

Fits and Related Systematics for the Hadronic Vacuum Polarization on the Lattice

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Based on coll. with **Christopher Aubin** (Fordham U.), **Tom Blum** (Connecticut U.),
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([arXiv:1405.2389](https://arxiv.org/abs/1405.2389) ; [arXiv:1309.2153](https://arxiv.org/abs/1309.2153) ; [arXiv:1307.4701](https://arxiv.org/abs/1307.4701) ; [arXiv:1205.3695](https://arxiv.org/abs/1205.3695))

Outline: How to get a_μ^{HVP} w. error less than 1% ?

- Low- Q^2 region, where data is scarce, dominates: **problem**.
- New Strategy: "divide and conquer".

$$\text{Split } 0 \leq Q^2 \leq Q_{min}^2 \quad \text{and} \quad Q_{min}^2 \leq Q^2 \leq \infty.$$

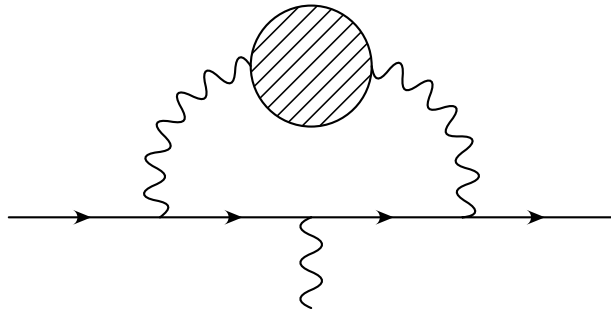
$$Q_{min}^2 \sim 0.1 \text{ GeV}^2$$

- Trapezoid rule approximation for $0.1 \text{ GeV}^2 \lesssim Q^2 \leq \infty$
- Use 3 (independent) methods for $0 \leq Q^2 \lesssim 0.1 \text{ GeV}^2$:
 1. Pades
 2. Polynomial in conformal variable
 3. $\text{N}^2\text{LO-ChPT}$ supplemented with $O(p^8)$ phenomenological LEC.

Introduction

Lautrup-de Rafael '69

Blum '02

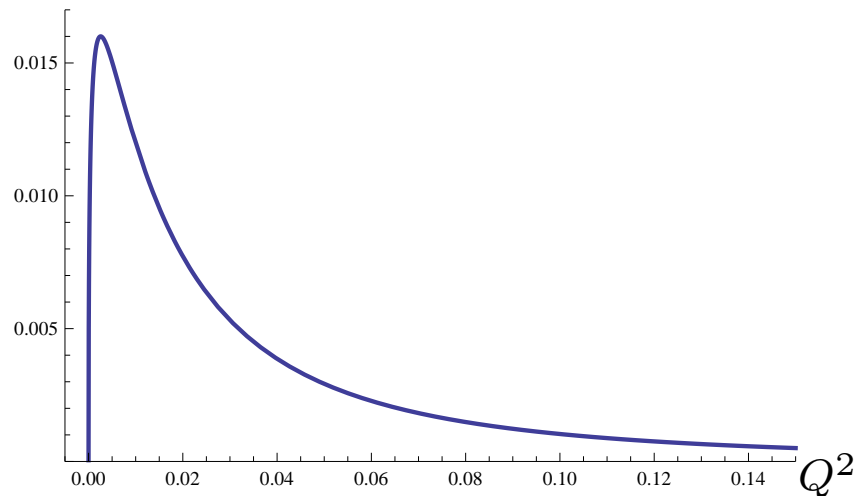


$$(g - 2)_\mu^{HVP} \sim \int_0^\infty dQ^2 \underbrace{f(Q^2)}_{\text{known}} \left[\Pi(Q^2) - \Pi(0) \right]$$

$$\hat{\Pi}(Q^2) \equiv \Pi(Q^2) - \Pi(0) \implies (g - 2)_\mu^{HVP}$$

★ integrand strongly peaked at $Q^2 \sim m_\mu^2/4 \sim 0.003 \text{ GeV}^2$

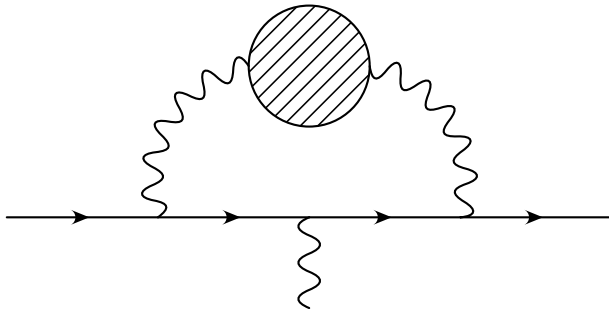
(g-2) integrand



Introduction

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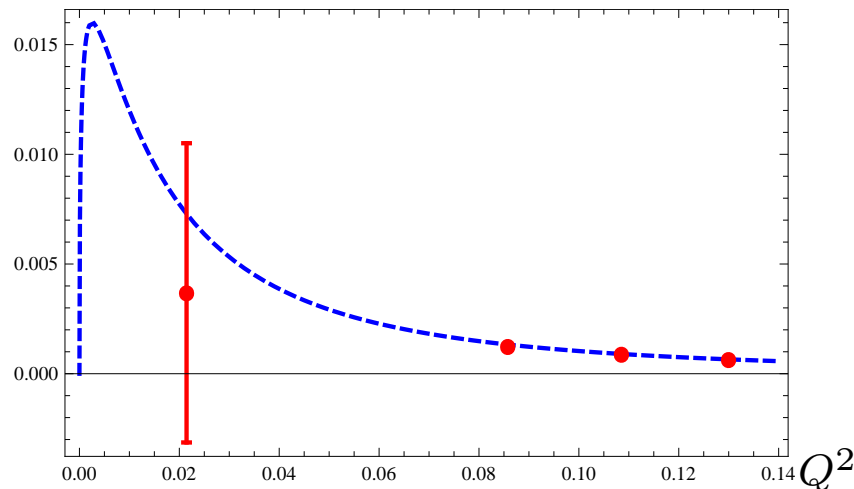
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if no good data in region of curvature \implies possibly wrong results !

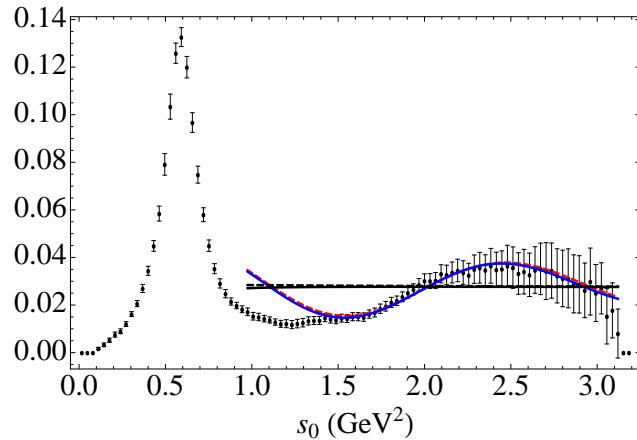
(g-2) integrand



Need reliable function !

how to test this theoretical error?

A τ -based model for $I = 1$ contributions



Boito, Cata, Golterman, Jamin, Mahdavi, Maltman, Osborne, SP '11 + '12

$t \leq t_{min} \text{ GeV}^2 \rightarrow \text{OPAL data.}$

$t \geq t_{min} \text{ GeV}^2:$

$$\text{Im}\Pi(t) = \rho_{\text{Pert.Th.}}(t) + e^{-\delta - \gamma t} \sin(\alpha + \beta t)$$

$$\Pi(Q^2) = -Q^2 \int_{4m_\pi^2}^{\infty} \frac{dt}{\pi} \frac{\text{Im}\Pi(t)}{t(t+Q^2)}$$

We take $t_{min} = 1.5 \text{ GeV}^2$.

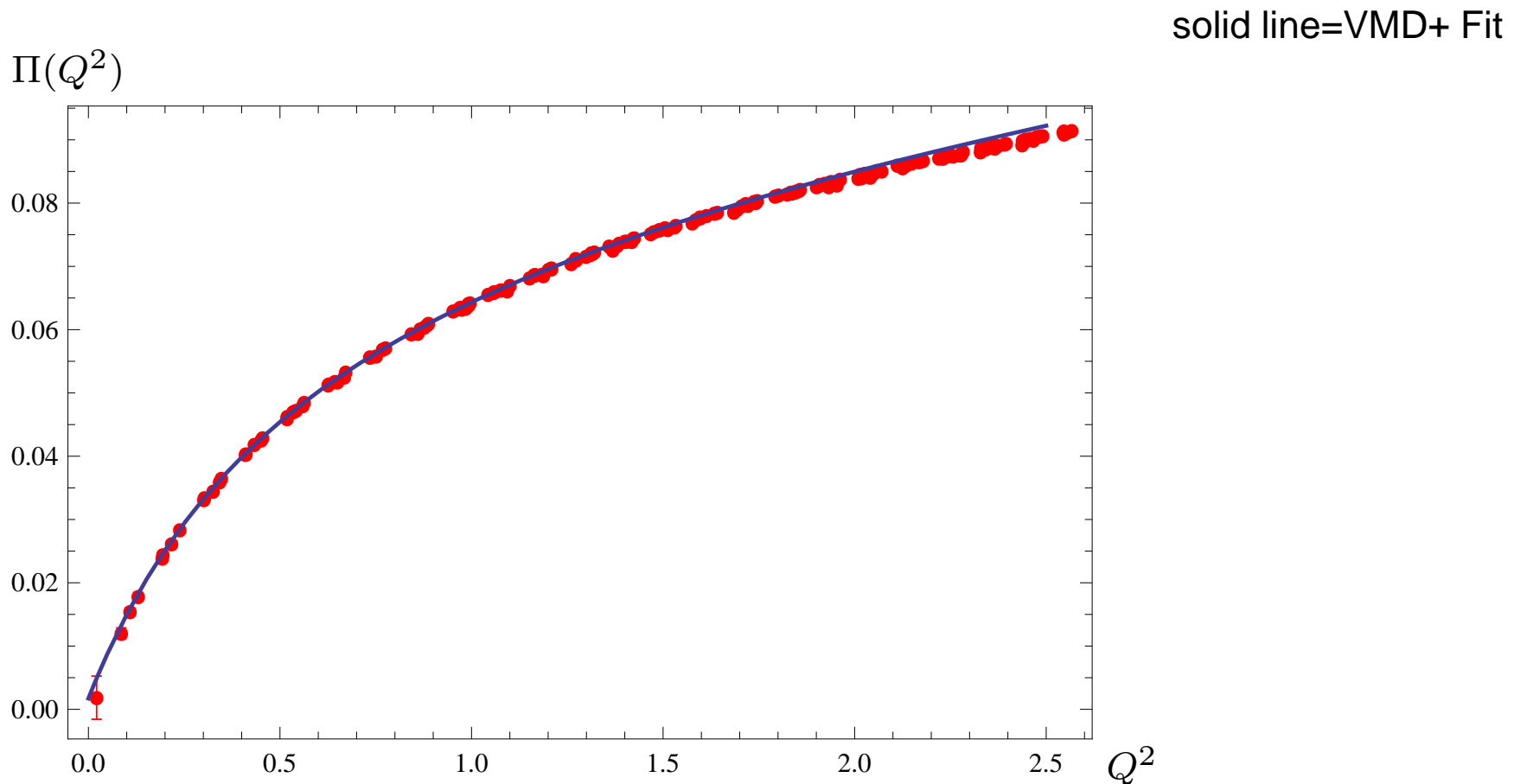
$$\Pi(0) = 0$$

- Take typical lattice Q^2 values + lattice covariance matrix
(e.g., Aubin et al. '12, $64^3 \times 144$ lattice, $a = 0.06 \text{ fm}$, $m_\pi = 220 \text{ MeV}$, periodic BCs)
- Generate fake lattice data for $\Pi(Q^2)$ and compare with true answer from model.
- You should try this model to check your systematics: it's very physically motivated !

The perils of using the wrong function

Take, e.g., the case of VMD+polynomial as fitting function

You may think this is a good fit for an accurate $(g - 2)_\mu$:

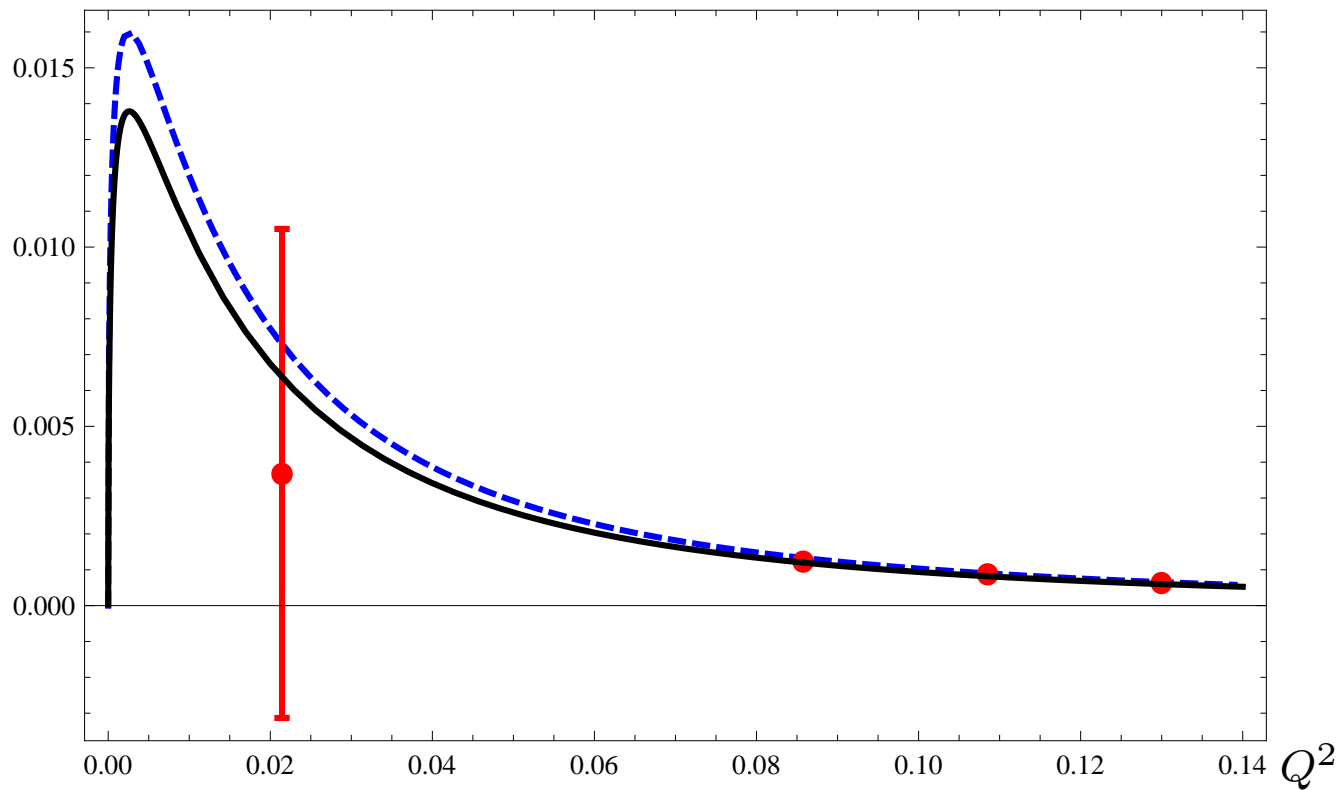


The perils of using the wrong function

while, in fact, this is what you should be looking at:

dashed blue=Truth
solid line=VMD+polynomial Fit

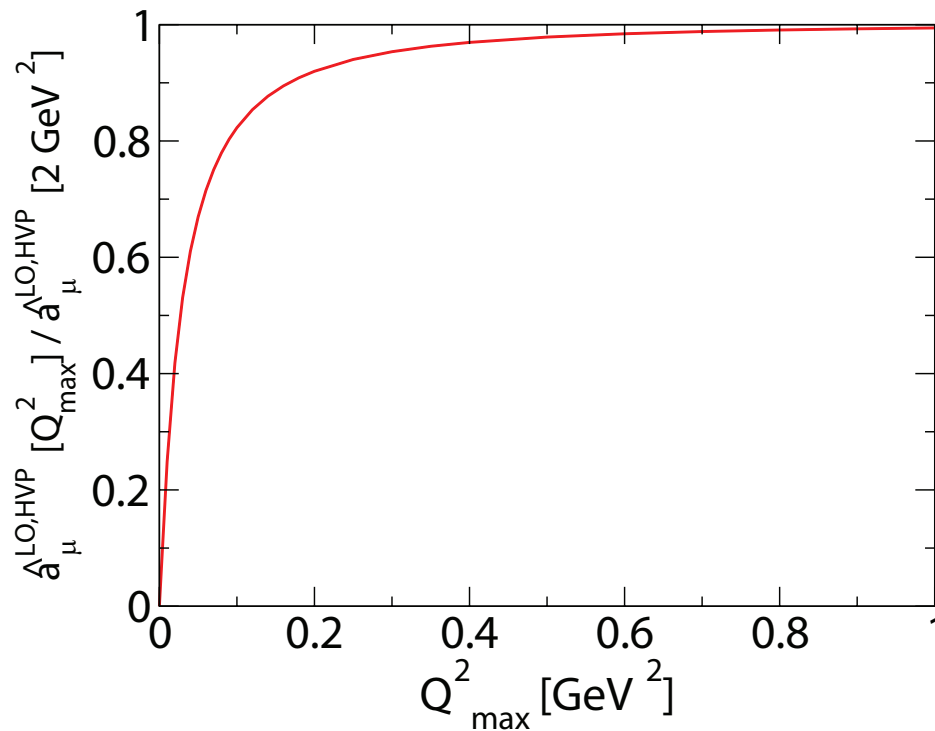
(g-2) Integrand



Systematic error in a_{μ}^{HVP} corresponds to 18 units of nominal fit error !

Hybrid strategy for a less than 1% error

More than 80% of a_μ^{HVP} accumulates for $Q^2 \lesssim 0.1 \text{ GeV}^2$.

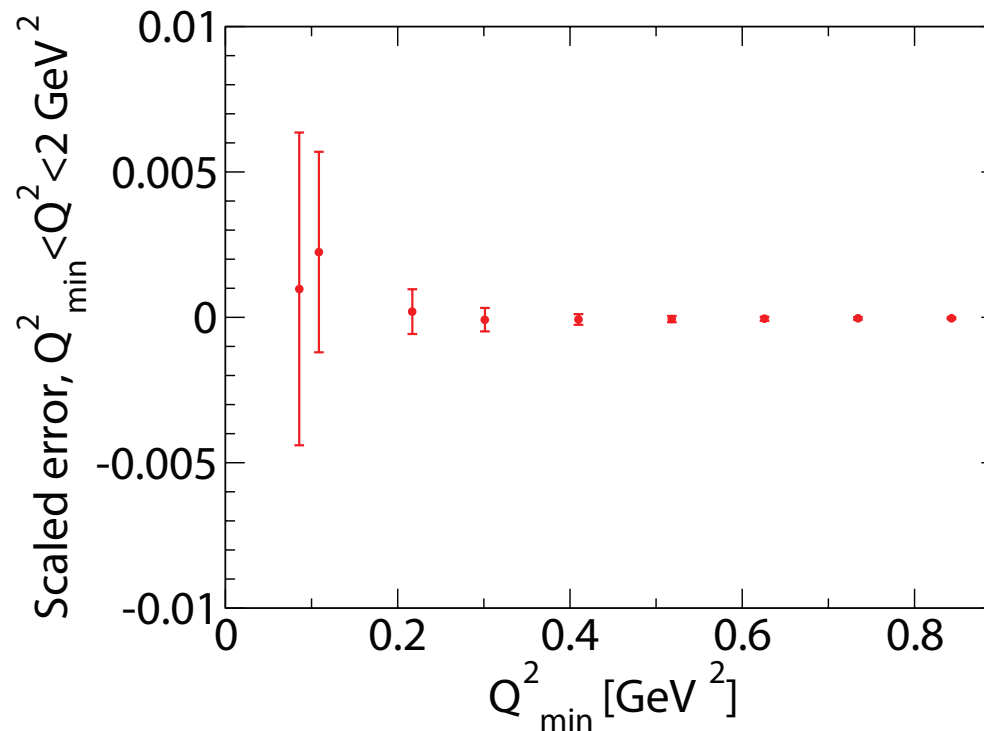


⇒ Split the contributions $0 \leq Q^2 \lesssim 0.1 \text{ GeV}^2$ and $0.1 \text{ GeV}^2 \lesssim Q^2 \leq \infty$

The $0.1 \text{ GeV}^2 \lesssim Q^2 \leq \infty$ interval (I)

Trapezoid rule approximation (use fake data for vac. pol. from $l=1$ model with Q^2 from MILC $64^3 \times 144$, $a = 0.06$ fm lattice data, including covariances, PBC).

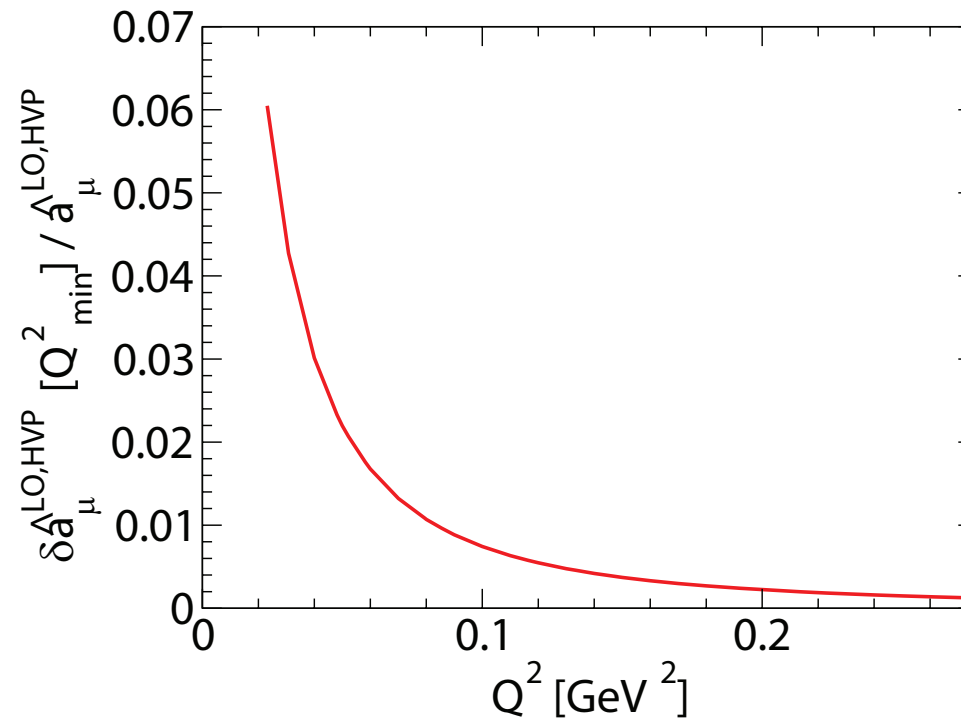
Systematic and statistical error due to trapezoid rule.



Should expect clear improvement w. AMA, twisting or analytic continuation (Feng et al. '13).

The $0.1 \text{ GeV}^2 \lesssim Q^2 \leq \infty$ interval (II)

Systematic error due to $\delta\Pi(0) = 0.001$ (obtained from $[1, 1]$ Pade fit to $Q^2 \leq 1 \text{ GeV}^2$).



The $0 \leq Q^2 \lesssim 0.1 \text{ GeV}^2$ interval

Three independent strategies:

- Pades .
- Polynomial in a conformal variable.
- N²LO- ChPT + C Q^4

Pades (I)

(Aubin, Blum, Golterman, S.P. '12; Chakraborty et al. '14)

Rational function matched to derivatives of $\Pi(Q^2)$ at $Q^2 = 0$, or matched/fitted to several Q^2 values.

Derivatives of $\Pi(Q^2)$ at $Q^2 = 0 \Leftrightarrow$ time moments

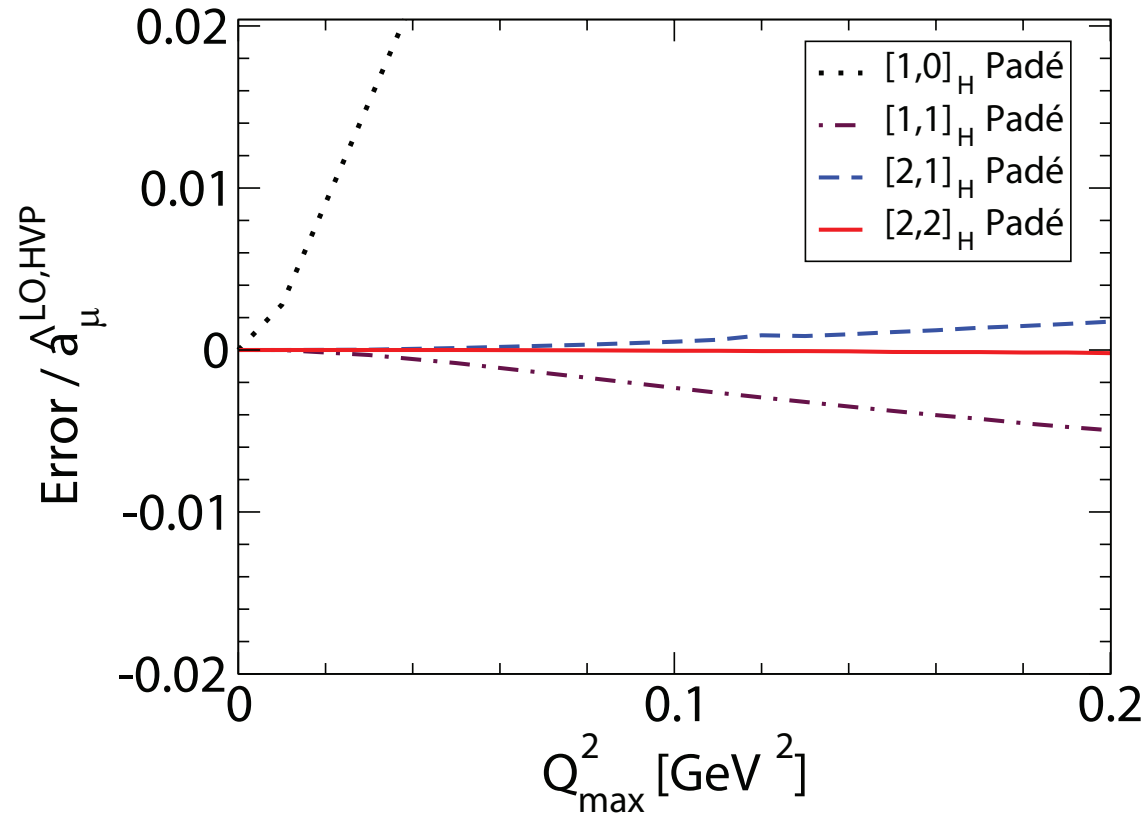
$$\hat{\Pi}(Q^2) = Q^2 (\Pi_1 + \Pi_2 Q^2 + \Pi_3 Q^4 + \Pi_4 Q^6 + \dots)$$

with $\Pi_j = \frac{(-1)^{j+1}}{(2j+2)!} \int d^4x t^{2j+2} \langle J_1(x) J_1(0) \rangle$

Interesting property (convergence+ alternating bounds):

$$[1, 0]_H \leq [2, 1]_H \leq \dots \leq [N + 1, N]_H \leq \hat{\Pi}(Q^2) \leq [N, N]_H \leq \dots \leq [2, 2]_H \leq [1, 1]_H$$

Padés (II)



$[1, 1]_H$ is good enough: 2 parameters i.e. up to t^6 moment.

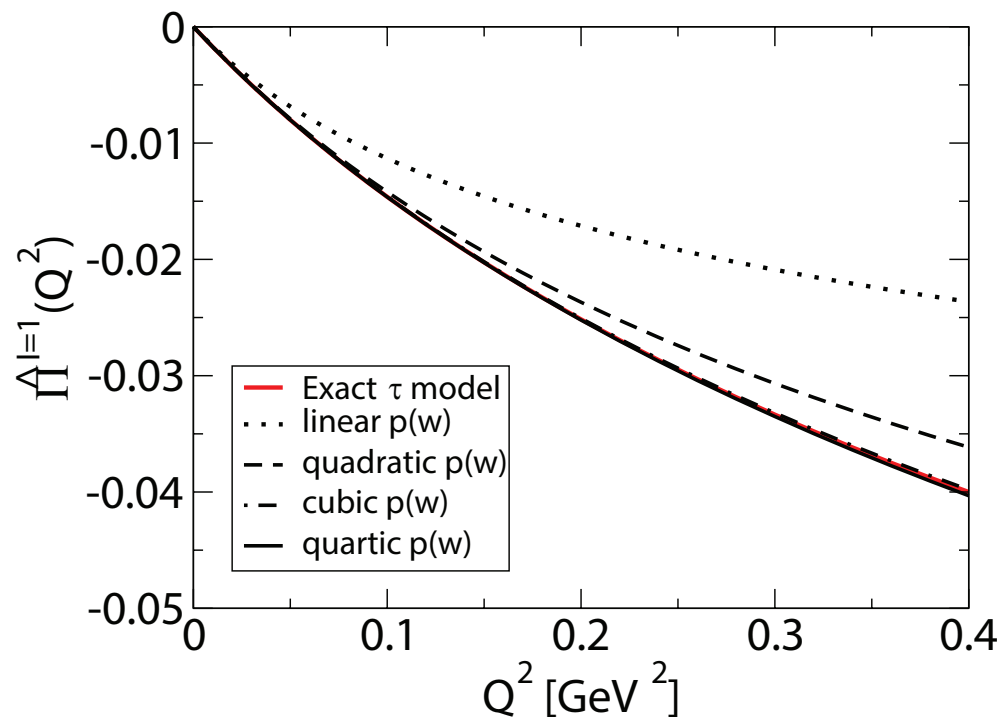
If, e.g., we fit Padé to data in interval $0.1\text{GeV}^2 \leq Q^2 \leq 0.2\text{GeV}^2$ then need $[2, 1]_H$, typically.

Polynomials in conformal variable

(Golterman, Maltman, S.P. '14)

$$w(Q^2) = \frac{1 - \sqrt{1+z}}{1 + \sqrt{1+z}} \quad , \quad z = \frac{Q^2}{4m_\pi^2} \quad , \quad \Pi(Q^2) = \sum_{n=0}^{\infty} p_n w^n$$

Also convergent (like Pades) but no alternating bounds.



Sub-1% systematic error \Rightarrow quadratic polynomial in w (2 parameters).

Similar to $[1, 1]_H$ Pade.

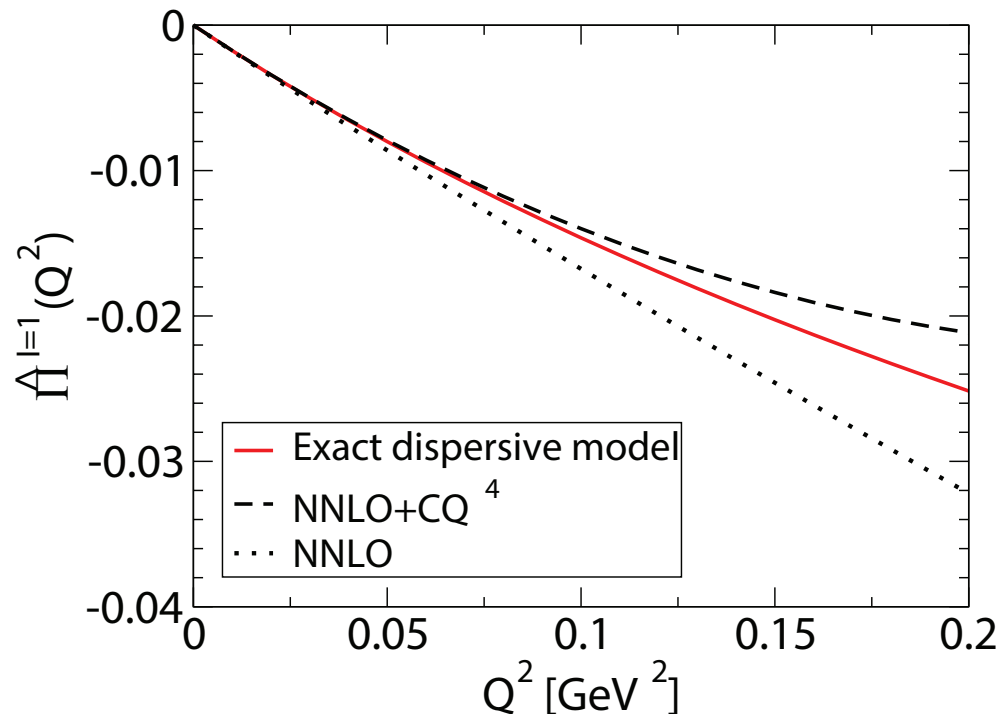
Again, if fit in an interval, then \Rightarrow cubic polynomial.

$N^2LO - ChPT + C Q^4$

(Golterman, Maltman, S.P. '14)

Phenomenologically motivated, but not fully model independent (no chiral log at $\mathcal{O}(p^8)$).

Best strategy is get LEC from time moments (as did with Pades and conf. polynomial).



However, even with the term CQ^4 , this ChPT functional form performs worse than Pades and/or conformal polynomial.

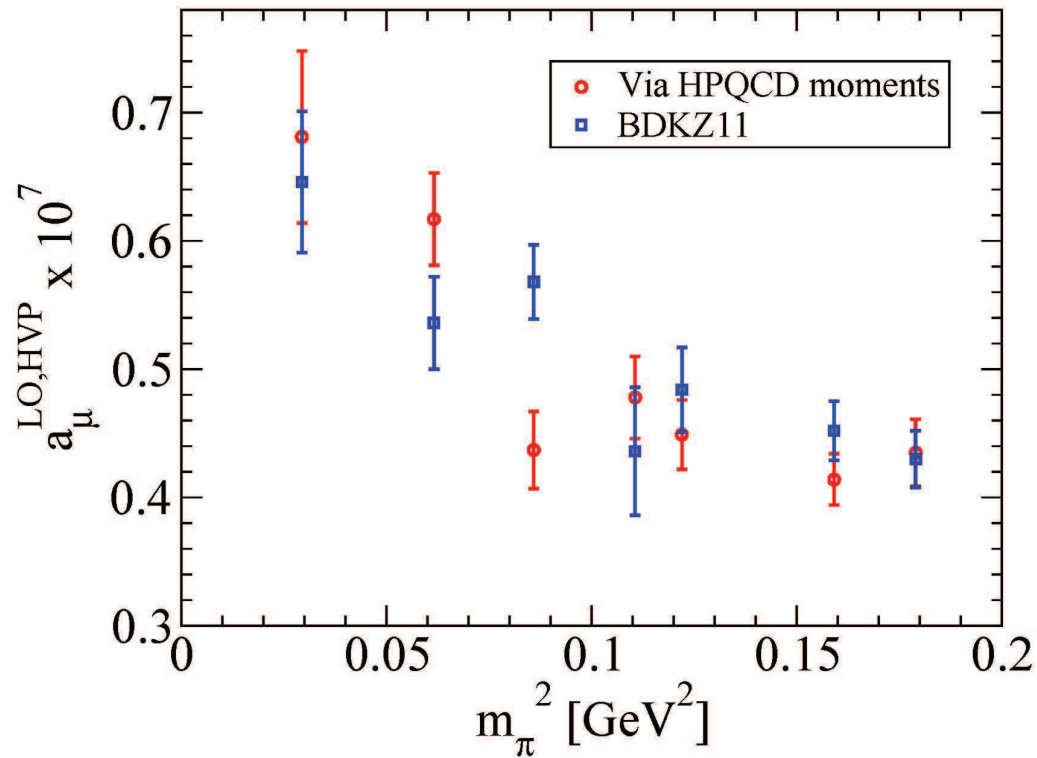
⇒ Mostly useful as a cross check.

Very Preliminary Results

Used time moments to construct Pades. u, d, s quarks.

RBC/UKQCD DWF data (Hudspith, Lewis, Maltman, Portelli '14)

RBC/UKQCD DWF ensemble $a_\mu^{\text{LO,HVP}}$ results



(BDKZ11: Boyle, Del Debbio, Kerrane, Zanotti '11)

Conclusions

- Fitting on a large Q^2 window very dangerous.
- Pointed out that a hybrid strategy with
 - trapezoid rule approximation for $0.1 \text{ GeV}^2 \lesssim Q^2 \leq \infty$
 - Pades, conformal polynomials and N²LO- ChPT + C Q^4 for $0 \leq Q^2 \lesssim 0.1 \text{ GeV}^2$

is capable of reaching the desired below 1% error in $a_{\mu\text{HVP}}$ provided time moments can be determined accurately enough (rule of thumb: $\delta\Pi_j/\Pi_j \lesssim j \%$)

∃ First attempt at time moments by Chakraborty et al. '14 for s, c quarks.

What about u, d ?

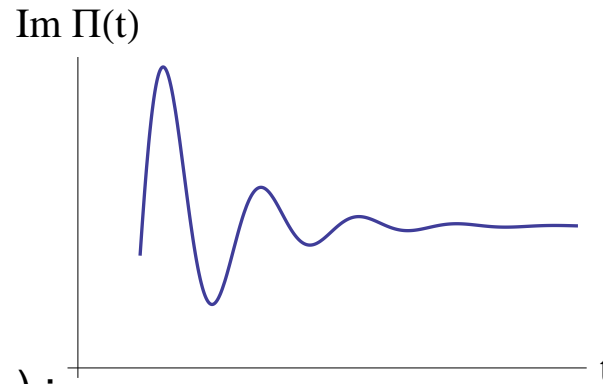
(encouraging results from Hudspith et al. RBC/UKQCD DWF data, still very preliminary...)

- Alternatively (and a good check!): get good data in the low $Q^2 \sim 0.1\text{GeV}^2$ interval (twisting, analytic continuation, etc...) and fit (Pades, Conf. polynomial,...)

BACK-UP SLIDES

Duality Violations

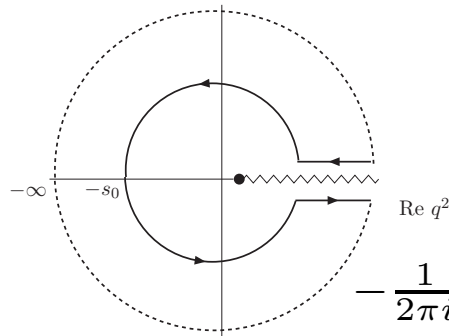
- OPE valid in euclidean, but not in minkowski. We know that spectrum \neq OPE



- We expect (@ large t) :

$$\text{Im}\Pi_{DV} \sim \text{Im}(\Pi - \Pi_{OPE}) \sim \underbrace{\kappa}_{\text{OPE asympt.}} \underbrace{e^{-\gamma t} \sin(\alpha + \beta t)}_{\text{Regge}}$$

- $\Pi_{DV}(s) \rightarrow 0$ as $|s| \rightarrow \infty$. Then:



$$-\frac{1}{2\pi i} \oint_{|z|=s_0} dz w(z) \Pi_{DV}(z) = - \underbrace{\int_{s_0}^{\infty} ds}_{\text{extrapolation!}} w(s) \frac{1}{\pi} \text{Im}\Pi_{DV}(s)$$

(Cata-Golterman-S.P. '05)

Fitting functions

Aubin, Blum, Golterman, SP '12

★ Padés, model independent, they enjoy a convergence theorem for $N \rightarrow \infty$:

$$\Pi(Q^2) = \underbrace{\Pi(0) + Q^2 \left(a_0 + \sum_{r=1}^N \frac{a_r}{Q^2 + b_r} \right)}_{\text{Pade}}$$

$\Pi(0)$, a 's and b 's are fitting parameters.

★ **VMD** is **not** a Pade, since you fix $b_1 = M_\rho^2$. (true $\Pi(Q^2)$ has cut starting at $4m_\pi^2 \dots$)

We have: $a_0 \neq 0 \implies [N, N]$ Pade; $a_0 = 0 \implies [N - 1, N]$ Pade.

For instance:

- $\frac{a_1}{Q^2 + b_1}$ is a $[0,1]$ Pade $\implies \Pi(Q^2) = \Pi(0) + Q^2 \left(\frac{a_1}{Q^2 + b_1} \right)$
- $a_0 + \frac{a_1}{Q^2 + b_1}$ is a $[1,1]$ Pade $\implies \Pi(Q^2) = \Pi(0) + Q^2 \left(a_0 + \frac{a_1}{Q^2 + b_1} \right)$

etc...

Example: Stieltjes. No errors.

Toy Vac. Pol. function ($\text{Im}\Pi \geq 0$), with a cut, $-\infty \leq Q^2 \leq -1$:

$$\begin{aligned} \frac{\Pi(0) - \Pi(Q^2)}{Q^2} &= \int_1^\infty \frac{dt}{(t + Q^2)} \frac{1}{t} \\ &= \frac{1}{Q^2} \log(1 + Q^2) \end{aligned}$$

Theorem: As $N \rightarrow \infty$, with a_i, b_i determined from the function (and derivatives) at $Q^2 = 0$, or at multiple points,

$$\frac{a_0 + a_1 Q^2 + a_2 Q^4 + \dots + a_N Q^{2N}}{1 + b_1 Q^2 + \dots + b_N Q^{2N}} \longrightarrow \frac{\Pi(0) - \Pi(Q^2)}{Q^2}$$

everywhere in a compact region in complex Q^2 , except on the cut.

