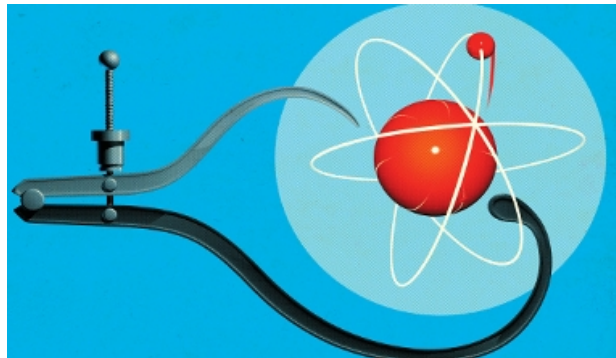


# The hadronic corrections to the Lamb shift in muonic hydrogen from ChPT and the proton radius



arxiv: 1403.3408

arxiv: 1406.4524

arxiv: 15xx.xxxx

work in collaboration with Antonio Pineda

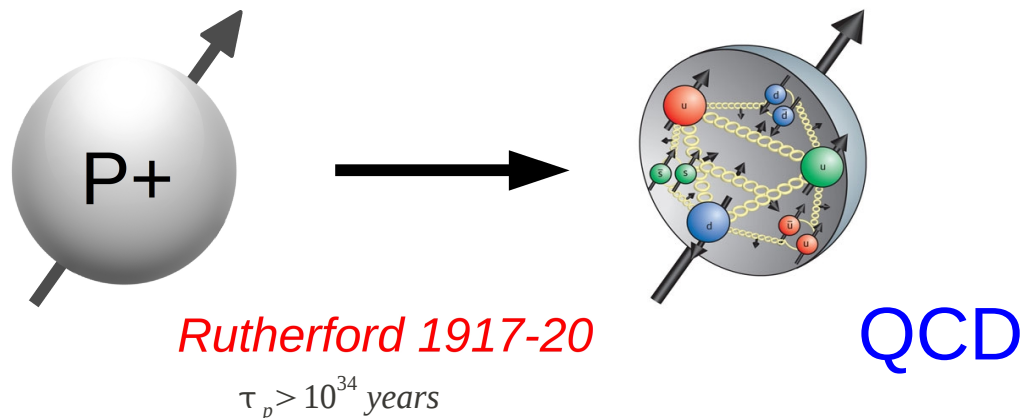
1st Hadron Spanish Network Days and Spanish-Japanese JSPS Workshop,  
IFIC, Valencia, 16th June 2015

**Clara Peset**

# Outline

- Model independent prediction of the proton radius
- Effective field theory framework: pNRQED
- Contributions to the Lamb shift
  - QED- like contributions
  - hadronic contributions (TPE)
- Summary

## Definition of the proton radius



$$r_p^2 = 6 \frac{dG_{p,E}(q^2)}{dq^2} \Big|_{q^2=0}$$

IR divergent!

NRQCD Lagrangian:

$$\delta \mathcal{L} = -e \frac{c_D^{(p)}}{m_p^2} N_p^\dagger \nabla \cdot \mathbf{E} N_p$$

Caswell, Lepage '86

$$\left[ \begin{array}{l} G_E(q^2) = F_1(q^2) + \frac{q^2}{4m_p^2} F_2(q^2) \\ c_D = 1 + 2F_2 + 8F_1' \end{array} \right.$$

$$r_{p,\overline{MS}}^2(\nu) = \frac{3}{4m_p^2} \left( c_D^{\overline{MS}}(\nu) - 1 \right)$$

General expression:

$$c_D^{(p)}(\nu) - c_{D,point-like}^{(p)}(\nu) \equiv \frac{4}{3} m_p^2 r_p^2$$

# The proton radius puzzle

Value from  $\mu$ -H:

$$r_p = 0.84087(39) \text{ fm (Science '13)}$$

7.1  $\sigma$

away!

CODATA value  
(2012):

$$r_p = 0.8775(51) \text{ fm}$$

from Hydrogen spectroscopy and electron-proton scattering

Probability of the lepton being  
*within* the volume of the proton

$$\left(\frac{r_p}{a_B}\right)^3 = (\alpha m_r r_p)^3$$

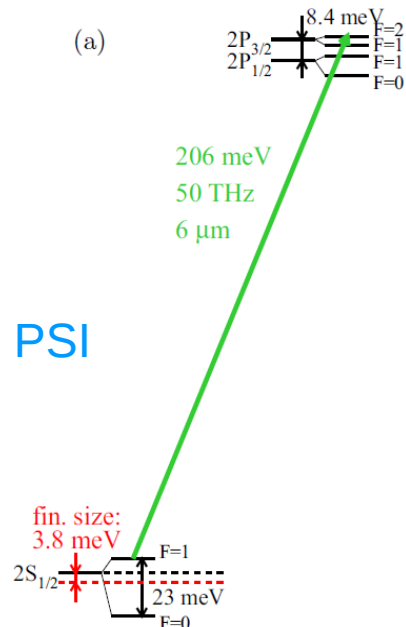
$$m_\mu \approx 200 m_e$$

8 million times  
larger for  $\mu$ H

But also larger UNKNOWN hadronic effects  $\rightarrow$  Two Photon Exchange

$$\text{LAMB SHIFT: } 2S_{1/2} \rightarrow 2P_{3/2}$$

displacement of nS energy levels



# EFT for bound states ( $\mu p$ )

Scales in bound states

Muonic hydrogen  
(Coulomb interaction)

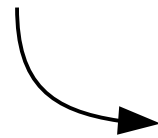
- Hard scale:  $m_r \longrightarrow m_r$
- Soft scale:  $|\mathbf{p}| \longrightarrow m_r \alpha$
- Ultrasoft scale:  $E \longrightarrow m_r \alpha^2$

Well separated  
scales!

pNRQED

Scales in  $\mu H$ :

$$m_p \sim m_\rho, \quad m_\mu \sim m_\pi \sim m_r, \quad m_r \alpha \sim m_e$$



Small expansion parameters:

$$\frac{m_\pi}{m_p} \sim \frac{m_\mu}{m_p} \approx \frac{1}{9}, \quad \frac{m_e}{m_r} \sim \frac{m_r \alpha}{m_r} \sim \alpha \approx \frac{1}{137}$$

Energy levels:

$$E(\mu p) = \frac{-m_r \alpha^2}{2n^2} (1 + c_1 \alpha + c_2 \alpha^2 + c_3 \alpha^3 \dots)$$

$$c_1 \sim c_1 \left( \frac{m_\mu \alpha}{m_e} \right) \text{ pure QED, and } c_n \sim \sum_{j=0}^{\infty} c_n^{(j)} \left( \frac{m_\pi}{m_p} \right)^j; \quad c_n^{(j)} \sim c_n^{(j)} \left( \frac{m_r}{m_\mu}, \frac{m_\mu}{m_\pi} \dots \right)$$

# pNRQED

pNRQED is a theory of ultrasoft photons

$$\text{HBET} \xrightarrow{m_\pi, \Delta} \text{NRQED} \longrightarrow \text{pNRQED}$$

The pNRQED lagrangian:

$$\mathcal{L}_{pNRQED} = \int d^3\mathbf{x} d^3\mathbf{X} S^\dagger(\mathbf{x}, \mathbf{X}, t) \left\{ i\partial_0 - \frac{\mathbf{p}^2}{m_r} + \frac{\mathbf{p}^4}{m_\mu^3} + \frac{\mathbf{p}^4}{m_p^3} - \frac{\mathbf{P}^2}{2M} \right. \\ \left. - V(\mathbf{x}, \mathbf{p}, \sigma_1, \sigma_2) + e\mathbf{x} \cdot \mathbf{E}(\mathbf{X}, t) \right\} S(\mathbf{x}, \mathbf{X}, t) - \int d^3\mathbf{x} \frac{d_1}{4} F_{\mu\nu} F^{\mu\nu}$$

The potential:

$$V = V^{(0)} + V^{(1)} + V^{(2)} + V^{(3)} + \dots$$

$$V^{(i)} \propto \left(\frac{1}{m}\right)^i + \text{expansions in the small parameters}$$

# Matching coefficients

Hadronic contributions:  $d_2^{had}$   $c_D^{had}$   $c_3^{had}$

EXAMPLE:

HBET:

$$\delta\mathcal{L}_{(N,\Delta)l_\mu} = \frac{1}{m_p^2} c_{3,R}^{pl_\mu} \bar{N}_p \gamma^0 N_p \bar{l}_\mu \gamma_0 l_\mu$$

NRQED:

$$\mathcal{L}_{Nl_\mu} = \frac{1}{m_p^2} c_{3,NR}^{pl_\mu} N_p^\dagger N_p l_\mu^\dagger l_\mu$$

pNRQED:

$$V^{(2)}(r) = c_{3,NR}^{pl_\mu} \frac{\delta^{(3)}(\mathbf{r})}{m_p m_\mu}$$

$$c_{3,NR}^{pl_\mu} = c_{3,R}^{pl_\mu} + c_{3,point-like}^{pl_\mu} + \boxed{c_3^{had}}$$

Theoretical equation for the Lamb Shift:

$$\Delta E_L = 206.0243(30) - 5.2271(7) \frac{r_p^2}{\text{fm}^2} + 0.0455(125) + \mathcal{O}(m_\mu \alpha^5 \frac{m_\mu^3}{m_\rho^3}, m_\mu \alpha^6) \text{meV}$$

CP, A. Pineda



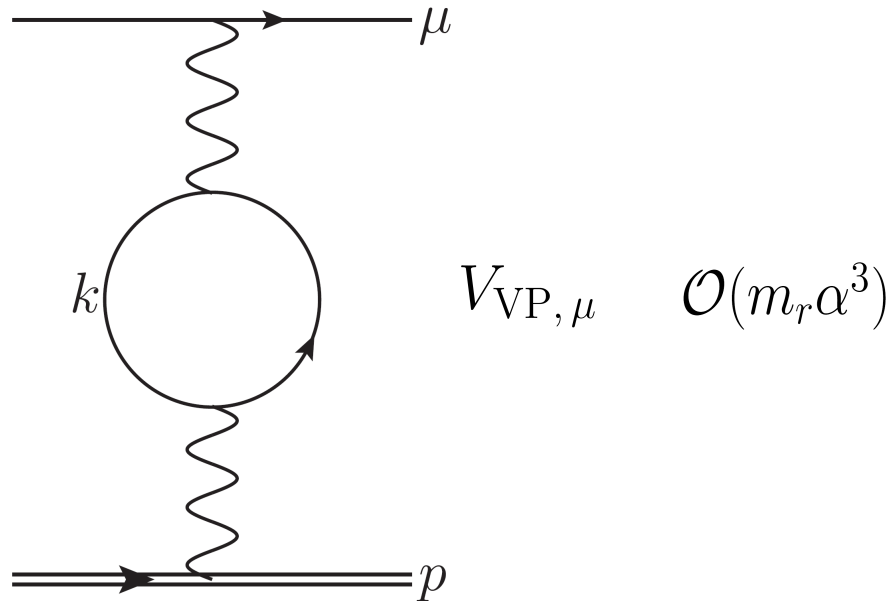
# “Pure” QED corrections

Theoretical equation for the Lamb Shift:

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CP, A. Pineda

$\mu$ H QED leading contribution: ELECTRON VACUUM POLARIZATION



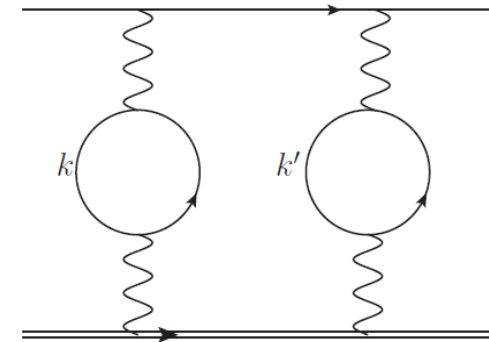
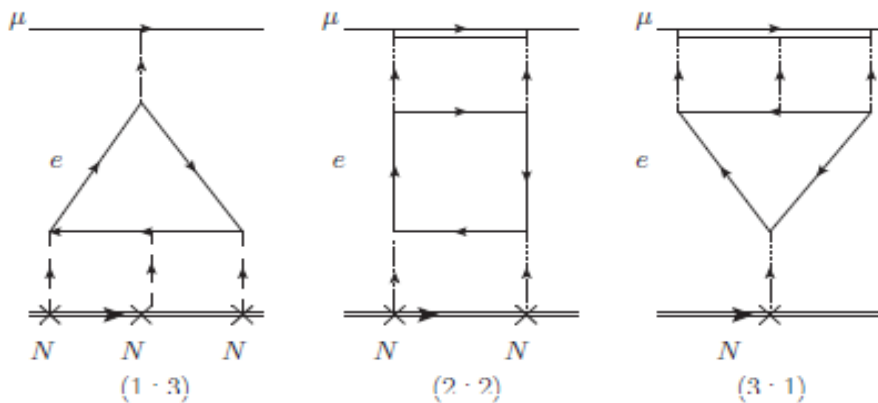
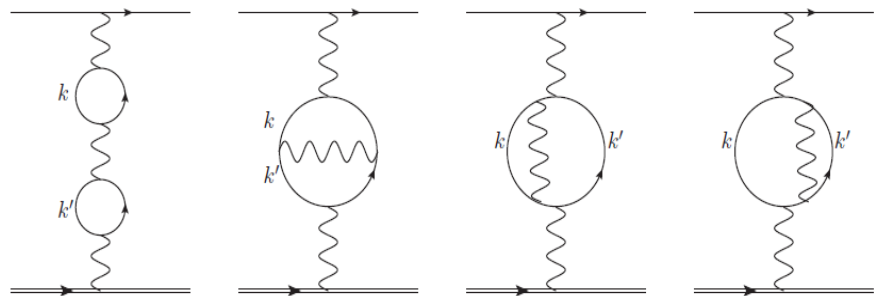
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CP, A. Pineda

Other contributions up to  $\mathcal{O}(m_r \alpha^5)$



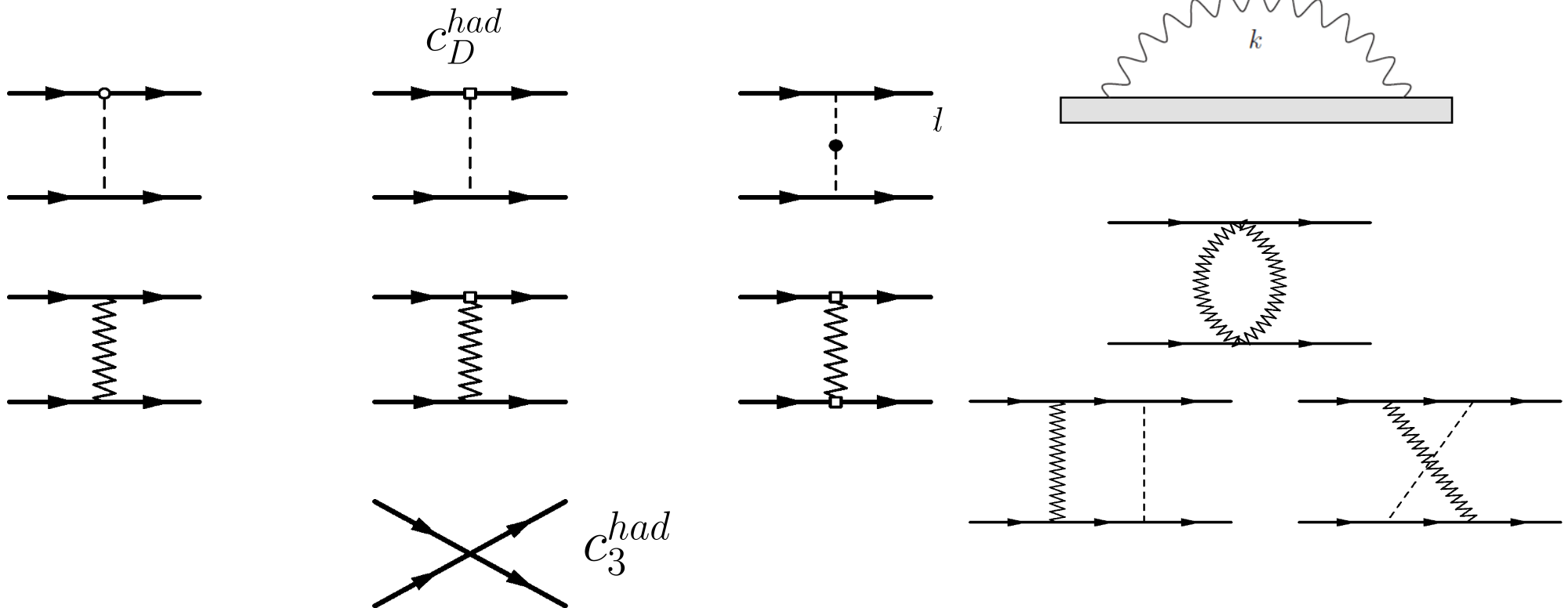
# “Pure” QED corrections

Theoretical equation for the Lamb Shift:

$$\Delta E_L = 206.0243(30) - 5.2271(7) \frac{r_p^2}{\text{fm}^2} + 0.0455(125) + \mathcal{O}(m_\mu \alpha^5 \frac{m_\mu^3}{m_\rho^3}, m_\mu \alpha^6) \text{meV}$$

CP, A. Pineda

Relativistic contributions up to  $\mathcal{O}(m_r \alpha^5)$



# Summary and error estimates

Theoretical equation for the Lamb Shift:

$$\Delta E_L = 206.0243(30) - 5.2271(7) \frac{r_p^2}{\text{fm}^2} + 0.0455(125) + \mathcal{O}(m_\mu \alpha^5 \frac{m_\mu^3}{m_\rho^3}, m_\mu \alpha^6) \text{meV}$$

CP, A. Pineda

$\mathcal{O}(m_r \alpha^3)$	$V_{\text{VP}}^{(0)}$	205.00745
$\mathcal{O}(m_r \alpha^4)$	$V_{\text{VP}}^{(0)}$	1.50795
$\mathcal{O}(m_r \alpha^4)$	$V_{\text{VP}}^{(0)}$	0.15090
$\mathcal{O}(m_r \alpha^5)$	$V_{\text{VP}}^{(0)}$	0.00752
$\mathcal{O}(m_r \alpha^5)$	$V_{\text{LbL}}^{(0)}$	-0.00089
$\mathcal{O}(m_r \alpha^4 \times \frac{m_\mu^2}{m_p^2})$	$V^{(2)} + V^{(3)}$	0.05747
$\mathcal{O}(m_r \alpha^5)$	$V_{\text{soft}}^{(2)}/\text{ultrasoft}$	-0.71903
$\mathcal{O}(m_r \alpha^5)$	$V_{\text{VP}}^{(2)}$	0.01876
$\mathcal{O}(m_\mu \alpha^6 \times \ln(\frac{m_\mu}{m_e}))$	$V^{(2)}; c_D^{(\mu)}$	-0.00127
$\mathcal{O}(m_\mu \alpha^6 \times \ln \alpha)$	$V_{\text{VP}}^{(2)}; c_D^{(\mu)}$	-0.00454

# Hadronic Effects

Theoretical equation for the Lamb Shift:

$$\Delta E_L = 206.0243(30) - 5.2271(7) \frac{r_p^2}{\text{fm}^2} + 0.0455(125) + \mathcal{O}(m_\mu \alpha^5 \frac{m_\mu^3}{m_\rho^3}, m_\mu \alpha^6) \text{meV}$$

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$\mathcal{O}(m_r \alpha^4 \times m_r^2 r_p^2)$	$V^{(2)}; c_D^{(p)}; r_p^2$	$-5.1975 \frac{r_p^2}{\text{fm}^2}$
$\mathcal{O}(m_r \alpha^5 \times m_r^2 r_p^2)$	$V_{\text{VP}}^{(2)}; c_D^{(p)}; r_p^2$	$-0.0283 \frac{r_p^2}{\text{fm}^2}$
$\mathcal{O}(m_r \alpha^6 \ln \alpha \times m_r^2 r_p^2)$	$V^{(2)}; c_D^{(p)}; r_p^2$	$-0.0014 \frac{r_p^2}{\text{fm}^2}$
$\mathcal{O}(m_r \alpha^5 \times \frac{m_r^2}{m_\rho^2})$	$V_{\text{VP}_{\text{had}}}^{(2)}; d_2^{\text{had}}$	$0.0111(2)$
$\mathcal{O}(m_r \alpha^5 \times \frac{m_r^2}{m_\rho^2} \frac{m_\mu}{m_\pi})$	$V^{(2)}; c_3^{\text{had}}$	$0.0344(125)$

$$\delta V^{\text{had}} = \frac{D_d^{\text{had}}}{m_p^2} \delta^{(3)}(\mathbf{r})$$

$$D_d^{\text{had}} \equiv -c_3^{\text{had}} - 16\pi\alpha d_2^{\text{had}} + \frac{\pi\alpha}{2} c_D^{\text{had}}$$

Definition of the proton radius

# Hadronic Effects

Theoretical equation for the Lamb Shift:

$$\Delta E_L = 206.0243(30) - 5.2271(7) \frac{r_p^2}{\text{fm}^2} + 0.0455(125) + \mathcal{O}(m_\mu \alpha^5 \frac{m_\mu^3}{m_\rho^3}, m_\mu \alpha^6) \text{meV}$$

CP, A. Pineda

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F. Jegerlehner

$$\delta V^{\text{had}} = \frac{D_d^{\text{had}}}{m_p^2} \delta^{(3)}(\mathbf{r})$$

$$D_d^{\text{had}} \equiv -c_3^{\text{had}} - 16\pi\alpha d_2^{\text{had}} + \frac{\pi\alpha}{2} c_D^{\text{had}}$$

Hadronic Vacuum polarization:  
Obtained from DR

# Hadronic Effects

Theoretical equation for the Lamb Shift:

$$\Delta E_L = 206.0243(30) - 5.2271(7) \frac{r_p^2}{\text{fm}^2} + 0.0455(125) + \mathcal{O}(m_\mu \alpha^5 \frac{m_\mu^3}{m_\rho^3}, m_\mu \alpha^6) \text{meV}$$

CP, A. Pineda

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$$\delta V^{\text{had}} = \frac{D_d^{\text{had}}}{m_p^2} \delta^{(3)}(\mathbf{r})$$

$$D_d^{\text{had}} \equiv -c_3^{\text{had}} - 16\pi\alpha d_2^{\text{had}} + \frac{\pi\alpha}{2} c_D^{\text{had}}$$

$$c_3^{\text{had}} = c_3^{\text{Born}} + c_3^{\text{pol}} \sim \alpha^2 \frac{m_\mu}{m_\pi} \left(1 + \# \frac{m_\pi}{\Delta} + \dots\right) + \mathcal{O}\left(\alpha^2 \frac{m_\mu}{\Lambda_{\text{QCD}}}\right)$$

Large-N<sub>c</sub> limit: we expect a large contribution from  $\Delta(1232)$

# $C_3^{\text{had}}$ : Born/Zemach contribution

$$c_{3,\text{Born}}^{pl_i} = 4(4\pi\alpha)^2 M_p^2 m_{l_i} \int \frac{d^{D-1}q}{(2\pi)^{D-1}} \frac{1}{q^6} G_E^{(0)} G_E^{(2)}(-\mathbf{q}^2) = \frac{\pi}{3} \alpha^2 M_p m_\mu \langle r^3 \rangle_{(2)} \quad \text{Zemach third momentum}$$

Zemach momenta:

		$\langle r^3 \rangle$	$\langle r^4 \rangle$	$\langle r^5 \rangle$	$\langle r^6 \rangle$	$\langle r^7 \rangle$	$\langle r^3 \rangle_{(2)}$
EFT: CP, Pineda  FITS:	$\pi$	0.4980	0.6877	1.619	5.203	20.92	0.9960
	$\pi \& \Delta$	0.4071	0.6228	1.522	4.978	20.22	0.8142
	Dipole	0.7706	1.083	1.775	3.325	7.006	2.023
	Kelly	0.9838	1.621	3.209	7.440	19.69	2.526
	Distler	1.16(4)	2.59(19)(04)	8.0(1.2)(1.0)	29.8(7.6)(12.6)	— — —	2.85(8)
	et al						

large dependence on the fitted function & large difference with EFT

Born energy shift:

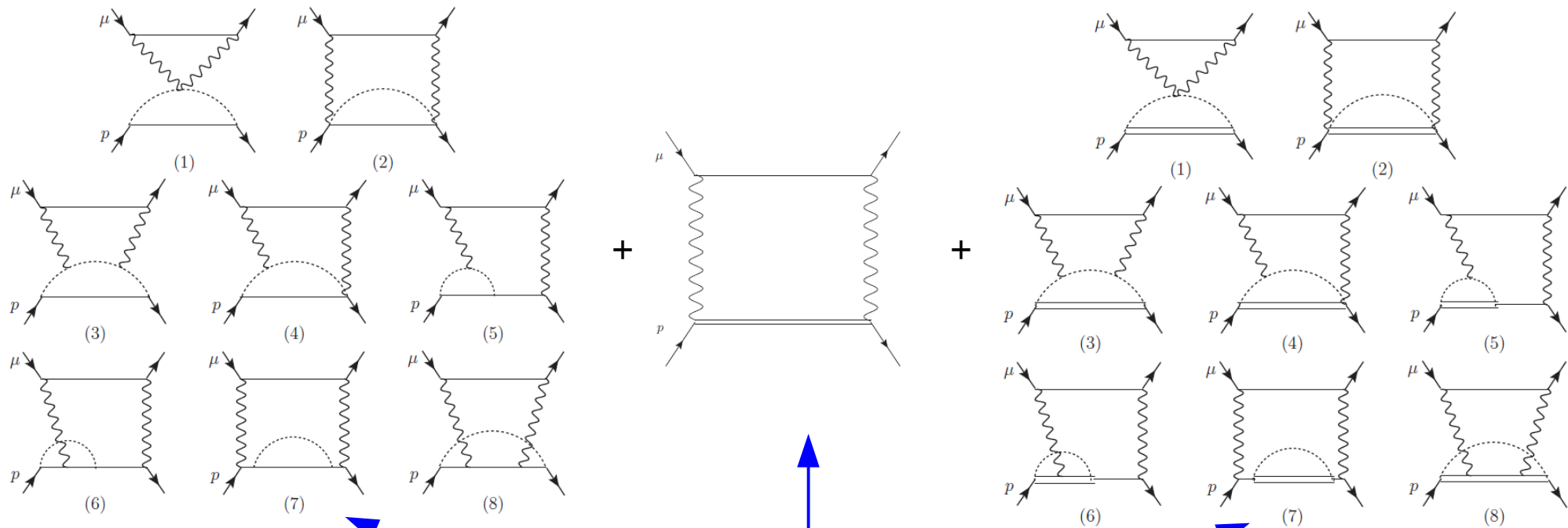
		CP, Pineda	
$\mu\text{eV}$	DR	Pachucki	Carlson et al
$\Delta E_{\text{Born}}$		23.2(1.0)	24.7(1.6)
	HBET	( $\pi$ )	( $\pi \& \Delta$ )
		10.1(5.1)	8.3(4.3)

we would expect less difference with the DR analysis



# $C_3^{\text{had}}$ : Polarizability

$$T_S^{\mu\nu} = \left( -g^{\mu\nu} + \frac{q^\mu q^\nu}{q^2} \right) S_1(\rho, q^2) + \frac{1}{M_p^2} \left( p^\mu - \frac{M_p \rho}{q^2} q^\mu \right) \left( p^\nu - \frac{M_p \rho}{q^2} q^\nu \right) S_2(\rho, q^2)$$



$$\Delta E_{\text{pol}} = \frac{C_{3,\text{pol}}^{pl_\mu}}{M_p^2} \frac{1}{\pi} \left( \frac{m_r \alpha}{2} \right)^3 = 18.51(\pi\text{-loop}) - 1.58(\Delta\text{-tree}) + 9.25(\pi\Delta\text{-loop}) = 26.2(10.0)\mu\text{eV}$$

# $c_3^{\text{had}}$ : Polarizability

Different results for the polarizability contribution:

( $\mu\text{eV}$ )	DR + Model				B $\chi$ PT ( $\pi$ )	HBET ( $\pi$ )	( $\pi$ & $\Delta$ )
$\Delta E_{\text{pol}}$	12(2)	11.5	7.4(2.4)	15.3(5.6)	8.2( $^{+1.2}_{-2.5}$ )	18.5(9.3)	26.2(10.0)
	Pachucki	Martynenko	Carlson et al	Gorchtein et al		Alarcon et al	CP, Pineda

corrections to HBET are suppressed by  $\frac{m_\mu}{m_\rho}$

The polarizability contribution from EFT is larger than the one computed using combinations of DR and different models

Total TPE energy shift:

$$\Delta E_{\text{TPE}} = \Delta E_{\text{Born}} + \Delta E_{\text{pol}} = 28.59(\pi) + 5.86(\pi \& \Delta) = 34.4(12.5)\mu\text{eV}$$

(LO)

(NLO)

CP, Pineda

The total contribution of TPE agrees with DR results better than when arbitrarily split into Born and polarizability pieces

# Summary

**We can theoretically predict the Lamb shift in a model independent way in an EFT framework**

$$\begin{array}{ll} \text{from } \mu\text{H:} & r_p = 0.8412(15) \text{ fm} \\ \text{CODATA value} & \\ \text{(from H):} & r_p = 0.8775(51) \text{ fm} \end{array} \left. \begin{array}{l} \text{CP, A. Pineda} \\ \end{array} \right\} 6.8\sigma \text{ away!}$$

- The hadronic contributions are the main source of uncertainty, although it is not enough to account for the discrepancy with ep
- The main radius-independent hadronic contribution is the **Two Photon Exchange**
  - The EFT approach gives Born & polarizability contributions which are quite different from the ones obtained by DR plus different models
  - Both the agreements and the disagreements should be further understood.
- In conclusion, **the proton radius puzzle survives the EFT analysis**





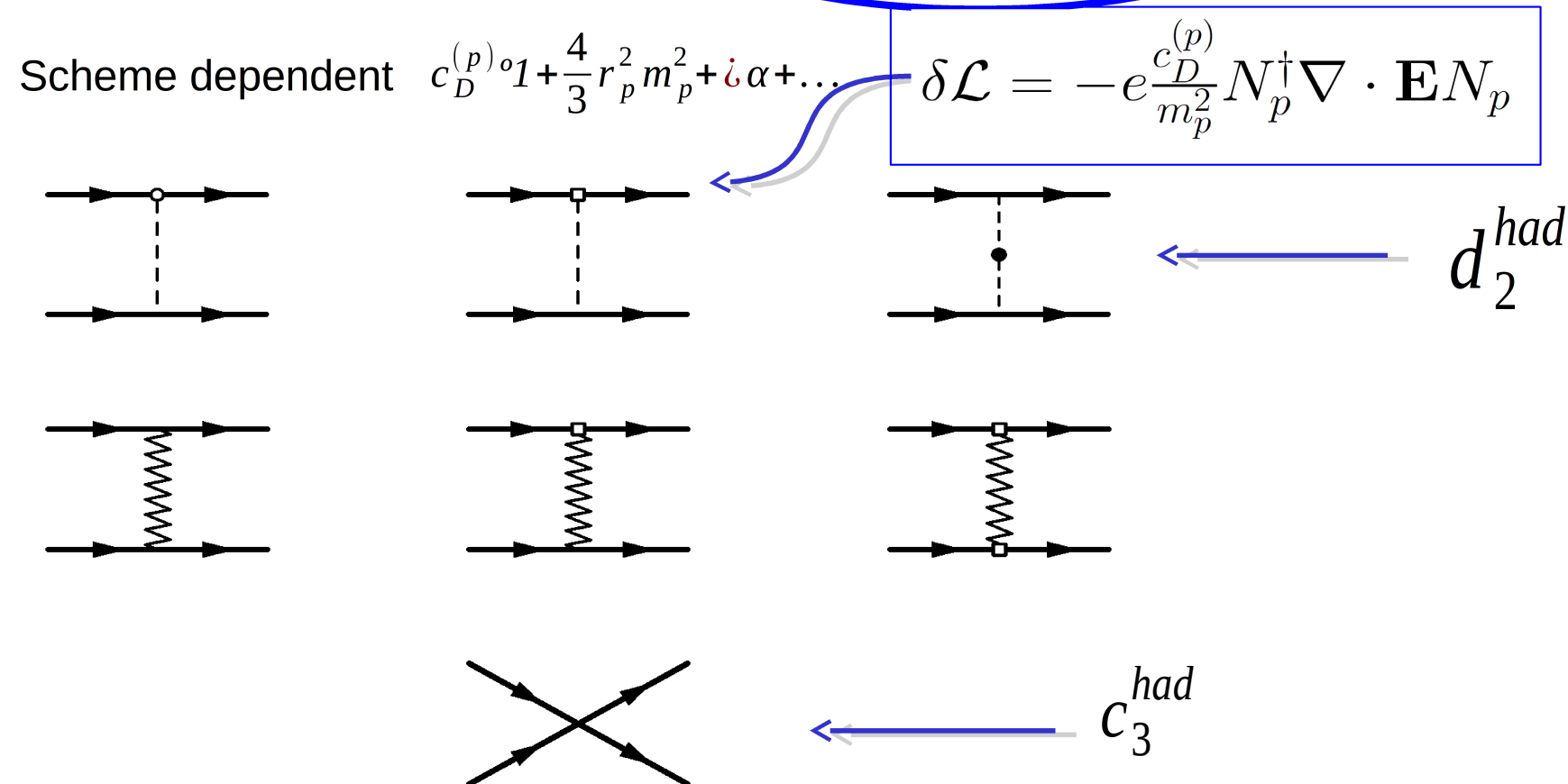
**THANK YOU!**

# Hadronic corrections

Theoretical equation for the Lamb Shift:

$$\Delta E_L = 206.0243(30) - 5.2271(7) \frac{r_p^2}{\text{fm}^2} + 0.0455(125) + \mathcal{O}(m_\mu \alpha^5 \frac{m_\mu^3}{m_\rho^3}, m_\mu \alpha^6) \text{meV}$$

CP, Pineda



# Zemach momenta

Even momenta:

$$G_E(-\mathbf{k}^2) = \sum_{n=0}^{\infty} \frac{(-1)^n}{(2n+1)!} \mathbf{k}^{2n} \int_0^{\infty} dr (4\pi) r^{2n} \rho_e(r) = \sum_{n=0}^{\infty} \frac{(-1)^n}{(2n+1)!} \mathbf{k}^{2n} \langle r^{2n} \rangle$$

Odd momenta:

$$\langle r^{2k+1} \rangle = \frac{\pi^{3/2} \Gamma[2+k]}{\Gamma[-1/2-k]} 2^{4+2k} \int \frac{d^3 q}{(2\pi)^3} \frac{1}{\mathbf{q}^{2(2+k)}} \left[ G_E(-\mathbf{q}^2) - \sum_{n=0}^k \frac{\mathbf{q}^{2n}}{n!} \left( \frac{d}{d\mathbf{q}^2} \right)^n G_E(-\mathbf{q}^2) \Big|_{\mathbf{q}^2=0} \right]$$

# Polarizability energy shift

$$\begin{aligned}
 c_3^{pl_i} = & -e^4 M_p m_{l_i} \int \frac{d^4 k_E}{(2\pi)^4} \frac{1}{k_E^4} \frac{1}{k_E^4 + 4m_{l_i}^2 k_{0,E}^2} \left\{ (3k_{0,E}^2 + \mathbf{k}^2) S_1(ik_{0,E}, -k_E^2) - \mathbf{k}^2 S_2(ik_{0,E}, -k_E^2) \right\} \\
 & + \mathcal{O}(\alpha^3) .
 \end{aligned}
 \tag{4.1}$$