

Universal geometrical scaling of the elliptic flow

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Outline

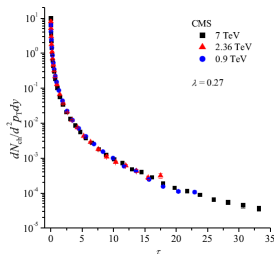
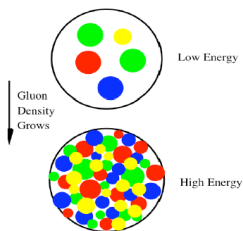
- 1 Motivation
- 2 The string percolation model
- 3 Geometrical scaling of the elliptic flow

Geometrical scaling in pp collisions. Color Glass Condensate.

- **CGC**: An effective theory of **QCD** at **high energies** and **high gluon occupation numbers**.
- Dynamical scale $Q_s \gg \Lambda_{QCD}$, called **saturation momentum**.
- This framework **predicts geometrical scaling in pp collisions**

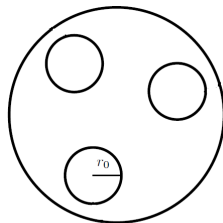
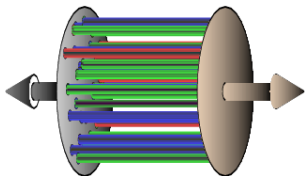
$$\frac{dN}{dydp_T^2} = \frac{1}{Q_0^2} F(\tau)$$

$$\tau = \frac{p_T^2}{Q_s^2} \quad Q_0^2 \approx 1 \text{ GeV}$$



GS in Heavy Ion Collisions?

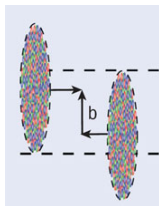
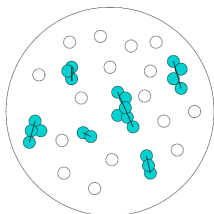
Physical picture



$$r_0 = 0.2 - 0.25 \text{ fm.}$$

- Projectile and target interact via **color field** created by the constituent **partons** of the nuclei.
- **Color field** is **confined** in a region with transverse size $r_0 \sim 0.2 \text{ fm}$.
- We can see them as **small areas** in transverse plane.
- These color “**strings**” break producing $q\bar{q}$ pairs (Schwinger mechanism) that subsequently lead to the observed hadrons.

Physical picture



- With growing **energy** and/or **mass number** of colliding particles, the number of **sources** grows → The **number of strings** grows with **energy** and/or **mass number**.
- The **number of strings** also increases with increasing **centrality**.
- The basic degree of freedom is the **density of strings**

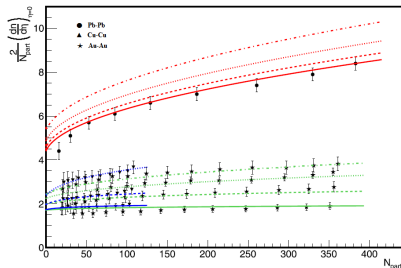
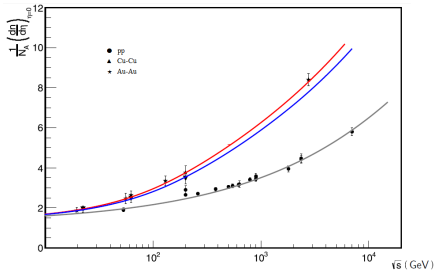
$$\rho = \frac{S_1}{S_A} N_s$$

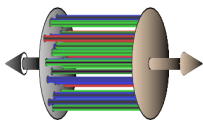
Multiplicity distributions

- In the framework of **percolation of strings**:

$$\rho_A = \rho_P A^{\alpha(s)} N_A^{1/3}$$

$$\alpha(s) = \frac{1}{3} \left(1 - \frac{1}{1 + \ln \left(\sqrt{s/s_0} + 1 \right)} \right) \quad N_A = N_{part}/2$$



Percolation of strings \leftrightarrow CGC

- Transverse size of a **string/flux tube**

Percolation

$$r_0^2 \sqrt{\frac{1 - e^{-\rho}}{\rho}}$$

CGC

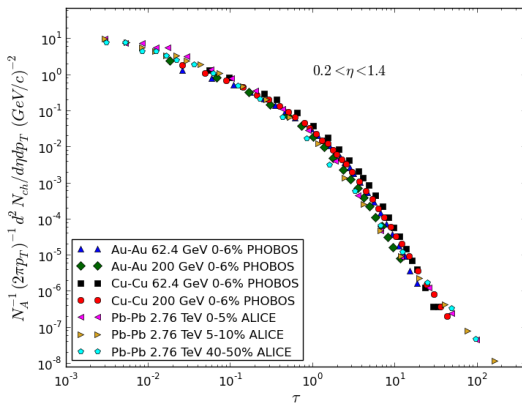
$$1/Q_s^2$$

- In the high density limit $\rightarrow \sqrt{\rho} \sim Q_s^2$.

$$\Rightarrow \left(Q_s^A \right)^2 = \left(Q_s^p \right)^2 A^{\alpha(s)/2} N_A^{1/6}$$

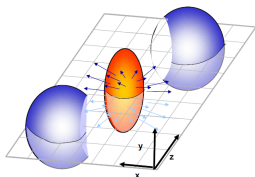
Geometrical scaling in heavy ion collisions

- Using the same approach as in pp collisions



Elliptic flow

- **Common interpretation of azimuthal anisotropy: Spatial anisotropy** \Leftrightarrow **momentum anisotropy** (but not the only one!)

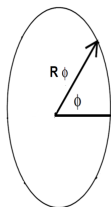


- **Fourier decomposition of momentum distribution** to measure this anisotropy

$$E \frac{d^3 N}{dp^3} = \frac{1}{2\pi} \frac{d^2 N}{p_T dp_T dy} \left(1 + \sum_{n=1}^{\infty} v_n(p_T, y) \cos(n\phi) \right)$$

- $v_2(p_T, y) \Rightarrow$ **Elliptic Flow**

Scaling of the elliptic flow



- Higher density for lower R_ϕ
- $\Rightarrow \phi$ -dependent density

$$\Rightarrow \left(Q_s^A \right)^2 \equiv \frac{L}{\lambda_{mfp}} \frac{1}{R_\phi^2} = \frac{Q_s^A L}{R_\phi^2}$$

- L is the length associated to the size of the collision area at a given impact parameter and energy $L \propto (1 + N_A^{1/3})/2$

Scaling of the elliptic flow

- Changing $(Q_s^A)^2 \leftrightarrow (Q_{s\phi}^A)^2$

$$v_2 = \frac{4 \int_0^{\pi/2} d\phi \cos 2\phi \frac{dN}{dp_T^2 d\phi}}{\frac{dN}{dp_T^2}} = \frac{4 \int_0^{\pi/2} d\phi \cos 2\phi F(\tau_\phi)}{F(\tau)}$$

- For non very peripheral collisions we can expand in powers of $R^2 - R_\phi^2$

$$v_2 = \frac{2}{\pi} \int_0^{\pi/2} d\phi \cos 2\phi \frac{R^2 - R_\phi^2}{R^2} \frac{4}{F(\tau)} \frac{dF}{d\tau} \tau Q_s^A L$$

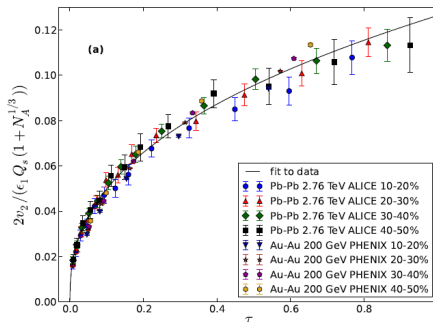
- Defining

$$\boxed{\frac{4}{F(\tau)} \frac{dF}{d\tau} = \varphi(\tau)}, \quad \boxed{\epsilon_1 = \frac{2}{\pi} \int_0^{\pi/2} d\phi \cos 2\phi \frac{R^2 - R_\phi^2}{R^2}}$$

Scaling of the elliptic flow in heavy ion collisions

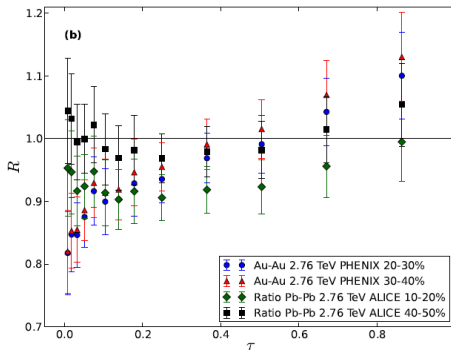
- We obtain the scaling law

$$\frac{v_2}{\epsilon_1 Q_s^A L} = \tau \varphi(\tau)$$



Scaling of the elliptic flow in heavy ion collisions

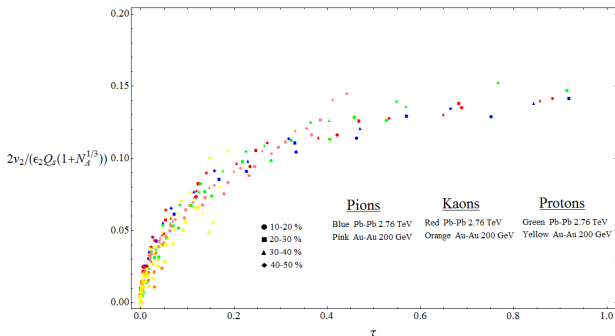
- Ratio over Pb-Pb 30-40% at 2.76 TeV



Scaling of the elliptic flow in heavy ion collisions: Identified hadrons

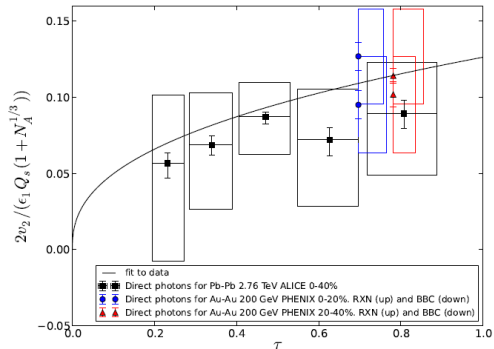
- In this framework, multiplicity distributions change:

$$p_T \rightarrow \sqrt{p_T^2 + m^2} - m$$



Scaling of the elliptic flow in heavy ion collisions: Photons

- Elliptic flow of photons is difficult to explain in final-state models because they interact weakly with QGP



Conclusions

- Elliptic flow of charged and identified particles satisfies a scaling law.
- Photon data are lying on the same scaling curve.
- Indication of possible initial-state effects in flow.
- Study of possible scaling of the rest of harmonics.