

1. NA48/2 and NA62 Beams and Detectors

NA48/2 [1] (2003 and 2004) and NA62 [2] (2007) share the same detectors but have different beam parameters. The spectrometer magnetic field was modified accordingly.

Both $K^\pm \rightarrow \pi^\pm \gamma \gamma$ data sets were collected with a minimum bias trigger condition:

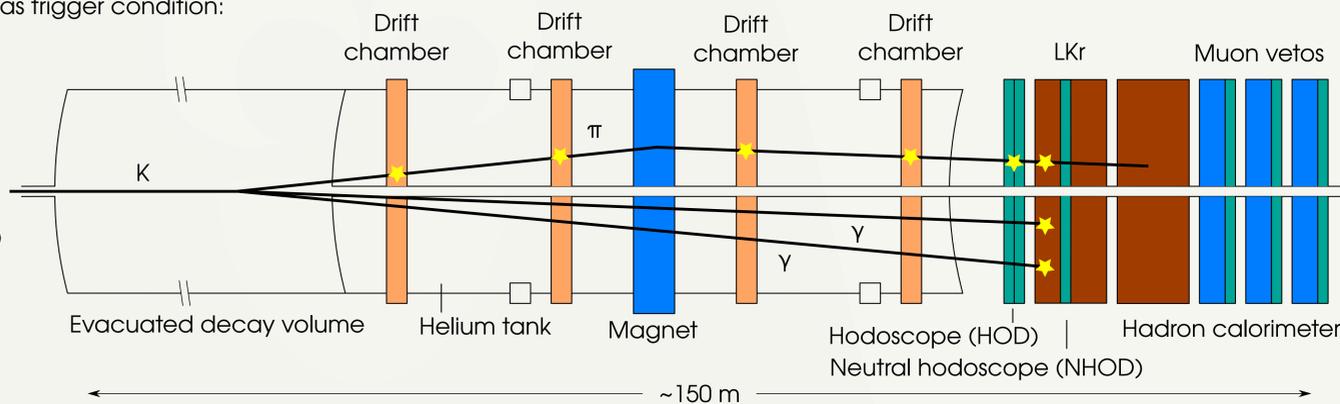
- Time coincidence of signals in the two HOD planes,
- Energy deposit > 10 GeV in LKr.

Unless otherwise stated the plots are from NA62 analysis.

	NA48/2	NA62
Data taking period	12 hours in 2003 54 hours in 2004	4 months in 2007 (Downscaled trigger)
Beam momentum	(60 ± 3) GeV/c	(74 ± 1.4) GeV/c

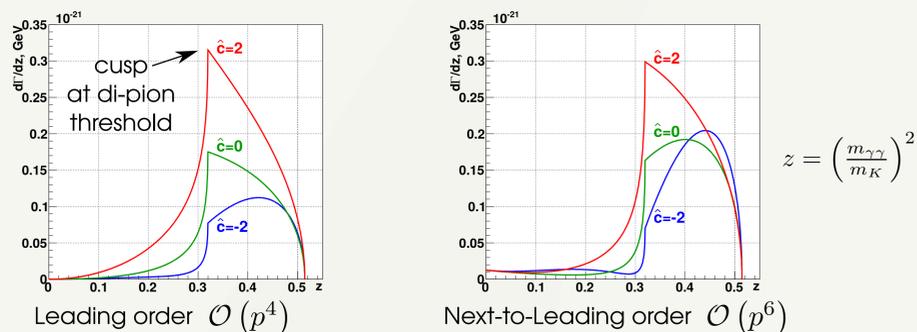
Pion momentum was reconstructed by a magnetic spectrometer while the energy of photons was measured by a liquid krypton calorimeter (LKr).

The detector resolutions are $\frac{\sigma_E}{E} = (3.2/\sqrt{E} \oplus 9/E \oplus 0.42)\%$ for LKr, $\frac{\sigma_p}{p} = (1.02 \oplus 0.044 \cdot p)\%$ for NA48/2 spectrometer and $\frac{\sigma_p}{p} = (0.48 \oplus 0.009 \cdot p)\%$ for NA62 spectrometer. (E): GeV, (p): GeV/c.



2. Chiral Perturbation Theory and $K^\pm \rightarrow \pi^\pm \gamma \gamma$

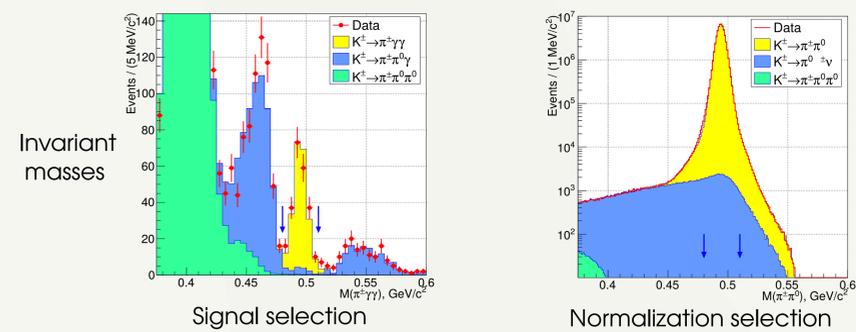
$K^\pm \rightarrow \pi^\pm \gamma \gamma$ is a good probe for testing the ChPT description of low energy weak processes. This channel has been extensively studied by theorists [3].



At leading order the predicted differential decay rate depends exclusively on \hat{c} . At next to leading order \hat{c} enter the decay rate as $\hat{c}^* = \hat{c} - 2(m_\pi/m_K)^2 \eta_1 - 2\eta_2 - 2\eta_3$. In this analysis η_i are fixed to 0.

3. Signal Selection and Normalization

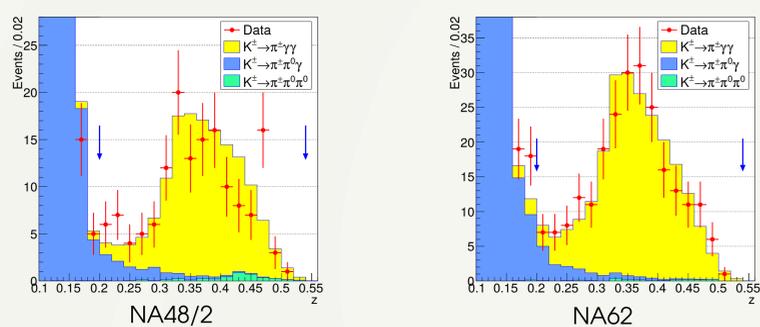
$K^\pm \rightarrow \pi^\pm \gamma \gamma$ and $K^\pm \rightarrow \pi^\pm \pi^0$ are collected with same trigger logic. This allows a first order cancellation of several systematic effects.



Branching ratio is given by $\mathcal{B}(K_{\pi\gamma\gamma}) = \frac{N'_{\pi\gamma\gamma}}{N'_{2\pi}} \cdot \frac{A_{2\pi}}{A_{\pi\gamma\gamma}} \cdot \frac{\epsilon_{2\pi}}{\epsilon_{\pi\gamma\gamma}} \cdot \mathcal{B}(K_{2\pi}) \cdot \mathcal{B}(\pi_{\gamma\gamma}^0)$ where N' are # events after background subtraction, A are acceptances and ϵ are trigger efficiencies. For $K^\pm \rightarrow \pi^\pm \pi^0$, di-photon mass has to be consistent with π^0 .

4. Reconstructed z Spectrum

Simulated signal with the best \hat{c} value in the $O(p^6)$ fit:



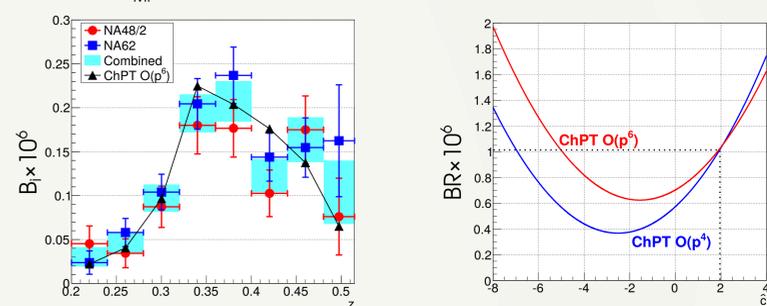
Main background to signal is due to $K^\pm \rightarrow \pi^\pm \pi^0 \gamma$ followed by a $\pi^0 \rightarrow \gamma \gamma$ decay, with two photons producing a merged LKr energy cluster.

5. Branching Ratio

Acceptances and differential decay rate are model dependent but for small width z bins these dependences become negligible. This allows to extract a model independent decay rate B_{MI} by summing over the z bins for $z > 0.2$.

Combined NA48/2 and NA62 model independent branching ratio:

$$B_{MI}(z > 0.2) = (0.965 \pm 0.063) \times 10^{-6}$$



Assuming an $O(p^6)$ description we can compute the total decay rate B_6 for a given \hat{c} .

6. Extraction of the \hat{c} Parameter

$O(p^6)$ differential decay rate can be written as [3]:

$$\frac{\partial \Gamma}{\partial y \partial z}(\hat{c}, y, z) = \frac{m_K}{2^9 \pi^3} \left[z^2 \left(|A(\hat{c}, z, y^2) + B(z)|^2 + |C(z)|^2 \right) + \left(y^2 - \frac{1}{4} \lambda(1, r_\pi^2, z) \right)^2 |B(z)|^2 \right]$$

where $\lambda(a, b, c) = a^2 + b^2 + c^2 - 2(ab + ac + bc)$. A and B are loop amplitudes and C pole amplitude.

Values of all the external parameters used in the NA48/2 and NA62 analyses are detailed in [1] and [2].

\hat{c} parameter is extracted by fitting the reconstructed z spectrum, maximizing the log-likelihood $\ln \mathcal{L} = \sum_{i=1}^N [n_i \ln m_i - m_i - \ln(n_i!)]$. The sum is running over N bins of the reconstructed z spectrum, m_i is the # of expected events for a given \hat{c} and n_i is the # of observed events.

8. References

- [1] Batley, J.R. et al. (NA48/2 collab.), Phys. Lett. B 730 (2014) 141
- [2] Lazzeroni, C. et al. (NA62 collab.), Phys. Lett. B 732 (2014) 65
- [3] D'Ambrosio, G. and Portolés, J. Phys. Lett. B 386 (1996) 403
- [4] Kitching, P. et al., Phys. Rev. Lett. 79 (1997) 4079

7. Take Home Message

Current experimental state of the art is the following:

	BNL E787 [4]	NA48/2	NA62
Candidates	31	149	232
Background Events	5.1 ± 3.3	15.5 ± 0.7	17.4 ± 1.1
$B_{MI}(z > 0.2) \times 10^6$	$0.6 \pm 0.15 \pm 0.07^*$	$0.877 \pm 0.087 \pm 0.017$	$1.088 \pm 0.093 \pm 0.027$
\hat{c}_4	1.6 ± 0.6	$1.37 \pm 0.33 \pm 0.14$	$1.93 \pm 0.26 \pm 0.08$
\hat{c}_6	1.8 ± 0.6	$1.41 \pm 0.38 \pm 0.11$	$2.10 \pm 0.28 \pm 0.18$
$B_6 \times 10^6$	$1.1 \pm 0.3 \pm 0.1$	$0.910 \pm 0.072 \pm 0.022$	$1.058 \pm 0.066 \pm 0.044$

* $0.157 < z < 0.384$

Combining the NA48/2 and NA62 results:
 $B_{MI}(z > 0.2) = (0.965 \pm 0.063) \times 10^{-6}$
 $\hat{c}_4 = 1.72 \pm 0.21$
 $\hat{c}_6 = 1.86 \pm 0.25$
 $B_6 = (1.003 \pm 0.056) \times 10^{-6}$

NA48/2 and NA62 results are in good agreement with the earlier E787 results and considerably improve our experimental knowledge of $K^\pm \rightarrow \pi^\pm \gamma \gamma$. We observed the predicted cusp at the di-pion threshold ($z = 0.320$).