

# Forward dispersion relations for $\pi k$ scattering.

A.Rodas Bilbao, J.R.Peláez

June 14, 2015

- 1 Motivation and Introduction
  - Forward dispersion relations (FDR)
- 2 Unconstrained Fits
  - FDR check
- 3 Constrained fits to data (CFD)
- 4 Conclusions

- $\pi K$  scattering appears as final state in many hadronic processes.
- Good description in many previous works with UChPT: **Oller,Oset (1999).Dobado,Peláez (1997). Oller,Oset,Peláez (1999). Jamin,Oller,Pich (2000). Nicola,Peláez (2002).**
- Roy-Steiner analysis: Bttiker,Descotes-Genon,Moussallam (2004).
- $\kappa(800)$  appears in these works and in other papers: **D.Bugg. Zheng,Zhou. M. Ablikim et al (BES). E.M. Aitala et al (E791). G. Bonvicini et al (CLEO). J.M. Link et al (FOCUS).**
- However  $\kappa(800)$  still needs confirmation according to PDG.
- We have been encouraged to perform a similar analysis for the  $\kappa(800)$  as done for the  $f_0(500)$  by our group.

# Introduction

- Simple fits with unitarity and analyticity. There is no dynamical input, pure data analysis.
- Check of the FDR (valid range of energy).
- Impose FDR to the fits.
- Important for other purposes (Roy-Steiner input, poles...)
- Data obtained from LASS experiments (Aston et al., Estabrooks et al.).
- First step in a long term project.

# Forward dispersion relations

- We form symmetric or antisymmetric amplitudes under  $s \leftrightarrow u$  exchange.

$$\begin{aligned} T^+ &= \frac{1}{3} T^{1/2} + \frac{2}{3} T^{3/2}, \\ T^- &= \frac{1}{3} T^{1/2} - \frac{1}{3} T^{3/2}. \end{aligned} \tag{1}$$

- $T^I$  is the amplitude of defined isospin  $I$ .

# Forward dispersion relations

- We will take for our analysis  $t = 0$ , they are called FDR.
- For the symmetric  $s \leftrightarrow u$  amplitude one subtraction is needed

$$\operatorname{Re}(T^+(s)) = T^+(s_{th}) + \frac{s(s - s_{th})}{\pi} P \int_{s_{th}}^{\infty} \frac{\operatorname{Im} T^+(s')(2s' - s_{th})}{s'(s' - s)(s' - s_{th})(s' + s - s_{th})} ds'. \quad (2)$$

- For the antisymmetric amplitude no subtraction is needed

$$\operatorname{Re}(T^-(s)) = \frac{(2s - s_{th})}{\pi} P \int_{s_{th}}^{\infty} \frac{\operatorname{Im} T^-(s')}{(s' - s)(s' + s - s_{th})} ds'. \quad (3)$$

- We use the unitary functional form for the partial waves

$$t_l^I(s) = \frac{1}{\sigma(s)} \frac{1}{\cot \delta_l^I(s) - i} \quad (4)$$

- Where

$$\cot \delta_l^I(s) = \frac{\sqrt{s}}{2q^{2l+1}} \sum B_n w(s)^n \quad (5)$$

- with  $w(s) = \frac{\sqrt{y(s)} - \alpha \sqrt{y(s_0) - y(s)}}{\sqrt{y(s)} + \alpha \sqrt{y(s_0) - y(s)}}$  as our new variable (conformal mapping).
- Here  $y(s) = (\frac{s-su}{s+su})^2$  defines the circular cut on the next figure.
- $w$  used to maximize the analyticity domain.

# Fits:Elastic region

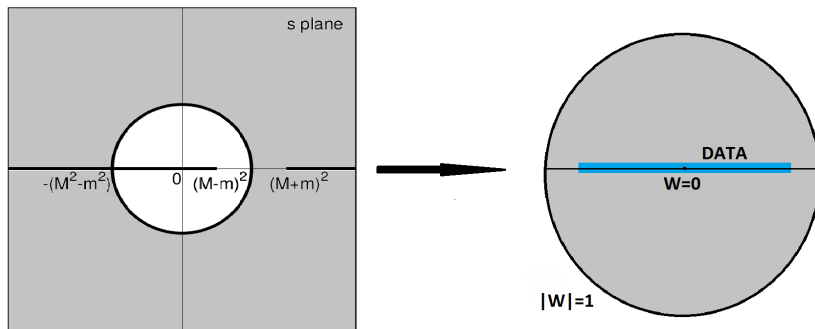


Figure: Structure of the PW.

- $\alpha$  is used to center the point of energy  $s_c$  for the expansion.



# Fits:Elastic region

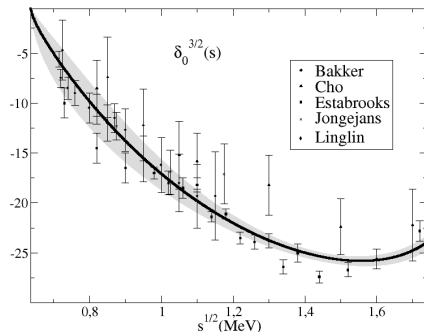
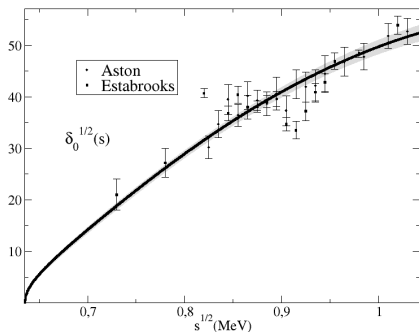


Figure:  $S^{1/2}$  and  $S^{3/2}$  phase shifts, the  $I=3/2$  is elastic in the entire region.

- $S^{3/2}$  is elastic in the whole region.

# Fits:Elastic region

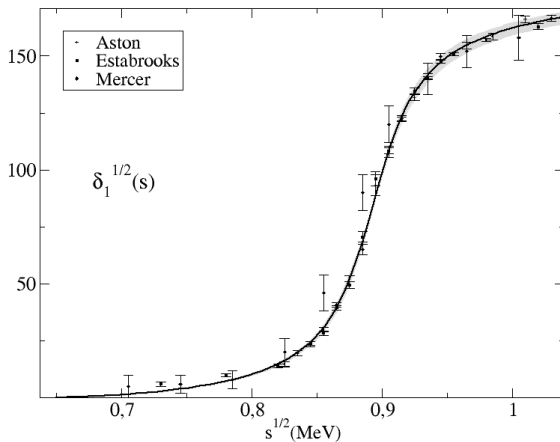


Figure:  $P^{1/2}$  phase shifts.

- Good description of the  $K^*(892)$ .

- In the inelastic region  $t_l^I = \frac{\eta_l^I(s)e^{2i\delta_l^I(s)} - 1}{2i} = |t_l^I|e^{i\phi_l^I}$ .
- We use complex rational functions that near each resonance look like BW.
- We impose matching conditions on the inelastic  $\eta k$  threshold.
- We neglect in this analysis the  $P^{3/2}$  and  $D^{3/2}$  PW, their phase shifts never reach 5 degrees.
- We use up to  $F^{1/2}$  which is very small and neglect  $G^{1/2}$  in the studied energy region.

# Fits: Inelastic region

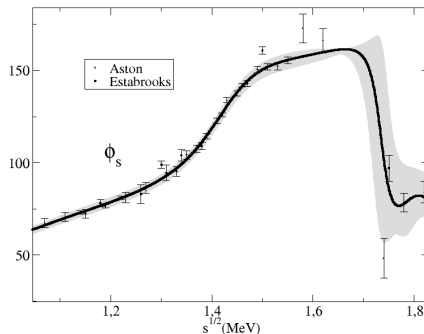
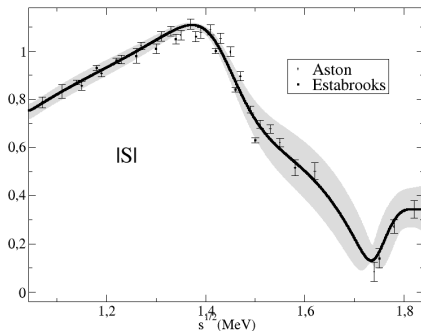


Figure: S-wave amplitude and total phase.

- Incompatibilities between Aston and Estabrooks sets of data.

# Fits: Inelastic region

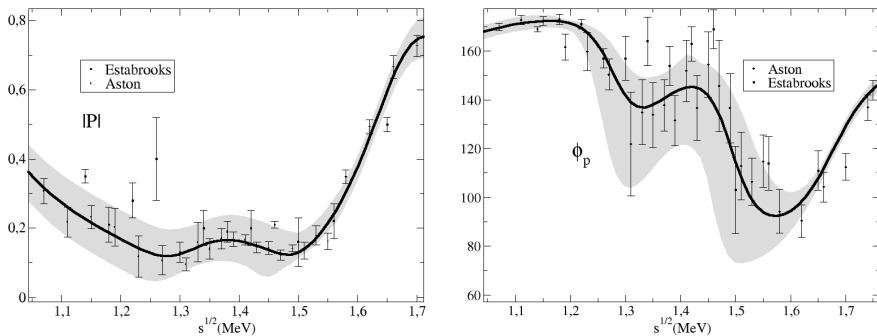


Figure:  $P^{1/2}$ -wave amplitude and total phase.

# Fits: Inelastic region

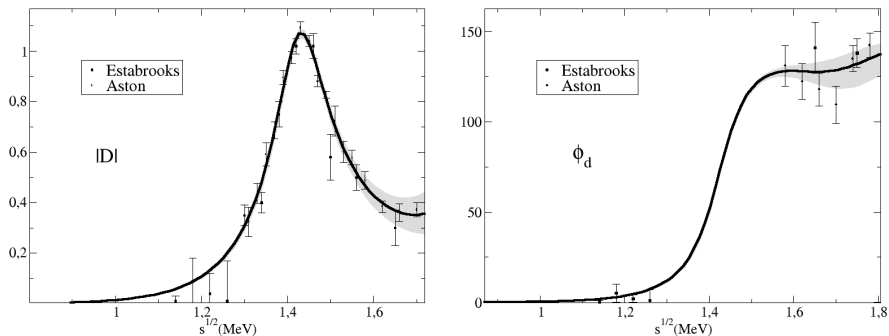
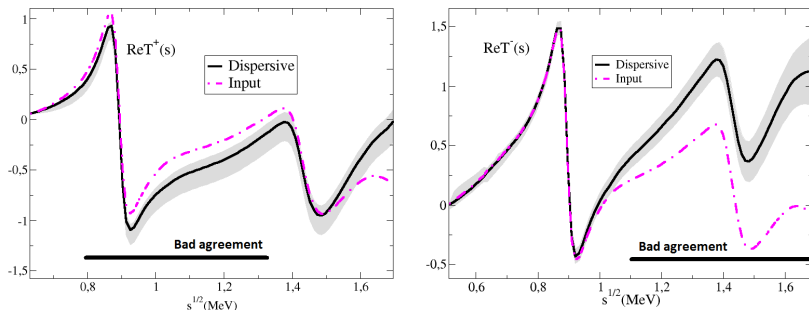


Figure:  $D^{1/2}$ -wave amplitude and total phase.

- No contribution below 1.2 – 1.4 GeV.
- Above 1.7 GeV we use Regge parametrizations. **Peláez, Yndurain (2004)**

# FDR check



**Figure:** FDR unconstrained, symmetric and antisymmetric. It is clear the huge difference between the input and output

- Symmetric incompatibilities caused by the  $P^{1/2}$  PW.
- Antisymmetric deviations due to Regge contribution.
- Room for improvement  $\rightarrow$  Constrained fits.
- Above 1.7 GeV the discrepancies are too huge to impose FDR.

# Constrained fits to data (CFD)

- We study the FDR up to  $1.6 - 1.7 \text{ GeV}$ .
- We define a  $\chi^2_1$  between the input and output.
- There is a  $\chi^2_2$  between the UFD parameters and the new ones.
- After the minimization of  $\chi^2_1 + \chi^2_2$  we obtain

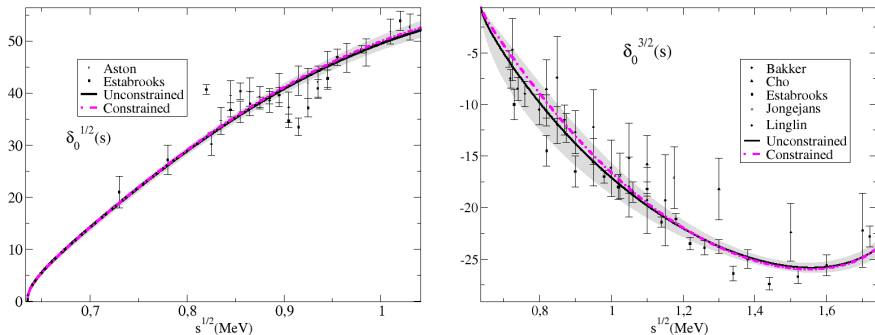


Figure:  $S^{1/2}$  and  $S^{3/2}$  phase shifts.



# Constrained fits to data (CFD)

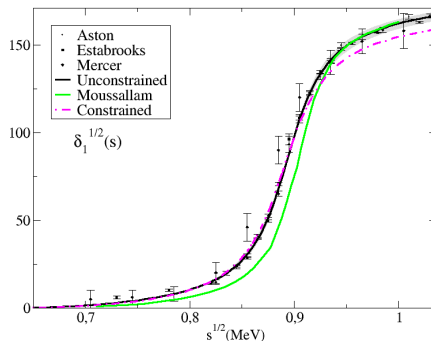


Figure:  $P^{1/2}$  phase shifts.

- If the  $K^*(892)$  is to be well described the phase shift must be lower at 1 GeV.
- Moussallam result  $\rightarrow$  solve Roy equations. They use the data at 0.935 GeV as the matching point.

# Constrained fits to data (CFD)

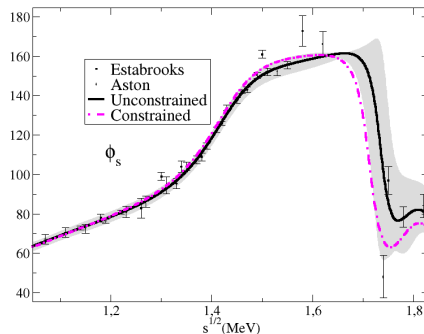
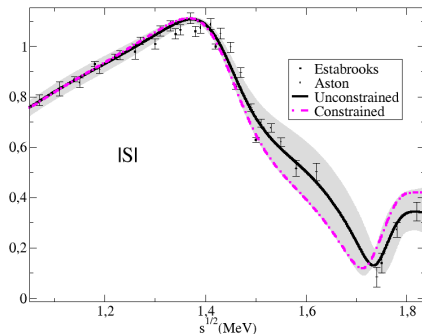


Figure: S-wave amplitude and total phase.

- Almost unchanged below 1.5 GeV.
- Constrained fit still compatible.

# Constrained fits to data (CFD)

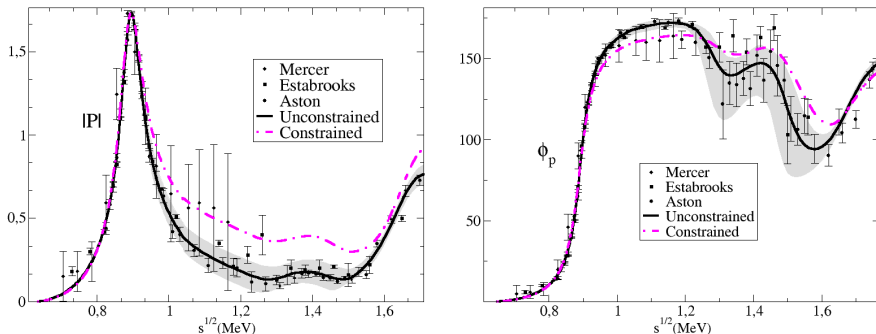


Figure:  $P^{1/2}$ -wave amplitude and total phase.

- We have tried to fit the data in the region 1 – 1.5 GeV but it spoils the  $T^+$  and also the  $K^*(892)$  after the minimization.
- The FDR demand a deviation from data in the P wave from 1 to 1.5 GeV.

# Constrained fits to data (CFD)

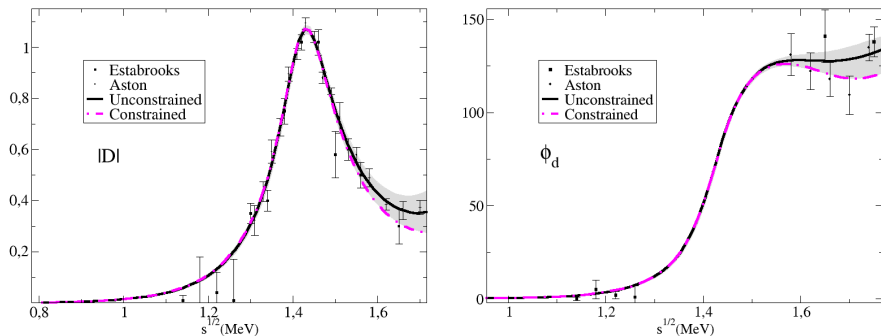


Figure:  $D^{1/2}$  wave amplitude and total phase.

- Little changes around 1.6 GeV.

# Constrained fits to data (CFD)

- The change in the symmetric amplitude around  $1 - 1.2\text{GeV}$  is caused by the change of the P-wave. The Regge contribution in this region is small.
- The huge change of the antisymmetric one is caused by the Regge  $\pi k$  factorization constant.

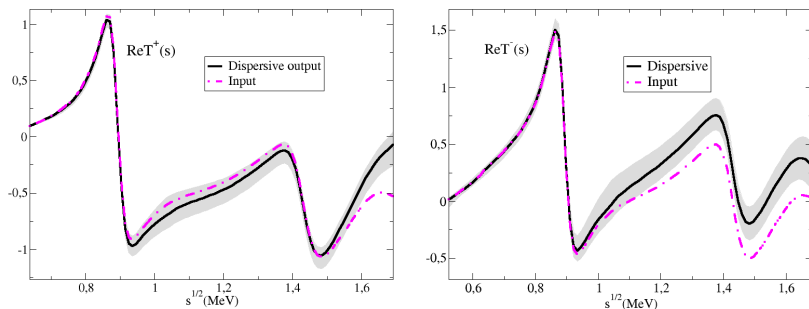
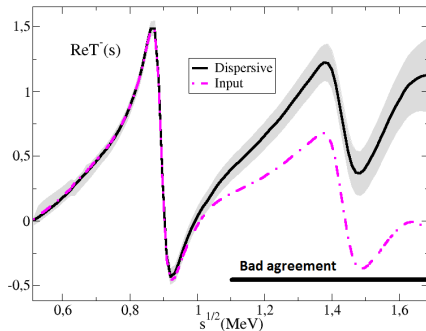
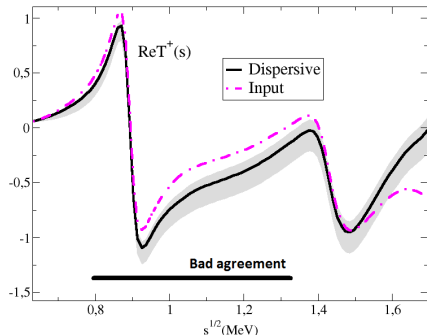


Figure: FDR constrained, symmetric and antisymmetric. They fairly compatible up to  $1.6\text{GeV}$

# Constrained fits to data (CFD)



**Figure:** FDR unconstrained, symmetric and antisymmetric. It is clear the huge difference between the input and output

# Scattering lengths

- Now we can obtain the threshold parameters for the most important partial waves

Table: Scattering lengths.

SL	UFD	CFD	Moussallam result
$m_\pi a_0^{1/2}$	0.217	$0.218 \pm 0.014$	$0.224 \pm 0.022$
$m_\pi a_0^{3/2}$	$-0.0790 \pm 0.03$	$-0.0614 \pm 0.025$	$-0.0448 \pm 0.0077$
$m_\pi^3 a_1^{1/2}$	$0.0259 \pm 0.008$	$0.0242 \pm 0.007$	$0.019 \pm 0.001$

- For the kappa resonance  $K_0(800)$  we obtain

Table:  $\kappa(800)$  parameters.

Group	Mass	Width
UFD	$680 \pm 19$	$677 \pm 24$
CFD	$681 \pm 19$	$674 \pm 23$
Moussallam et al.	$658 \pm 13$	$557 \pm 24$
D.Bugg	$663 \pm 34$	$658 \pm 44$
Zheng,Zhou	$694 \pm 53$	$606 \pm 59$

- The values of the masses are compatible. However we expect to obtain a lower value of the partial width in a future analysis with the help of  $\pi K$  Roy equations.



- We have studied FDR in  $\pi k$  up to  $1.7\text{GeV}$ .
- Imposing FDR yields significant changes in  $|P^{1/2}|$  above  $1\text{GeV}$ .
- Our solution still describes the data of the experiments in the energy region for the rest of the PW.
- $\kappa(800)$  mass compatible with PDG and Roy-Steiner equations but width comes wider.
- This is first step on our program to obtain fits consistent with Roy-Steiner equations and rigorous determination of  $\kappa(800)$ .

¡Thank you for your attention!

- Too huge differences in both  $T^+$  and  $T^-$ .
- We can't correct only the PW or the Regge, too much deviation from the old parameters in the PW.

# Regge parametrization

- The problem arises with the factorization constants  $f_{k/\pi}$  (Pomeron) and  $g_{k/\pi}$  (Rho).
- The problem in the antisymmetric amplitude is caused by  $g_{k/\pi}$ .
- The previous works in  $\pi\pi$  scattering together with the paper of the Regge parametrizations suggest a lower value (30 percent) of this constant.