

## Abstract

Models of electroweak symmetry breaking with extended Higgs sectors are theoretically well motivated. In this study, we investigate the phenomenology of the new Yukawa couplings in generic two-Higgs-doublet models. We find that a heavy charged Higgs together with  $\alpha, \tan\beta \sim O(1)$ , type-II and III could enhance the two-photon production cross section; however, with large  $\tan\beta$  scenario, only type-III could match the LHC data. Additionally, we study the implications of LHC data on the production cross sections for the channel  $h \rightarrow Z\gamma$  and branching ratio for  $t \rightarrow ch$  decay.

## Motivation -2HDM-

- Tree-level flavour changing neutral currents (FCNC's) due to non diagonal Yukawa interactions .
- New sources of CP violation in scalar-scalar interactions .
- An increase in the scalar spectrum with one charged and three neutral scalar .

The scalar potential can be written as :

$$V(\Phi_1, \Phi_2) = m_{11}^2(\Phi_1^\dagger \Phi_1) + m_{22}^2(\Phi_2^\dagger \Phi_2) - (m_{12}^2(\Phi_1^\dagger \Phi_2) + h.c.) + \frac{\lambda_1}{2}(\Phi_1^\dagger \Phi_1)^2 + \frac{\lambda_2}{2}(\Phi_2^\dagger \Phi_2)^2 + \lambda_3(\Phi_1^\dagger \Phi_1)(\Phi_2^\dagger \Phi_2) + \lambda_4(\Phi_1^\dagger \Phi_2)(\Phi_1^\dagger \Phi_1) + (\frac{\lambda_5}{2}(\Phi_1^\dagger \Phi_2)^2 + h.c.) + \{(\lambda_6(\Phi_1^\dagger \Phi_1) + \lambda_7(\Phi_2^\dagger \Phi_2))(\Phi_1^\dagger \Phi_2) + h.c.\}$$

The lightest higgs coupling to fermions are written as :

$$\mathcal{L}_Y^h = \bar{u}_i \left( \frac{\cos \alpha}{\sin \beta} \frac{m_{u_i}}{v} \delta_{ij} - \frac{\cos(\alpha - \beta)}{\sqrt{2} \sin \beta} X_{ij}^u \right) u_j h + \bar{d}_i \left( -\frac{\sin \alpha}{\cos \beta} \frac{m_{d_i}}{v} \delta_{ij} + \frac{\cos(\alpha - \beta)}{\sqrt{2} \cos \beta} X_{ij}^d \right) d_j h,$$

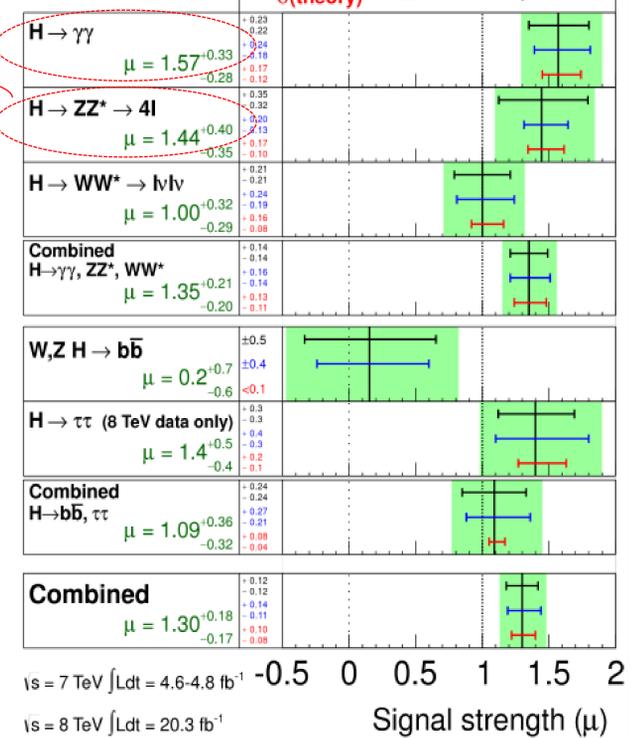
The overall signal strength for diphoton at ATLAS may be explained by one of Standard model's extensions like **Two-Higgs-Doublet Models** (2HDMs) .

So , in order to study the new physics effects on the LHC measurements for some specific XY finale state, we define the quantities  $R_{XY}$

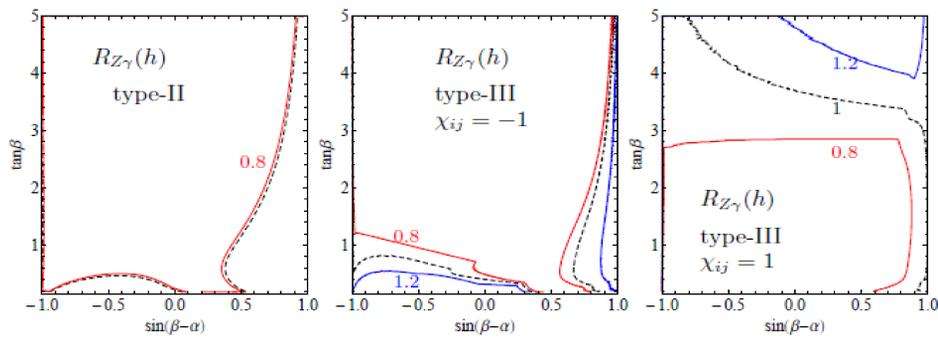
$$R_{XY} = \frac{\sigma(pp \rightarrow h)B(h \rightarrow XY)}{\sigma(pp \rightarrow h)_{SM}B(h \rightarrow XY)_{SM}}$$

ATLAS Prelim.

$m_H = 125.5 \text{ GeV}$

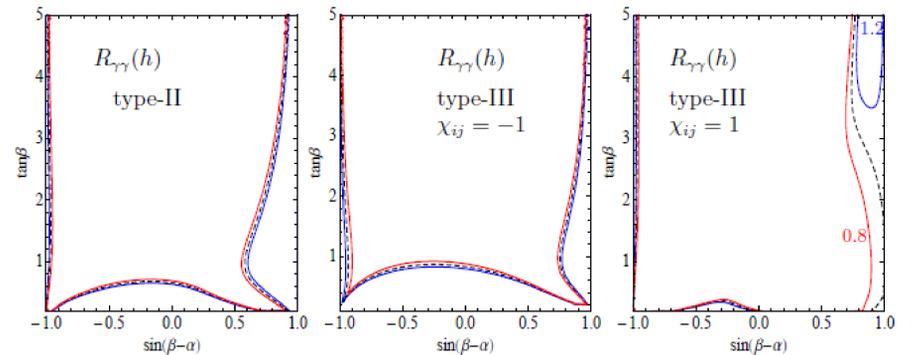


## Numerical Results



The  $R_{Z\gamma}(h)$  for  $m_h = 125 \text{ GeV}$  in the type-II (left panel), type-III with  $\chi_{ij} = -1$  (middle panel) and type-III with  $\chi_{ij} = 1$  (right panel). The contour lines are 0.8 (red), 1 (dashed) and 1.2 (blue).

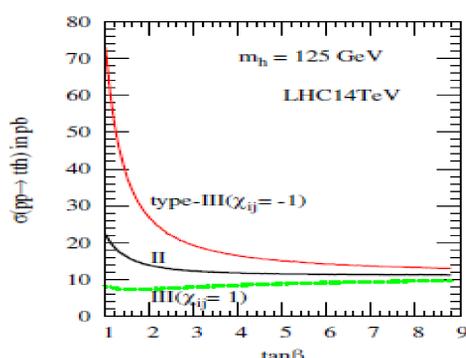
If  $R_{Z\gamma}(h)$  requires to be between 1 and 2, then type-II is almost excluded. Both type-III with  $\chi_{ij} = 1$  provides a wider allowed parameter space especially type-III with  $\chi_{ij} = -1$  in the region where  $\sin(\beta - \alpha)$  is negative .



The  $R_{\gamma\gamma}(h)$  for  $m_h = 125 \text{ GeV}$  in the type-II (left panel), type-III with  $\chi_{ij} = -1$  (middle panel) and type-III with  $\chi_{ij} = 1$  (right panel). The contour lines are 0.8 (red), 1 (dashed) and 1.2 (blue).

The new fermion couplings in type-II and type-III could enhance  $R_{\gamma\gamma}$ . Type-III could fit the data for both  $\chi_{ij}$  at different region of  $\sin(\beta - \alpha)$  and provides a wider allowed parameter space .

As we already know, flavor changing neutral currents (FCNCs) can be induced at tree level in type-III model, so if a large BR for  $t \rightarrow ch$  is measured, the type-III model could be the best extension of the SM .



Production cross section of  $pp \rightarrow t\bar{t}h$  at the LO as a function of  $\tan\beta$  in the general THDM where  $m_H = 200 \text{ GeV}$ ,  $m_{H^\pm} = 480 \text{ GeV}$ ,  $m_{A0} = 420 \text{ GeV}$  and  $m_h = 125 \text{ GeV}$  are chosen .

## Conclusion:

- ✓ The phenomenologies of the three types of yukawa interactions are completely different from each other.
- ✓ The differences between the types of the yukawa interactions largely affect the production and the decay of the Higgs boson.

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