independent operators)
 $\mathcal{O}_1 = |\phi^\dagger D_\mu \phi|^2 \quad \mathcal{O}_2 = \phi^\dagger \phi \partial^2 \phi^\dagger \phi \quad \mathcal{R} = \phi^\dagger \phi \phi^\dagger D^2 \phi$
 - Compute $\langle H_1 H_1^* H_2 H_2^* \rangle$ in full and effective theories

$$\alpha_1^{(1l)} = \frac{N_C |\lambda_T|^2}{16\pi^2 M^2} \left(\frac{1}{2} \lambda_t^2 - \frac{1}{2} |\lambda_T|^2 \right),$$
$$\alpha_2^{(1l)} = \frac{N_C |\lambda_T|^2}{16\pi^2 M^2} \left(\frac{3}{2} \lambda_t^2 - \frac{1}{3} |\lambda_T|^2 \right),$$
$$\alpha_{\mathcal{R}}^{(1l)} = \frac{N_C |\lambda_T|^2}{16\pi^2 M^2} \left(-\frac{1}{2} \lambda_t^2 + \frac{1}{2} |\lambda_T|^2 \right),$$

$$\Delta \hat{T} = -v^2 \alpha_1$$

$$\Delta \hat{T} = \frac{N_C}{32\pi^2} \frac{v^2}{M^2} \left[|\lambda_T|^4 + 2\lambda_t^2 |\lambda_T|^2 \left(\log \frac{M^2}{m_t^2} - 1 \right) \right]$$

How to use EFTs (from the top-down) at one loop

- Sample result: T parameter from charge 2/3 vector-like quark singlet (Alonso), Jenkins, Manohar, Trott ('13); Elias-Miró, Espinosa, Masso, Pomarol ('13); Elias-Miró, Grojean, Gupta, Marzocca ('13)
 - Running to m_t : tree-level operators relevant

$$\Delta\hat{T} = -v^2\alpha_1$$

$$16\pi^2 \frac{d\alpha_1}{d\log\mu} = 8N_C\lambda_t^2\alpha_{\phi q}^{(1)} + \dots,$$

$$\Delta\hat{T} = \frac{N_C}{32\pi^2} \frac{v^2}{M^2} \left[|\lambda_T|^4 + 2\lambda_t^2|\lambda_T|^2 \left(\log \frac{M^2}{m_t^2} - 1 \right) \right]$$

$$\mathcal{O}_{\phi q}^{(1)} = i\phi^\dagger D_\mu\phi\bar{q}\gamma^\mu q \qquad \alpha_{\phi q}^{(1)} = \frac{|\lambda_T|^2}{4M^2}$$

$$\begin{aligned} \alpha_1(m_t) &= \alpha_1(M) - \frac{N_C\lambda_t^2\alpha_{\phi q}^{(1)}(M)}{2\pi^2} \log\left(\frac{M}{m_t}\right) \\ &= \frac{N_C}{32\pi^2 M^2} \left[\lambda_t^2|\lambda_T|^2 - |\lambda_T|^4 - 2\lambda_t^2|\lambda_T|^2 \log\left(\frac{M^2}{m_t^2}\right) \right]. \end{aligned}$$

How to use EFTs (from the top-down) at one loop

- Sample result: T parameter from charge 2/3 vector-like quark singlet
 - Matching at m_t : top contribution with anomalous tree-level couplings

$$g_{W_3 t_L t_L} = g_{W_3 t_L t_L}^{\text{SM}} [1 - 2v^2(\alpha_{\phi q}^{(1)} - \alpha_{\phi q}^{(3)})] = g_{W_3 t_L t_L}^{\text{SM}} \left(1 - \frac{|\lambda_T|^2 v^2}{M^2}\right),$$
$$g_{W_1 t_L b_L} = g_{W_1 t_L b_L}^{\text{SM}} [1 + 2v^2 \alpha_{\phi q}^{(3)}] = g_{W_1 t_L b_L}^{\text{SM}} \left(1 - \frac{|\lambda_T|^2 v^2}{2M^2}\right).$$

$$\hat{T}(m_t^+) = \frac{N_C}{32\pi^2} \lambda_t^2 \left(1 - \frac{|\lambda_T|^2 v^2}{M^2}\right) = \hat{T}_{\text{SM}} + \Delta \hat{T}_{\text{SM}}$$

$$\Delta \hat{T}(m_t^-) = -v^2 \alpha_1(m_t) + \Delta \hat{T}(m_t^+) = \frac{N_C}{32\pi^2} \frac{v^2}{M^2} \left[|\lambda_T|^4 + 2\lambda_t^2 |\lambda_T|^2 \left(\log \frac{M^2}{m_t^2} - 1 \right) \right]$$

Summary

- The goal is to complete a UV/IR dictionary that maps arbitrary UV completions to experimental observables
- The tree-level, dimension-6 dictionary is (almost) finished
- The required automation for the one-loop dictionary is well advanced
- MatchMaker: General, fully automated and flexible code to compute tree-level and one-loop matching conditions

