

BPS Skyrmions, a unified approach to nuclear matter: from nuclei to neutron stars

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VI CPAN Days

October 20th, 2014 - Sevilla, Spain



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- 1 Introduction
- 2 The BPS Skyrme Model
- 3 Coupling to gravity: Neutron Stars

[based on work in collaboration with C. Adam, J. Sanchez-Guillen, R. Vazquez and A. Wereszczynski]

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Skyrme Model

T.H.R. Skyrme: Low-energy field theory of QCD.

- Primary fields are mesons (pions).
- Baryons and nuclei are described by collective nonlinear excitations of the fundamental degrees of freedom.
- Topological charge = baryon number.

Simplest case (two flavors): target space = $SU(2)$, Skyrme field U :

$$x^\mu \rightarrow U(\mathbf{x}) : \quad \mathbb{R}^3 \times \mathbb{R} \rightarrow SU(2)$$

Original Skyrme Model

Lagrangian of the Standard Skyrme Model:

$$\mathcal{L} = \mathcal{L}_2 + \mathcal{L}_4$$

- Sigma model term:

$$\mathcal{L}_2 = -\frac{f_\pi^2}{4} \text{Tr} (U^\dagger \partial_\mu U U^\dagger \partial^\mu U)$$

Pion kinetic term

- Quartic Skyrme term:

$$\mathcal{L}_4 = \frac{1}{32e^2} \text{Tr} ([U^\dagger \partial_\mu U, U^\dagger \partial_\nu U]^2)$$

Sufficient to circumvent Derrick stability. No BPS saturated solutions.

Generalized Skyrme Model

Effective field theory \sim derivative expansion \Rightarrow higher powers of derivatives expected.

$$\mathcal{L} = \mathcal{L}_2 + \mathcal{L}_4 + \mathcal{L}_6 + \mathcal{L}_0$$

- Sextic Term:

$$\mathcal{L}_6 = -\lambda^2 \pi^4 \mathcal{B}_\mu^2$$

where

$$\mathcal{B}^\mu = \frac{1}{24\pi^2} \text{Tr}(\epsilon^{\mu\nu\rho\sigma} U^\dagger \partial_\nu U U^\dagger \partial_\rho U U^\dagger \partial_\sigma U).$$

(Square of the topological current) Quadratic in time derivatives
 \Rightarrow standard hamiltonian formulation

- Potential term

$$\mathcal{L}_0 = -\mu^2 U$$

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The model

Limit of generalized Skyrme Model:

$$\mathcal{L}_{06} = \mathcal{L}_6 + \mathcal{L}_0$$

- Parametrization for U

$$U = e^{i\xi\vec{n}\cdot\vec{\sigma}} = \cos \xi + i \sin \xi \vec{n} \cdot \vec{\sigma} \quad \vec{n}^2 = 1$$

and stereographic projection

$$\vec{n} = \frac{1}{1 + |u|^2} (u + \bar{u}, -i(u - \bar{u}), 1 - |u|^2)$$

$$\Rightarrow \mathcal{L}_{06} = \frac{\lambda^2 \sin^4 \xi}{(1 + |u|^2)^4} (\epsilon^{\mu\nu\rho\sigma} \xi_\nu u_\rho \bar{u}_\sigma)^2 - \mu^2 \mathcal{U}(\xi)$$

The model

- Symmetries: Area-preserving diffeomorphisms on target space S^2 spanned by u :

$$\xi \rightarrow \xi, \quad u \rightarrow \tilde{u}(u, \bar{u}, \xi)$$

$$(1 + |\tilde{u}|^2)^{-2} d\xi d\tilde{u} d\bar{\tilde{u}} = (1 + |u|^2)^{-2} d\xi u d\bar{u}$$

- BPS (Bogomolny) bound

$$E = \int d^3x \left(\frac{\lambda^2 \sin^4 \xi}{(1 + |u|^2)^4} (\epsilon^{mnl} i \xi_m u_n \bar{u}_l)^2 + \mu^2 \mathcal{U}(\xi) \right)$$

$$\geq 2\lambda\mu\pi^2 |B| \frac{1}{2\pi^2} \int_{S^3} d\Omega \sqrt{\mathcal{U}(\xi)} \equiv 2\lambda\mu\pi^2 \langle \sqrt{\mathcal{U}} \rangle_{S^3} |B|$$

- BPS equation:

$$\frac{\lambda \sin^2 \xi}{(1 + |u|^2)^2} \epsilon^{mnl} i \xi_m u_n \bar{u}_l = \mp \mu \sqrt{\mathcal{U}}.$$

Example: Skyrme potential

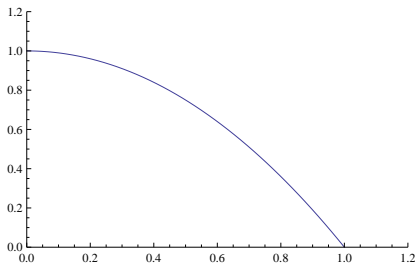
We assume the standard Skyrme potential

$$U = \frac{1}{2} \text{Tr}(1 - U) \rightarrow U(\xi) = 1 - \cos \xi$$

With the ansatz: $\xi = \xi(r)$, $u(\theta, \phi) = \tan \frac{\theta}{2} e^{in\phi}$

- Compact solution with radius proportional to $n^{1/3}$.
- Energy linear in the baryon number:

$$E = \frac{64\sqrt{2}\pi}{15} \mu\lambda |n|$$



C.Adam, J. Sanchez-Guillen, A. Wereszczynski, Phys. Lett. B **691**, 105 (2010)

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BPS Skyrme model coupled to gravity

In general, the corresponding action is

$$S_{06} = \int d^4x |g|^{\frac{1}{2}} (-\lambda^2 \pi^4 |g|^{-1} g_{\rho\sigma} B^\rho B^\sigma - \mu^2 \mathcal{U})$$

- Energy-momentum tensor (perfect fluid)

$$T^{\rho\sigma} = -2|g|^{-\frac{1}{2}} \frac{\delta}{\delta g_{\rho\sigma}} S_{06} = (p + \rho) u^\rho u^\sigma - p g^{\rho\sigma}$$

where 4-velocity

$$u^\rho = B^\rho / \sqrt{g_{\sigma\pi} B^\sigma B^\pi}$$

- Energy density and pressure

$$\rho = \lambda^2 \pi^4 |g|^{-1} g_{\rho\sigma} B^\rho B^\sigma + \mu^2 \mathcal{U}$$

$$p = \lambda^2 \pi^4 |g|^{-1} g_{\rho\sigma} B^\rho B^\sigma - \mu^2 \mathcal{U}$$

BPS Skyrme model coupled to gravity

- Axially symmetric ansatz for Skyrme field

$$\xi = \xi(r), \quad \vec{n} = (\sin \theta \cos B\phi, \sin \theta \sin B\phi, \cos \theta)$$

- Metric (Schwarzschild coordinates)

$$ds^2 = \mathbf{A}(r)dt^2 - \mathbf{B}(r)dr^2 - r^2(d\theta^2 + \sin^2 \theta d\phi^2)$$

- Compatible with Einstein eq. $G_{\rho\sigma} = \frac{\kappa^2}{2} T_{\rho\sigma} \Rightarrow$

$$\frac{1}{r} \frac{\mathbf{B}'}{\mathbf{B}} = -\frac{1}{r^2}(\mathbf{B} - 1) + \frac{\kappa^2}{2} \mathbf{B}\rho$$

$$r(\mathbf{B}\rho)' = \frac{1}{2}(1 - \mathbf{B})\mathbf{B}(\rho + 3p) + \frac{\kappa^2}{2}\mu^2 r^2 \mathbf{B}^2 \mathcal{U}(h)\rho$$

$$\frac{\mathbf{A}'}{\mathbf{A}} = \frac{1}{r}(\mathbf{B} - 1) + \frac{\kappa^2}{2} r\mathbf{B}\rho$$

where $(h \equiv (1/2)(1 - \cos \xi))$

$$\rho = \frac{4\mathbf{B}^2\lambda^2}{\mathbf{B}r^4} h(1 - h)h'^2 + \mu^2 \mathcal{U}(h), \quad p = \rho - 2\mu^2 \mathcal{U}(h)$$

How to proceed with the calculations

- Fit λ, μ to nucleon mass $m_N = m_{\text{He}}/4 = 931.75 \text{ MeV}$, nucleon radius $r_N = 1.25 \text{ fm}$

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- Integrate via shooting from the center:
at $r = 0$: $h(0) = 1$ (anti-vacuum), $\mathbf{B}(0) = 1$ (no enclosed matter), free constant $\rho(0)$
till R where $h(R) = 0$ (vacuum) and $p'(R) \equiv 0$

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- Concretely, we will assume two different potentials:

$$\mathcal{U} = \mathcal{U}_\pi = 1 - \cos \xi,$$

$$\mathcal{U} = \mathcal{U}_\pi^2 = (1 - \cos \xi)^2,$$

Solving the system

Solutions exist up to maximum value of B or $n = B/B_\odot$ where

$$\mathcal{U}_\pi : \quad n_{\max} = 5.005, \quad M_{\max} = 3.734M_\odot, \quad R_{\max} = 18.458 \text{ km},$$

$$\mathcal{U}_\pi^2 : \quad n_{\max} = 3.271, \quad M_{\max} = 2.4388M_\odot, \quad R_{\max} = 16.44 \text{ km}$$

Recent observations: $M_{\max} \sim 2M_\odot$ established, indications for $M_{\max} \sim 2.5M_\odot$

$$10 \text{ km} < R_{\max} < 20 \text{ km}$$

Neutron star mass vs baryon number

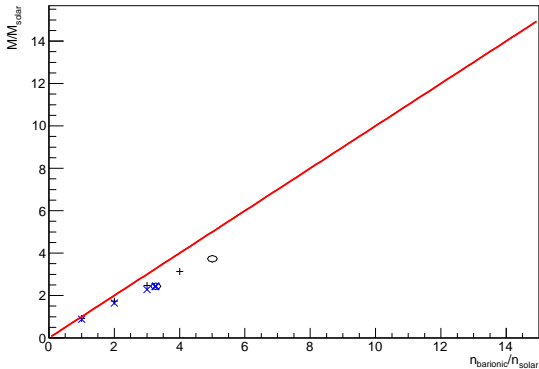


Figure: Neutron star mass as function of baryon number, both in solar units. Symbol plus (+): potential \mathcal{U}_π . Symbol cross (\times): potential \mathcal{U}_π^2 .

Neutron star mass vs. neutron star radius

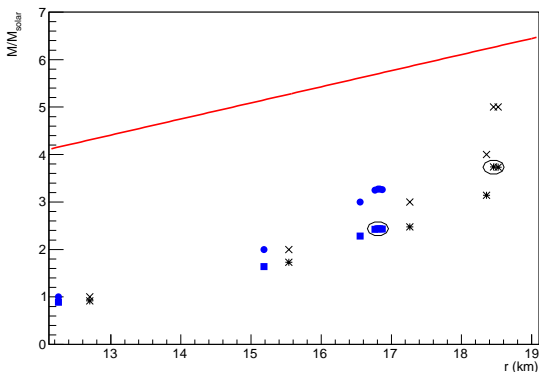


Figure: Neutron star mass as a function of the neutron star radius. Symbol asterisk (*): potential \mathcal{U}_π . Symbol square: potential \mathcal{U}_π^2 . Maximum values are indicated by circles. Also the value of n is plotted. The cross (\times) corresponds to \mathcal{U}_π and the circle to \mathcal{U}_π^2 .

Equation of State for $n = 1$

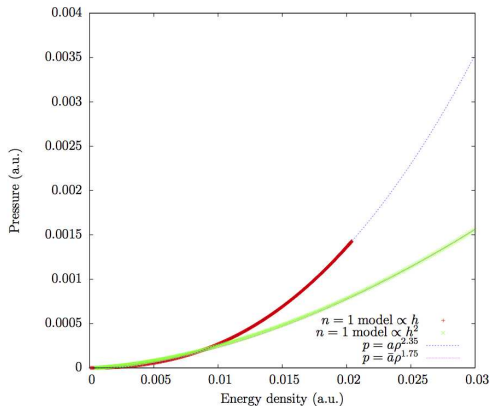


Figure: Symbol plus (+): potential \mathcal{U}_π . Symbol cross (×): potential \mathcal{U}_π^2 .
 Dotted lines: corresponding fit functions.

Equation of State for n_{max}

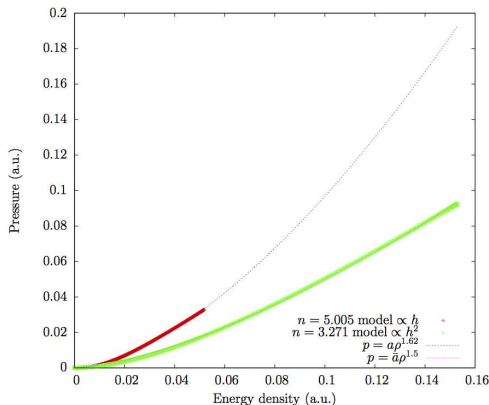


Figure: Symbol plus (+): potential \mathcal{U}_π . Symbol cross (\times): potential \mathcal{U}_π^2 .
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Results

- $M_{\max} \sim 2.5 - 3.5M_{\odot}$ and $R_{\max} \sim 10 - 20\text{km}$ precisely in right range

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$$\rho(\rho) = a(B)\rho^{b(B)}$$

EoS of polytropic star but $a = a(B)$, $b = b(B)$

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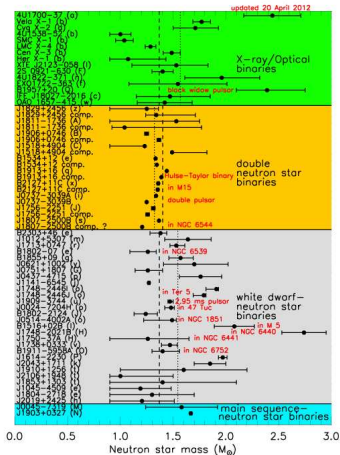
EoS of polytropic star but $a = a(B)$, $b = b(B)$

- Concretely EoS stiffer for larger B (larger M)

C.Adam, C. Naya, J. Sanchez-Guillen, R. Vazquez, A. Wereszczynski, arXiv:1407:3799

Comparing to Observed Neutron Star Masses

J. M. Lattimer, Ann. Rev. Nucl. Part. Sci. 62 (2012) 485



Conclusions

- Novel Skyrme model as a limit of generalized ones with analytical solutions.
- Linear mass - baryon number relation (BPS property).
- BPS Skyrme model produces excellent results also for neutron stars (concretely M_{\max} and R_{\max}).
- Extrapolation from $B = 1$ to $B \sim 10^{57}$ within one and the same model.
- Equation of state (EoS) NOT universal
 $\Rightarrow M(R)$ is monotonously growing function.
- Completely compatible with (not very precise) observational data.
- Work in progress: Better crust corrections and new potentials.

