

NLO corrections to vector-boson scattering at the LHC

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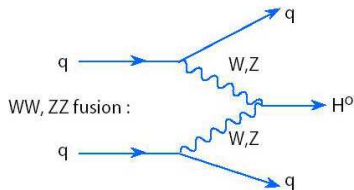
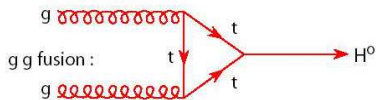
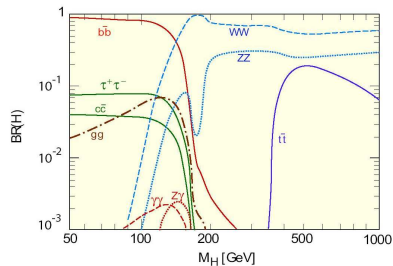
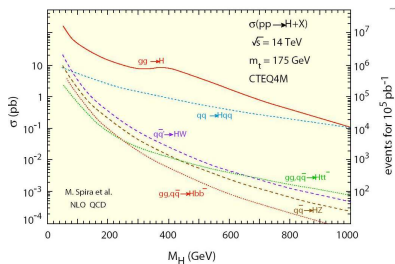
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Tools and Precision Calculations for Physics Discoveries at Colliders

- 1 Motivation
 - Search for a light Higgs boson
 - Study of Electroweak Symmetry Breaking
- 2 What kinds of diagrams are involved?
- 3 What has been done so far?
 - Tree-level
 - NLO QCD corrections
- 4 Our project
 - Many subprocesses, many diagrams
 - Block Structure
 - Elements of calculation
 - Monte Carlo and cuts
- 5 Summary

SM Higgs Search

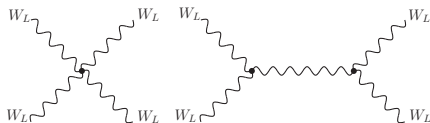




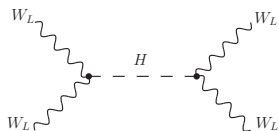
- vector boson fusion ($qq \rightarrow qqH$)
 - possible discovery mode for a light SM Higgs boson
 - second largest cross-section for the light Higgs
 - statistically significant thanks to forward jet tagging and suppressed QCD in the central region
 - relatively low luminosity needed for discovery
- vector boson scattering ($qq \rightarrow qqVV$) - background to $H \rightarrow VV$ decay mode via VBF
- VBF can be used to determine Higgs properties and constraint its couplings to gauge bosons

What if there is no light Higgs?

- the scattering of longitudinal W's grows with energy and violates unitarity



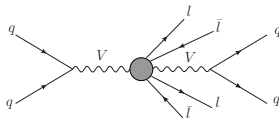
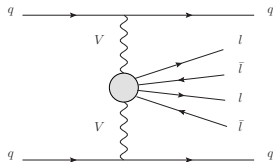
$$\propto g_W^2 \frac{E^2}{M_W^2}$$



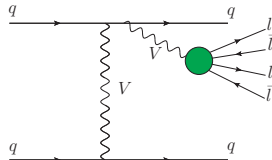
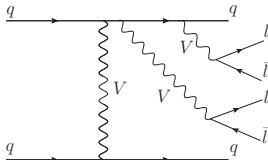
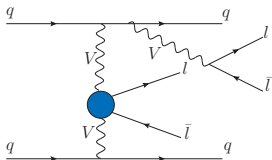
$$\propto -g_W^2 \frac{E^2}{M_W^2}$$

- without Higgs, new mechanism of EW symmetry breaking must be considered
- new physics (composite Higgs, extra dimensions,...) predicts new resonances and modify VBF
- $qq \rightarrow qqWW$ very sensitive channel to new interaction
- minimizes the background from transversely polarized WW

Diagrams - Leading Order



- following types of diagrams have to be included (to preserve gauge invariance)

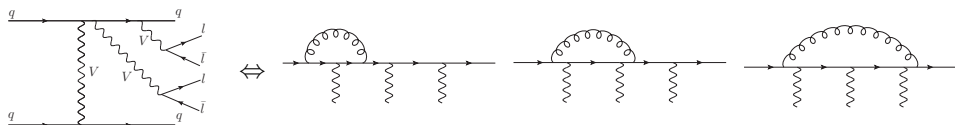
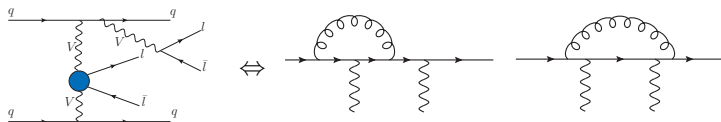
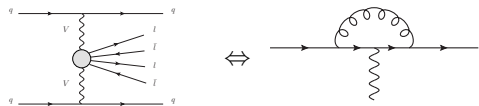


Tree-level studies

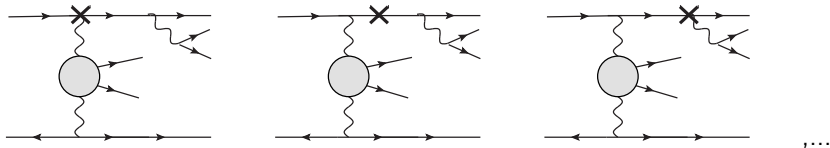
- first partial results [Cahn, Dawson \(1984\)](#)
- $pp \rightarrow qqWW$ in effective gauge boson approximation, only for longitudinal polarization [Duncan, Kane, Repko \(1986\)](#)
- exact calculation of $pp \rightarrow qqWW$, all polarizations [Dicus, Vega \(1986\)](#)
- $pp \rightarrow qqZZ$, effective gauge boson approximation [Abbasabadi, Repko \(1988\)](#)
- $pp \rightarrow qqZZ \rightarrow qqllll$, narrow width approximation [Baur, Glover \(1990\)](#)
- $pp \rightarrow (qqZW \rightarrow qqZW) + X$, effective gauge boson approximation, longitudinal polarization [Dobado, Herrero, Terron \(1991\)](#)
- $pp \rightarrow qqZW$, full tree-level, leptonic decay correlations [Barger, Cheung, Han, Stange, Zeppenfeld \(1992\)](#)
- $pp \rightarrow qqWW$ - electroweak chiral lagrangian formalism, semileptonic decay [Butterworth, Cox, Forshaw \(2002\)](#)
- $pp \rightarrow qqllll$ - complete parton level analysis, SM and SILH [Ballestrero, Accomando, Bevilacqua, Franzosi, Maina \(2006-2010\)](#)
- multiple BSM studies for the LHC [Han, Krohn, Wang, Zhu \(2009\)](#), [Cheung, Chiang, Yuan \(2008\)](#)...

Diagrams - NLO QCD Contributions

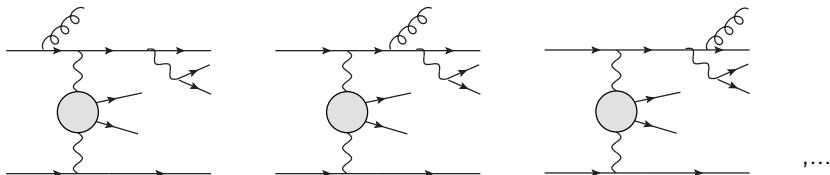
Virtual corrections



Counterterms

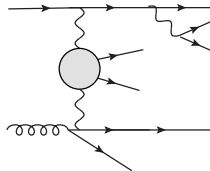
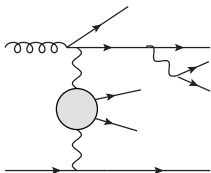


Real corrections

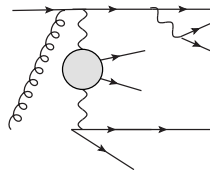
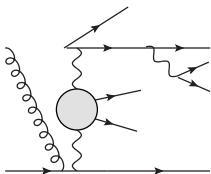


Initial state gluon

- spacelike gauge bosons

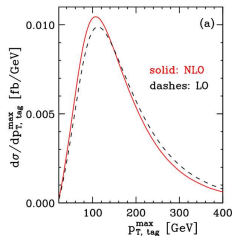
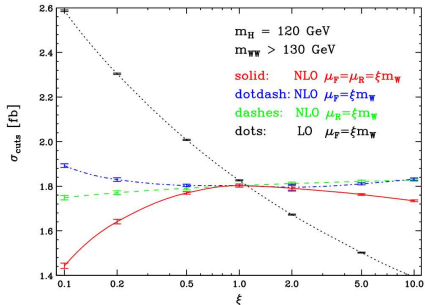


- timelike gauge bosons



NLO QCD calculation

- full tree-level calculation and NLO QCD corrections (real and virtual contributions)
 - 2006 - Jäger, Oleari, Zeppenfeld: $qq \rightarrow jjW^+W^- \rightarrow jjllll$
 - 2006 - Jäger, Oleari, Zeppenfeld: $qq \rightarrow jjZZ \rightarrow jjllll$
 - 2007 - Bozzi, Jäger, Oleari, Zeppenfeld: $qq \rightarrow jjWZ \rightarrow jjllll$
 - 2009 - Jäger, Oleari, Zeppenfeld: $qq \rightarrow jjW^\pm W^\pm \rightarrow jjllll$



Organizing the Subprocesses

$qq \rightarrow qq W^+ W^+$	$\left\{ \begin{array}{l} uc \rightarrow ds e^+ \nu_e \mu^+ \nu_\mu \\ us \rightarrow dc e^- \bar{\nu}_e \mu^- \bar{\nu}_\mu \\ uc \rightarrow dc e^+ \nu_e \mu^+ \mu^- \\ us \rightarrow ds e^+ \nu_e \mu^+ \mu^- \\ ds \rightarrow ds e^+ \nu_e \mu^- \bar{\nu}_\mu \\ uc \rightarrow uc e^+ \nu_e \mu^- \bar{\nu}_\mu \\ us \rightarrow dc e^+ \nu_e \mu^- \bar{\nu}_\mu \\ us \rightarrow us e^+ \nu_e \mu^- \bar{\nu}_\mu \end{array} \right.$	$qq \rightarrow qq ZZ \left\{ \begin{array}{l} ds \rightarrow ds e^+ e^- \mu^+ \mu^- \\ uc \rightarrow uc e^+ e^- \mu^+ \mu^- \\ us \rightarrow dc e^+ e^- \mu^+ \mu^- \\ us \rightarrow us e^+ e^- \mu^+ \mu^- \\ ds \rightarrow ds e^+ e^- \nu_\mu \bar{\nu}_\mu \\ uc \rightarrow uc e^+ e^- \nu_\mu \bar{\nu}_\mu \\ us \rightarrow dc e^+ e^- \nu_\mu \bar{\nu}_\mu \\ us \rightarrow us e^+ e^- \nu_\mu \bar{\nu}_\mu \end{array} \right.$
$qq \rightarrow qq W^- W^-$		
$qq \rightarrow qq W^+ Z$		
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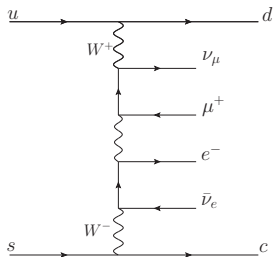
- one subprocess involves between 300 and 2000 Feynman diagrams
- u-channel and s-channel obtained via crossing
- interferences between channels at LO

Example - $uc \rightarrow ds W^+ W^+ \rightarrow ds e^+ \nu_e \mu^+ \nu_\mu$

$QQ \rightarrow QQe^\pm \nu_e \mu^\pm \nu_\mu$	$\bar{Q}Q \rightarrow \bar{Q}Qe^\pm \nu_e \mu^\pm \nu_\mu$
$uc \rightarrow dse^+ \nu_e \mu^+ \nu_\mu \quad t$	$\bar{d}c \rightarrow \bar{u}se^+ \nu_e \mu^+ \nu_\mu \quad t$
$cu \rightarrow dse^+ \nu_e \mu^+ \nu_\mu \quad u$	$\bar{c}\bar{d} \rightarrow \bar{u}se^+ \nu_e \mu^+ \nu_\mu \quad u$
$uu \rightarrow dde^+ \nu_e \mu^+ \nu_\mu \quad \frac{t+u}{2}$	$\bar{s}u \rightarrow \bar{c}de^+ \nu_e \mu^+ \nu_\mu \quad t$
$cs \rightarrow sse^+ \nu_e \mu^+ \nu_\mu \quad \frac{t+u}{2}$	$u\bar{s} \rightarrow \bar{c}de^+ \nu_e \mu^+ \nu_\mu \quad u$
$\bar{Q}\bar{Q} \rightarrow \bar{Q}\bar{Q}e^\pm \nu_e \mu^\pm \nu_\mu$	$\bar{d}u \rightarrow \bar{u}de^+ \nu_e \mu^+ \nu_\mu \quad t+s$
$\bar{d}\bar{s} \rightarrow \bar{u}\bar{c}e^+ \nu_e \mu^+ \nu_\mu \quad t$	$u\bar{d} \rightarrow \bar{u}de^+ \nu_e \mu^+ \nu_\mu \quad t+s$
$\bar{s}\bar{d} \rightarrow \bar{u}\bar{c}e^+ \nu_e \mu^+ \nu_\mu \quad u$	$\bar{s}c \rightarrow \bar{c}se^+ \nu_e \mu^+ \nu_\mu \quad t+s$
$\bar{d}\bar{d} \rightarrow \bar{u}\bar{u}e^+ \nu_e \mu^+ \nu_\mu \quad \frac{t+u}{2}$	$c\bar{s} \rightarrow \bar{c}se^+ \nu_e \mu^+ \nu_\mu \quad t+s$
$\bar{c}\bar{c} \rightarrow \bar{s}\bar{s}e^+ \nu_e \mu^+ \nu_\mu \quad \frac{t+u}{2}$	$\bar{d}u \rightarrow \bar{c}se^+ \nu_e \mu^+ \nu_\mu \quad s$
	$u\bar{d} \rightarrow \bar{c}se^+ \nu_e \mu^+ \nu_\mu \quad s$

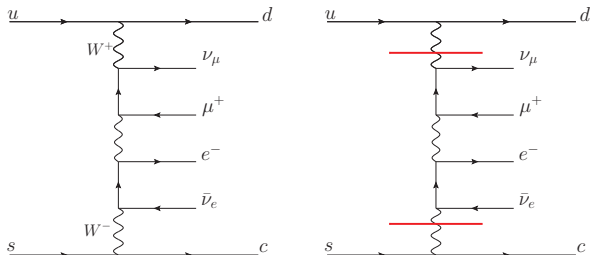
Block structure

- EW and QCD parts are completely independent and can be evaluated separately and reused
- introducing so called "leptonic tensors"
- separating QCD and EW blocks
 - simplifies calculation
 - speeds up Monte Carlo simulations



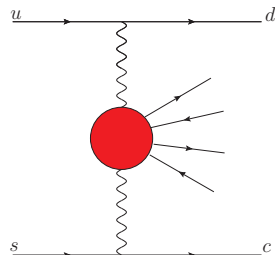
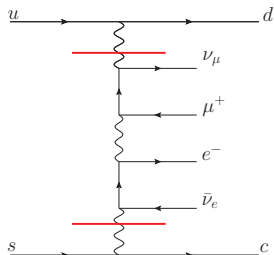
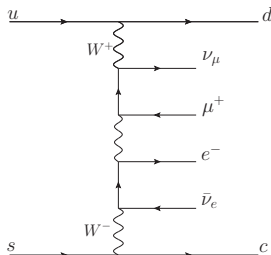
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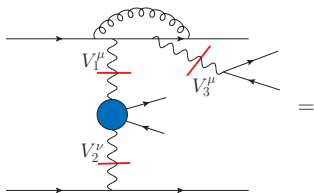


How it works - polarization sums

- relatively small number of building blocks required to construct large number of diagrams

$$\mathcal{M} = \mathcal{M}_{QCD\mu} \mathcal{A}^\mu = \mathcal{M}_{QCD\mu} g^{\mu\nu} \mathcal{A}_\nu \quad \text{and} \quad g_{\mu\nu} = - \sum_i \varepsilon(k)_{i\mu} \varepsilon(k)_{i\nu} + \frac{k_\mu k_\nu}{k^2}$$

$$\begin{aligned} \mathcal{M} &= -(\mathcal{M}_{QCD} \cdot \varepsilon_+)(\mathcal{A} \cdot \varepsilon_+) - (\mathcal{M}_{QCD} \cdot \varepsilon_-)(\mathcal{A} \cdot \varepsilon_-) \\ \Rightarrow & -(\mathcal{M}_{QCD} \cdot \varepsilon_0)(\mathcal{A} \cdot \varepsilon_0) + \frac{1}{k^2} (\mathcal{M}_{QCD} \cdot k)(\mathcal{A} \cdot k) \end{aligned}$$

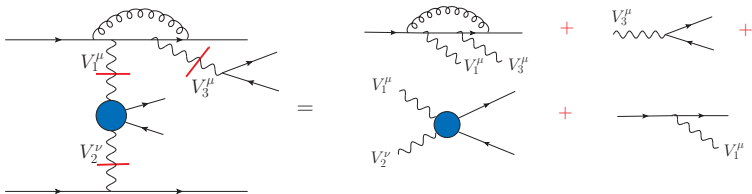


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Elements of Calculation

- diagrams generated with FeynArts
 - EW and QCD blocks are generated independently
- analytical expressions generated with FormCalc and modified in Mathematica and exported to Fortran
- Weyl-van-der-Waerden formalism - translates all kinematic objects into two-component WvdW spinors in chiral representation

$$\Psi = \begin{pmatrix} \phi_A \\ \psi^{\dot{A}} \end{pmatrix} \quad \psi_A \phi^A = (\psi\phi) \quad \psi_{\dot{A}} \phi^{\dot{A}} = \langle \psi\phi \rangle \quad 2k_\mu p^\mu = (k p) \langle k p \rangle$$

- s- and u-channel obtained via crossing which amounts to sign reversal of certain spinors

Dipole subtraction method

$$\sigma = \int_m d\sigma^B + \int_{m+1} d\sigma^R + \int_m d\sigma^V$$

- $d\sigma^B$: Born cross-section - no singularities
- $d\sigma^R$: real cross-section - collinear and soft singularities
- $d\sigma^V$: virtual cross-section - IR singularities
- KLN-theorem: singularities in σ^R and σ^V cancel each other
- analytical integration not possible
- dipole subtraction

$$\sigma^{NLO} = \int_{m+1} (d\sigma^R - d\sigma^A) + \int_m (d\sigma^V + \int_1 d\sigma^A)$$

- $d\sigma^A$: subtraction term containing all IR singularities and analytically integrable over 1-particle phase-space causing collinear and soft singularities

$$\int_{m+1} d\sigma^A = \sum_{\text{dipoles}} \int_m d\sigma^{B'} \otimes \int_1 dV_{\text{dipole}}$$

- $d\sigma^{B'}$: colour correlated Born matrix elements
- dV_{dipole} : process independent dipole terms implemented in the MC program
- $d\sigma^B$, $d\sigma^R$, $d\sigma^V$, $d\sigma^{B'}$ calculated using helicity formalism
- external software (Coli, LoopTools) used to perform tensor reduction
- UV singularities are dealt with by generating counterterm blocks
- IR, soft and collinear singularities
 - regularized in dimensional regularization (LoopTools, Coli) or mass regularization scheme (Coli)
 - pole structure of the virtual blocks

$$\mathcal{M}_V = \mathcal{M}_B \frac{\alpha_s(\mu_R)}{3\pi} \left(\frac{\mu_R^2}{Q^2} \right)^\epsilon \left(-\frac{2}{\epsilon^2} - \frac{3}{\epsilon} \right) + \text{const.} + o(\epsilon)$$

Monte Carlo and cuts

- custom-made multi-channel Monte Carlo is being developed
- choosing proper cuts is essential for distinguishing VBF from the background
- 'typical' VBF cuts include
 - tagging jets - two hard reconstructed jets with $p_T \geq 20$ GeV and large rapidity separation $\Delta y_{jj} > 4$ and invariant mass $M_{jj} > 600$ GeV
 - parallel to the beam (within 1°) - $|\eta_j| < 4.9$
 - separation of jets and leptons $\Delta R_{ll} > 0.2$, $\Delta R_{jl} > 0.4$
 - jets in opposite hemispheres $y_{j1} \times y_{j2} < 0$
 - cuts on invariant masses of the leptons

Progress so far...

- leading order MEs ($qq \rightarrow jj4l$) incl. t-, s- and u-channel and their interferences - completed and compared with MadGraph and FormCalc results
- total LO cross section comparison with existing results (Zeppenfeld, Ballestrero) with different sets of cuts
- LO analysis with and without approximations (s-channel, interferences, mass scheme) in progress
- next-to-leading order - implemented, interfacing with MC and comparison with existing results (Zeppenfeld) in progress
- to be done - gluon exchange at LO

Summary

- vector boson scattering might turn out to be principal for
 - search for a light Higgs at the LHC
 - probing for new strong effects in case of no Higgs is found
- full NLO corrections are fairly complicated (many legs, large number of diagrams) and require modular approach
- despite long history of studies NLO calculations started to emerge only recently (and only QCD corrections)
- our project - first independent verification of Zeppenfeld's calculation
- possibility of incorporating EW corrections and BSM physics in the future