

α_s from τ decays: higher orders and perturbative behaviour

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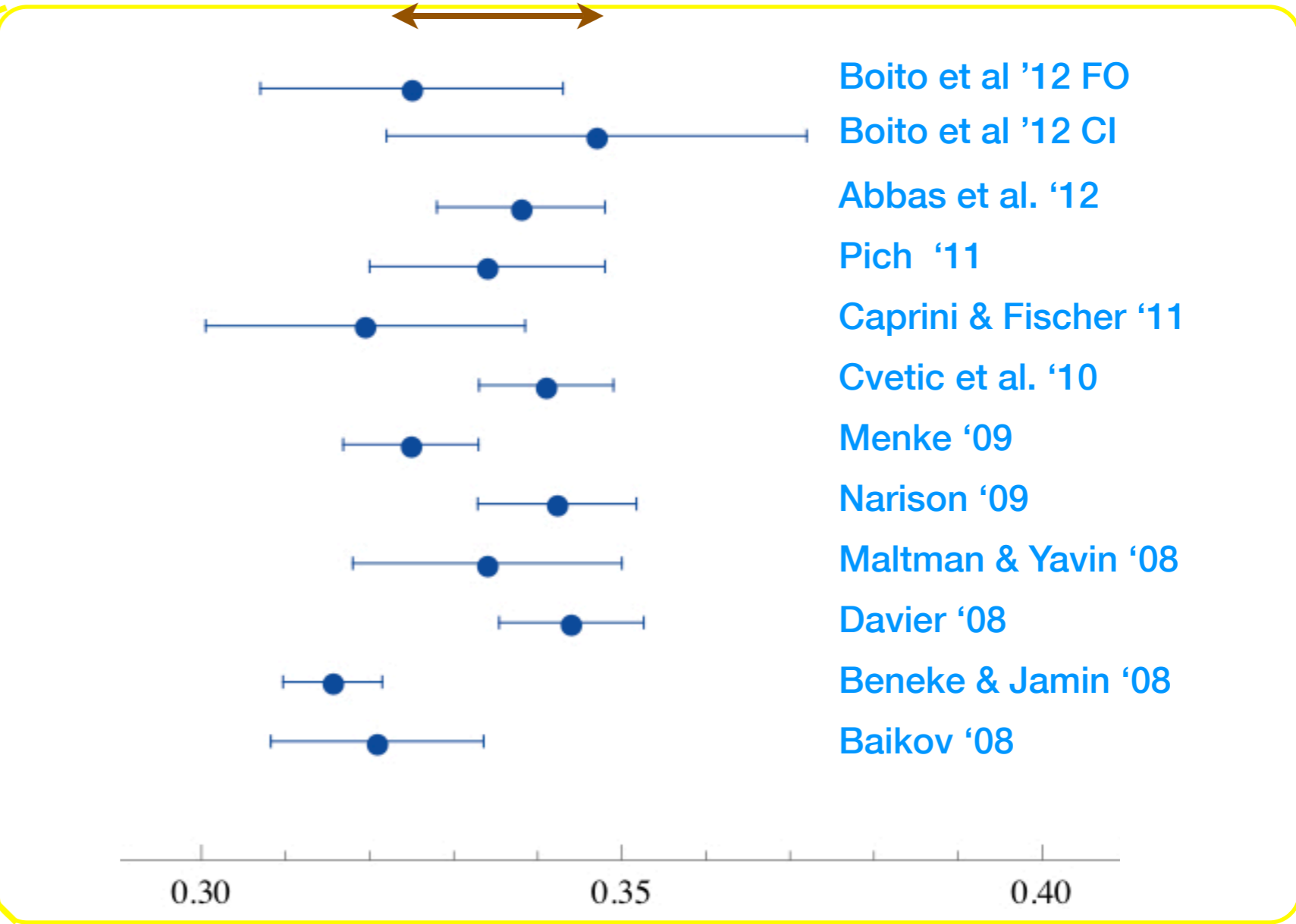
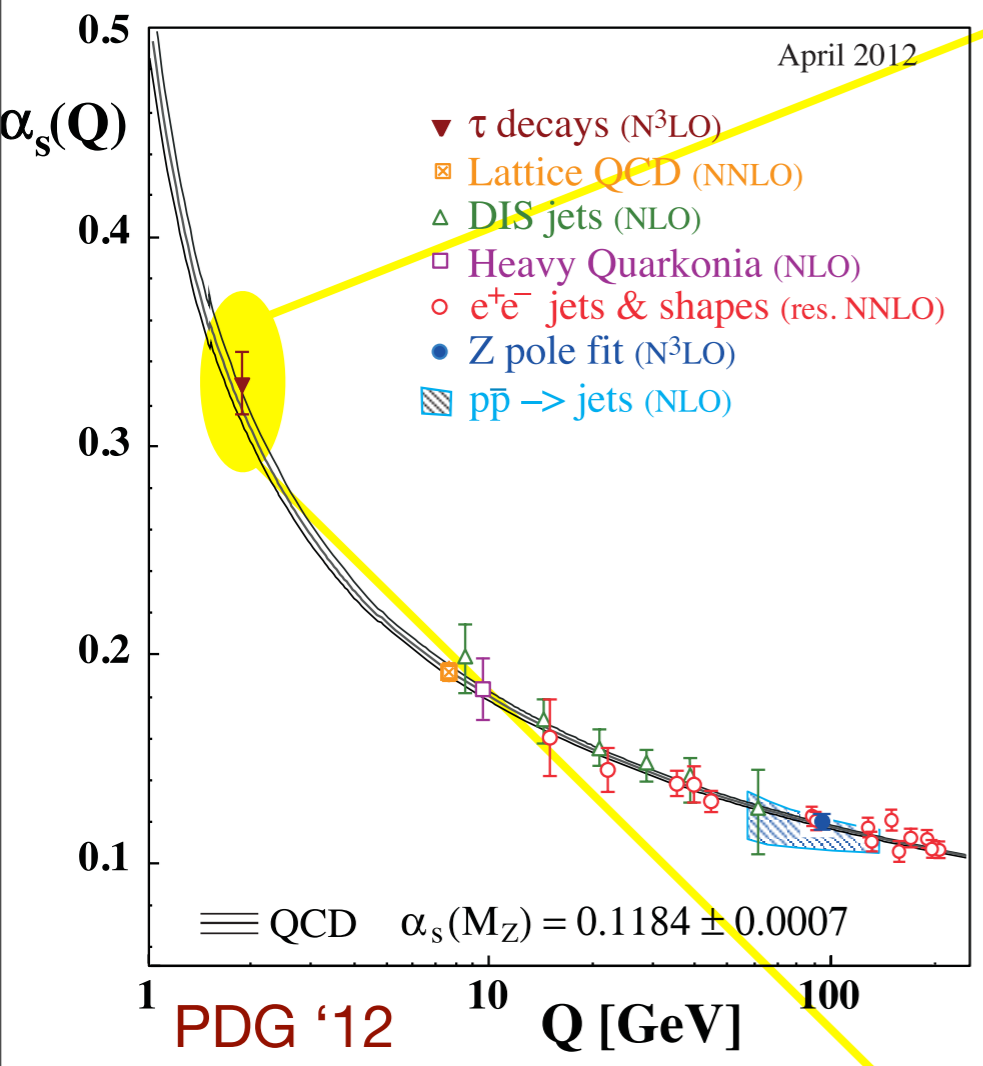
- M Beneke, DB, M Jamin, *JHEP* **1301** (2013)
- DB, M Golterman, M Jamin, A Mahdavi, K Maltman, J Osborne, S Peris, *Phys Rev D* **85** (2012)
- DB, O Catà, M Golterman, M Jamin, K Maltman, J Osborne, S Peris, *Phys Rev D* **84** (2011)
- DB, PoS Confinement X (2012)

- **Introduction**

- **Recent developments in α_s from tau decays**
 - Moments, higher orders and FOPT vs CIPT
 - Duality Violations
 - Updated ALEPH data (correct correlation matrices)

- **Prospects**
 - New analysis of updated ALEPH data (in progress, no new α_s value)

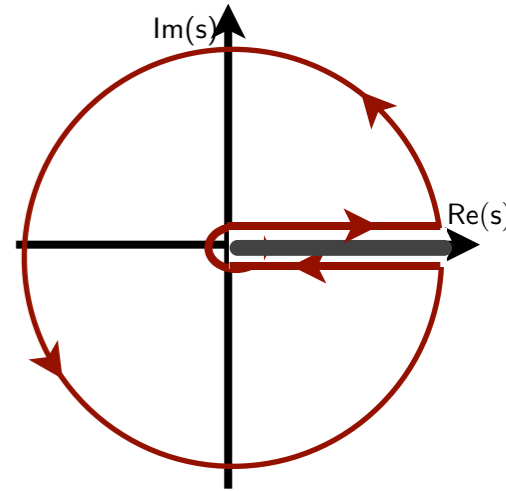
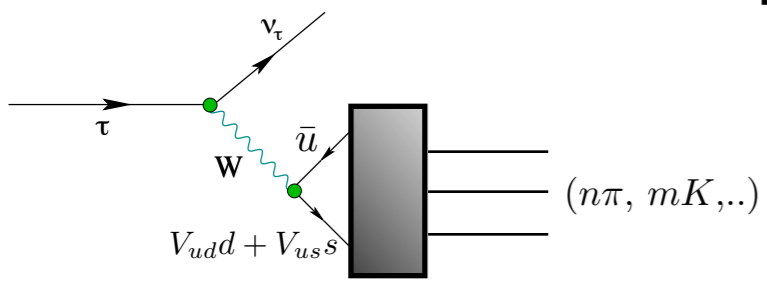
Prescription for the RG improvement ($\sim 7\%$ error)



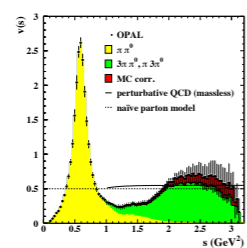
- Spread in the results reflect (mainly) details of the theoretical input.
- There are still open questions (Renormalization Group Improvement, duality violations, ...)

● Sum rules for the spectral functions

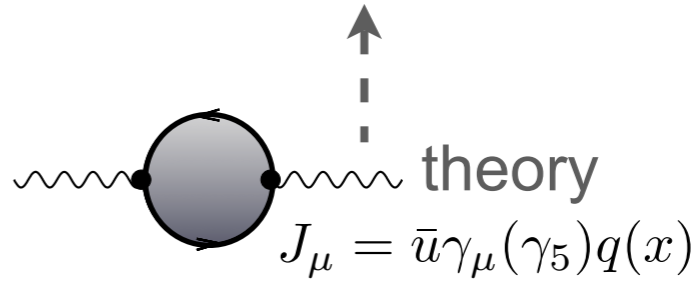
(in tau decays) Braaten, Narison, and Pich, 1992



$$\int_0^{s_0} ds w(s) \frac{1}{\pi} \text{Im} \tilde{\Pi}(s) = \frac{-1}{2\pi i} \oint_{|z|=s_0} dz w(z) \tilde{\Pi}(z)$$



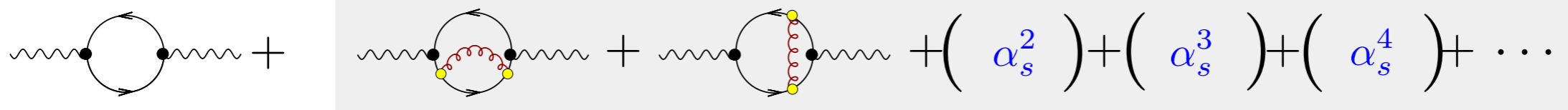
experiment (OPAL and ALEPH)



● Contributions to the sum rule (theory side)

$$R_{V/A}^{w_i}(s_0) = \frac{N_c}{2} S_{EW} |V_{ud}|^2 \left[\delta_{w_i}^{\text{tree}} + \delta_{w_i}^{(0)}(s_0) + \sum_{D \geq 2} \delta_{w_i, V/A}^{(D)}(s_0) + \delta_{w_i, V/A}^{\text{DV}}(s_0) \right]$$

OPE DVs



α_s^4 : Baikov, Chetyrkin, Kühn 2008

Recent developments in α_s from tau decays

● Systematic study of the weight function dependence

Beneke, DB, Jamin, '13

$$\int_0^{s_0} ds w(s) \frac{1}{\pi} \text{Im} \tilde{\Pi}(s) = \frac{-1}{2\pi i} \oint_{|z|=s_0} dz w(z) \tilde{\Pi}(z)$$

- Convergence of pt. series is weight function dependent. What are *ideal* weight functions?
- Several moments used in literature. Renormalization group: what's best CIPT or FOPT?
(**C**ontour **I**mproved **P**erturbation **T**heory or **F**ixed **O**rders **P**erturbation **T**heory)

● Duality violations

$$R_{V/A}^{w_i}(s_0) = \frac{N_c}{2} S_{EW} |V_{ud}|^2 \left[\delta_{w_i}^{\text{tree}} + \delta_{w_i}^{(0)}(s_0) + \sum_{D \geq 2} \delta_{w_i, V/A}^{(D)}(s_0) + \delta_{w_i, V/A}^{\text{DV}}(s_0) \right]$$

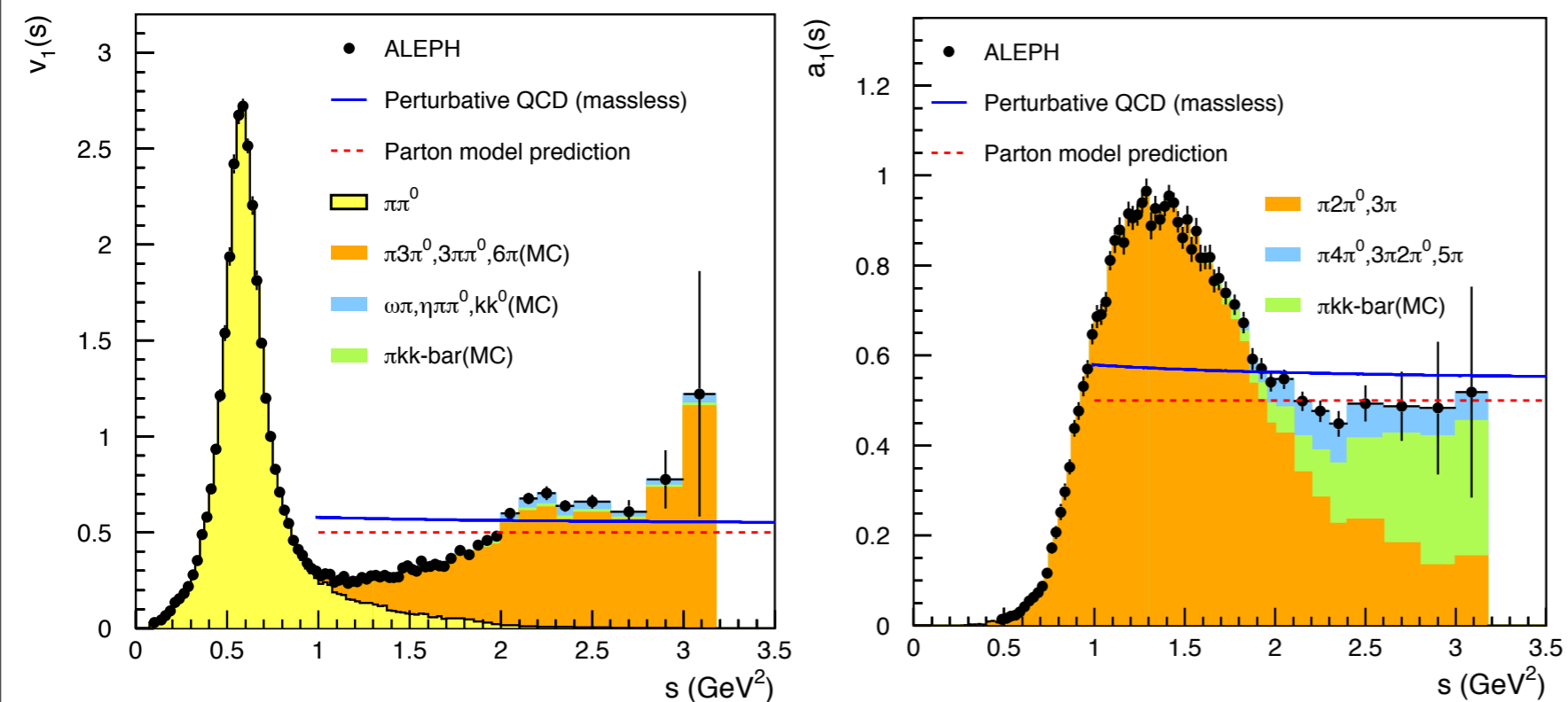
Catà, Golterman, & Peris, '05, '09; Boito et al '11, '12.

- Resonance effects close to the real axis.
- Must take them into account (at the very least for error calculation).
- One must resort to reasonable models based on QCD.
- Source of uncertainty not considered in the “standard” analysis.

Davier et al 2013

● Updated ALEPH data for spectral functions

Davier, Höcker, Malaescu, Yuan, Zhang, 1312.1501v2



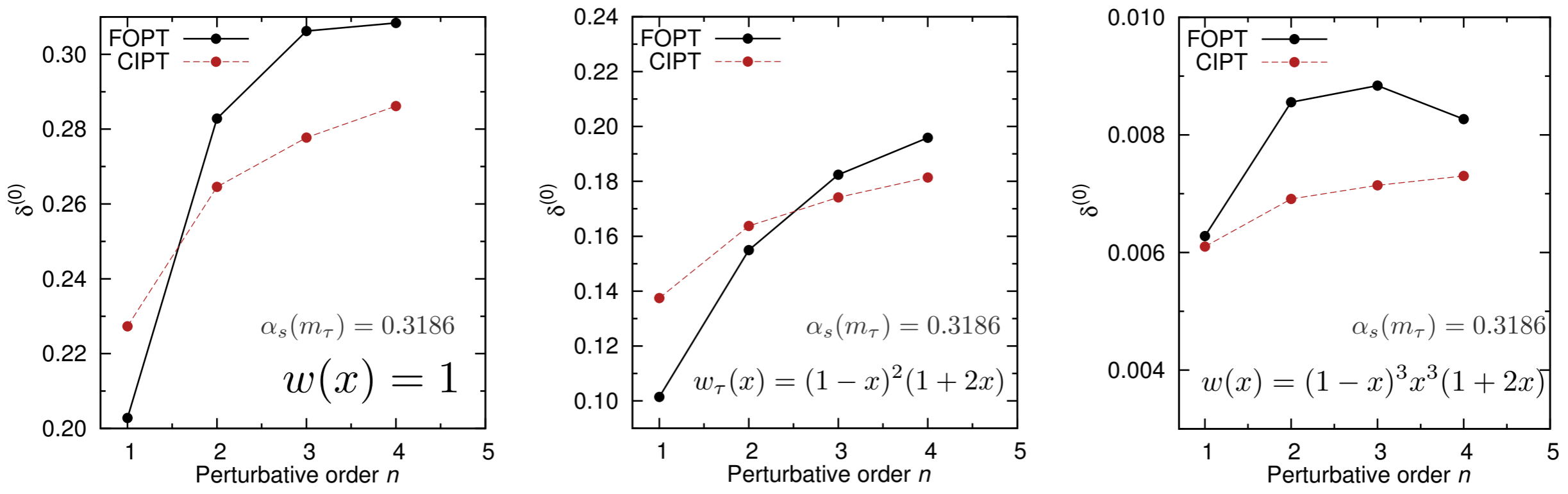
- Corrected covariances matrix (due to unfolding).
- Corrects a problem discovered the previous version (05/08) of the ALEPH data.
DB, Catà, Golterman, Jamin, Maltman, Osborne and Peris, '11
- Smaller uncertainties than the OPAL data.

Perturbative series, weight functions, and RG improvement

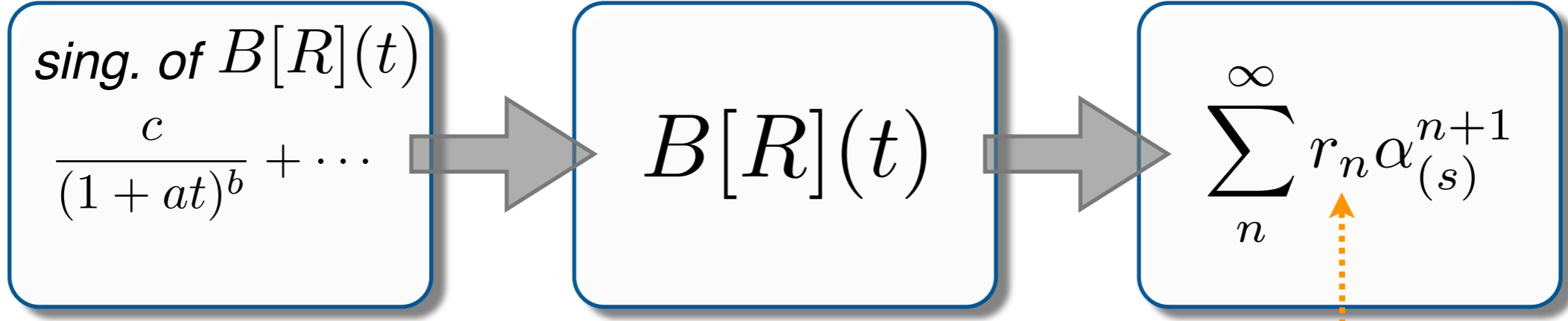
Beneke, DB, Jamin, '13

FOPT vs CIPT

- RG improvement: main theoretical uncertainty. Weight function dependent.



- Strategy: reasonable model for the leading singularities of the Borel transformed perturbative series.

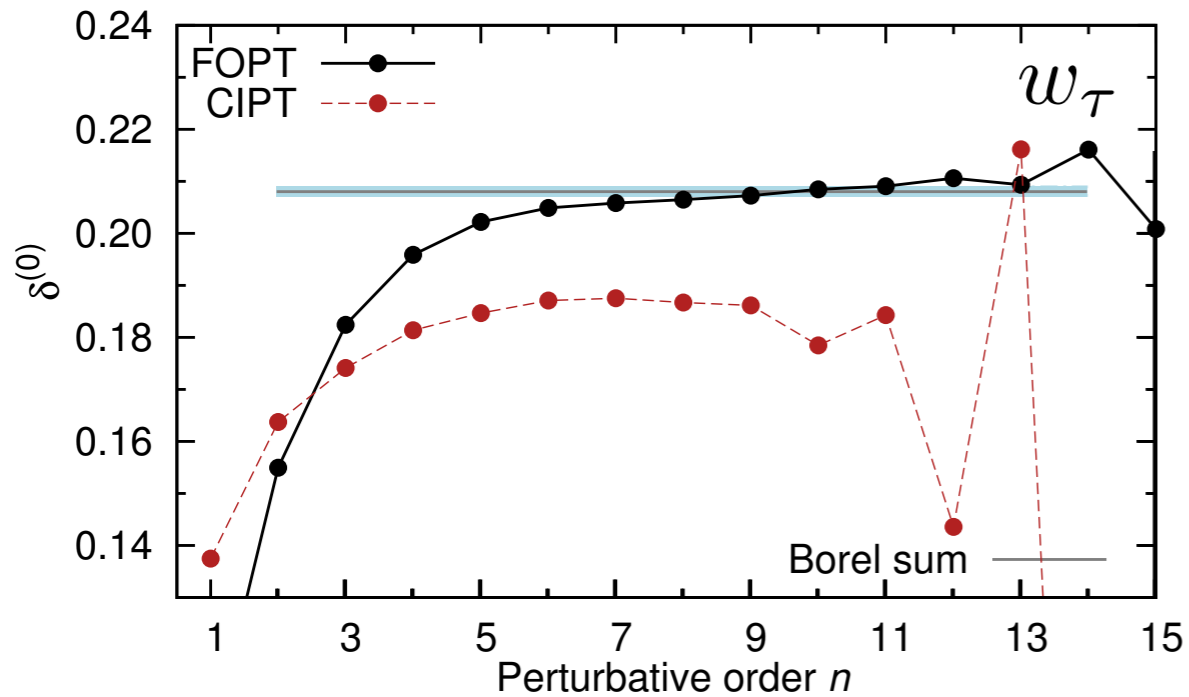


First 4 coefficients are fixed from loop computations

Reference model (RM)

Beneke, Jamin '08

$$B[\hat{D}](u) = B[\hat{D}_1^{UV}](u) + B[\hat{D}_2^{IR}](u) + B[\hat{D}_3^{IR}](u) + d_0^{PO} + d_1^{PO} u$$

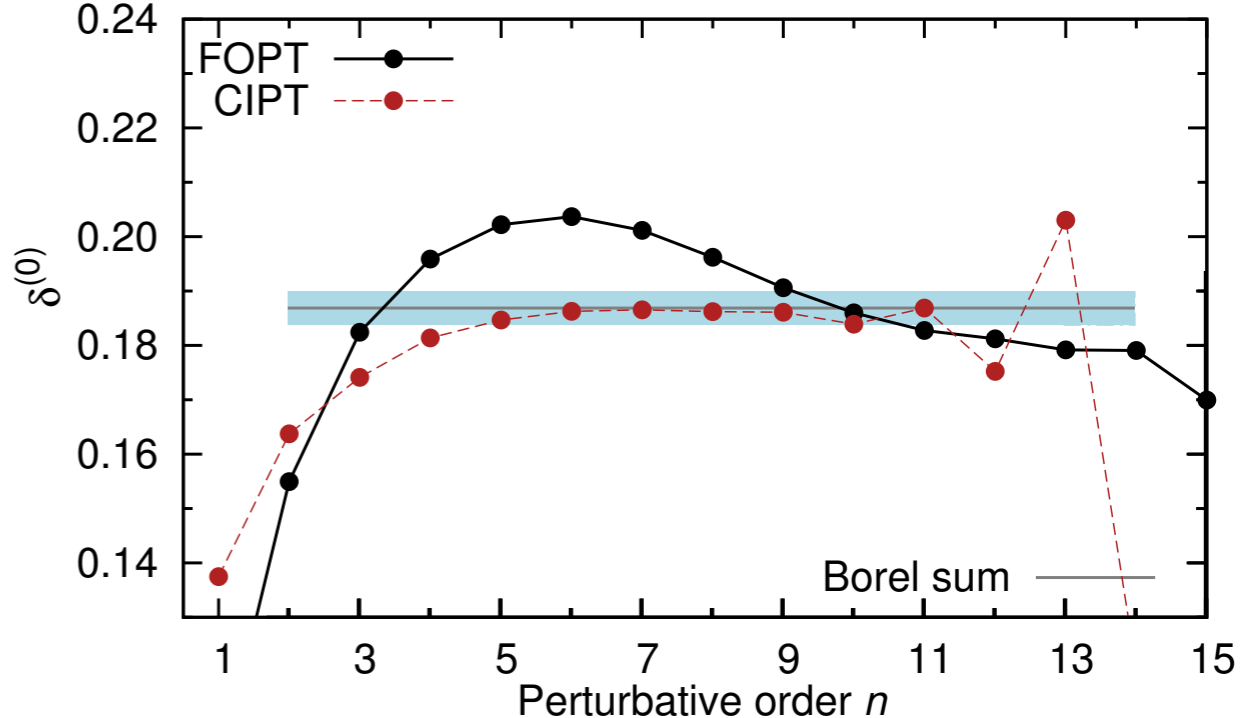


- Model with the leading UV and the first two IR singularities.
- Favors FOPT (related to the presence of $u = 2$ sing).

Alternative Model (AM)

Beneke, Boito, Jamin '13

$$B[\hat{D}](u) = B[\hat{D}_1^{UV}](u) + B[\hat{D}_3^{IR}](u) + B[\hat{D}_4^{IR}](u) + d_0^{PO} + d_1^{PO} u$$

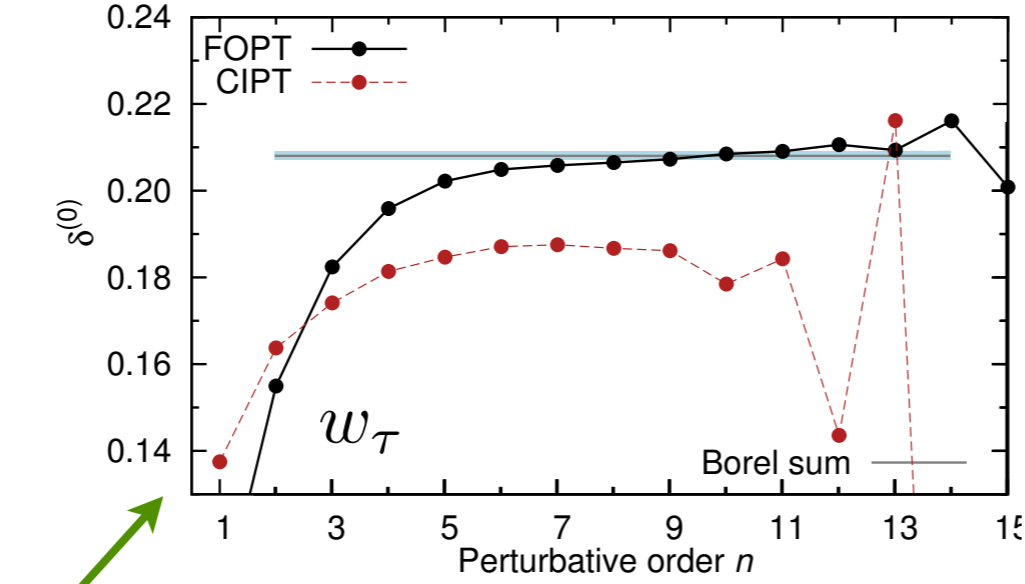


- No IR singularity at $u = 2$.
- Favors CIPT.

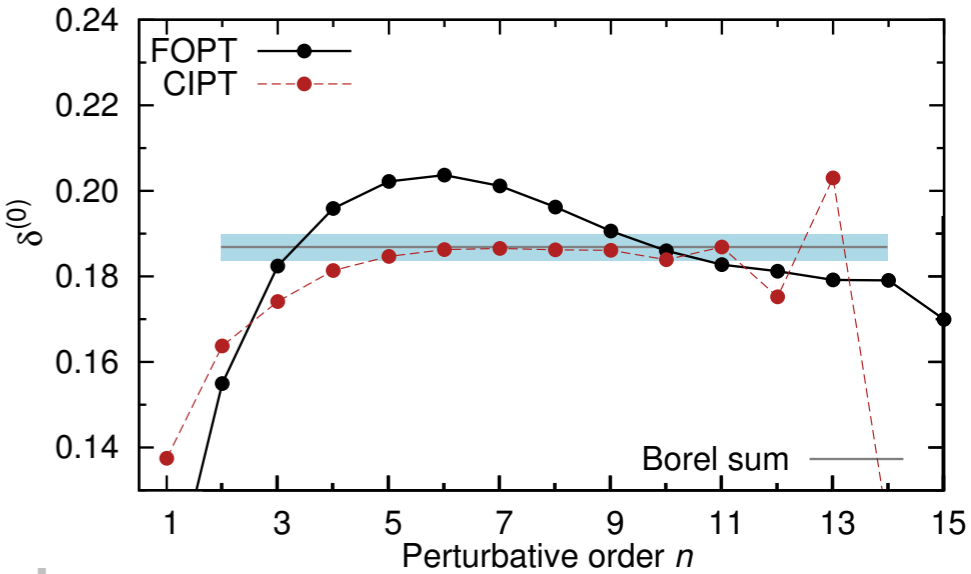
Conclusions depend on assumptions about higher orders

● Systematic study of different weight functions used in literature

k	$w_k(x)$
1	1
2	x
3	x^2
4	x^3
5	x^4
6	$1 - x$
7	$1 - x^2$
8	$1 - x^3$
9	$1 - \frac{3x}{2} + \frac{x^3}{2}$
10	$(1 - x)^2$
11	$(1 - x)^3$
w_τ	$(1 - x)^2(1 + 2x)$
13	$(1 - x)^3(1 + 2x)$
14	$(1 - x)^2x$
15	$(1 - x)^3x(1 + 2x)$
16	$(1 - x)^3x^2(1 + 2x)$
17	$(1 - x)^3x^3(1 + 2x)$

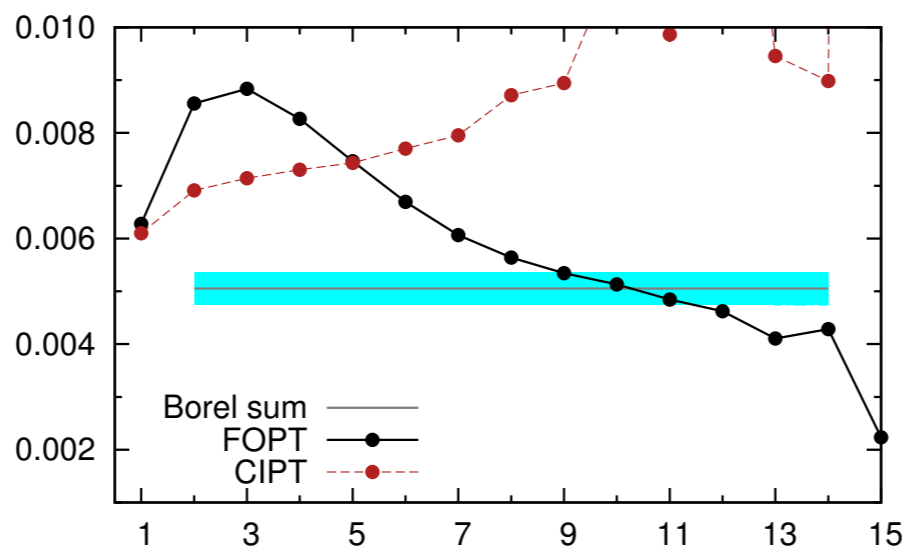


Good pt. behaviour.

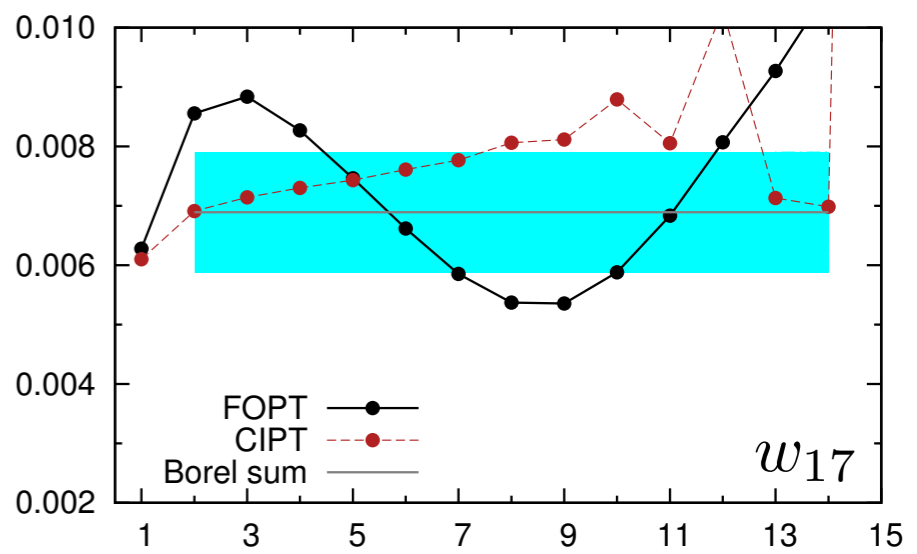


RM

AM



Bad pt. behaviour.



- A decision in favor of FOPT or CIPT depends on the higher order coefficients.
- Within reasonable assumptions **FOPT is the preferred** framework for RGI.
- Some moments are more suitable for the extraction of α_s .
- Some of the recent extractions of α_s employed moments that are not optimal.
- Interesting alternatives to FO-/CIPT using knowledge about Borel plane sing.:
 - Borel sum [Beneke and Jamin '08](#)
 - Conformal mapping [Abbas, Ananthanarayan, Caprini, and Fischer, '13](#)

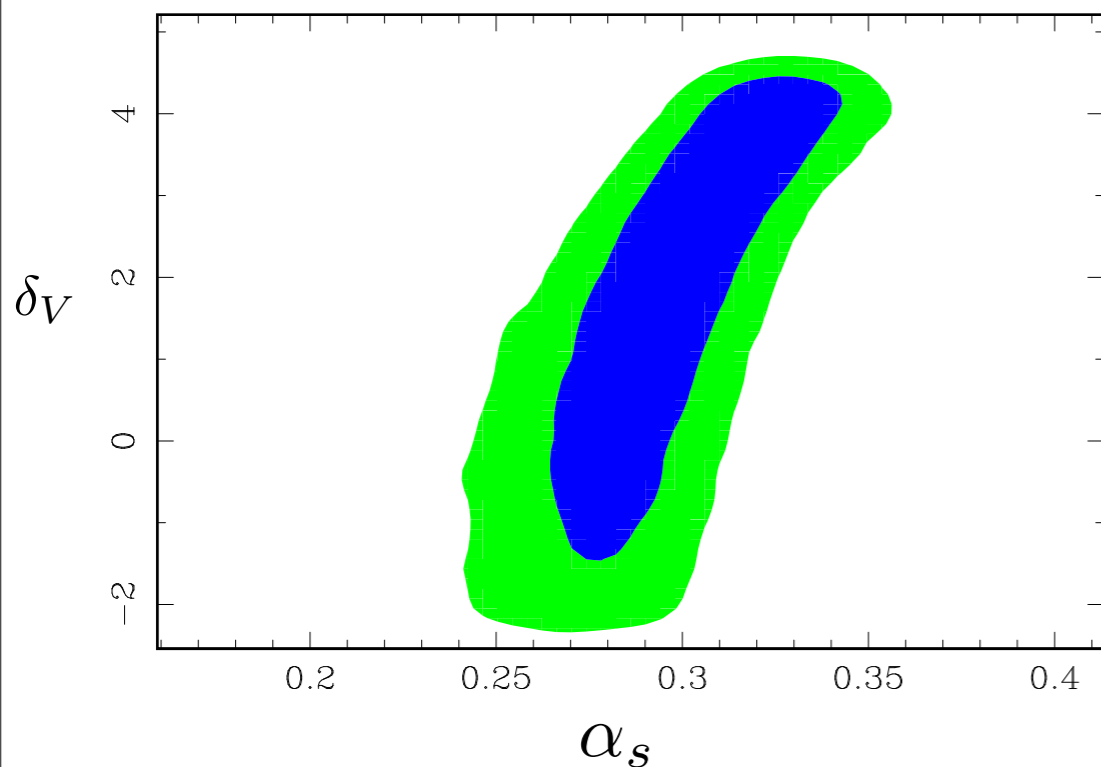
Updated ALEPH data: prospects (preliminary)

- Smaller uncertainties and correct covariances allow for a better determination of α_s and the non-perturbative contributions
- Example: simultaneous determination of DVs and α_s from the V channel.

Model for DVs:
$$\frac{1}{\pi} \text{Im}\Delta(s) = e^{-\delta - \gamma s} \sin(\alpha + \beta s)$$

● OPAL data

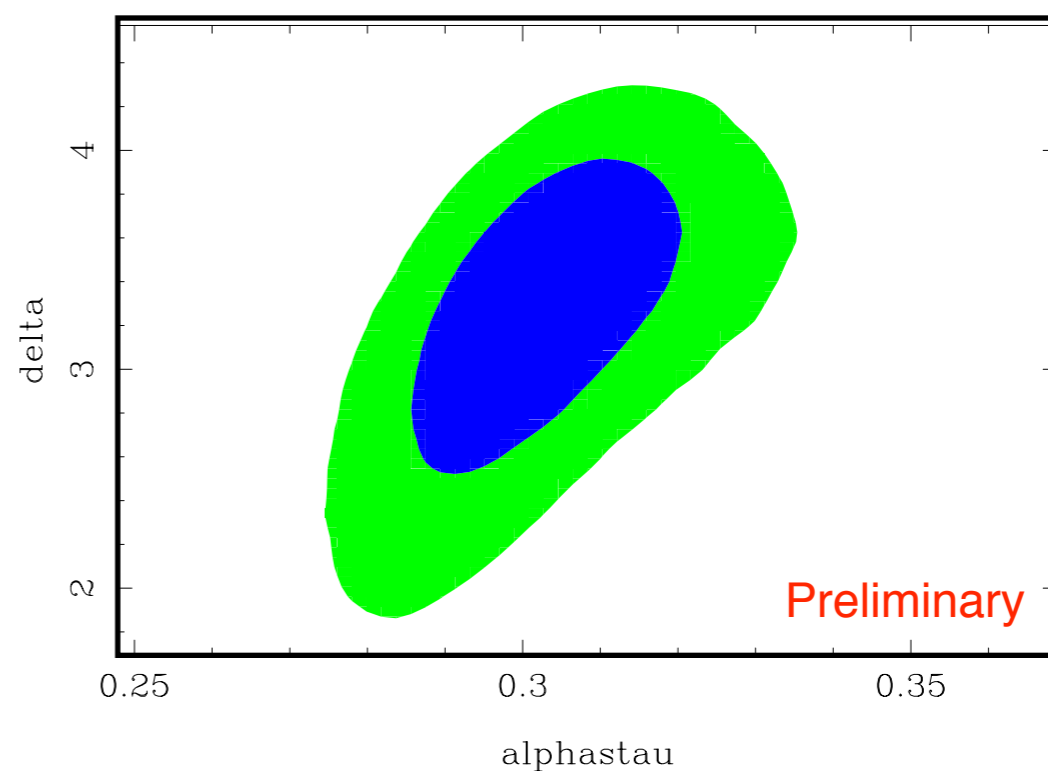
DB, Golterman, Jamin, Mahdavi, Maltman, Osborne, S Peris, '12



$$\sigma_{\alpha_s(m_\tau^2)} \sim 5.5\%$$

● Updated ALEPH data

DB, Golterman, Maltman, Osborne, S Peris, (in preparation)



$$\sigma_{\alpha_s(m_\tau^2)} \sim 3.4\%$$

- Smaller uncertainties and correct covariances allow for a better determination of α_s and the non-perturbative contributions
- Example: simultaneous determination of DVs and α_s from the V channel.

Model for DVs: $\frac{1}{\pi} \text{Im}\Delta(s) = e^{-\delta - \gamma s} \sin(\alpha + \beta s)$

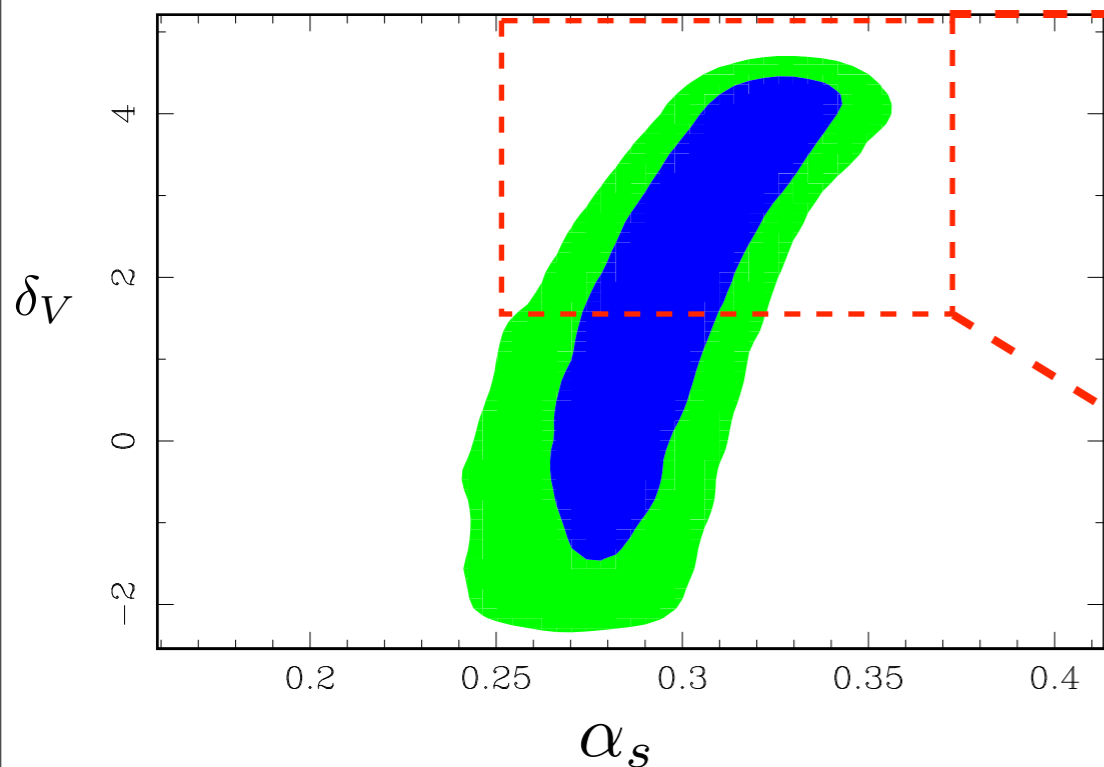
● OPAL data

DB, Golterman, Jamin, Mahdavi, Maltman, Osborne, S Peris, '12

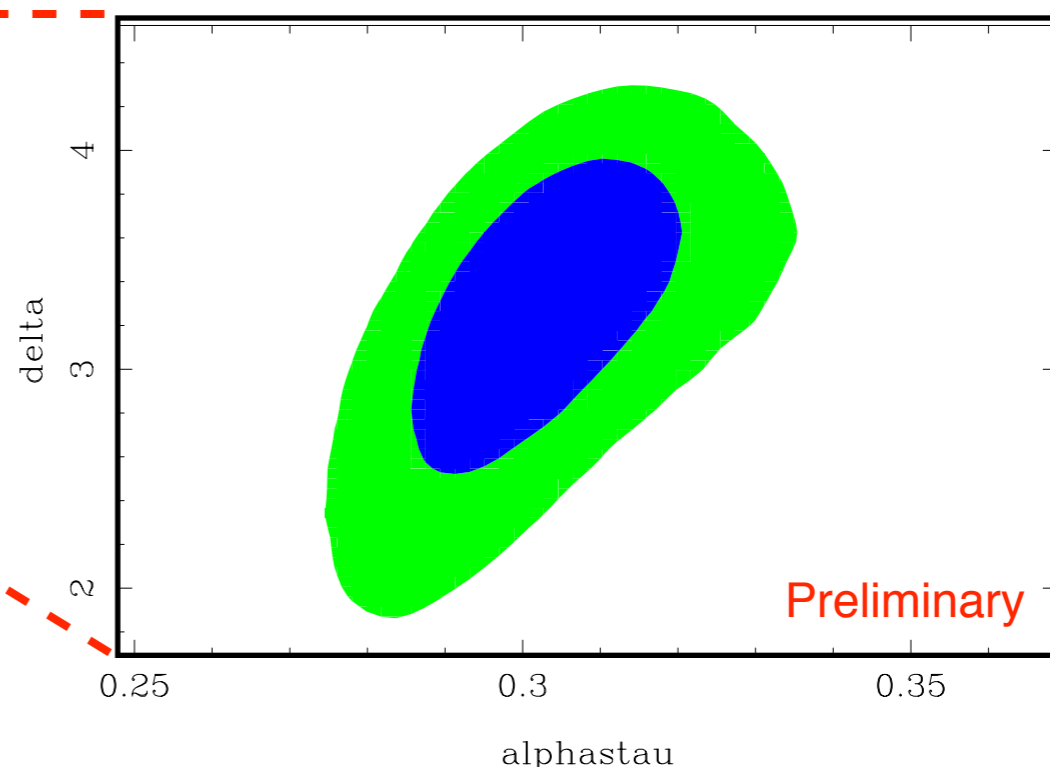
● Updated ALEPH data

DB, Golterman, Maltman, Osborne, S Peris, (in preparation)

Note the different scale



$\sigma_{\alpha_s(m_\tau^2)} \sim 5.5\%$



$\sigma_{\alpha_s(m_\tau^2)} \sim 3.4\%$

Good prospects for an independent α_s analysis from the updated ALEPH data

- New analysis incorporating recent knowledge.
 - Duality Violations
 - Perturbative behaviour of different moments
- Better control of the non-perturbative contribution.
- New results from the α_s analysis this autumn (TAU14/Aachen, September).

Back-up slides

RGImprovement

- Description in terms of the Adler function (derivative of )

$$D_{\text{pert}}^{(1+0)}(s) = \frac{N_c}{12\pi^2} \sum_{n=0}^{\infty} a_{\mu}^n \sum_{k=1}^{n+1} k c_{n,k} \left(\log \frac{-s}{\mu^2} \right)^{k-1} \quad a_{\mu} = \alpha(\mu)/\pi$$

- only $c_{n,1}$ are independent (known up to $c_{4,1}$). $c_{n,k}$ depend on $c_{n,1}$ and β_m .

- Prescriptions for the RG improvement

FOPT
 $\mu = s_0$

$$\delta_{\text{FO},w_i}^{(0)} = \sum_{n=1}^{\infty} a(s_0)^n \sum_{k=1}^n k c_{n,k} J_{k-1}^{\text{FO},w_i}$$

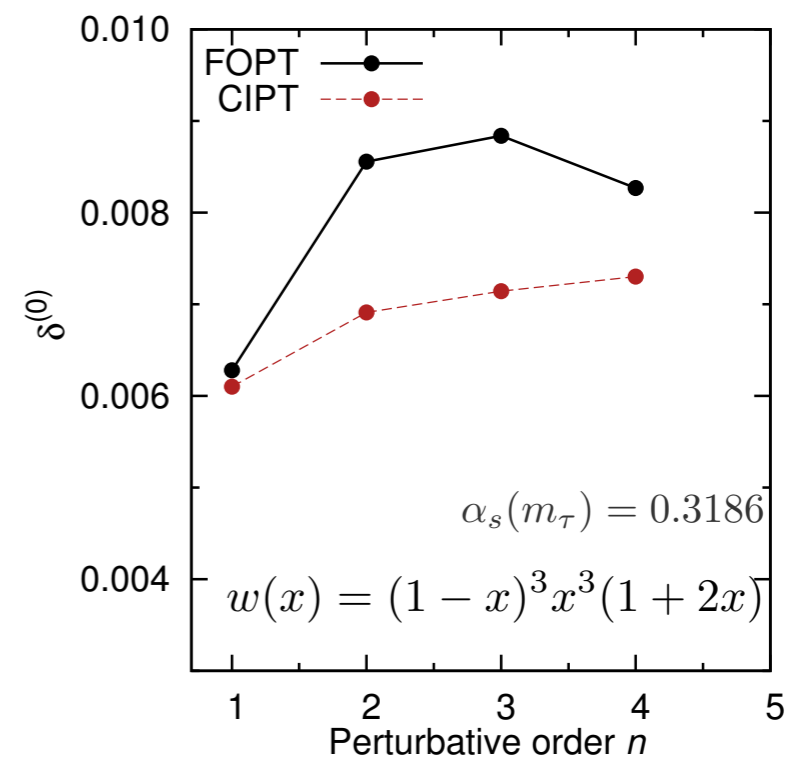
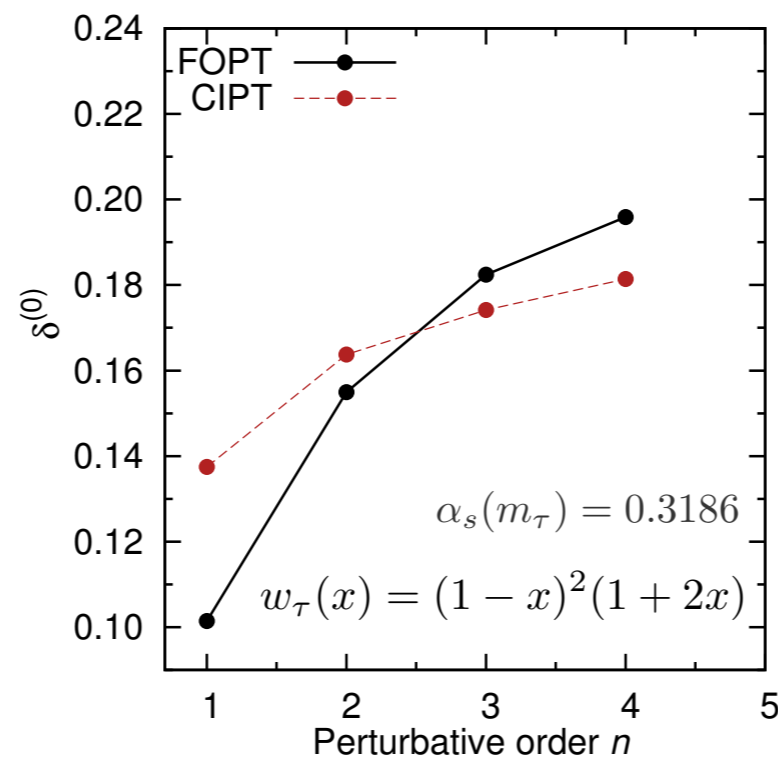
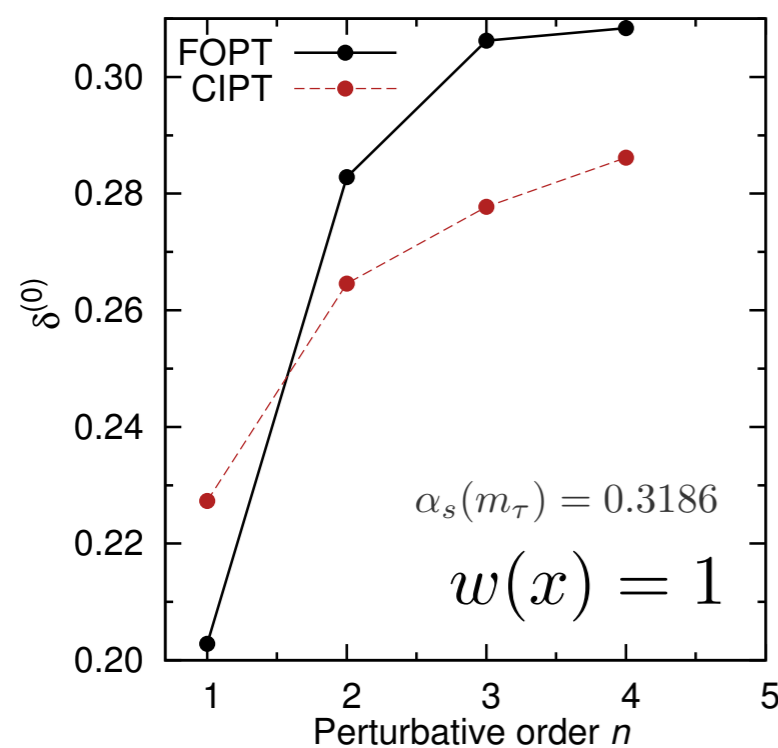
$$J_n^{\text{FO},w_i} \equiv \frac{1}{2\pi i} \oint_{|x|=1} \frac{dx}{x} W_i(x) \log^n(-x)$$

CIPT
 $\mu = -s_0 x$

$$\delta_{\text{CI},w_i}^{(0)} = \sum_{n=1}^{\infty} c_{n,1} J_n^{\text{CI},w_i}(s_0)$$

$$J_n^{\text{CI},w_i}(s_0) \equiv \frac{1}{2\pi i} \oint_{|x|=1} \frac{dx}{x} W_i(x) a^n(-s_0 x)$$

Le Diberder and Pich '92



Borel transform

$$R \sim \sum_n r_n \alpha_{(s)}^{n+1} \quad \text{divergent but (hopefully) **asymptotic**}$$

Dyson 1952

↓ ↓

? *in QFT we only know the expansion*

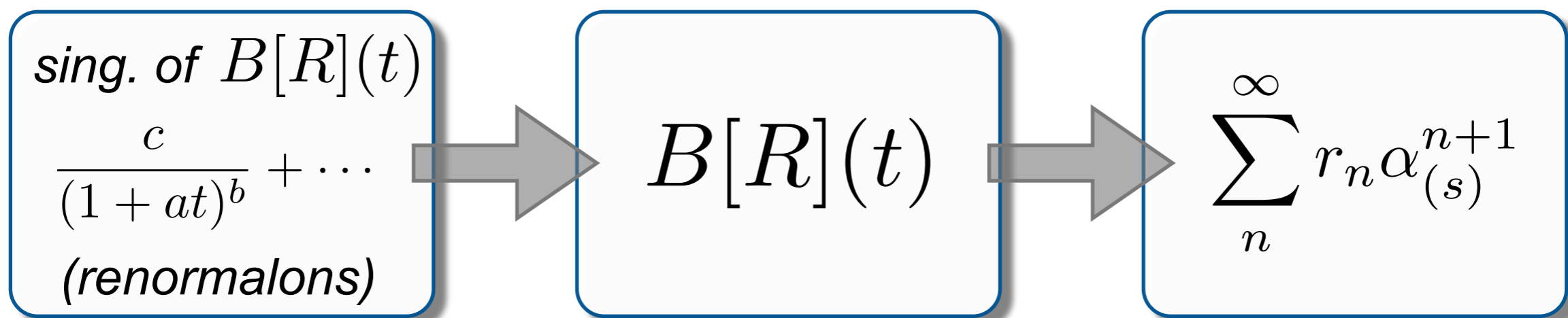
- Define the Borel transformed series

$$B[R](t) \equiv \sum_{n=0}^{\infty} r_n \frac{t^n}{n!} \quad \text{which can be "summed"} \implies \tilde{R} \equiv \int_0^{\infty} dt e^{-t/\alpha} B[R](t)$$

- Divergent behaviour encoded in the singularities of $B[R](t)$

(review) Beneke 1999

Strategy:



$$D_{\text{pert}}^{(1+0)}(Q^2) = \frac{N_c}{12\pi^2} \sum_{n=0}^{\infty} \left(\frac{c_{n,1}}{\pi^n} \right) \alpha_Q^n$$

singularities in the Borel plane

- General structure of large-order behavior (believed to be) known

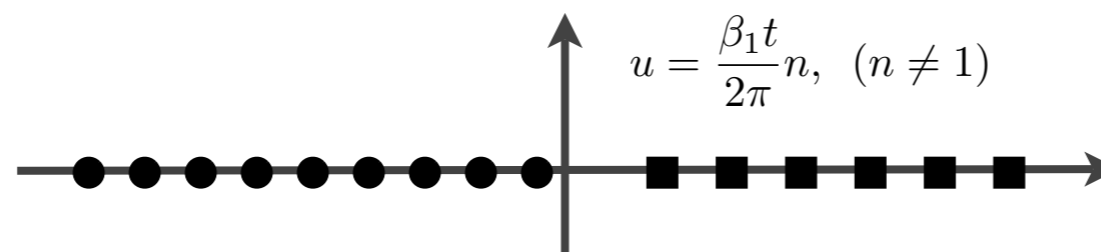
(review) Beneke 1999

Borel transformed Adler function

$$B[\hat{D}](t) \equiv \sum_{n=0}^{\infty} \frac{c_{n,1}}{\pi^n} \frac{t^n}{n!}$$

Borel sum: $\hat{D}(\alpha) \equiv \int_0^{\infty} dt e^{-t/\alpha} B[\hat{D}](t)$

- Singularities in the t plane



UV renormalons

- sign alternating
- leading sing. in the Adler function at $u = -1$
- no-sign alternation in known coeff.: small residue for the leading UV pole

IR renormalons

- fixed sign
- sing. at $u = 2, 3, 4 \dots$ related to dim-4, dim-6, dim-8... contributions
- $u = 2$ related to the gluon condensate

$$B[D_p] = \frac{c_p}{(p-u)^\gamma} \left[1 + \tilde{b}_1(p-u) + \dots \right]$$

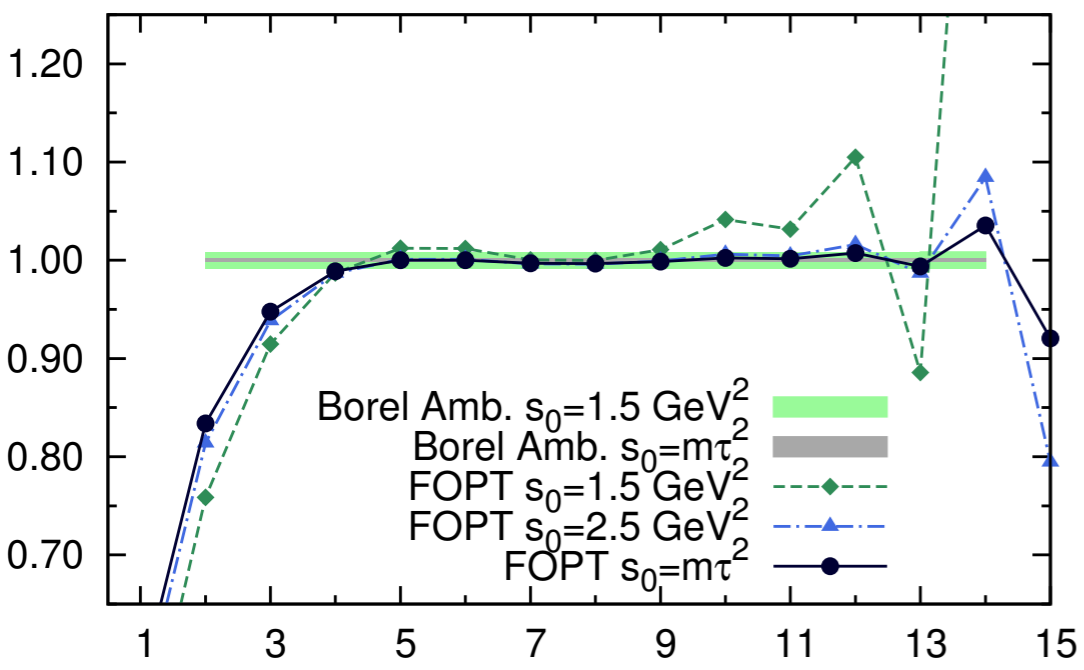
Structure of each singularity in principle calculable (up to)

energy dependence

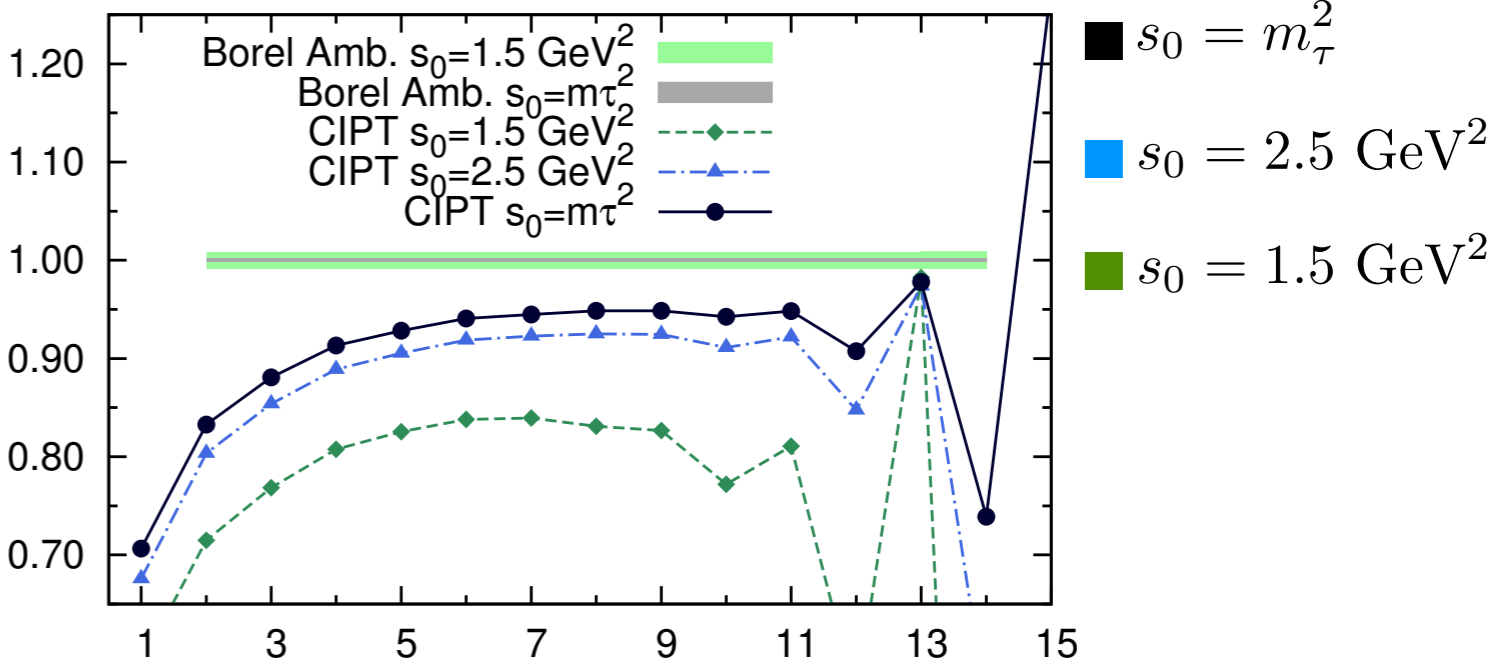
Reference model

$w(x) = 1 - x^2$

FOPT

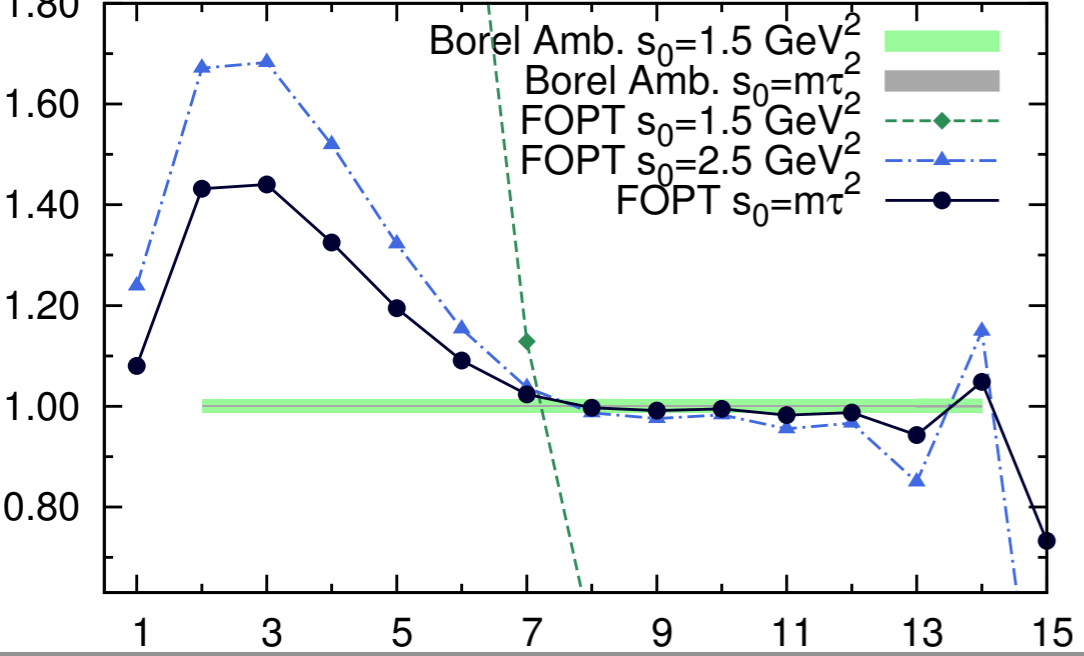


CIPT

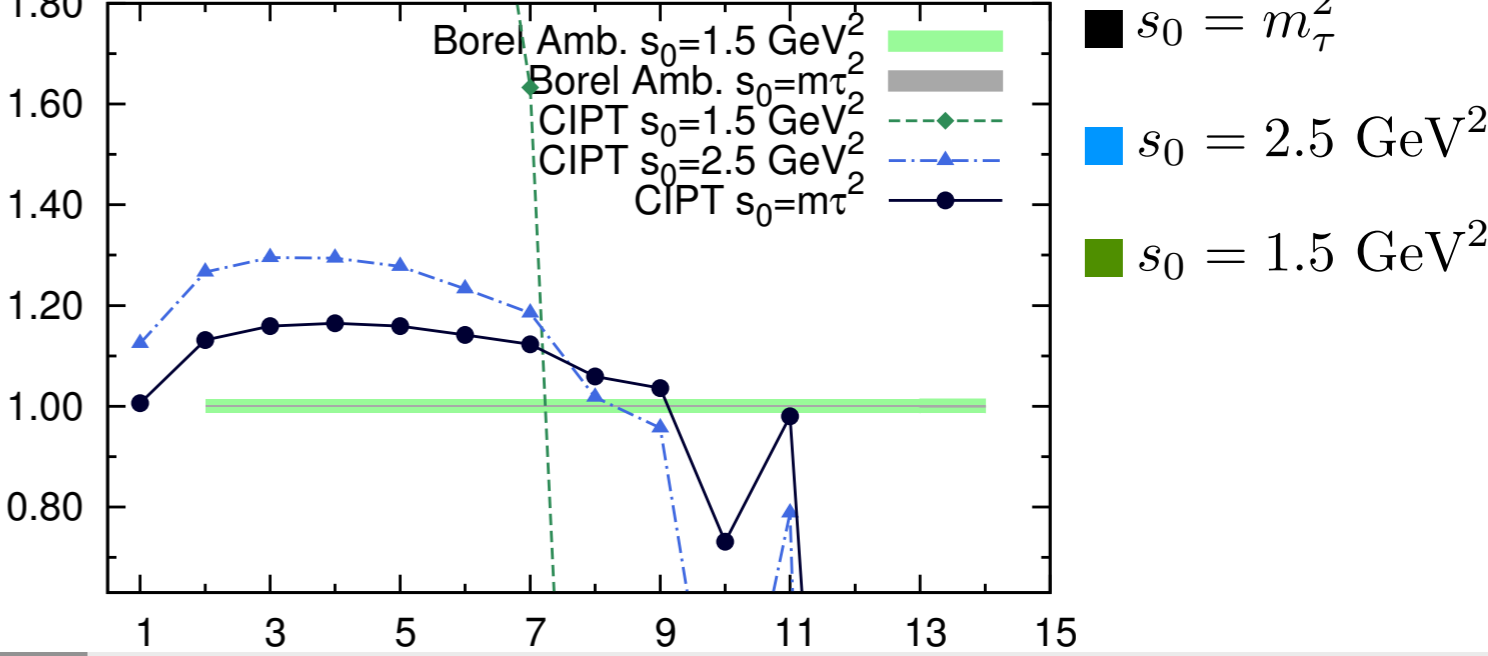


$w(x) = (1 - x)^3 x^3 (1 + 2x)$

FOPT



CIPT



Why FOPT is better in the RM

Reference model

Beneke & Jamin '08

- Separating the contributions in FOPT

$$\delta_{\text{FO},w_i}^{(0)} = \sum_{n=1}^{\infty} \left[c_{n,1} \delta_{w_i}^{\text{tree}} + g_n^{[w_i]} \right] a(s_0)^n$$

$$g_n^{[w_i]} = \sum_{k=2}^n k c_{n,k} J_{k-1}^{\text{FO},w_i}$$

- Result at α_s^n . FOPT sums the first n rows. **Important cancellations.**

w_τ

α_s^n	$c_{1,1}$	$c_{2,1}$	$c_{3,1}$	$c_{4,1}$	$c_{5,1}$	$c_{6,1}$	$c_{7,1}$	$c_{8,1}$	g_n	$\frac{c_n + g_n}{c_n}$
1	1									1
2	g_2 3.56	+ 1.64							3.56	3.17
3	g_3 8.31	+ 11.7	+ 6.37						20.0	4.14
4	g_4 -20.6	+ 30.5	+ 68.1	+ 49.1					78	2.59
⋮	→									⋮
6	g_6 -2924	-2858	-2280	2214	5041	3275			-807	0.754
⋮	→									⋮
8	g_8 14652	-29552	-145846	-502719	-393887	260511	467787	388442	-329054	0.153

Fixed Order

- CIPT sums the first n columns to all orders. Misses the cancellations.